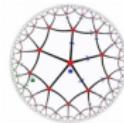


Entanglement in QM and QFT

5/5 -Holographic entanglement entropy

Pablo Bueno



It from Qubit

Simons Collaboration on
Quantum Fields, Gravity and Information



— II School of Holography and Entanglement Entropy —
December, 2020

OUTLINE

- ① THE HOLOGRAPHIC PRINCIPLE AND AdS/CFT
- ② THE RYU-TAKAYANAGI PRESCRIPTION
- ③ CORRECTIONS TO THE RT FORMULA
- ④ SOME EXPLICIT CALCULATIONS
- ⑤ GRAVITY FROM ENTANGLEMENT

SOME REFERENCES

- A nice review of the holographic principle is <https://arxiv.org/pdf/hep-th/0203101.pdf>. The original insight is due to 't Hooft.
- The original AdS/CFT paper by Maldacena is <https://arxiv.org/pdf/hep-th/9711200.pdf> (most cited paper in the history of physics). Seminal papers by Witten: <https://arxiv.org/pdf/hep-th/9802150.pdf> and Gubser, Klebanov and Polyakov: <https://arxiv.org/pdf/hep-th/9802109.pdf>.
- The subject of holographic entanglement entropy started with the seminal papers of Ryu and Takayanagi <https://arxiv.org/pdf/hep-th/0603001.pdf> and <https://arxiv.org/pdf/hep-th/0605073.pdf>, where they presented their prescription. The covariant version, due to Hubeny, Rangamani and Takayanagi was proposed in <https://arxiv.org/pdf/0705.0016.pdf>.
- Other relevant papers on the subject including various generalizations are: <https://arxiv.org/pdf/1310.5713.pdf>, <https://arxiv.org/pdf/1101.5813.pdf>, <https://arxiv.org/pdf/1307.2892.pdf>, <https://arxiv.org/pdf/1304.4926.pdf>, <https://arxiv.org/pdf/1102.0440.pdf>. There are many more...
- Various reviews on the subject of holographic entanglement entropy can be found in <https://arxiv.org/pdf/0905.0932.pdf>, <https://arxiv.org/pdf/1609.01287.pdf>, <http://www2.yukawa.kyoto-u.ac.jp/~tadashi.takayanagi/CERNEE.pdf>, <https://arxiv.org/pdf/1609.00026.pdf>
- The derivation of the linearized gravity equations from the entanglement first law is from <https://arxiv.org/pdf/1312.7856.pdf>.

The holographic principle and AdS/CFT

Black hole entropy and holographic principle

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The idea that all things happening within a volume A should be describable in terms of things happening at its boundary ∂A is called the **holographic principle**.

AdS/CFT correspondence

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When the CFT is in a strongly-coupled regime and the number of components of the fields is sufficiently large, the equivalent quantum gravity theory is in the regime where stringy and quantum effects are small and physics is described by (semi)classical Einstein gravity:

$$I_{\text{grav.}} = \frac{1}{16\pi G} \int d^{d+1}x \sqrt{-g} \left[\frac{d(d-1)}{L^2} + R + \dots \right]$$

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We can use the tools of GR to learn quantum things!

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(For today, we mostly restrict our discussion to the vacuum state.)

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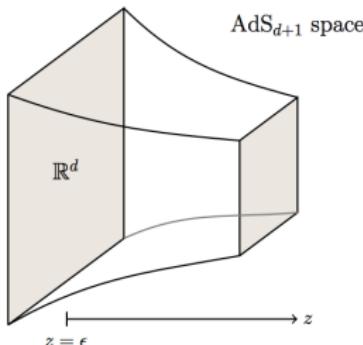
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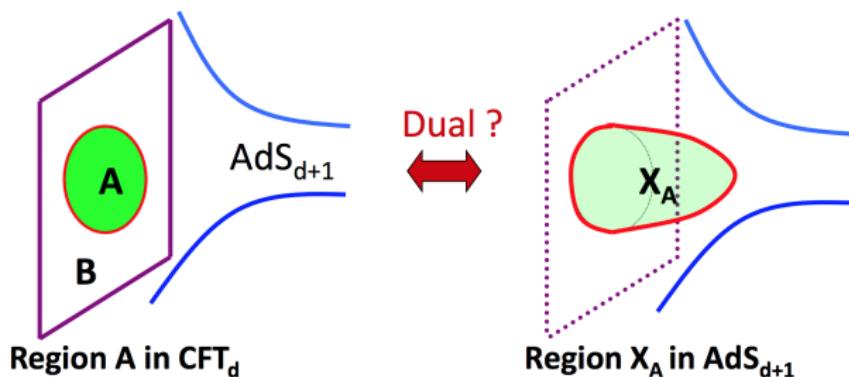
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- Greater values of z probe the “bulk region” and corresponds to low energies. Features which distinguish different states become apparent in the geometry (*e.g.*, by the presence of a horizon).



