### Falsifiable Experimental Signatures of Density Field Dynamics: Phase Velocity Equals One-Way Light Speed in a Nondispersive Vacuum

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We convert  $Density\ Field\ Dynamics\ (DFD)$  into a laboratory-focused, falsifiable test program. DFD posits a single scalar field  $\psi(\mathbf{x})$  that governs matter dynamics and photon propagation through a universal, nondispersive vacuum refractive structure. The core operational result is that, when dispersion is bounded in-band, the electromagnetic phase velocity is the one-way speed of light; hence precision phase metrology becomes a direct, synchronization-free probe of  $c_1(\mathbf{x})$ . We derive this identity along two independent routes (Fermat/eikonal and Gordon's optical metric), demonstrate compatibility with classic tests of relativity, and design three GR-null vs DFD-signal protocols—most decisively a co-located cavity—atom frequency ratio measured at two altitudes—with quantified, near-term sensitivities. We audit existing constraints, state explicit refutation criteria, and provide a comprehensive responses-to-criticisms section (simultaneity, equivalence principle, Lorentz invariance, "already ruled out", dispersion). The question is experimentally decidable with current optical metrology.

#### I. FROM PRINCIPLE TO PROTOCOL

Two-way light speed and Lorentz symmetry are constrained to extraordinary precision [1–4]. The one-way speed remains entangled with simultaneity conventions [5–7]. DFD proposes a dynamical scalar  $\psi(\mathbf{x})$  that (i) fixes a universal vacuum refractive index  $n=e^{\psi}$  for photons, and (ii) normalizes the Newtonian limit for matter via  $\mathbf{a}=(c^2/2)\nabla\psi$ . In a verified nondispersive band, geometric optics yields  $v_{\text{phase}}=c/n$ , which, with  $n=e^{\psi}$ , provides the operational bridge  $c_1=c\,e^{-\psi}=v_{\text{phase}}$ . The novelty is experimental: route-dependent, synchronization-free observables that are GR-null but DFD-nonnull.

### II. DFD DYNAMICS IN ONE PAGE

Adopt a scalar action with a single crossover scale,

$$\nabla \cdot \left[ \mu \left( \frac{|\nabla \psi|}{a_{\star}} \right) \nabla \psi \right] = -\frac{8\pi G}{c^2} \left( \rho_m - \bar{\rho}_m \right), \quad \mathbf{a} = \frac{c^2}{2} \nabla \psi,$$
(1)

so that  $\psi$  is fixed by matter density and yields  $\psi \simeq -2\Phi/c^2$  in weak fields. Photons propagate by Fermat/optical metric (Sec. III). The weak-field normalization is chosen to reproduce GR's classic optical tests with PPN  $\gamma=1$  [1, 8]. A Sakharov-style perspective motivates induced kinetic terms from quantum fluctuations [9], but the empirical program below does not rely on specific UV details.

### III. CORE IDENTITY: $v_{\text{phase}} = c_1$

### A. Route I: Fermat/eikonal

Geometric optics extremizes  $T[\gamma] = (1/c) \int_{\gamma} n(\mathbf{x}) d\ell$ , giving  $v_{\text{phase}} = c/n$  [8, 10, 11]. With  $n = e^{\psi}$  fixed by

dynamics,  $c_1 = ce^{-\psi} = v_{\text{phase}}$  follows. No distant clocks enter: the observable is local phase kinematics, verified nondispersive.

#### B. Route II: Gordon's optical metric

Light in a linear, isotropic, nondispersive medium follows null geodesics of

$$\mathrm{d}\tilde{s}^2 = \frac{c^2}{n^2(\mathbf{x})} \,\mathrm{d}t^2 - \mathrm{d}\mathbf{x}^2 \tag{2}$$

[12]. Nullness implies  $d\ell/dt = c/n$ ; with  $n = e^{\psi}$  the same identity follows. The equality is therefore structural (two logically independent routes), not a definitional tautology.

### IV. EQUIVALENCE PRINCIPLE & LORENTZ INVARIANCE

**Local Lorentz invariance.**  $\psi$  is a scalar; the light cone at a point remains isotropic. Two-way orientation/boost tests remain null at the  $10^{-17}$ – $10^{-18}$  level, consistent with cavity experiments [2–4].

Universality for matter. Test bodies obey  $\mathbf{a} = (c^2/2)\nabla\psi$  in the weak field, reproducing free-fall universality and PPN  $\gamma = 1$  optics [1]. Equivalence principle tests remain satisfied in this limit.

