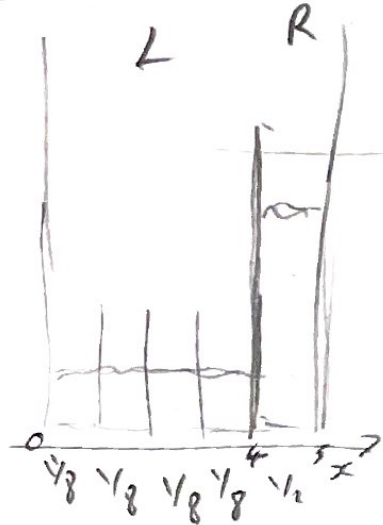


# Lecture # 5



$$S = 4 \times \frac{1}{8} \times 3 + \frac{1}{2}$$

2 bits

$$-\log_2 2 = 1$$

$$S = 4 \times \frac{1}{4} \times 2 =$$

= 2 bits

$$-\int_0^4 \frac{1}{8} \log \frac{1}{8} \frac{1}{1/4} dx = \int_0^5 \frac{1}{2} \log \frac{1}{2} \frac{1}{1/2} dx$$

$$= 2 \text{ bits} - \log 2$$

↑

Adrian

Renew entropy of derivatives



Continuum

$$-\int p \ln \frac{p}{p_0} dx$$

$$S(A, B) = S(A) + S(B|A)$$

$$= S(A) + \sum P(a, b) \log P(b|a)$$

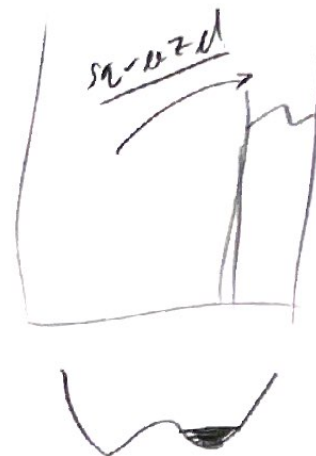
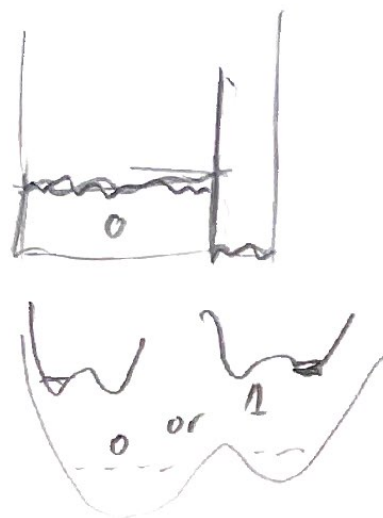
$$\sum_a P(a) \sum_b P(b|a) \log P(b|a)$$

$\langle S(B|a) \rangle$

$$S(A) = 1 \text{ bit} + \frac{1}{2} S(B|a) + \frac{1}{2} S(B|r)$$

2      0

$$= 2 \text{ bits}$$



(2)

NOTE: Quantum Different! No Cloning Theorem  
 Can't copy Quantum information.

# Intro to Jarzynski

③

CLAUSIUS 1864

The energy of the world is constant

The entropy of the world tends to a maximum.

$$\Delta S^{\text{tot}} \geq 0$$

$$\Delta S^{\text{S}} + \Delta S^{\text{env}}$$

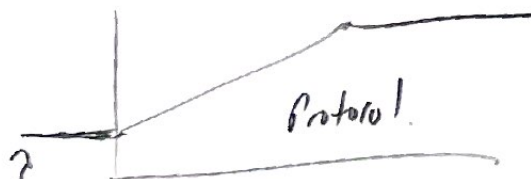
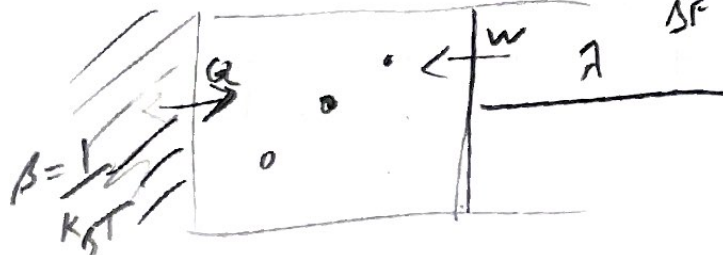
$$- \beta Q + \beta W - \beta \Delta E$$

$$\Delta S - \beta \Delta E$$

$$- \beta \Delta E + \beta W \geq 0$$

$$\Delta F \leq W$$

$$\Delta E = Q + W$$



Jarzynski equality 1997

$$\langle e^{-\beta W} \rangle = e^{-\beta \Delta F} !$$

↓ Jensen

$$\Delta F \leq \langle W \rangle$$

[Surprise in 1997!  
Not a near-equilibrium result]