The Ryu-Takayanagi prescription

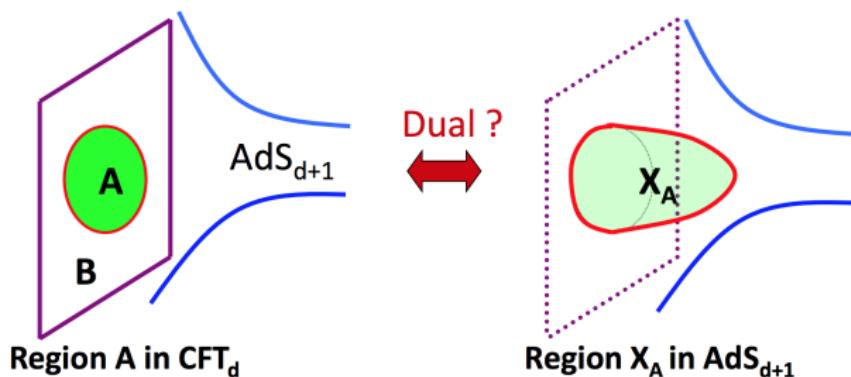
Holographic entanglement entropy

One may wonder: is there a region in AdS_{d+1} which encodes the information corresponding to a given region A in the CFT_d ?



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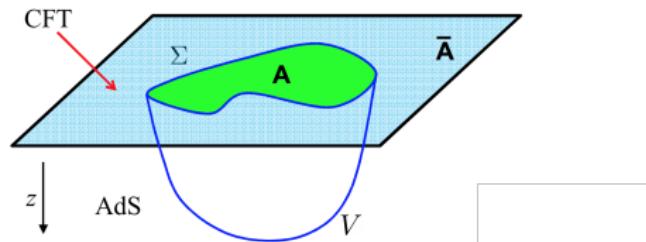
This question was partially answered by Ryu and Takayanagi with their sensational proposal for the computation of entanglement entropy of CFTs with a holographic dual...

Holographic entanglement entropy

The **Ryu-Takayanagi prescription** tells us that the EE for a region A in the CFT [dual to Einstein (super)gravity in the bulk] can be computed as

$$S_{\text{HEE}}(A) = \min_{V \sim A} \left[\frac{\text{Area}(V)}{4G} \right]$$

We need to minimize amongst all bulk surfaces V whose boundary coincides with the one of A

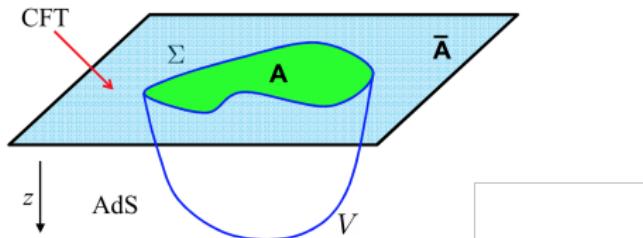


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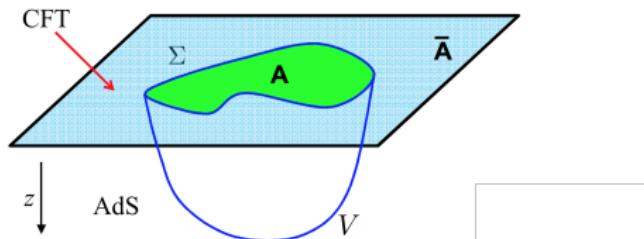
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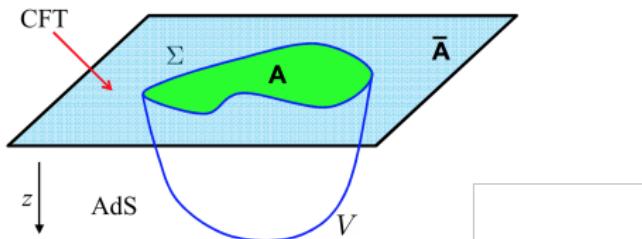
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Nice new entry in the holographic dictionary! Very powerful computationally. Remarkably similar to Bekenstein-Hawking formula for black hole entropy!