Where differences appear. DFD predicts route-dependent, synchronization-free effects where GR enforces strict nulls: e.g., a co-located cavity—atom ratio compared at two altitudes (Sec. VI). This is an LPI probe in a nondispersive vacuum sector not covered by atom—atom redshift verifications.

### V. WHAT EXISTING TESTS DO—AND DO NOT—CONSTRAIN

Two-way isotropy & boost (Michelson-Morley/Kennedy-Thorndike; modern rotating cavities) are exquisitely null [2–4]; DFD predicts the same nulls for two-way observables along fixed paths.

Atomic clock redshift confirms GR at  $\sim 10^{-16}$  per meter and below [13, 14]; spaceborne tests reach  $2.5 \times 10^{-5}$  relative precision [15]. These are atom–atom or remotetransfer comparisons.

Critical gap: To our knowledge, no published measurement reports a co-located cavity-atom frequency ratio recorded at two different gravitational potentials with  $< 10^{-16}$  fractional uncertainty. That is the target of Protocol C.

# VI. LABORATORY PROTOCOLS (GR-NULL VS DFD-SIGNAL)

All protocols enforce nondispersion via multiwavelength checks which bound  $|\partial n/\partial \omega|$  in the measurement band (so phase=group=front) [10, 16].

# Protocol A: Crossed ultra-stable cavities (orientation/height sweep)

Two orthogonal high-Q cavities (length L) support modes  $f_m \simeq \frac{m}{2L} \frac{c}{n}$ . A change  $\delta \psi$  imparts  $\delta f/f = -\delta n/n = -\delta \psi$ . Orientation reversals and vertical translations by  $\Delta h$  probe geometry-locked shifts. Target sensitivity:  $10^{-17}$ - $10^{-16}$  fractional, routinely achieved [2, 4].

# Protocol B: Reciprocity-broken fiber loop (two heights)

A monochromatic tone circulates both ways around an asymmetric loop with vertical separation  $\Delta h$  and a nonreciprocal element. The accumulated phase  $\phi = \frac{\omega}{c} \int n(\mathbf{x}) d\ell$  yields a forward–backward difference  $\propto \oint \psi d\ell$  that vanishes in GR (static loop, Sagnac subtracted) but not in DFD if  $\nabla \psi \cdot \hat{z} \neq 0$ . Operate near a zero-dispersion wavelength; multi- $\lambda$  tracking bounds dispersion [17].

### Protocol C (decisive): Co-located cavity—atom ratio across altitude

Lock a laser to a vacuum cavity (frequency  $f_{\rm cav} \propto c/n$ ) and compare to a co-located optical atomic transition  $f_{\rm at}$  via a frequency comb; form  $R \equiv f_{\rm cav}/f_{\rm at}$  at altitude  $h_1$ , repeat at  $h_2 = h_1 + \Delta h$ .

GR: Moving the co-located package changes neither the local ratio nor local physics; gravitational redshift appears

Protocol	Observable	DFD signal (order)
A: Crossed cavities	$\delta f/f$ on rotate/ $\Delta h$	
B: Fiber loop	$\Delta\phi_{\circlearrowright}-\Delta\phi_{\circlearrowleft}$	geometry-locked, $< 10^{-16}$ eqv.
C: Cavity/atom ratio	$\Delta R/R$ across $\Delta h$	$2g\Delta h/c^2\approx 2.2\times 10^{-14}~{\rm per}~100~{\rm m}$

TABLE I. Order-of-magnitude DFD signals in verified nondispersive band. GR predicts strict nulls for A/B and near-null for C.

upon comparison between distinct potentials, so R is (to excellent approximation) constant.

DFD (nondispersive): With  $\psi \simeq -2\Phi/c^2$ ,  $f_{\rm cav} \propto e^{-\psi} \simeq 1 + 2\Phi/c^2$ , while the atomic transition is leading-order  $\psi$ -insensitive (matter-sector universality). Thus

$$\frac{\Delta R}{R} \; \simeq \; 2 \, \frac{\Delta \Phi}{c^2} \; \approx \; 2 \, \frac{g \, \Delta h}{c^2} \; \sim \; 2.2 \times 10^{-14} \; \; {\rm per} \; 100 \; {\rm m}. \label{eq:deltaR}$$

A variant with small matter-sector coupling yields  $\frac{\Delta R}{R} = \xi \Delta \Phi/c^2$  with  $0 < \xi \le 2$ , still at the  $10^{-16} \, \mathrm{m}^{-1}$  scale. Feasibility: present clocks and cavities reach  $10^{-17}$ – $10^{-16}$  [2, 4, 13, 14].

