

Postulates of QM, Part I

- Up til now we have studied the Schr eqn, but we haven't discussed the fundamentals of QM

- We have seen that the following associations

Schr w.f. \leftrightarrow config. or location of particle

eigenvalues of energy \leftrightarrow definite energies of system

$|y(x)|^2 \leftrightarrow$ probability of finding particle @ x

Provide sensible physical interp + correct predictions

- But we have been following a semi-historical approach, discovering as we go.

- Really QM is on a completely different foundation from CM, + ~~modern thinking~~ ^{entire QM} has a framework built up from a few postulates

Addition :

$$v, w \in V \Rightarrow v + w \in V$$

Examples of complex vector spaces :

① n -tuple

$$a = \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_n \end{pmatrix}$$

$$\Rightarrow \alpha a = \begin{pmatrix} \alpha a_1 \\ \alpha a_2 \\ \vdots \\ \alpha a_n \end{pmatrix}$$

(mult criteria obviously
met ~~if~~ if mult
works this way)

$$a + b = \begin{pmatrix} a_1 + b_1 \\ a_2 + b_2 \\ \vdots \\ a_n + b_n \end{pmatrix}$$

(Similar for addition
criterion)

- We can define linear independence :

b is linearly independent of $\{a_k\}_{k=1 \dots m}$,
if there is no choice of ~~any~~ α_k

$$\text{s.t. } b = \sum_{k=1}^m \alpha_k a_k$$

(Impossible to construct b from linear combos
of a_k)

- In this representation of the n -dimensional vector space, each

$$e_k = \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 1 \\ \vdots \\ 0 \end{pmatrix} \leftarrow k\text{th entry}$$

~~is~~ is linearly independent from set of all the others.

- Also every vector in space can be written

$$a = \sum_{k=1}^n a_k e_k$$

~~MA set of vectors~~

- $\Rightarrow e_k$ form a "basis" for V .

- Also it's a "minimal" basis, ~~n~~ⁿ dimensional vector space requires a basis of at least ~~n~~ⁿ elements - with ~~n~~ⁿ-element ones being "minimal"

- Many possible choices of minimal basis for V , all containing n elements

Example 2 of Complex vector space:

Space of complex-valued functions $\psi(x)$

- $\alpha \psi(x) + \psi_1(x) + \psi_2(x)$ still all within space, so criteria are met ✓
- Can think of $\psi(x)$ as ∞ -length n-tuple:

$$\begin{pmatrix} \psi(x_1) \\ \psi(x_2) \\ \vdots \end{pmatrix}$$

- For space of functions with same B.C.s,

Sturm-Liouville tells us that ~~eigen~~ eigenfunctions of any positive, self-adjoint operator form

complete basis for that space of fns.

- E.g. ~~$\psi \rightarrow 0$~~ $\psi \rightarrow 0$ @ $\pm \infty$. ~~$\{\psi_i(x)\}$~~ $\{\psi_i(x)\}$ eigenfunctions of H form complete basis.
- * - New step here is to consider that ~~this~~

~~we~~ we are ~~not~~ talking about a vector space, complex \mathbb{V} to be more precise.

Next def: Positive Norm:

- First: Inner product:

↗

Takes vectors v_1, v_2 to a complex #: \mathbb{C}

$$\langle v_1, v_2 \rangle = \langle v_1, v_2 \rangle = \alpha \in \mathbb{C}$$

such that $\langle \cdot, \cdot \rangle$ is linear in 2nd arg:

$$\langle v_1, \alpha v_2 + \beta v_3 \rangle = \alpha \langle v_1, v_2 \rangle + \beta \langle v_1, v_3 \rangle$$

+ "antilinear" in 1st arg:

$$\langle \alpha v_1 + \beta v_2, v_3 \rangle = \alpha^* \langle v_1, v_3 \rangle + \beta^* \langle v_2, v_3 \rangle$$

where α^* is complex conj. of α .

- We say that the inner product is a positive norm if

$\langle v, v \rangle$ is real + positive for any $v \in V$ s.t. $v \neq 0$

- Both of our vector space examples have natural inner products with positive norm:

~~Let's think: you define the inner product and the norm when you define Hilbert space. Maybe there are more than 1 option?~~

- For ^{vector} ~~inner~~ space of ~~real~~ n -tuples:

$$N(a, b) = \sum_{i=1}^n a_i^* b_i$$

$$\Rightarrow N(a, a) = \sum_{i=1}^n |a_i|^2 > 0 \quad \text{if } a \neq 0 \quad \checkmark$$

- For a vector space of functions:

$$N(\varphi_1, \varphi_2) = \int dx \varphi_1^*(x) \varphi_2(x)$$

$$\Rightarrow N(\varphi_1, \varphi_1) = \int dx |\varphi_1(x)|^2 > 0 \quad \text{if } \varphi_1 \neq 0$$

* - $\psi(t)$ describes motion of vector throughout \mathcal{H}

- We have now defined all terms in Postulate 1.

- A few more concepts are helpful.

