Numerical Methods in Steady State 1D and 2D Heat Equations

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1 Introduction

The steady-state heat equation with a constant coefficient in two dimensions is given by:

 $-k\nabla^2 u(x,y) = q(x,y)$

where k is the thermal conductivity, u is the material temperature, and q is a heat source term. But in the code, I use f instead of q, it is just a notation.

In this Documentation, we first list some assumptions that help us to simplify and reformulate the 1D and 2D problems. Second, we derive the 2nd and 4th order finite difference expression using node-based central difference. Third, we transform the PDE into linear systems with certain matrices by flattening. Then, we can use Jacobi and Gauss-Seidel iterative methods to solve PDE numerically by solving linear systems.

2 Assumptions

• Dirichlet Boundary Condition: The solution is known on the boundary, i.e.

$$u(x)|_{x=a,b} = u_0(x)$$

$$u(x,y)|_{\partial\Omega} = u_0(x,y)$$

for 1D and 2D cases, respectively, hence the scheme is node-base.

- Central Finite Difference Method: use central finite difference for both 2nd order and 4th order approximations.
- Special Points in 4th Order: We assume that we know the boundary points in the case also for i = 1, N 1; j = 1, N 1 as inputs.
- **Domain Size:** In 2D case, the domain is rectangular, i.e $\Omega = [a_1, b_1] \times [a_2, b_2]$. Moreover, if $b_1 a_1 = b_2 a_2$, it's square, we consider this case.
- Mesh Size: $\Delta x = \Delta y = h$, i.e. using square mesh.
- **Scheme:** The computational scheme is node-based.
- Smoothness: u is smooth enough that we can do the Taylor expansion.

3 Steady State Heat Equation Numerical Formulation

3.1 Governing Equations

3.1.1 1D Case

Using assumptions above, we can reformulate 1D steady-state heat equation with dirichlet boundary condition into:

$$-k\nabla^2 u(x) = q(x), \quad \forall x \in \Omega = [a, b]$$
 (1)

$$u(a) = u_0(a), \quad u(b) = u_0(b)$$
 (2)

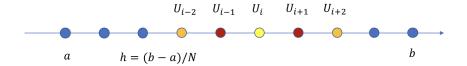


Figure 1: Node-based 1D discretized mesh

3.1.2 2D Case

Using assumptions above, we can reformulate 2D steady-state heat equation with dirichlet boundary condition into:

$$-k\nabla^2 u(x,y) = q(x,y), \quad (x,y) \in \Omega = [a_1, b_1] \times [a_2, b_2]$$
 (3)

$$u(x,y)|_{\partial\Omega} = u_0(x,y) \tag{4}$$

3.2 Generate Grid Points

For 1D case, assume we have N+1 points, and let $h=\Delta x$, then

$$x_i = a + i * h, i = 0, 1, 2, ..., N$$

Let $\Delta x = \Delta y = h = \frac{b_i - a_i}{N_i}$ be the distance between two grid points. Indeed, for rectangular domain, we only have to choose different N_1 and N_2 to get the same h. Then

$$x_i = a_1 + i * h, i = 0, 1, 2, ..., N$$

 $y_i = a_2 + j * h, i = 0, 1, 2, ..., N$

We want to get approximations of $U_{ij} = u(x_i, y_j)$ at grid points (x_i, y_j) . Since by Dirichlet Boundary Condition, the solution on boundaries are known, we only have $(N-1) \times (N-1)$ unknowns to solve.

3.3 2nd-order Finite Difference Approximation

3.3.1 1D Case

Derived from Taylor expansion of u(x,y), we have

$$\nabla^{2} u(x_{i}) = \frac{u(x_{i-1}) - 2u(x_{i}) + u(x_{i+1})}{h^{2}} - \frac{2h^{2}}{4!} \frac{d^{4}u}{dx_{i}^{4}} + O(h^{4})$$

$$= \frac{U_{i-1} - 2U_{i} + U_{i+1}}{h^{2}} + T_{i}$$
(5)

