

Time Evolution

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H is the time translation operator

$$T(dt) \equiv 1 - \frac{i dt}{\hbar} H$$

$$T(dt) |\psi(t)\rangle = |\psi(t+dt)\rangle$$

$$|\psi(t)\rangle - \frac{i}{\hbar} dt H |\psi(t)\rangle = |\psi(t+dt)\rangle$$

$$\frac{|\psi(t+dt)\rangle - |\psi(t)\rangle}{dt} = \frac{1}{i\hbar} \frac{dt}{dt} H |\psi(t)\rangle$$

$$\boxed{i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = H |\psi(t)\rangle} \quad \text{S.E.}$$

Time evolution operator $U(t, t_0)$

$$|\psi(t)\rangle = U(t, t_0) |\psi(t_0)\rangle$$

$$i\hbar \frac{\partial}{\partial t} U(t, t_0) |\psi(t_0)\rangle = H U(t, t_0) |\psi(t_0)\rangle$$

$$\boxed{i\hbar \frac{\partial U(t, t_0)}{\partial t} = H U(t, t_0)} \quad \rightarrow -i\hbar \frac{\partial U^\dagger}{\partial t} = U^\dagger H$$

Solution for H ind. of time

$$\underline{U(t, t_0) = e^{-\frac{i}{\hbar} H(t-t_0)}}$$

Two pictures (for \mathcal{O} ind. of time) \rightarrow Review

Schrodinger: $|\psi(t_0)\rangle \rightarrow |\psi(t)\rangle = U(t, t_0)|\psi(t_0)\rangle$
 $\mathcal{O} \rightarrow \mathcal{O}$

Heisenberg: $|\psi(t_0)\rangle \rightarrow |\psi(t_0)\rangle$
 $\mathcal{O}_H(t_0) \rightarrow \mathcal{O}_H(t)$

Insist that observables are the same

$t=t_0$: $\langle \psi(t_0) | \mathcal{O}_H(t_0) | \psi(t_0) \rangle = \langle \psi(t_0) | \mathcal{O} | \psi(t_0) \rangle$
 $\therefore \boxed{\mathcal{O}_H(t_0) = \mathcal{O}}$

$t=t$: $\langle \psi(t_0) | \mathcal{O}_H(t) | \psi(t_0) \rangle = \langle \psi(t) | \mathcal{O} | \psi(t) \rangle$
 $= \langle \psi(t_0) | U^\dagger(t, t_0) \mathcal{O} U(t, t_0) | \psi(t_0) \rangle$
 $\boxed{\mathcal{O}_H(t) = U^\dagger(t, t_0) \mathcal{O} U(t, t_0)}$

Equation of motion for $\mathcal{O}_H(t)$

$$\begin{aligned} i\hbar \frac{\partial \mathcal{O}_H(t)}{\partial t} &= i\hbar \frac{\partial}{\partial t} [U^\dagger \mathcal{O} U] \\ &= i\hbar \frac{\partial U^\dagger}{\partial t} \mathcal{O} U + U \mathcal{O} i\hbar \frac{\partial U}{\partial t} \\ &= -U^\dagger H \mathcal{O} U + U \mathcal{O} H U \\ &= -\underbrace{U^\dagger H U}_{H_H} \underbrace{U^\dagger \mathcal{O} U}_{\mathcal{O}_H} + \underbrace{U \mathcal{O} U^\dagger}_{\mathcal{O}_H} \underbrace{U H U^\dagger}_{H_H} \end{aligned}$$

$$\boxed{i\hbar \frac{\partial \mathcal{O}_H(t)}{\partial t} = [\mathcal{O}_H, H_H]}$$

For H ind. of $t \Rightarrow [U, H] = 0 \rightarrow$
 Heisenberg eq. of motion

$$\boxed{i\hbar \frac{\partial \mathcal{O}_H}{\partial t} = [\mathcal{O}_H, H]}$$