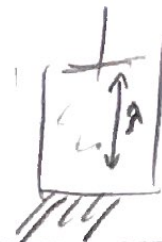
[Young → more important the test than  
not going to follow this original derivation.  
→ come back to that later]

"Fluctuation Theorems" (then notes, experiments, consequences)

(4)

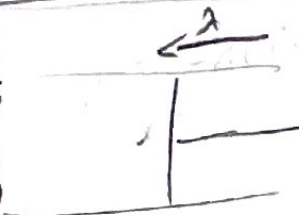
# Canonical Ensemble

"Boltzmann distribution"



Gibbs 1902  
~~Boltzmann 1877~~

$$P_A(x) = \frac{e^{-\beta E(x, \lambda)}}{Z}$$



$$= e^{\beta F - \beta E(x, \lambda)}$$

$$Z = \sum_x e^{-\beta E(x, \lambda)}$$

$$\beta F = \ln Z$$

$$\beta F = -S + \beta \langle E \rangle$$

$$= \sum P(x) \ln P(x) + \beta \langle E \rangle$$

$$= \sum P(x) \ln \frac{e^{-\beta E(x)}}{Z} + \beta \langle E \rangle$$

$$= -\ln Z + \underbrace{\sum P(x) (-\beta E(x))}_{0} + \beta \langle E \rangle$$

$$\beta F \stackrel{or}{=} -\ln Z$$

$$\frac{\partial \beta F}{\partial \beta} = -\frac{\partial \ln Z}{\partial \beta} = -\frac{1}{Z} \frac{\partial Z}{\partial \beta}$$

$$= -\frac{1}{Z} \sum [-E(x)] e^{-\beta E(x)}$$

$$= \langle E \rangle$$

(5) (4)

## Maximum Entropy

Entropy is maximized at Thermodynamic equilibrium.

Calculus of Variations / Lagrange multipliers

Maximize  $S = -\sum P(x) \ln P(x)$  subject to constraints  $\sum P_x = 1$  and  $\sum P_x E(x) = \langle E \rangle$

$$J(P) = S - \lambda_0 \sum P_x - \lambda_1 \sum P_x E(x)$$

Variation

$$\frac{\delta J(P)}{\delta P(x)} = 0$$

$$= -1 - \ln P(x) - \lambda_0 - \lambda_1 E(x)$$

$$P(x) = e^{-1 - \lambda_0 - \lambda_1 E(x)} = \frac{e^{-\beta E(x)}}{Z}$$

Max entropy

Jaynes [What happens?

1963 → Not a general appeal to maximize  
Max entropy property of equilibrium  
Not to know constraints!]



Tsallis Entropy

$$S_q = \frac{-\sum p_i^q - 1}{q-1}$$

$$\langle E \rangle_q = \frac{\sum p_i^q E_i}{\sum p_i^q}$$

Maximize  $\rightarrow p_i = \frac{1}{Z} [1 + \beta(q-1) E_i]^{-\frac{1}{q-1}}$

"Power laws"

JUNK

Don't waste Your Time!

"Maximum Caliber"

# Temperature.

7

Thermodynamics  $S$  function of  $E, V, N, \dots$  extensive.

$$cS(E, V, N) \stackrel{\text{T.L.}}{=} S(cE, cV, cN) \quad | \quad S \sim$$

Euler's Theorem for 1st order Homogeneous function.