where the local truncation error $T_i = -\frac{2h^2}{4!} \frac{d^4u}{dx_i^4} + O(h^4)$ and we have $\lim_{h\to 0} T_i = 0, \forall i, j$, we can get the following equation:

$$\frac{-k}{h^2}[U_{i-1} - 2U_i + U_{i+1}] = q_i \tag{6}$$

where $q_i = q(x_i), i = 1, 2, ..., N - 1.$

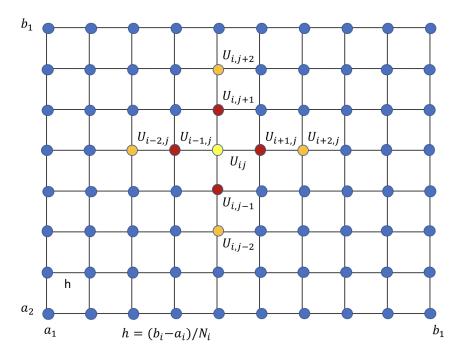


Figure 2: Node-based 2D discretized mesh

3.3.2 2D Case

Derived from Taylor expansion of u(x,y), we have

$$\nabla^{2}u(x_{i}, y_{j}) = \frac{u(x_{i-1}, y_{j}) - 2u(x_{i}, y_{j}) + u(x_{i+1}, y_{j})}{h^{2}} + \frac{u(x_{i}, y_{j-1}) - 2u(x_{i}, y_{j}) + u(x_{i}, y_{j+1})}{h^{2}}$$

$$- \frac{2h^{2}}{4!} \left(\frac{\partial^{4}u}{\partial x_{i}^{4}} + \frac{\partial^{4}u}{\partial y_{j}^{4}}\right) + O(h^{4})$$

$$= \frac{U_{i-1, j} + U_{i+1, j} + U_{i, j-1} + U_{i, j+1} - 4U_{ij}}{h^{2}} + T_{ij}$$
(7)

where the local truncation error $T_{ij} = -\frac{2h^2}{4!} (\frac{\partial^4 u}{\partial x_i^4} + \frac{\partial^4 u}{\partial y_j^4}) + O(h^4)$ and we have $\lim_{h\to 0} T_{ij} = 0, \forall i, j$, we can get the following equation:

$$\frac{-k}{h^2}[U_{i,j-1}+U_{i-1,j}-4U_{ij}+U_{i+1,j}+U_{i,j+1}]=q_{ij}$$
 where $i=1,2,...,N-1,j=1,2,...,N-1$ and $q_{ij}=q(x_i,y_j)$

3.4 4th-order Finite Difference Approximation

3.4.1 1D Case

$$\nabla^2 u(x_i) = \frac{-U_{i-2,j} + 16U_{i-1,j} - 30U_{ij} + 16U_{i+1,j} - U_{i+2,j}}{12h^2} - \frac{8h^4}{6!} \frac{d^6 u}{dx_i^6} + O(h^6)$$
(9)

where $\lim_{h\to 0} -\frac{8h^4}{6!} \frac{d^6u}{dx^6} + O(h^6) = 0, \forall i$. Then we get:

$$\frac{-k}{h^2} \left[-\frac{1}{12} U_{i-2} + \frac{4}{3} U_{i-1} - \frac{5}{2} U_i + \frac{4}{3} U_{i+1} - \frac{1}{12} U_{i+2} \right] = q_i \tag{10}$$

where i = 2, ..., N - 2 and $q_i = q(x_i)$.

3.4.2 2D Case

Similar as above, ignore the high order term of h, we have

$$\nabla^{2}u(x_{i}, y_{j}) = \frac{-U_{i-2,j} + 16U_{i-1,j} - 30U_{ij} + 16U_{i+1,j} - U_{i+2,j}}{12(\Delta x)^{2}} + \frac{-U_{i,j-2} + 16U_{i,j-1} - 30U_{i,j} + 16U_{i,j+1} - U_{i,j+2}}{12(\Delta y)^{2}} + T_{ij}$$
(11)

where the local truncation error $T_{ij} = -\frac{8h^4}{6!} (\frac{\partial^6 u}{\partial x^6} + \frac{\partial^6 u}{\partial y^6}) + O(h^6) \lim_{h\to 0} T_{ij} = 0, \forall i, j \text{ as well, we can get the equations:}$

