

#15

zeta zeta zeta zeta zeta

#1



friction coefficient

x_i

Longvin

$$m_x \ddot{x}$$

$=$

$f(x, v, t)$
external force

$$- \frac{\partial E(x)}{\partial x}$$

✓

$$- \gamma \dot{x}$$

+

$$\sum \xi$$

$$\sigma_\xi^2 = \frac{2\gamma}{\beta} \delta(t)$$

overdamped Langevin
(Brown)

Δt : conditional
white noise.

$$0 =$$

[Free energy estimate of work is Potential & Mean Force]

#2



F_{xy}

$E(x, y)$

Joint Free energy

Mean force on x = $\left\langle \frac{\partial E}{\partial x} \right\rangle = \int dy P(y|x) \frac{\partial E}{\partial x}$

$$P(y|x) = \frac{P(x, y)}{P(x)} = \frac{e^{-\beta E(x, y) + \beta F_{xy}}}{\int dy e^{-\beta E(x, y) + \beta F_{xy}}}$$

Point wise Free energy

define $F'_{xy}(x) = -\frac{1}{\beta} \ln \int dy e^{-\beta E(x, y)}$

$$\frac{\partial F'_{xy}(x)}{\partial x} = \frac{\int dy e^{-\beta E(x, y)} \frac{\partial E}{\partial x}}{\int dy e^{-\beta E(x, y)}} = \int dy P(y|x) \frac{\partial E}{\partial x} = \left\langle \frac{\partial E}{\partial x} \right\rangle = E^{MF}$$



#3

Berezikovskii & Szabo 2011

(Klein - Kramers = Fokker-Planck equation
= Smoluchowski eq)

$$m \ddot{x} = f_x(t) = - \frac{\partial}{\partial x} E^{MF}(x) + \delta f_x$$

$$- \frac{\partial E(x, y)}{\partial x} - \frac{\partial E(x, y)}{\partial x} + \left\langle \frac{\partial E(x, y)}{\partial x} \right\rangle_y$$

Approximate $\delta f_x = 0$

Friction not to displacement

or δf_x is delta correlated white noise

$$\text{force } \gamma = \beta \int_0^\infty \langle \delta f_x(u) \delta f_x(t) \rangle dt$$

$$\delta f_x \approx -\gamma \ddot{x} + c \xi$$

$$\downarrow$$

$$\sqrt{\frac{2\gamma}{\beta}}$$

Fluctuation dissipation theorem

Sivak et al 2012

$$\delta \left(\frac{\partial F(x)}{\partial x} \right)$$

#4

$$D = \frac{\mu}{\beta} \leftarrow \text{mobility} = \frac{\langle v \rangle}{F_{ex}} = \frac{1}{\beta \gamma}$$

Einstein Relation

diffusion coeff

$$(\beta \gamma)^{-1} = D \quad (\gamma \text{ could be matrix})$$



Stokes-Einstein eq.

$$D = \frac{1}{\beta 6\pi\eta r} \leftarrow \text{Rotus}$$

\uparrow Viscosity

low Reynolds number Re

(Ratio of inertial to viscous forces)

(Micro hydro. is at low Reynolds #)

$$\gamma = 6\pi\eta r \leftarrow \text{friction scale \& Rotus.}$$

effective Mass

$$P(v) \propto \frac{e^{-\frac{\beta m_e v^2}{2}}}{\sqrt{2\pi/\beta m_e}}$$

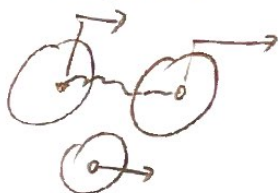
#5

Problems → Aggregation of N particles $R_{rms} \sim N^{1/3}$

$$D \sim \frac{1}{N^{1/3}}$$



But if Langevin does not work



even more to move in same direction! $D \stackrel{?}{\sim} \frac{1}{e^N}$

Why? Langevin breaks Galilean Symmetries No absolute frame at rest.

Potential Solution Dissipative Particle Dynamics



dissipative-fluctuating force between particles.

Longvin Path Action

Huhl & Ross 1981 ²¹

$$m \ddot{x}(t) = f(x) - \gamma(x) \dot{x}(t) + \xi(t) \quad \langle \xi_i \rangle = 0 \quad \langle \xi_i(t) \xi_j(t') \rangle = \frac{2\gamma}{\beta} \delta(t-t')$$

$$P[x|x(0)] \propto e^{-A_{\text{active}}(x)} \propto e^{-\int \mathcal{L}(t) dt}$$

$$P[\xi(t)] = \left(\frac{\beta}{\pi}\right)^{\frac{d}{2}} |\gamma|^{-\frac{d}{2}} e^{-\frac{\beta}{4} \xi^T(t) \gamma^{-1} \xi(t)} \quad [\text{Multidimensional Normal}]$$

$$P[\xi] \propto \exp\left(-\frac{\beta}{4} \int \xi^T \gamma^{-1} \xi dt\right)$$

$$P[x|x(0)] \propto \int e^{-\int dt \mathcal{L}(t)}$$

Jacobian is path independent.

$$\propto \frac{1}{4} \beta \left(m \ddot{x} - f(x) + \gamma \dot{x} \right)^T \gamma^{-1} \left(m \ddot{x} - f(x) + \gamma \dot{x}(t) \right) + C + \frac{1}{2} \ln |\beta \gamma^{-1}|$$

Se Kimoto 1998

$$A = \int \mathcal{L} dt, \quad \mathcal{L} = \frac{1}{2} \beta (m \ddot{x} - f(x)) \dot{x} \quad (1)$$

$$+ \frac{1}{4} \beta [m \ddot{x} - f(x)]^T \gamma^{-1} [m \ddot{x} - f(x)] \quad (2) \quad \text{kinetic}$$

$$+ \frac{1}{4} \beta \dot{x}^T \gamma \dot{x} \quad (3) \quad \text{kinetic}$$

$$+ \frac{1}{2} \ln |\beta \gamma^{-1}| \quad (4)$$

①

$$Q = \frac{\delta E}{\delta x} \frac{\partial x}{\partial t} \frac{\partial}{\partial \lambda} E(x, \dot{x}, \lambda) = \frac{\partial E}{\partial x} \frac{\partial x}{\partial t} + \frac{\partial (\frac{1}{2} m \dot{x}^2)}{\partial \dot{x}} \frac{d\dot{x}}{dt} + \underbrace{\frac{\partial E}{\partial \lambda} \frac{\partial \lambda}{\partial t}}_W$$

$$- f \dot{x} \quad m \dot{x} \ddot{x}$$

$$\underbrace{(m \ddot{x} - f) \dot{x}}_Q + W$$

$$\frac{A - \tilde{A}}{2} = \beta Q$$

kinetic friction, the ② dominates
 $m \ddot{x} - f(x) = 0$
 (classical)