$$f(x, y, z) = ax + by + cz$$
$$a = \partial_x f = \partial_y f + \partial_z f$$

$$\lambda f(x, y, z) = f(\lambda x, \lambda y, \lambda z)$$

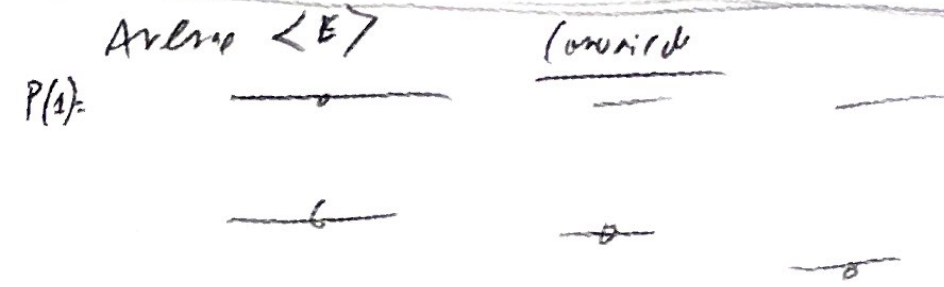
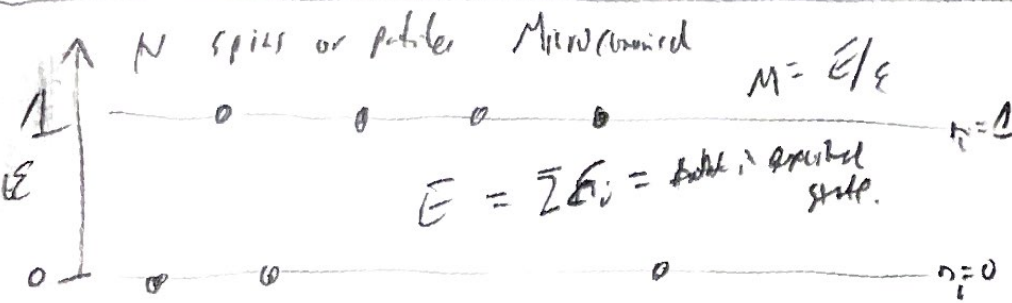
$$\frac{\partial}{\partial \lambda} f(x, y, z) = \frac{\partial f}{\partial \lambda x} \underbrace{\frac{\partial \lambda x}{\partial \lambda}}_x + \frac{\partial f}{\partial \lambda y} \underbrace{\frac{\partial \lambda y}{\partial \lambda}}_y + \frac{\partial f}{\partial \lambda z} \underbrace{\frac{\partial \lambda z}{\partial \lambda}}_z$$

$$\text{let } \lambda = 1$$

MC  $\rightarrow$  Canonical.

(9)

19505



$N$  spins, Energy  $E$

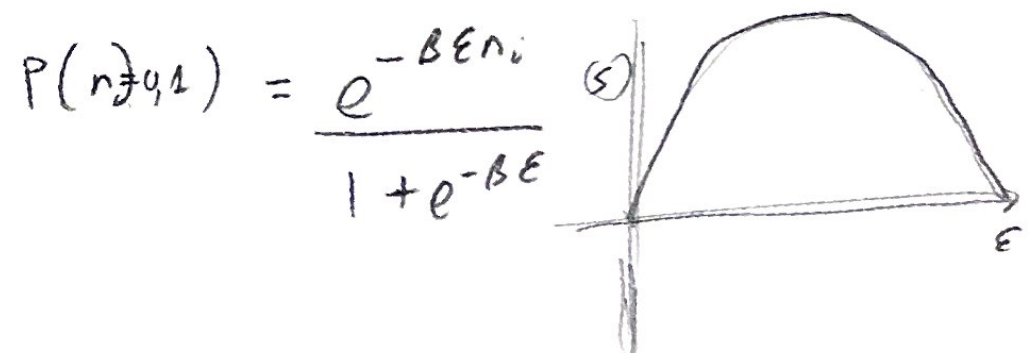
$$S = k_B \ln \Omega = k_B \ln \frac{N!}{E/\epsilon! (N - E/\epsilon)!}$$

