

# BLACK HOLE THERMODYNAMICS

## Lecture 4

### Euclidean Methods

#### BH Mechanics

$K$ : surface gravity  
constant on  $\mathbb{H}^+$

$M, J$  conserved  
quantities.

$$dM = \frac{k}{8\pi} dA_{++} + \sum_i dJ_i + \bar{\Phi}_{++} dQ$$

$dA_{++} \gg 0$  area increase

#### BH Thermodynamics

$k_+ \sim \text{temp}$   
canonical ensemble  
(temp constant in thermal eq.)

$E \sim M$  energy system

$$dE = T dS + \sum_i \mu_i dN_i$$

$dS \geq 0$  entropy increase

We will find that

$$T_+ = \frac{k}{2\pi} K \rightarrow \text{Hawking Temperature}$$

$$S_{BH} = \frac{A_+}{4 l_p^2} \rightarrow \text{Bekenstein-Hawking entropy formula.}$$

Einstein-Hilbert

$$(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi G T_{\mu\nu})$$

## Euc Schw BH:

Schw:

$$ds^2 = -\left(1 - \frac{2M}{r}\right) dt^2 + \frac{dr^2}{\left(1 - \frac{2M}{r}\right)} + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$

↓ EUC Schw:  $t = -it_E$  Wick Rotation

$$ds_E^2 = \left(1 - \frac{2M}{r}\right) dt_E^2 + \frac{dr^2}{\left(1 - \frac{2M}{r}\right)} + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$

Zoom near  $r = 2M$ , define new coords

$$r = \frac{\rho^2}{8M} + 2M \quad dr = \frac{\rho}{4M} d\rho$$

look at behaviour as  $\rho \rightarrow 0$  ( $r \rightarrow 2M$ )

$$1 - \frac{2M}{r} = 1 - \frac{1}{1 + \frac{\rho^2}{16M^2}} = \frac{\rho^2}{16M^2} + O(\rho^4)$$

$$ds_E^2 = \left.\frac{\rho^2}{16M^2} dt_E^2\right|_{\rho=0} + \frac{1}{\frac{\rho^3}{16M^2}} \frac{\rho^2}{16M^2} d\rho^2 + (2M)^2(d\theta^2 + \sin^2\theta d\phi^2) + \dots$$

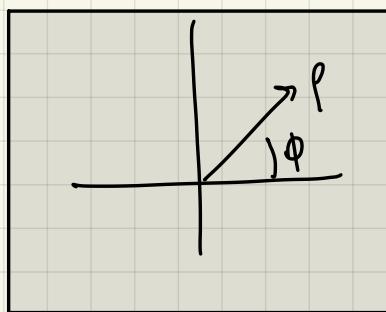
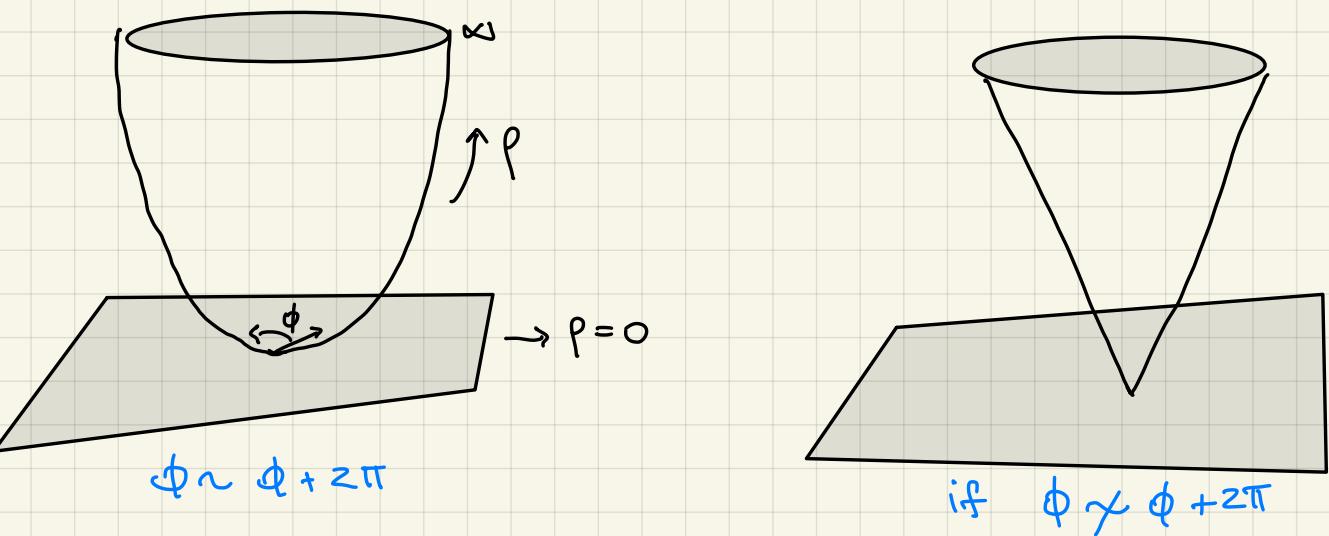
$$= \underbrace{\frac{\rho^2}{16M^2} dt_E^2}_{\mathbb{R}^2 \text{ in polar coords}} + d\rho^2 + (2M)^2(d\theta^2 + \sin^2\theta d\phi^2) + \dots$$