- Vector is normalized :A

$$N(\psi, \psi) = 1$$

a) n -tuple: ~~(or n -vector model)~~

$$\sum_i |a_i|^2 = 1$$

b) function ~~(or function space model)~~

$$\int dx |\psi(x)|^2 = 1$$

- In general v can be "normalized" as follows; assuming $v \neq 0$

$$v_1 = \frac{1}{\sqrt{N(v, v)}} v$$

- So we have generalized the concept of unit vectors. Don't usually think of $f(x)$ as unit vector, but if normalized it really ~~is~~ ^{can} be a unit vector for ~~some~~ ^{some corresponding} Hilbert spaces

- ^{Normalized} $\psi_i(x)$ are unit vectors in Hilbert space of $f_{ns} \rightarrow 0$ at ∞

- Now that we have generalized concept of unit vectors, let's also generalize orthogonality

- Orthogonal if $N(v_1, v_2) = 0$

- For n -tuple $\sum_{i=1}^n a_i^* b_i = 0$

- For functions

$$\int dx \psi_1^*(x) \psi_2(x) = 0$$

- It is convenient to choose basis for \mathcal{H} in which basis vectors are normalized & mutually orthogonal

- For basis $\{v_k\}$ of \mathcal{H} , $N(v_i, v_j) = \delta_{ij}$

δ_{ij} is Kronecker delta,

$$\delta_{ij} = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases}$$

- Such a basis is "orthonormal"

- In 3D real space

$$\hat{x} = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad \hat{y} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad \hat{z} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

form an orthonormal basis.

- Alternatively,

$$\hat{r} = (\sin\theta \cos\phi, \sin\theta \sin\phi, \cos\theta)$$

$$\hat{\theta} = (\cos\theta \cos\phi, \cos\theta \sin\phi, -\sin\theta)$$

$$\hat{\phi} = (-\sin\phi, \cos\phi, 0)$$

is also an orthonormal basis

- For our \mathcal{H} , it is ^{possible} 1 states of the particle which form the "axes". Just as for real 3D space, many choices exist for which axes to use.

- Once we have orthonormal basis $\{u_k\}$, we can get very explicit representation of inner-product:

Let a, b be any vectors in H . Then

$$a = \sum a_k u_k \quad b = \sum \beta_k u_k$$

$$\begin{aligned} \Rightarrow N(a, b) &= N\left(\sum_k a_k u_k, \sum_l \beta_l u_l\right) \\ &= \sum_{k, l} a_k^* \beta_l \overbrace{N(u_k, u_l)}^{\delta_{kl}} = \sum_k a_k^* \beta_k \end{aligned}$$

- Recall before we had for n -tuple

$$N(a, b) = \sum a_i^* b_i \text{ which looks same b/c unit vectors}$$

$\begin{pmatrix} 1 \\ 0 \\ \vdots \\ 0 \end{pmatrix}$ are one ~~the~~ good choice for orthonormal basis

vector space of

- But for n functions, we had

~~$$N(\psi_1, \psi_2) = \int dx \psi_1^*(x) \psi_2(x)$$~~

$$N(\psi_1, \psi_2) = \int dx \psi_1^*(x) \psi_2(x)$$

- That looks at first pretty different. But using ^{orthonormal} n basis to represent any $\psi(x)$, it really does behave same way as for n -type.
 \Rightarrow "Vector space" description seems ^{more} reasonable ~~than~~

- Finally "Dual Space"

- For fixed v_1 , $N(v_1, v_2)$ is a mapping from H to \mathbb{C} :
 \downarrow
denotes entire space

$$N(v_1, \cdot) : H \rightarrow \mathbb{C}$$

- (A complex # gets assigned to each vector in Hilbert space. I.e. mapped. Mapping changes with v_1 .)

- Set of such maps (for diff v_1) is called the "dual space" H^*

- We can think of $N(v_1, v_2)$ as describing action of vector v_1 on vector v_2 .

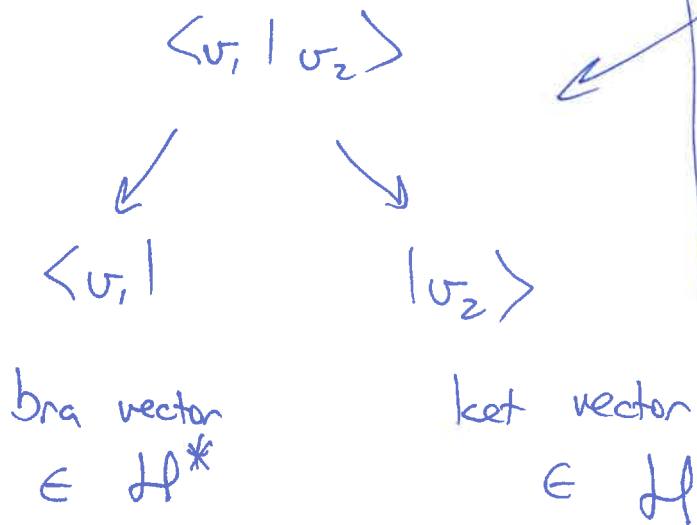
- Dirac expressed it in elegant notation:

$$N(v_1, v_2) = \langle v_1 | v_2 \rangle$$

"bracket of v_1 & v_2 "

- So we have 2 spaces. A space of vectors & a space of maps. Latter is dual space.

- Notation suggests we can pull bracket apart into pieces:



If $|u_1\rangle$ is $\psi_1(x)$
 $\langle u_1 |$ is not just $\psi_1^*(x)$. it is that $\times \int \rightarrow$ to make a ~~spatial~~ map

- Next: Results of measurements ... about to get less abstract.