$$\frac{-k}{h^2} \left[-\frac{1}{12} U_{i-2,j} + \frac{4}{3} U_{i-1,j} - 5 U_{ij} + \frac{4}{3} U_{i+1,j} - \frac{1}{12} U_{i+2,j} - \frac{1}{12} U_{i,j-2} + \frac{4}{3} U_{i,j-1} + \frac{4}{3} U_{i,j+1} - \frac{1}{12} U_{i,j+2} \right] = q_{ij}$$
where $i = 2, ..., N-2, j = 2, ..., N-2$ and $q_{ij} = q(x_i, y_j)$.

Linear System of Heat Equation and Matrix Form 3.5

3.5.1 1D Case

For 1D case, just let $z = [U_0, U_1, ..., U_{N-1}, U_N]^T$, where $U_0 = u_0(a), U_N = u_0(b)$, then we can get the Az = b form, which can be solved by iterative methods. For 2nd order approximation, with Dirichlet boundary conditions known, we have the following form:

$$b_{2nd}^{1D} = [U_0, q_1, ..., q_{N-1}, U_N]^T$$

$$A_{2nd}^{1D} = \frac{-k}{h^2} \begin{bmatrix} \frac{-h^2}{k} \\ 1 & -2 & 1 \\ & 1 & -2 & 1 \\ & & \ddots & \ddots & \ddots \\ & & & 1 & -2 & 1 \\ & & & & \frac{-h^2}{k} \end{bmatrix}$$

#nonzeros of interior: 3, except first and last is 1.

Note: Since for boundary points, we only need 1 for it, to make the matrix simple, I add $\frac{-h^2}{k}$ to recover it to 1. For 4th order approximation,

$$b_{4th}^{1D} = [U_0, U_1, q_2, ..., q_{N-2}, U_N - 1, U_N]^T$$

$$A_{4th}^{1D} = \frac{-k}{h^2} \begin{bmatrix} \frac{-h^2}{k} & & & & & & & \\ & \frac{-h^2}{k} & & & & & \\ -\frac{1}{12} & \frac{4}{3} & -\frac{5}{2} & \frac{4}{3} & -\frac{1}{12} & & & & \\ & \ddots & \ddots & \ddots & \ddots & \ddots & \ddots & \\ & & -\frac{1}{12} & \frac{4}{3} & -\frac{5}{2} & \frac{4}{3} & -\frac{1}{12} & & \\ & & & \frac{-h^2}{k} & & & \\ & & & & \frac{-h^2}{k} & & \end{bmatrix}$$

#nonzeros of interior row: 5, except first, second, first last and second last is 1.

3.5.2 2D Case

We want to transform the 2D problem into a linear system Az = b as well and use iterative methods like Jacobi and Gauss-Seidel to solve it. First, assume $N_1 = N_2 = N$, flatten U_{ij} into a vector z using rule:

$$z_{i(N+1)+j} = U_{i,j}, \forall i = 0, 1, 2, ..., N, j = 0, 1, 2, ..., N$$

Then we have $z = [U_{00}, U_{01}, ..., U_{0N}, U_{10}, ..., U_{NN}]^T$.

Second, we use the same rule of adding boundary variables as 1D case to flatten f into b:

$$b_{2nd}^{2D} = \left[U_{00}, U_{01}, ..., U_{0,N}, U_{10}, q_{11}, ..., q_{1,N-1}, ..., q_{N-1,N-1}, U_{N-1,N}, ..., U_{N,N}\right]^T$$

Third, we can get the form of corresponding A:

$$A_{2nd}^{2D} = \begin{bmatrix} I \\ I_{2nd} & B_{2nd} & I_{2nd} \\ & \ddots & \ddots & \ddots \\ & & & I_{2nd} & B_{2nd} & I_{2nd} \\ & & & & I \end{bmatrix}$$

where I_{2nd} and B_{2nd} are $N+1\times N+1$ matrix. #nonzeros of interior row: 5, except special cases with 1.