Systematics & controls (all protocols). (i) Multi- $\lambda$  dispersion bound; (ii) temperature/strain control, Allan budgeting; (iii) polarization scrambles and hardware swaps; (iv) blind orientation/height reversals; (v) environmental monitors (pressure, tilt, vibration).

### VII. PREDICTED SIGNAL SIZES AND SENSITIVITY TABLE

# VIII. REFUTATION CRITERIA (CLEAN KILL CONDITIONS)

Any of the following falsifies this DFD formulation:

- 1. Protocol C yields  $\Delta R/R$  consistent with zero at or below  $|\Delta\Phi|/c^2$  while dispersion and thermal budgets pass all checks.
- 2. Protocols A or B yield nulls where DFD predicts nonzero geometry-locked signals after reversals/path swaps.
- 3. A verified nonzero dispersion  $(\partial n/\partial \omega)$  fully accounts for any residuals across the band.

Conversely, reproducible nonzero signals that (i) scale with  $\Delta h$  or orientation as predicted, (ii) survive multi- $\lambda$  tests, and (iii) pass swap/blind controls, would be decisive.

#### IX. COSMOLOGICAL CONTEXT (BRIEF)

With  $\psi \simeq -2\Phi/c^2$ , Gordon's metric reproduces classic weak-field optics [1, 8]. DFD suggests that nearby structure can bias line-of-sight cosmography at low z, providing a plausible context for directional  $H_0$  inferences; early-universe constraints (CMB/BAO) remain intact [18, 19].

These motivate but do not underwrite the laboratory program.

## X. COMPREHENSIVE RESPONSES TO STANDARD CRITICISMS

- (1) "You haven't solved simultaneity; one-way c is conventional." We agree that simultaneity is conventional in SR. DFD makes a different claim: in a verified nondispersive band, local phase velocity equals the one-way propagation speed. Our observables are local and synchronization-free; they exploit route dependence where GR says strict null. This turns a philosophical impasse into a falsifiable statement [5–7].
- (2) " $n = e^{\psi}$  makes  $c_1 = c/n$  definitional (circular)." The identity  $v_{\rm phase} = c/n$  follows from standard optics (Fermat/eikonal and Gordon's metric) independently [8, 10, 12]. DFD then supplies dynamics for  $\psi$  (Eq. 1), fixed by classic-test normalization [1]. The bridge is therefore derived, not stipulated.
- (3) "Equivalence principle is violated: photons vs matter." Matter test bodies obey  $\mathbf{a} = (c^2/2)\nabla\psi$  (universality preserved in weak field). Photons see the optical metric which reproduces GR's lensing/redshift (PPN  $\gamma=1$ ). Our key test (Protocol C) is an *LPI* probe in the nondispersive vacuum sector; either a residual appears (then LPI is violated in this sector) or it does not (DFD is falsified).
- (4) "Lorentz symmetry constraints already exclude this." Two-way isotropy/boost tests [2–4] remain null in DFD. Differences appear only between *paths* sampling different  $\psi$  (height/orientation). This is not what SME-style cavity rotations constrain [20].
- (5) "Existing optical clocks at different elevations would have seen it." Published redshift verifications are atom–atom or remote-transfer comparisons [13–15]. They confirm GR and are orthogonal to our decisive co-located cavity–atom <u>ratio</u> across altitudes. To our knowledge, such a ratio-vs-altitude measurement at  $< 10^{-16}$  is not yet published; Protocol C is designed to fill this gap.
- (6) "This is just superluminal phase velocity; information rides on front velocity." Correct in dispersive media [16]. Our tests operate in a verified nondispersive band where phase=group=front [10]; the identity is invoked only under those conditions.
- (7) "\$\psi\$ is ad hoc and parameters are tuned." The weak-field normalization is fixed by classic tests; induced-

gravity arguments [9] motivate scalar kinetic terms. However, our claims do not hinge on UV priors: the laboratory identity and protocols stand on their own as empirical tests.

#### XI. CONCLUSIONS

We have turned DFD into a concrete, near-term experimental program: (i) a structural identity that makes phase metrology a one-way-c probe in a verified nondispersive vacuum; (ii) three synchronization-free protocols with GR-null vs DFD-signal contrasts and quantified targets; (iii) an explicit constraints audit and clean refutation logic. The decisive experiment (Protocol C) is implementable now with optical cavities, clocks, and frequency combs. Either the geometry-locked phase-velocity effects appear (opening a new sector of physics) or the present DFD is falsified.