$\underbrace{\hspace{10em}}_{S^2}$

$\mathbb{R}^2$  in polar coords

$$\rho^2 d\phi^2 + d\rho^2 \Rightarrow \phi = \frac{t_E}{4M}$$

$$\rho \in [0, \infty) \quad \phi \in [0, 2\pi]$$



If we want a regular Euclidean geometry for Schw  
 $\phi \sim \phi + 2\pi$

$$\Rightarrow t_E \sim t_E + 2\pi (4M)$$

### QFT interpretation

What does it mean for Euc time to be periodic ?

$$t_E \sim t_E + \beta$$

From a path integral we would have

$$Z[\bar{\Phi}_i, \bar{\Phi}_f] = \int \mathcal{D}\bar{\Phi} e^{\frac{iS[\bar{\Phi}]}{\hbar}} = \langle \bar{\Phi}_i | e^{-\frac{iHt_f}{\hbar}} | \bar{\Phi}_f \rangle$$

$\bar{\Phi}(0) = \bar{\Phi}_i$

$$iS[\bar{\Phi}] = i \int dt d^3x L = -S_E = - \int dt_E d^3x L$$

$t = i t_E$

$$e^{-i\epsilon t + \frac{h}{\hbar}} \rightarrow e^{-\epsilon t + \frac{h}{\hbar}}$$

$$\Rightarrow t_E \sim t_E + \beta$$

$$Z[\beta] = \text{Tr}(e^{-\beta H}) \quad : \text{partition function at finite temp}$$

for angular momentum

$$Z[\beta, \Omega] = \text{Tr}(e^{-\beta(H - \Omega J)})$$

$$T = \frac{1}{\beta}$$

Interpretation implemented for Euc Schw

$$t_B \sim t_E + 2\pi(4M)$$

$$\Rightarrow \beta = 2\pi \cdot \frac{4M}{K} \Rightarrow T_4 = \frac{1}{8\pi M}$$

$$T_4 = \frac{K}{2\pi}$$

Rindler spacetime: back to Lorentzian

$$ds_E = \underbrace{\frac{r^2}{16M^2} dt_E^2}_{\text{TR}^2 \text{ in polar coords}} + dp^2 + (2M)^2 (d\theta^2 + \sin^2\theta d\phi^2) + \dots$$

↓ back to Lorentzian

$$= -\underbrace{\frac{r^2}{16M^2} dt^2}_{\text{Rindler}} + dp^2 + (2M)^2 (d\theta^2 + \sin^2\theta d\phi^2)$$

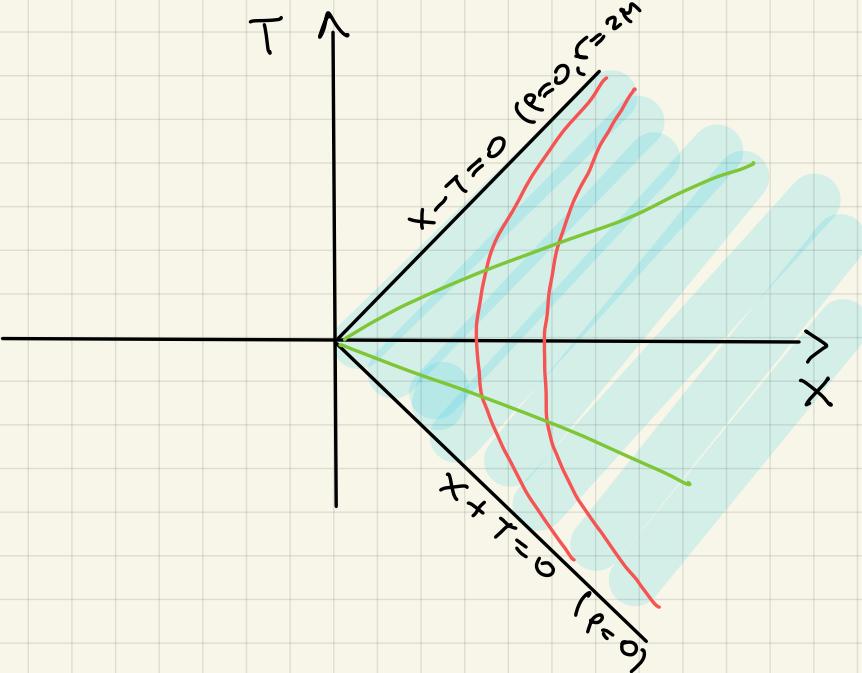
$$T = r \sinh\left(\frac{t}{4M}\right)$$

$$X = r \cosh\left(\frac{t}{4M}\right)$$

$$= -dT^2 + dx^2 + \dots$$

$$x^2 - T^2 = p^2 \quad -$$

$$\frac{T}{X} = \tanh\left(\frac{t}{4M}\right) \quad -$$



Rindler: acc obs  
in Minkowski

$$\text{acc : } \frac{1}{4M} = k$$