$$B_{2nd} = \frac{-k}{h^2} \begin{bmatrix} \frac{-h^2}{k} \\ 1 & -4 & 1 \\ & \ddots & \ddots & \ddots \\ & & 1 & -4 & 1 \\ & & & \frac{-h^2}{k} \end{bmatrix}$$

$$I_{2nd} = \frac{-k}{h^2} \begin{bmatrix} 0 & & & & \\ & 1 & & & \\ & & \ddots & & \\ & & & 1 & \\ & & & & 0 \end{bmatrix}$$

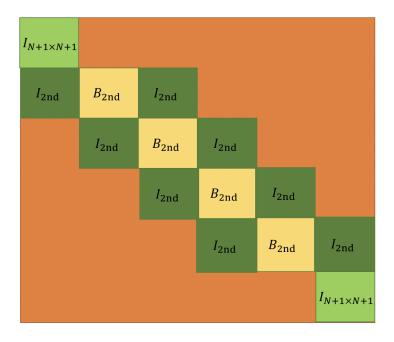


Figure 3: Structure of 2nd order difference matrix

For 4th order approximation, we know more boundary conditions including i = 0, 1, N - 1, N; j = 0, 1, N - 1, N, we have:

$$b_{4th}^{2D} = [U_{00}, U_{01}, ..., U_{0N}, U_{10}, U_{11}, ..., U_{1N}, U_{20}, q_{21}, ..., q_{N-2,N-2}, U_{N-1,N}, ..., U_{N,N}]^T$$

where $I = I_{N+1 \times N+1}$ and B is also $N+1 \times N+1$ matrix. #nonzeros of interior row: 9, except boundary cases are 1.

$$B_{4th} = \frac{-k}{h^2} \begin{bmatrix} \frac{-h^2}{k} & & & & & & & \\ & \frac{-h^2}{k} & & & & & \\ -\frac{1}{12} & \frac{4}{3} & -5 & \frac{4}{3} & -\frac{1}{12} & & & & \\ & \ddots & \ddots & \ddots & \ddots & \ddots & \ddots & \\ & & -\frac{1}{12} & \frac{4}{3} & -5 & \frac{4}{3} & -\frac{1}{12} & & \\ & & & \frac{-h^2}{k} & & & \\ & & & & \frac{-h^2}{k} & \end{bmatrix}$$

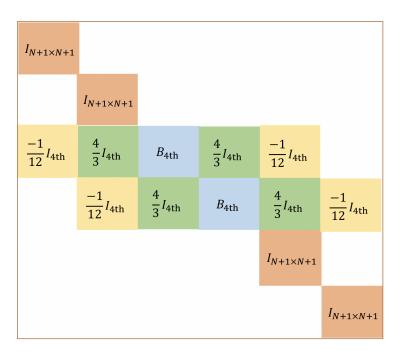


Figure 4: Structure of 4th order difference matrix

$$I_{4th} = \frac{-k}{h^2} \begin{bmatrix} 0 & & & & & \\ & 0 & & & & \\ & & 1 & & & \\ & & & \ddots & & \\ & & & & 1 & \\ & & & & 0 & \\ & & & & 0 \end{bmatrix}$$

3.6 Iterative Methods to Solve Linear Systems

3.6.1 Jacobi Iterative Method

Given an initial guess \mathbf{z}^0 , here we use 0 for initialization, then the rule of Jacobi iterative method is

$$z_i^{k+1} = \frac{1}{a_{ii}} (b_i - \sum_{j=1, j \neq i}^n a_{ij} z_j^k)$$
(13)

3.6.2 Gauss-Seidel Iterative Method

Given an initial guess \mathbf{z}^0 , here we use 0 for initialization, then the rule of Gauss-Seidel iterative method is

$$z_i^{k+1} = \frac{1}{a_{ii}} \left(b_i - \sum_{j=1}^{i-1} a_{ij} z_j^{k+1} - \sum_{j=i+1}^{n} a_{ij} z_j^k \right)$$
 (14)

iter_method/dim	1D	2D
2nd order	$(11 \times (N+1) + 20) \times 8B$	$(15 \times (N+1)^2 + 20) \times 8B$
4th order	$(15 \times (N+1) + 20) \times 8B$	$(23 \times (N+1)^2 + 20) \times 8B$

Table 1: Memory Estimate Summary

Gauss-Seidel method use the newest to solve linear system, while Jacobi use the old one. For performance, Gauss-Seidel usually converges and use less memory, but not for Jacobi. In practice, we can see from the results, Jacobi does not converge for 4th order finite difference.

