

Physics 412-1 Problem Set 4 Solutions

Problem 1

We must solve the equation

$$\left(-\frac{\hbar}{2m} \frac{d^2}{dx^2} - V_0 \delta(x)\right) \phi(x) = E \phi(x) \quad (1)$$

for $E < 0$.

For $x > 0$, the equation is

$$-\frac{\hbar}{2m} \frac{d^2}{dx^2} \phi(x) = E \phi(x). \quad (2)$$

Therefore the solution is

$$\phi_{>}(x) = e^{-\kappa x}, \quad (3)$$

where $\kappa = \sqrt{\frac{-2mE}{\hbar^2}}$. Similarly for $x < 0$, the solution is

$$\phi_{<}(x) = e^{\kappa x}. \quad (4)$$

To see how to match these solutions, integrate the full equation over a small interval around $x = 0$, we have

$$-\frac{\hbar}{2m} \frac{d\phi}{dx} \Big|_{-\epsilon}^{+\epsilon} - V_0 \phi(0) = O(\epsilon). \quad (5)$$

So as $\epsilon \rightarrow 0$,

$$-\frac{\hbar}{2m} \left(\frac{d\phi_{>}}{dx}(0) - \frac{d\phi_{<}}{dx}(0) \right) = V_0 \phi(0). \quad (6)$$

And therefore $\kappa = \frac{mV_0}{\hbar^2}$ and $E = -\frac{mV_0^2}{2\hbar^2}$.

To sum up, $\phi(x) = \frac{\sqrt{mV_0}}{\hbar} e^{-mV_0|x|/\hbar^2}$ and energy $E = -\frac{mV_0^2}{2\hbar^2}$, which is a unique answer by the procedure.

Problem 2

(a)

Let the wavefunction to be an eigenstate over all sapce with $E = \frac{\hbar^2 k^2}{2m}$. Inside the well, we have $\phi_a(x) = T_0 e^{ik_0 x} + R_0 e^{-ik_0 x}$, with energy $\frac{\hbar^2 k_0^2}{2m} - V_0$ Therefore

$$k_0 = \sqrt{k^2 + \frac{2mV_0}{\hbar^2}} \quad (7)$$

Consider the boundary conditions, at $x = 0$ we have

$$\phi_L = \phi_R \rightarrow 1 + R = R_0 + T_0, \quad (8)$$

$$\phi'_R = \phi'_L \rightarrow k(1 - R) = k_0(T_0 - R_0). \quad (9)$$

At $x = a$, we have

$$R_0 e^{-ik_0 a} + T_0 e^{ik_0 a} = T e^{ika}, \quad (10)$$

$$k_0(T_0 e^{ik_0 a} - R_0 e^{-ik_0 a}) = k T e^{ika}. \quad (11)$$

Solve the above equations, we have

$$T_0 = \frac{2(1 + c)}{(1 + c)^2 - (1 - c)^2 e^{i\theta}} \quad (12)$$

$$R_0 = \frac{-2e^{-i\theta}(1 - c)}{(1 + c)^2 - (1 - c)^2 e^{i\theta}} \quad (13)$$

$$R = \frac{(e^{i\theta} - 1)(c^2 - 1)}{(1 + c)^2 - (1 - c)^2 e^{i\theta}} \quad (14)$$

$$T = \frac{4e^{i(k_0 - k)a}}{(1 + c)^2 - (1 - c)^2 e^{i\theta}} \quad (15)$$

where $c = \frac{k_0}{k}$ and $\theta = 2k_0 a$

(b)

As $E \rightarrow \infty$, $c \rightarrow 1$, so $R \rightarrow \frac{(e^{i\theta} - 1)(0)}{1} = 0$.

(c)

As $V_0 \rightarrow \infty$, $c \rightarrow \infty$. So

$$R \rightarrow \frac{(e^{i\theta} - 1)c^2}{c^2(1 - e^{i\theta})} \quad (16)$$

If $e^{i\theta} - 1 \neq 0$, $R \rightarrow -1$.

But if $e^{i\theta} = 1$, we need to go back and calculate

$$R = T_0 + R_0 - 1 = \frac{2(1+c) - 2(1-c)}{1+2c+c^2 - 1+2c-c^2} - 1 = \frac{4c}{4c} - 1 = 0 \quad (17)$$

To sum up, as $V_0 \rightarrow \infty$, $R \rightarrow -1$ (totally reflected with π phase shift), unless $k_0 a = n\pi$, in which case $R \rightarrow 0$. This phenomenon is called a resonance, and has no analog in motion of classical particles.