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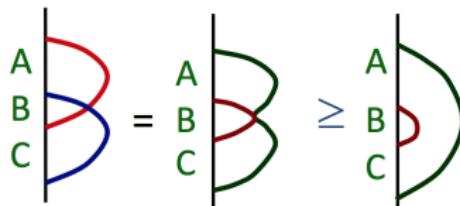
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- However, the formula passed numerous tests, including reproducing all expected properties of EE such as SSA, structure of universal terms, etc. in various dimensions.
- Using the AdS/CFT dictionary, the replica trick and the interpretation of EE as the $n \rightarrow 1$ limit of Rényi entropies, Lewkowycz and Maldacena finally rigorously proved it.

Inequalities

Strong subadditivity can be seen to hold rather easily.

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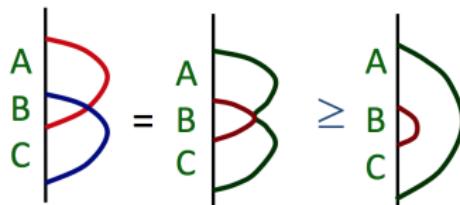
Strong subadditivity can be seen to hold rather easily. For instance, for a CFT₂, the RT surfaces for intervals are simply semicircles, and the EE is proportional to their length so:



$$S_{\text{EE}}(AB) + S_{\text{EE}}(BC) \geq S_{\text{EE}}(ABC) + S_{\text{EE}}(B)$$

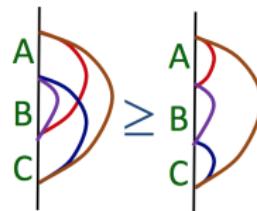
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Holographic entanglement entropy satisfies an additional property called “monogamy”, not satisfied by general CFTs: $I_3(A, B, C) \equiv I(A, B) + I(A, C) - I(A, BC) \leq 0$



$$S_{EE}(AB) + S_{EE}(BC) + S_{EE}(AC) \geq S_{EE}(A) + S_{EE}(B) + S_{EE}(C) + S_{EE}(ABC)$$

Corrections to the RT formula

Higher-curvature corrections

When the gravitational action contains higher-curvature (“stringy”) corrections to the Einstein-Hilbert action, the RT gets modified (similarly to Bekenstein-Hawking’s formula for black holes).

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- $f(R)$ gravity: $\mathcal{L} = \frac{d(d-1)}{L^2} + R + f(R)$

$$S_{\text{EE}}^{f(R)} = \frac{\mathcal{A}(V)}{4G} + \frac{1}{4G} \int_V d^{d-1}y \sqrt{h} f'(R)$$

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$$S_{\text{EE}}^{\text{Riem}^2} = \frac{\mathcal{A}(V)}{4G} + \frac{L^2}{4G} \int_V d^{d-1}y \sqrt{h} \left[2\alpha_1 R + \alpha_2 (R_a^a - \frac{1}{2} K^a K_a) + 2\alpha_3 (R_{ab}^{ab} - K_{aij} K^{aij}) \right]$$

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- Lovelock gravity: $\mathcal{L} = \frac{d(d-1)}{L^2} + R + \sum_{n=2}^{\lfloor \frac{d+1}{2} \rfloor} \lambda_n L^{2(n-1)} \mathcal{X}_{2n}(R)$, where $\lfloor x \rfloor$ integer part of x and the order- n invariants: $\mathcal{X}_{2n}(R) \equiv \frac{1}{2^n} \delta_{\nu_1 \nu_2 \dots \nu_{2n-1} \nu_{2n}}^{\mu_1 \mu_2 \dots \mu_{2n-1} \mu_{2n}} R^{\nu_1 \nu_2}_{\mu_1 \mu_2} \dots R^{\nu_{2n-1}}_{\mu_{2n-1} \mu_{2n}}$.

