

## Postulates of QM, Part III

Postulate 2 covers case where system is prepared in an eigenstate of measurement operator  $\hat{O}$ . But what if it is not in an eigenstate of  $\hat{O}$ ?

### Postulate 3:

The only possible results of a measurement of  $\hat{O}$  is one of the eigenvalues of  $\hat{O}$ .

### Postulate 4:

The probability that, in a normalized state  $|\psi\rangle$ , we measure a particular eigenvalue  $\lambda_i$ , is given by

$$P_i = |\langle \psi_i | \psi \rangle|^2, \text{ where } \hat{O}|\psi_i\rangle = \lambda_i |\psi_i\rangle$$

### Postulate 5:

The state of the system immediately after the measurement of the value  $\lambda_i$  is  $|\psi_i\rangle$

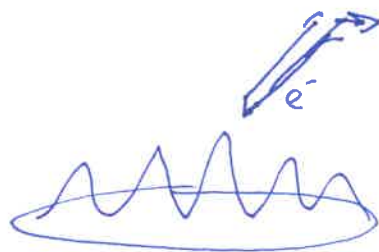
- In other words, the measurement of  $\hat{O}$  is a catastrophic operation on  $|\psi\rangle$ . "Wave function collapse."
- We can expand  $|\psi\rangle$  in basis of eigenstates of  $\hat{O}$ :

$$|\psi\rangle = \sum_i \alpha_i |\psi_i\rangle, \text{ where } \alpha_i = \langle \psi_i | \psi \rangle$$

From here, take  $|\psi\rangle$  to be the basis to be normalized &  $|\psi_i\rangle$  to be orthonormal,

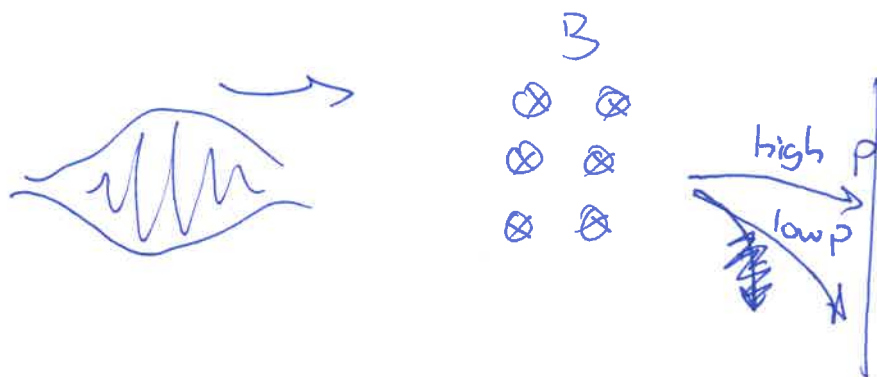
$$1 = \langle \psi | \psi \rangle = \sum_i |\alpha_i|^2$$

- Postulates 3-5 say that the result of measuring  $\hat{O}$  is to force  $|\psi\rangle$  into one of the eigenstates  $|\psi_i\rangle$ , and that occurs with a probability  $|\alpha_i|^2$
- At first this makes some sense. Measurements of microscopic things generally are very invasive to those delicate systems.
- Consider measuring the position of an electron in a "quantum corral", using a scanning tunneling microscope



The electron gets removed in order to measure its position. Certainly seems like a ~~small~~<sup>big</sup> change <sup>in  $|\psi\rangle$</sup>  is necessary there to measure.

- On measuring the momentum of an electron using deflection in a  $\vec{B}$  field.



Position it strikes on screen tells you momentum, but again we have done something catastrophic to the state.

- In some sense, the measurement postulates reflect the fact that there is no way to measure the properties ~~of~~ of an atomic-scale system without changing them in some essential way

- We will return to this question in a moment, to look at it from an updated perspective.
- This idea of wave function collapse, is built into the Copenhagen interpretation of QM.
- Regardless of whether we believe that interpretation of what is <sup>really</sup> going on when a measurement is made, Postulates 3-5 are agreed upon.
- And we need to distinguish between a given measurement (always one of the  $\lambda_i$ ) and the average result of many repeated measurements. The latter is called the expectation value  $\langle O \rangle$  & it is what we get for averaging over many experiments where the system was prepared in the same way at the beginning of each.