Interaction Picture

$$H = H_0 + V(t)$$

↑ independent of t

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H_0 has eigenstates: $H_0 |n\rangle = \hbar \omega_n |n\rangle$

Time evolution from H_0 alone:

$$i\hbar \frac{\partial U_0(t, t_0)}{\partial t} = H_0 U_0(t, t_0) \rightarrow U_0(t, t_0) = e^{-\frac{iH_0(t-t_0)}{\hbar}}$$

$V(t)$ makes transitions between $|n\rangle$

$$|\psi(t)\rangle = \sum_n \underbrace{c_n(t)}_{\text{transitions between states}} \underbrace{e^{-i\omega_n(t-t_0)}}_{\text{time evolution of state } |n\rangle \text{ due to } H_0} |n\rangle$$

transitions
between
states

time evolution of
state $|n\rangle$ due to H_0

Focus on the time evolution due to interaction

$$|\psi(t)\rangle_I \equiv U_0^\dagger(t, t_0) |\psi(t)\rangle$$

undoes the evolution due to H_0

$$= \sum_n c_n(t) |n\rangle$$

$$= U_0^\dagger \sum_n c_n e^{-i\omega_n(t-t_0)} |n\rangle$$

$$= \sum_n c_n e^{-i\omega_n(t-t_0)} e^{i\omega_n(t-t_0)} |n\rangle$$

$$= \sum_n c_n(t) |n\rangle$$

wavefunction in the
interaction representation

$$\text{At } t=t_0: |\psi(t_0)\rangle_I = \underbrace{U_0^\dagger(t_0, t_0)}_1 |\psi(t_0)\rangle$$

$$= |\psi(t_0)\rangle$$

same in Schr.
and Int.
rep.

Observables in the Interaction Representation

Require
Want

$$\underbrace{\langle \psi(t) |}_{\text{I}} \underbrace{\theta}_{\text{I}} \underbrace{|\psi(t)\rangle}_{\text{I}} = \langle \psi(t) | \theta | \psi(t) \rangle$$

$$\langle \psi(t) | U_0 \quad U_0^\dagger | \psi(t) \rangle$$

$$\therefore U_0 \theta_{\text{I}} U_0^\dagger = \theta$$

Sch. rep.

$$\theta_{\text{I}}(t) = U_0^\dagger(t, t_0) \theta U_0(t, t_0)$$

Evolution of θ_{I} (for θ ind. of time)

and H_0 ind. of time

$$i\hbar \frac{\partial \theta_{\text{I}}}{\partial t} = i\hbar \frac{\partial}{\partial t} U_0^\dagger \theta U_0$$

$$= \underbrace{i\hbar \frac{\partial U_0^\dagger}{\partial t} \theta U_0}_{-U_0^\dagger H_0} + U_0^\dagger \theta \underbrace{i\hbar \frac{\partial U_0}{\partial t}}_{H_0 U_0}$$

$$= \underbrace{U_0^\dagger \theta U_0}_{\theta_{\text{I}}} H_0 - H_0 \underbrace{U_0^\dagger \theta U_0}_{\theta_{\text{I}}}$$

$$i\hbar \frac{\partial \theta_{\text{I}}}{\partial t} = [\theta_{\text{I}}, H_0]$$

not for av. values

$$i\hbar \frac{\partial \langle \theta \rangle}{\partial t} = i\hbar \frac{\partial}{\partial t} \langle \psi(t) | \theta_{\text{I}}(t) | \psi(t) \rangle_{\text{I}}$$

$$= i\hbar \frac{\partial}{\partial t} \langle \psi(t) | \theta_{\text{I}}(t) | \psi(t) \rangle_{\text{I}} + i\hbar \langle \psi(t) | \theta_{\text{I}}(t) \frac{\partial}{\partial t} | \psi(t) \rangle_{\text{I}} + \langle \psi(t) | i\hbar \frac{\partial \theta_{\text{I}}}{\partial t} | \psi(t) \rangle_{\text{I}}$$