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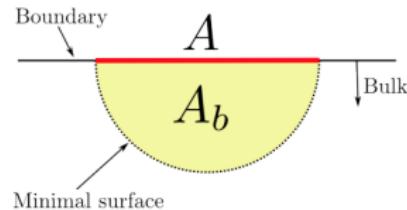
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General formula exists for arbitrary theories, but there are some obstructions to its application beyond the present cases.

Quantum corrections

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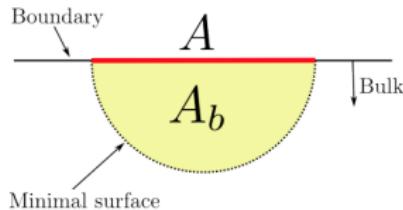
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The leading correction (order zero in $1/G$) comes from the entanglement entropy of bulk modes subject to the bipartitioning across the RT surface.

Some explicit calculations

Interval in a CFT₂

In Poincaré coordinates, $ds_{\text{AdS}_3}^2 = L_\star^2/z^2[-dt^2 + dz^2 + dx^2]$, we consider the EE of an interval $x \in [-\ell/2, \ell/2]$ in a time slice.

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$$ds_h^2 = \frac{L_\star^2}{Z(x)^2}[1 + Z'(x)^2]dx^2 \Rightarrow \sqrt{h} = \frac{L_\star}{Z(x)} \sqrt{1 + Z'(x)^2}$$

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where we introduced a UV cutoff at $Z = \delta$. This is exactly the result expected for a CFT₂ with central charge $c = 3L_\star/(2G)$. The value of c for holographic Einstein gravity can be computed via alternative methods and perfectly matches this result!

Spheres in a CFT_d

We write now the AdS_{d+1} metric as: $ds_{\text{AdS}_{d+1}}^2 = \frac{L^2}{z^2}[-dt^2 + dz^2 + dr^2 + r^2 d\Omega_{d-2}^2]$.

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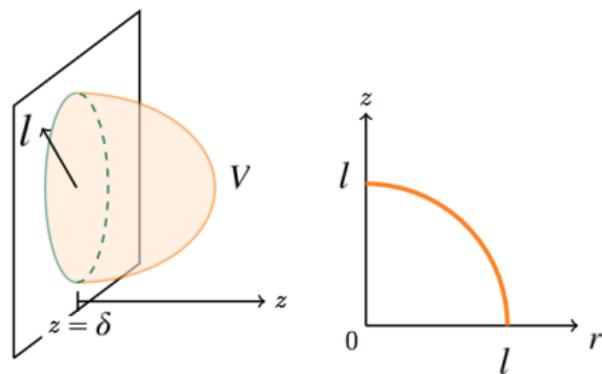
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where I already integrated over the angular directions. The Euler-Lagrange equation that we need to solve in order to find the minimal surface reads now:

$$rZZ'' + (d-2)ZZ'(1 + Z'^2) + (d-1)r(1 + Z'^2) = 0$$

Spheres in a CFT_d

The solution for this differential equation satisfying the boundary condition $Z(\ell) = 0$ is $Z = \sqrt{\ell^2 - r^2}$.



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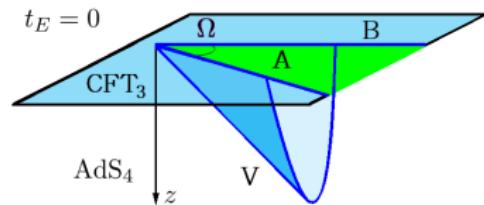
where

$$a^* = \frac{\pi^{\frac{d-2}{2}}}{8\Gamma\left[\frac{d}{2}\right]} \frac{L_\star^{d-1}}{G}.$$

This is the expression expected for a spherical entangling surface in a CFT_d. The coefficient a^* can also be computed using alternative methods and it exactly matches the one obtained here.