4 Algorithm

5 Memory Estimate

Memory that we need includes the following things:

- Iteration Variables: matrix A, vector b, vector z, k, error
- Inputs: tolerance ϵ , max_iter, iter_method, dim, fd_method, domain, N, output_mode, verification_mode, output_file
- 1. For matrix A, because it has special diagonal structures, we can use sparse structure to store it. Use 1D array to save number of non-zeros in each row, let n = N+1 for 1D case and n = (N+1)*(N+1) for 2D case. Then we need n for non-zeros. Use 2D array to save non-zero column index and values, we need 3n for 1D 2nd, 5n for 1D 4th and 2D 2nd, 9n for 2D 4th.
- 2. For inputs and iteration variables k, error, we only need 1 double to store them, we can estimate 20 doubles for them.
- 3. For z and b, we need N+1 for 1D case and $(N+1)^2$ for 2D case, respectively.
- 4. For verification mode, we need n for exact value.

The difference between Jacobi and Gauss-Seidel iterations is that we update the variable in Gauss-Seidel iteration in the loop, so we only need one z to store it. But for Jacobi, we update them until we get all the new values for a new iteration, so we need at least two z to store them. But we need to compare the l2 norm for current and previous solution, so here we need one more n doubles for "old z".

And we need 8B to store a double, then the total memory we need is as follows in Table 2. And we allocate dynamic memory to these variables.

Algorithm 1 Numerical Methods to Solve Steady State Heat Equations

Input: tolerance $\epsilon > 0$, max_iter, iter_method, dim, fd_method, domain, N, dim, output_mode, verification_mode, output_file

Output: $u(x_i, y_j), i = 1, 2, ..., N - 1, j = 1, 2, ..., N - 1$

Initialize: k = 0, error = 1000, $f(x_i, y_j)$

if $\dim == 1$ then

$$z = [U_0, U_1, ..., U_{N-1}, U_N]^T$$

 $if \text{ fd_method} == 2nd then$

$$b = b_{2nd}^{1D}$$

$$A = A_{2nd}^{1D}$$

else

$$b = b_{4th}^{1D}$$

$$A = A_{4th}^{1D}$$

if $\dim == 2$ then

$$z = [U_{00}, U_{01}, ..., U_{0N}, U_{10}, ..., U_{NN}]^T$$

if $fd_method == 2nd$ then

$$b = b_{2nd}^{2D}$$

$$A = A_{2nd}^{2D}$$

else

$$b = b_{2nd}^{2D}$$

$$A = A_{Ath}^{2D}$$

while error $> \epsilon$ and k <= max_iter do

 $\mathbf{if} \text{ iter_method} == \text{Jacobi } \mathbf{then}$

for
$$i = 0, 1, 2, ..., N$$
 do sum = 0

for
$$j = 0, 1, 2, ..., N$$
 do

if
$$j = 0, 1, 2, ..., N$$
 do
if $j \neq i$ then $sum = sum + a_{ij}z_j^k$
 $z_i^{k+1} = \frac{1}{a_{ii}}(b_i - sum)$

 $\mathbf{if} \ \mathrm{iter_method} == \mathrm{Gauss_Seidel} \ \mathbf{then}$

for
$$i = 0, 1, 2, ..., N$$
 do sum = 0

for
$$j = 0, 1, 2, ..., i - 1$$
 do $sum = sum + a_{ij}z_j^{k+1}$

for
$$j = i + 1, i + 2, ..., N$$
 do $sum = sum + a_{ij}z_j^k$
 $z_i^{k+1} = \frac{1}{a_{ii}}(b_i - sum)$

error =
$$||z^{k+1} - z^k||_2$$

 $k \leftarrow k + 1$
return z^k

$$k \leftarrow k + 1$$

6 User Instructions

6.1 Build Procedures

Please follow the procedures to build heat equation solving systems on Stampede2, and the command is also include in "command.sh".