Average value of many measurements:

5

$$\langle \psi | \hat{O} | \psi \rangle = \sum_i \lambda_i P_i$$

$$= \sum_i \lambda_i |\alpha_i|^2$$

$$= \sum_i \lambda_i \cancel{|\alpha_i|^2} |\langle \psi_i | \psi \rangle|^2$$

$$= \sum_i \lambda_i \langle \psi | \psi_i \rangle \langle \psi_i | \psi \rangle$$

$$= \sum_i \langle \psi | \psi_i \rangle \langle \psi_i | \hat{O} | \psi \rangle$$

(Now using the completeness of  $|\psi_i\rangle$ :  
 $\mathbb{I} = \sum_i |\psi_i\rangle \langle \psi_i|$ )

$$= \langle \psi | \hat{O} | \psi \rangle \text{ or } \cancel{\langle \psi | \hat{O} | \psi \rangle} \\ \langle \psi | \hat{O} | \psi \rangle$$



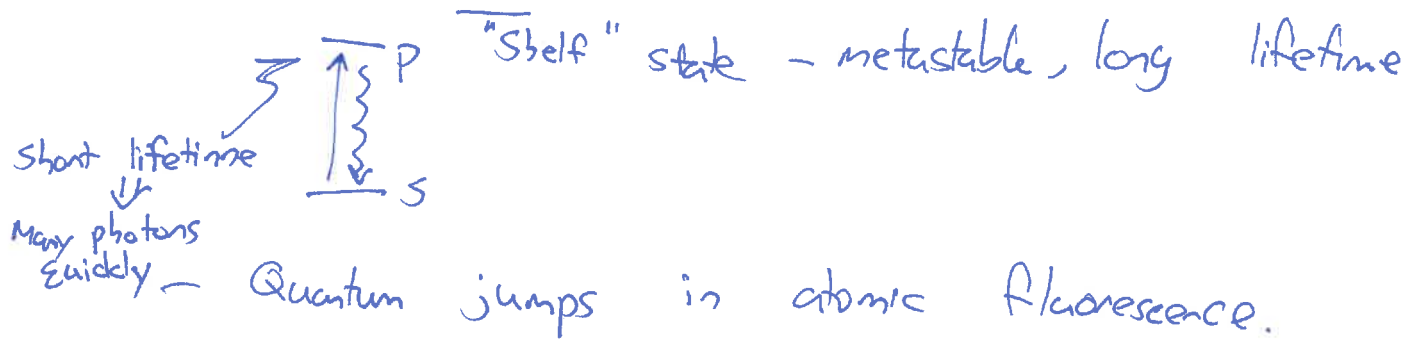
$\langle \hat{O} \rangle$ , the "expectation value of  $\hat{O}$  in  $\psi$ "  
 $= \langle \psi | \hat{O} | \psi \rangle$

Another quantity  $\langle \psi | \hat{O} | \chi \rangle$  is "the matrix element of  $\hat{O}$  between  $\psi$  &  $\chi$ "

# QND Measurements

Quantum non-demolition

## Ex. 1 Atomic level structure (AMO state readout technique)



- We detect single atom state without changing it, or at least its eigenvalue of an energy measurement.
- If electron was in shelf state, we seemingly didn't even disturb its phase.