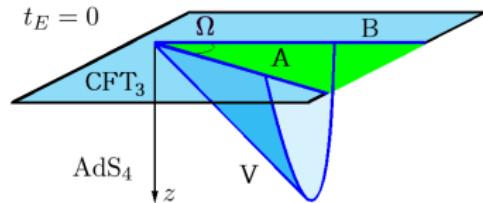
Corners in a CFT₃

What about a corner region?



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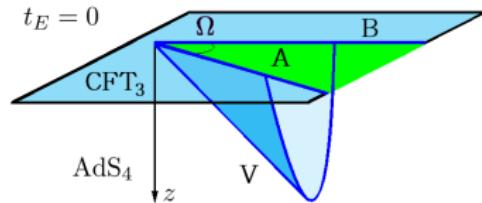


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where the holographic corner function is given by

$$a_E(\Omega) = \frac{L_\star^2}{2G} \int_0^\infty dy \left[1 - \sqrt{\frac{1 + h_0(\Omega)^2(1 + y^2)}{2 + h_0(\Omega)^2(1 + y^2)}} \right],$$

$$\Omega(h_0) = \int_0^{h_0} dh \frac{2h^2 \sqrt{1 + h_0^2}}{\sqrt{1 + h^2} \sqrt{(h_0^2 - h^2)(h_0^2 + (1 + h_0^2)h^2)}}.$$

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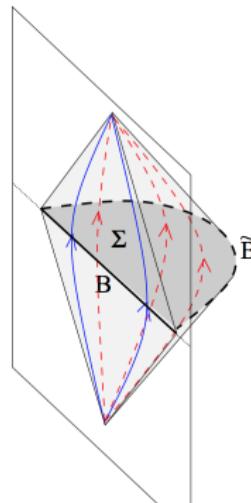
- The modular Hamiltonian is a complicated and non-local object in general. However, for a general CFT $_d$ with stress tensor T_{ab} , for a spherical ball $B(R, \vec{x}_0)$ of radius R centered at \vec{x}_0 , it has the following useful representation

$$\langle H_{\text{ball}} \rangle = 2\pi \int_{B(R, \vec{x}_0)} d^{d-1}x \left[\frac{R^2 - |\vec{x} - \vec{x}_0|^2}{2R} \right] \langle T_{tt} \rangle$$

Gravity from entanglement

- Holographically, entanglement entropy is computed via the RT prescription, and the surface corresponding to a ball-shaped region B is given by $\tilde{B} = \{t = t_0, |x^i - x_0^i|^2 + z^2 = R^2\}$. We have then for holographic theories

$$\delta S_{\text{EE}}(B) = \delta \langle H_B \rangle \stackrel{\text{AdS/CFT}}{\implies} \frac{\delta \mathcal{A}(\tilde{B})}{4G} = \delta \langle H_B^{\text{grav.}} \rangle$$



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- On the other hand, using the CHM map, we know that the EE for a ball-shaped region can be obtained as the thermal entropy of the CFT in hyperbolic space, $S_{\text{EE}}(B) = S_{\text{therm.}}(\mathbb{H}^{d-1})$.

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- Holographically, this means that the EE can be obtained from the Bekenstein-Hawking entropy of a hyperbolic black hole.

$$\frac{\delta \mathcal{A}(\tilde{B})}{4G} = \delta \langle H_B^{\text{grav.}} \rangle \xrightleftharpoons{\text{CHM}} \delta S_{\text{BH}} = \delta \langle H_B^{\text{grav.}} \rangle$$

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$$\delta \left[R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R - 8\pi G T_{\mu\nu} \right] = 0 \quad \xrightarrow{\text{IW}} \quad \delta S_{\text{BH}} = \delta \langle H_B^{\text{grav.}} \rangle$$

$$\xrightleftharpoons[\text{CHM+AdS/CFT}]{\text{EE}} \delta S_{\text{EE}}(B) = \delta \langle H_B \rangle$$

Namely, the linearized Einstein equations imply the quantum first law of entanglement entropy for ball-shaped regions.

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- **Classical (linearized) gravitational dynamics in AdS spaces follows from a fundamental quantum principle!**

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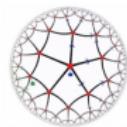
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It from Qubit
Simons Collaboration on
Quantum Fields, Gravity and Information



You can reach me at:
pablo.bueno@cab.cnea.gov.ar