

Method of Partial Waves

Assume a central potential $V(r)$.

\vec{L} is a constant of the motion.

There are eigenstates common to H, L^2, L_z .

the wave functions of these states are partial waves.

$$H_0 = \frac{p^2}{2\mu} \text{ e.s. for a free particle.}$$

a complete set of commuting observables would be

$$H_0, p_x, p_y, p_z$$

this is like being in a central potential w/ $V=0$.

H, L^2, L_z form a complete set of commuting observables also.

There are 2 separate bases: $[\vec{p}, \vec{L}] \neq 0$

For example spin-0 particle: $\hat{p}|p\rangle = p|p\rangle$

$$H_0|p\rangle = \frac{p^2}{2\mu}|p\rangle$$

$|p| = \sqrt{2\mu E}$ an infinite # of kets satisfy this
each energy E is infinitely degenerate.

$$\langle r|p\rangle = \left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{i\vec{p}\cdot\vec{r}/\hbar} \quad \text{w.f. of momentum}$$

ket. (plane waves)

$$H_0 |k\rangle = \frac{\hbar^2 k^2}{2m} |k\rangle$$

$$p|k\rangle = \hbar k |k\rangle$$

$$\langle k|k'\rangle = \delta(k-k')$$

$$\int d^3k |k\rangle \langle k| = 1$$

$$\langle r|k\rangle = \left(\frac{1}{2\pi}\right)^{3/2} e^{i\vec{k}\cdot\vec{r}}$$

The stationary states with well-defined angular momentum

$$\psi_{k,\ell,m}^{(0)} = \sqrt{\frac{2k^2}{\pi}} j_\ell(kr) Y_\ell^m(\theta, \varphi)$$

$j_\ell(\rho)$ is a spherical Bessel function

$$j_\ell(\rho) = (-1)^\ell \rho^\ell \left(\frac{1}{\rho} \frac{\partial}{\partial \rho} \right)^\ell \left(\frac{\sin \rho}{\rho} \right)$$

The eigenvalues of

$$H_0, L^2, \text{ and } L_z \text{ are } \frac{\hbar^2 k^2}{2m}, \ell(\ell+1)\hbar^2, m\hbar$$

these free spherical waves are orthonormal in the sense

$$\langle \psi_{k,\ell,m}^{(0)} | \psi_{k',\ell',m'}^{(0)} \rangle = \frac{2}{\pi} k k' \int_0^\infty j_\ell(kr) j_{\ell'}(k'r) r^2 dr \int d\Omega Y_\ell^m Y_{\ell'}^{m'*}$$

$$= \delta(k-k') \delta_{\ell\ell'} \delta_{mm'}$$

$$\int_0^\infty dk \sum_{\ell=0}^\infty \sum_{m=-\ell}^\ell | \psi_{k,\ell,m}^{(0)} \rangle \langle \psi_{k,\ell,m}^{(0)} | = 1$$