- 1. Download the tar of codes, "proj01" from Github, untar the files;
- 2. Use "autoreconf -f -i" to do the bootstrap;
- 3. Use the following command to do the configuration;

```
./configure --with-masa=$PKGPATH/masa-gnu7-0.50 \
--with-grvy=$PKGPATH/grvy-gnu7-0.34
```

- 4. Use "make" to automatically build the system and "make check" to check it runs correctly;
- 5. Use "cd src" to enter in the src file and "./solver input.dat" to solve heat equation, where you can change input.dat following the instructions in the file, which will also be illustrated in the "Input Options".

6.2 Input Options

Input options can be changed in input.dat, which shows in Figure 5. Note that you can also change the name and the file name on the command line.

- k: thermal conductivity, which will be used as parameter for MASA
- verify_mode: 1 will enable verification mode with MASA and output error norm
- output_ mode: 0 = silent, 1 = standard, 2 = debug
- output_file: name of solution output file, for "make check", please set it as "sol.dat"
- dimensions: choose 1 for 1D and 2 for 2D
- xmin, xmax, ymin, ymax: set range of each axis
- N: number of intervals in one axis, points are N+1 for 1D, $(N+1)^2$ for 2D.
- fd_method: 2 = second order, 4 = fourth order
- \bullet iter_method: choose 1 for Jacobi or 2 for Gauss-Seidel
- \bullet eps: iterative solver tolerance, default is 1.0e-12
- max_iter: max solver iterations

```
# -*-sh-*-
# input file for solving heat conduct equation
                                    # thermal conductivity [W/mK]
                                    # enable verification mode with MASA
# output mode (0 = silent, 1 = standard, 2 = debug)
# name of solution output file
verify_mode = 1
output_mode = 2
output_file = 'sol.dat'
[mesh]
dimensions = 2
                                     # 1 or 2 dimensions?
xmin = 0
                                     # min x location [m]
                                     # max x location [m]
ymin = 0
                                     # min y location [m]
ymax = 1
N = 16
                                     # number of intervals in one axis, points are N+1 for 1D, (N+1)^2 for 2D,
                                     \# 2 = second order, 4 = fourth order
iter\_method = 2
                                     # choose 1 for Jacobi or 2 for Gauss-Seidel
           = 1.0e-12
= 250000
                                      # iterative solver tolerance
max_iter
                                     # max solver iterations
```

Figure 5: Input options example

7 Verification Procedures and Exercise

Just change the input.dat with verify_mode to run in verification mode, since the silent output mode will show nothing, please change the output_mode to 1 or 2 to see the l_2 norm for the difference between the solution and the reference result generated from MASA. You can find example output for standard output in Figure 7, for debug output in Figure 8.

For verification exercise, we can run with N=8,16,32,64,128,256, and use "loglog" to plot and to illustrate the convergence rate with reference line. The expected slope for 2nd order is -2 and for 4th order is -4. We can see the result of Gauss-Seidel iterative method for 1D and 2D in Figure 9 and Figure 10, respectively; and Jacobi iterative method for 1D and 2D with 2nd order in Figure 11. They are all almost parallel to the reference line respectively.

From the results we can get slope as follows, Gauss and Jacobi are close in 2nd order.