## Ex. 2

### Readout of electron spin in Penning trap



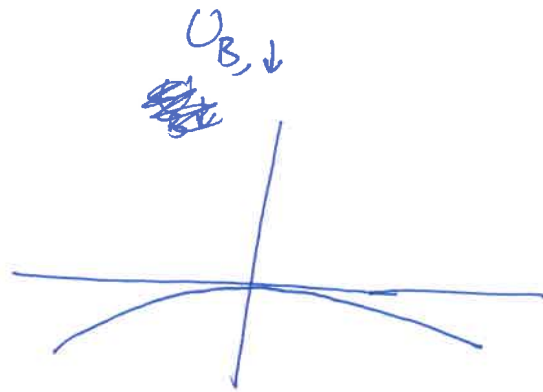
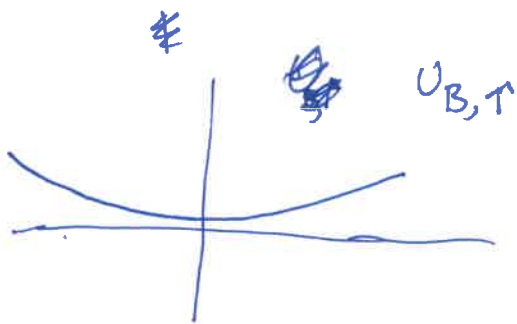
Add to this

$$B(\vec{z}) = \beta z^2, \quad U_B = -\vec{\mu} \cdot \vec{B}$$

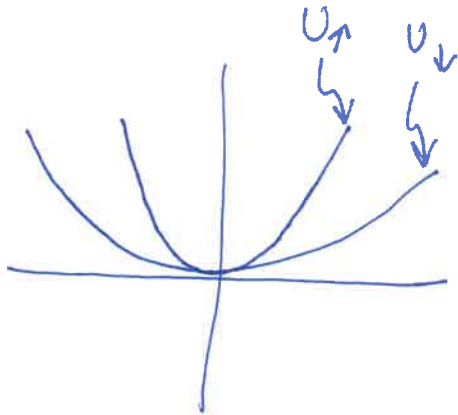
Then if spin is  $\uparrow$ , the electron

sees a different quadratic potential than if

it is down:



Now overall potential looks like



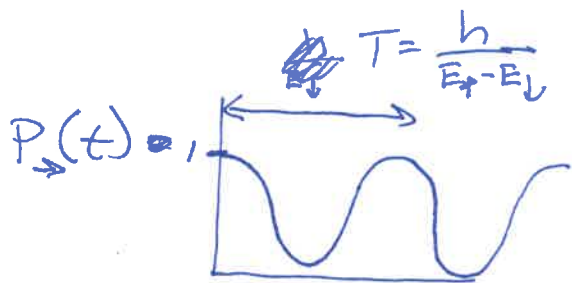
If we detect ~~res~~ harmonic oscillations of electron in trap, without destroying it or changing its spin orientation, we can read out its quantum state!

- These modern AMO measurement techniques are much more gentle than what <sup>(most?)</sup> people thought was possible in the early days of quantum mechanics.

- And yet <sup>the picture of</sup> wave function collapse still seems to give the right results.

- Signatures for electron spin.

① Say ~~the state~~  $|\psi(t=0)\rangle = |\rightarrow\rangle = \frac{1}{\sqrt{2}}|\uparrow\rangle + \frac{1}{\sqrt{2}}|\downarrow\rangle$   
 $= |\rightarrow\rangle_x$



But if at  $t = t'$ , we measure the spin along the  $z$  axis, we get  $P_{\rightarrow}(t) = \frac{1}{2}$  for all subsequent times. ~~(That's true whether we found  $|\uparrow\rangle$  or  $|\downarrow\rangle$  in our measurement)~~

(If we measured  $|\uparrow\rangle$ , we prepared the system in <sup>that</sup> energy eigenstate, so there is no time dependence for any subsequent measurement)



And ~~the~~  $|\uparrow\rangle = \frac{1}{\sqrt{2}} |\rightarrow\rangle + \frac{1}{\sqrt{2}} |\leftarrow\rangle$ ,

thus the  $P_{\rightarrow} = \frac{1}{2}$ .

② At all subsequent times after  $t'$ ,  
we measured the spin along  $z$ , we would  
get with 100% confidence the same <sup>result</sup> ~~thing~~ we  
got at  $t'$ .

⇒ Evidently even the act of "looking" as  
gently as we can really does fundamentally disturb  
the system.

~~Many worlds~~  
~~What if~~