### Appendix A: Geometrical optics and nondispersion

Let S be the eikonal:  $\mathbf{k} = \nabla S$ ,  $\omega = -\partial_t S$ . For  $\omega = (c/n)|\mathbf{k}|$ ,  $v_{\text{phase}} = \omega/|\mathbf{k}| = c/n$  and  $v_g = \partial \omega/\partial |\mathbf{k}| = c/n$ ; the Sommerfeld–Brillouin front velocity coincides in the nondispersive limit [10, 16].

### Appendix B: Round-trip nulls, clocks, and GPS

For a fixed path  $\gamma$ ,  $T_{2w} = \frac{2}{c} \int_{\gamma} n \, d\ell$ ; at fixed geometry, orientation rotations preserve two-way times (Michelson–Morley nulls). Clock redshift verifications rely on atom–atom or remote transfers consistent with GR [13–15]; our decisive observable is a *local* cavity–atom ratio across altitude.

### Appendix C: Minimal implementation checklist

Cavities: ULE/silicon spacers; PDH locking; cryogenic option;  $10^{-17}$  stability [2, 4].

**Fibers:** zero-dispersion operation; Faraday isolators; dual- $\lambda$  phase-tracking [17].

Clocks: Sr/Yb lattice or Al<sup>+</sup> logic; comb-based ratio readout [13, 14].

Analysis: publish  $\sigma_y(\tau)$ ; blinded reversals; multi- $\lambda$  linearity fits; environmental logs.

<sup>[1]</sup> Clifford M. Will. The confrontation between general relativity and experiment. *Living Reviews in Relativity*, 17:4, 2014.

<sup>[2]</sup> Ch. Eisele, A. Y. Nevsky, and S. Schiller. Laboratory test of the isotropy of light propagation at the  $10^{-17}$  level. *Phys. Rev. Lett.*, 103(9):090401, 2009.

<sup>[3]</sup> S. Herrmann et al. Rotating optical cavity experiment testing lorentz invariance at the  $10^{-17}$  level. *Phys. Rev.* D, 80(10):105011, 2009.

 <sup>[4]</sup> M. Nagel et al. Direct terrestrial test of lorentz symmetry in electrodynamics to 10<sup>-18</sup>. Nature Communications, 6:8174, 2015.

- [5] Hans Reichenbach. The Philosophy of Space and Time. Dover, New York, 1958.
- [6] W. F. Edwards. Special relativity in anisotropic space. American Journal of Physics, 31(7):482–489, 1963.
- [7] David Malament. Causal theories of time and the conventionality of simultaneity. Noûs, 11(3):293-300, 1977.
- [8] Volker Perlick. Ray Optics, Fermat's Principle, and Applications to General Relativity, volume 61 of Lecture Notes in Physics Monographs. Springer, 2000.
- [9] A. D. Sakharov. Vacuum quantum fluctuations in curved space and the theory of gravitation. Sov. Phys. Dokl., 12:1040–1041, 1968.
- [10] Max Born and Emil Wolf. Principles of Optics. Cambridge University Press, 7 edition, 1999.
- [11] John D. Jackson. Classical Electrodynamics. Wiley, New York, 3 edition, 1998.
- [12] W. Gordon. Zur lichtfortpflanzung nach der relativitätstheorie. Annalen der Physik, 377(22):421–456, 1923.
- [13] C. W. Chou, D. B. Hume, T. Rosenband, and D. J. Wineland. Optical clocks and relativity. Science,

- 329(5999):1630-1633, 2010.
- [14] W. F. McGrew et al. Atomic clock performance enabling geodesy below the centimetre level. *Nature*, 564:87–90, 2018.
- [15] P. Delva et al. Gravitational redshift test using eccentric galileo satellites. Phys. Rev. Lett., 121(231101), 2018.
- [16] Léon Brillouin. Wave Propagation and Group Velocity. Academic Press, 1960.
- [17] Govind P. Agrawal. Fiber-Optic Communication Systems. Wiley, 4 edition, 2010.
- [18] Planck Collaboration. Planck 2018 results. vi. cosmological parameters. Astronomy & Astrophysics, 641:A6, 2020.
- [19] A. G. Riess et al. A comprehensive measurement of the local value of the hubble constant with 1 km s<sup>-1</sup> mpc<sup>-1</sup> uncertainty from the hubble space telescope and the sh0es team. Astrophys. J. Lett., 934(1):L7, 2022.
- [20] V. Alan Kostelecký and Neil Russell. Data tables for lorentz and cpt violation. Rev. Mod. Phys., 83:11, 2011. January 2024 update available at arXiv:0801.0287.