Microcanonical  
Spins not independent

$$\langle E \rangle$$

$$Z = \sum_{E=0}^N \binom{N}{E} e^{-\beta E} = \sum_{n_1, n_2, \dots, n_N} e^{-\beta E}$$

$$S = N k_B \ln \left( 1 + e^{-\beta \epsilon} \right)$$





$$\ln \frac{N!}{E! (N-E)!}$$

$$\ln N! =$$

$$N \ln N - N$$

$$S = N \ln N - E \ln E - (N-E) \ln (N-E)$$

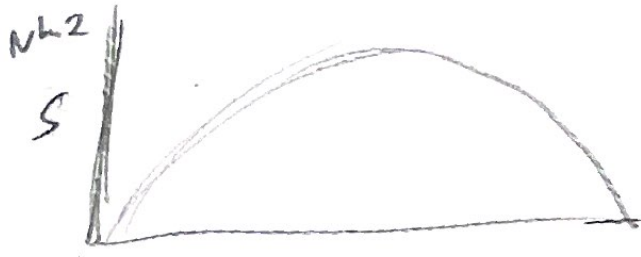
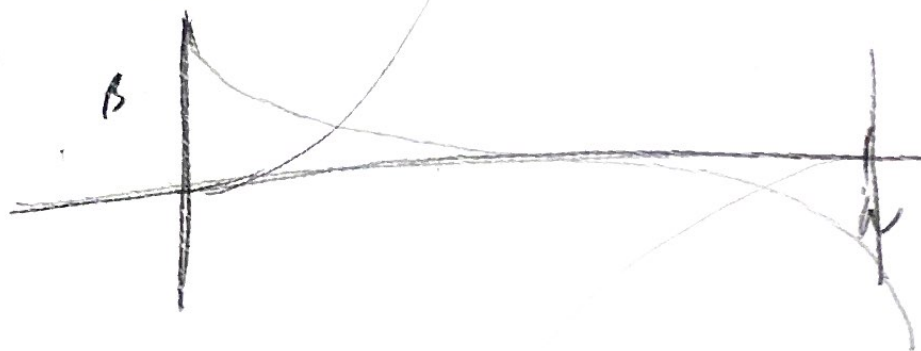
$$\frac{\partial S}{\partial E} = -1 - \ln E + 1 + \ln (N-E)$$

$$\frac{\partial S}{\partial E} = \ln \left( \frac{N-E}{E} \right)$$

$$N \ln (N-E) + E \ln (N-E)$$

$$\frac{N}{N-E}$$

$$-\left( \frac{N-E}{N-E} \right) + \ln (N-E)$$



(Dunkel et al 2014)

Temp cont

Inventin, Temperature, Chong 2004

why not reg? Even not generally has upper bound

not stable.Temperature  $\rightarrow$  well defined for Macroscopic eq. Systems.

$$P(x) = \frac{e^{-\beta E}}{Z}$$

Temperature of environment

Reserve, But convert energy to entropy

e.g. bit error  $W = kT \ln 2$ 

Does not generalize! to real.

e. Kinetic Temperature

even Harmonic Degree of freedom  $\langle E \rangle = \frac{1}{2} k_B T$ 

$$\langle K \rangle = \frac{3}{2} k_B T$$

$$k_B T_{\text{kinetic}} \neq \frac{\langle K \rangle}{N \frac{3}{2}}$$

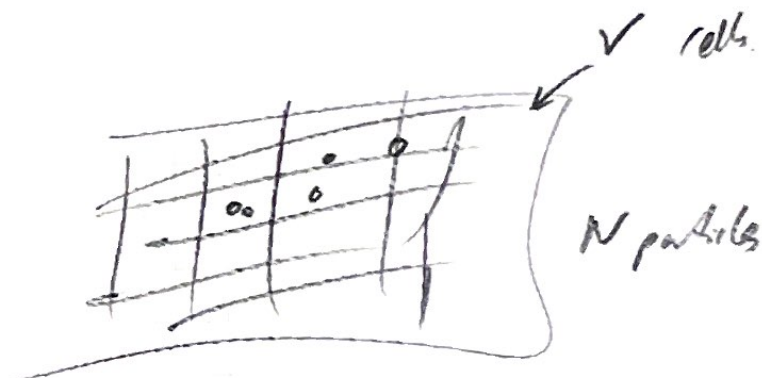
10% opt

Ideal Gas

$$pV = nRT$$

$$pV = N k_B T$$

$$\beta p = \frac{N}{V}$$
$$\left[ \frac{\partial S}{\partial V} \right]$$



$$S = k_B \Omega = k_B \ln \frac{V^N}{N!}$$

$$\boxed{\frac{\partial S}{\partial V} = \frac{N}{V}}$$