• 1d 2nd gauss: -2.0238

• 1d 4th gauss: -3.8642

• 2d 2nd gauss: -1.9933

• 2d 4th gauss: -3.8663

• 1d 2nd jacobi: -2.0038

• 2d 2nd jacobi: -1.9933

```
** Finite-difference based Heat Equation Solver (steady-state)
--> Parsing runtime options from input.dat
--> h = 0.062500

** Runtime mesh settings (1D):
--> nx = 16 (xmin,xmax) = (0.000000, 1.000000)

** Runtime solver settings:
--> finite difference method = CENTRAL2
--> iterative method = GAUSS.SEIDEL
--> max iterations = 250000
--> thermal conductivity = 1.0000000
--> solution output file = sol.dat

** Initializing data structures...

** Building linear system...
--> Enforcing analytic Dirichlet BCs using MASA (1D)

** Solving linear system...
--> Converged at iter: 526
--> Writing output to sol.dat

** Computing 12 error norm
--> 12 error norm = 2.460876e-02

Steady Heat Equation Solver - Performance Timings: | Mean Variance Count
--> output : 2.66288e-03 secs (36.2762 %) | [2.65288e-03 0.00000e+00 1]
--> parse_input : 2.60910e-03 secs (35.6763 %) | [2.60010e-03 0.00000e+00 1]
--> init : 1.92595e-03 secs (26.3359 %) | [1.92595e-03 0.00000e+00 1]
--> solve_system : 2.00272e-05 secs (1.2715 %) | p3.29832e-05 0.00000e+00 1]
--> solve_system : 2.00272e-05 secs (0.2739 %) | [2.00272e-05 0.00000e+00 1]
--> GRYY_Unassigned : 1.21593e-05 secs (0.1663 %)
Total Measured Time = 7.31301e-03 secs (100.0000 %)
```

Figure 6: Example of standard output for verification

```
** Solving linear system...
--> Converged at iter: 92
--> Writing output to sol.dat

** Computing l2 error norm.
--> 12 error norm = 6.948365e-03
```

Figure 7: Standard for verification

```
[debug]: solve_system - function end
[debug]: output - function begin
--> Writing output to sol.dat
[debug]: output - function end
[debug]: error_norm - function begin

** Computing 12 error norm.
--> 12 error norm = 6.948365e-03

[debug]: error_norm - function end
[debug]: ~Laplacian_FD - function begin
```

Figure 8: Debug for verification

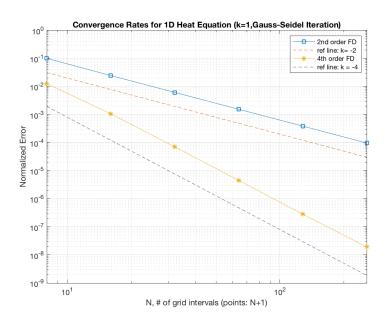


Figure 9: Convergence Rates for 1D Heat Equation with Gauss-Seidel

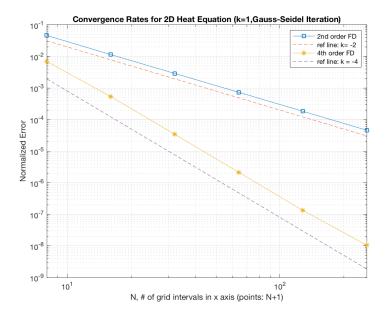


Figure 10: Convergence Rates for 2D Heat Equation with Gauss-Seidel

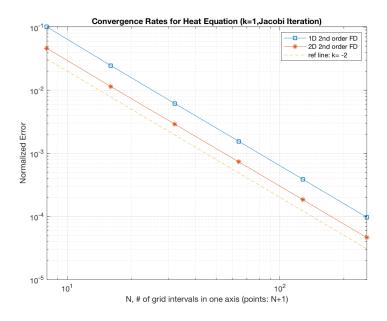


Figure 11: Convergence Rates for Heat Equation with Jacobi

8 Runtime Performance

We use "GRVY" to get the summary of application time. As we can see from Figure 12, the 3 sections accounting for at least 90% are: parse_input, init and output, since the number of points are very small and solving process is very fast, the time spends more on read and output files. From Figure 13, the top three are solver_system parse_input and output, as the number of points rises, more time will spend on solver system procedure. When the number of points arrives at 256, over 99% time spends on solving procedure.

Figure 12: Runtime summary for 8 points

N / process %	solve_system	parse_input	output
8	0.91	46.87	25.16
32	71.54	11.29	9.19
256	99.88	0.0078	0.0667

Table 2: Runtime Summary

Figure 13: Runtime summary for 32 points

Figure 14: Runtime summary for 256 points