$$i\hbar \frac{\partial U_0}{\partial t} = H_0 U_0$$

$$-i\hbar \frac{\partial U_0^\dagger}{\partial t} = U_0^\dagger H_0$$

not so useful

Time evolution of w.f. in interaction rep. 5

$$\underbrace{|\psi(t)\rangle_I}_{U_0^\dagger |\psi(t)\rangle} = U_I(t, t_0) \underbrace{|\psi(t_0)\rangle_I}_{|\psi(t_0)\rangle}$$

$$U |\psi(t_0)\rangle$$

$$\therefore \boxed{U_0^\dagger U = U_I} \rightarrow \boxed{U = U_0 U_I}$$

$$\begin{aligned} i\hbar \frac{\partial U_I}{\partial t} &= i\hbar \frac{\partial}{\partial t} U_0^\dagger U \\ &= \underbrace{i\hbar \frac{\partial U_0^\dagger}{\partial t} U}_{-U_0^\dagger H_0} + U_0^\dagger \underbrace{i\hbar \frac{\partial U}{\partial t}}_{HU} \end{aligned}$$

$$i\hbar \frac{\partial U}{\partial t} = HU$$

$$-i\hbar \frac{\partial U^\dagger}{\partial t} = U^\dagger H$$

$$= U_0^\dagger \underbrace{(H - H_0)}_V U_0 U_0^\dagger U$$

$$\underbrace{\quad}_V \underbrace{\quad}_{U_I}$$

$$\boxed{i\hbar \frac{\partial U_I}{\partial t} = V_I U_I} \xrightarrow{\text{exact}} \boxed{i\hbar \frac{\partial |\psi(t)\rangle_I}{\partial t} = V_I(t) |\psi(t)\rangle_I}$$

In terms of $c_n(t)$: $|\psi(t)\rangle_I = \sum_n c_n(t) |n\rangle$

$$V_I = U_0^\dagger V U_0$$

$$i\hbar \frac{\partial}{\partial t} \sum_m c_m |m\rangle = V_I \sum_m c_m |m\rangle$$

$\langle n| \rightarrow$

$$i\hbar \dot{c}_n = \sum_m \langle n| U_0^\dagger V U_0 |m\rangle c_m$$

$$\boxed{i\hbar \dot{c}_n = \sum_m c_m V_{nm} e^{i\omega_{nm}t}} \quad \text{exact}$$

Time Dependent Perturbation Theory

⑥
quick first
order

$$H = H_0 + V(t)$$

\uparrow
 $H_0 |n\rangle = \hbar \omega_n |n\rangle$ causes transitions
"small"

Exact Schrodinger equation: (interaction picture)

$$i\hbar \dot{C}_n = \sum_m C_m V_{nm}(t) e^{-i\omega_{nm}t}$$

First order perturbation theory

$$t=0: C_n(0) = \delta_{ni}$$

$$t=t: C_i(t) \approx 1 \text{ and } C_{f \neq i}(t) \approx 0$$

Plug into Schrodinger equation

$$n=i: i\hbar \dot{C}_i \approx 0 \leadsto C_i(t) \approx 1$$

$$n=f \neq i: i\hbar \dot{C}_f \approx \sum_i C_i V_{fi}(t) e^{-i\omega_{if}t} \approx 1$$

$$= V_{fi}(t) e^{-i\omega_{if}t}$$

$$C_f(t) - C_f(0) \approx -\frac{i}{\hbar} \int_0^t V_{fi}(t') e^{-i\omega_{if}t'} dt'$$

$$C_f^{(1)}(t) = -\frac{i}{\hbar} \int_0^t V_{fi}(t') e^{-i\omega_{if}t'} dt'$$

1st order in the interaction

$$C_f(t) = C_f^{(1)}(t) + C_f^{(2)}(t) + \dots$$