

Radar Cross Section Engineering Homework 3

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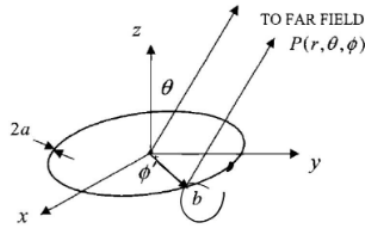
ECE-5973-004

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Part 1:

Circular Wire Example



P_1 : Loop made from resistive film
 & surfaces resistivity R_s .
 Find MoM expression for
 lead impedances $(Z_L)_{mn}$.

\mathbf{Z} : (3.68) (3.69)

\mathbf{R} Or \mathbf{V} : (3.72)

$$\hat{k} \cdot \hat{r}' = b \sin \theta \cos(\phi - \phi') \quad \left. \begin{array}{l} \text{Impedance} \downarrow \text{Excitation} \\ \text{Matrices} \end{array} \right\}$$

$$\hat{\phi} \cdot \hat{\phi}' = \cos(\phi - \phi')$$

$$\vec{J}_s(\phi) = \hat{\phi} \sum_{n=-\infty}^{+\infty} I_n \cdot \frac{e^{jn\phi}}{2\pi a} \quad \left. \begin{array}{l} \text{Entire Domain Base} \\ \text{Function (Mode)} \end{array} \right\}$$

$$\vec{W}_n(\phi) = \hat{\phi} \cdot \frac{e^{-jn\phi}}{2\pi a} \quad \left. \begin{array}{l} \text{Conjugate of origin} \\ \text{base function.} \end{array} \right\}$$

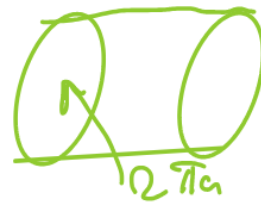
Impedance and Excitation Matrices

$$Z_{mn} = \iint_{S_m} ds \iint_{S_n} ds' \left[j\omega\mu_0 \vec{W}_m(\vec{r}) \cdot \vec{J}_n(\vec{r}') - \frac{j}{\omega\epsilon_0} \nabla' \cdot \vec{J}_n(\vec{r}') \nabla \cdot \vec{W}_m(\vec{r}') \right] G(\vec{r}, \vec{r}')$$

$$V_m = \iint_{S_m} \vec{W}_m(\vec{r}) \cdot \vec{E}_i(\vec{r})|_{\text{tan}} ds \rightarrow \begin{cases} V_m^\theta, & \vec{E}_i = \hat{\theta} E_{i\theta} \text{ (TM)} \\ V_m^\phi, & \vec{E}_i = \hat{\phi} E_{i\phi} \text{ (TE)} \end{cases}$$

$$Z_m = \iint_{S_m} ds \iint_{S_n} ds' \left[j\omega\mu_0 \vec{W}_m(\vec{r}) \cdot \vec{J}_n(\vec{r}') - \frac{j}{\omega\epsilon_0} (\nabla' \cdot \vec{J}_n(\vec{r}')) (\nabla \cdot \vec{W}_m(\vec{r}')) \right] G(\vec{r}, \vec{r}')$$

$$= \frac{\hat{\phi} \cdot \hat{\phi}'}{4\pi a^2} \cdot e^{j(n\phi' - m\phi)}$$



$$= \iint_S ds = 2\pi a \int_0^{2\pi} b d\phi$$

$$\nabla = \hat{\phi} \frac{1}{b} \frac{\partial}{\partial \phi}$$

$$\vec{J}_n(\phi') = \hat{\phi} \frac{e^{jn\phi'}}{2\pi a} \quad \nabla' \cdot \vec{J}_n(\phi') = \frac{n \cdot e^{jn\phi'}}{2\pi \cdot a \cdot b}$$

$$\vec{W}_m(\phi) = \hat{\phi} \cdot \frac{e^{-jm\phi}}{2\pi a} \quad \nabla \cdot \vec{W}_m(\phi) = -m \cdot \frac{e^{-jm\phi}}{2\pi ab}$$

$$R = b \sqrt{4 \sin^2\left(\frac{\Delta}{2}\right) + \frac{a^2}{b^2}} \quad \Delta = \phi - \phi' \therefore \text{Not zero only when } m=n$$

$$\begin{aligned} Z_{mn} &= \int_0^{2\pi} 2\pi a \cdot b d\phi \cdot \int_0^{2\pi} 2\pi a b d\phi' \\ &= \left[\frac{\hat{\phi} \cdot \hat{\phi}}{4\pi^2 a^2} \cdot e^{-j(m\phi - n\phi')} \cdot j\omega\mu + \frac{jn \cdot n}{\omega\epsilon b^2} \cdot e^{-j(m\phi - n\phi')} \right. \\ &\quad \left. \cdot \frac{1}{4\pi^2 a^2} \right] \cdot \frac{e^{-jkR}}{4\pi R} \end{aligned}$$

$$Z_{mn} = Z_{nm} = \frac{1}{2\pi} \int_0^{2\pi} d\Delta \left[j\omega\mu b \cos \Delta - \frac{jn^2}{\omega\epsilon b} e^{jn\Delta} \right] \cdot \frac{e^{-jkR}}{R}$$

$$R_n^\phi = 2\pi a \int_0^{2\pi} \vec{E}_r \cdot \vec{J}_n b d\phi = \pi b j^{n+1} e^{jn\phi} [\vec{J}_{n+1}(kb \sin \theta) - \vec{J}_{n-1}(kb \sin \theta)]$$

Can now calculate MoM solution using looped or straight wires.

$$(Z_L)_{mn} = \iint_S \vec{W}_m \cdot [\vec{R}_S(\vec{r}) \vec{J}_n(\vec{r})] dS$$

Thus:

$$Z_{Lm} = (2\pi a) R_s \int_{-1/2}^{1/2} \left(\frac{1}{z} \frac{P_m(z)}{2\pi a} \right) \cdot \left(\frac{1}{z} \frac{P_n(z)}{2\pi a} \right) dz$$
$$= \begin{cases} R_s \Delta / (2\pi a), & m=n \\ 0, & m \neq n \end{cases}$$

Part 2:

- a. Installed FEKO
- b. Using FEKO standard MoM solver, compute mono-static RCS of the rectangular PEC plate. Assume $aa=bb=3\lambda\lambda$. Generating results similar to Fig.2.13 in text book.

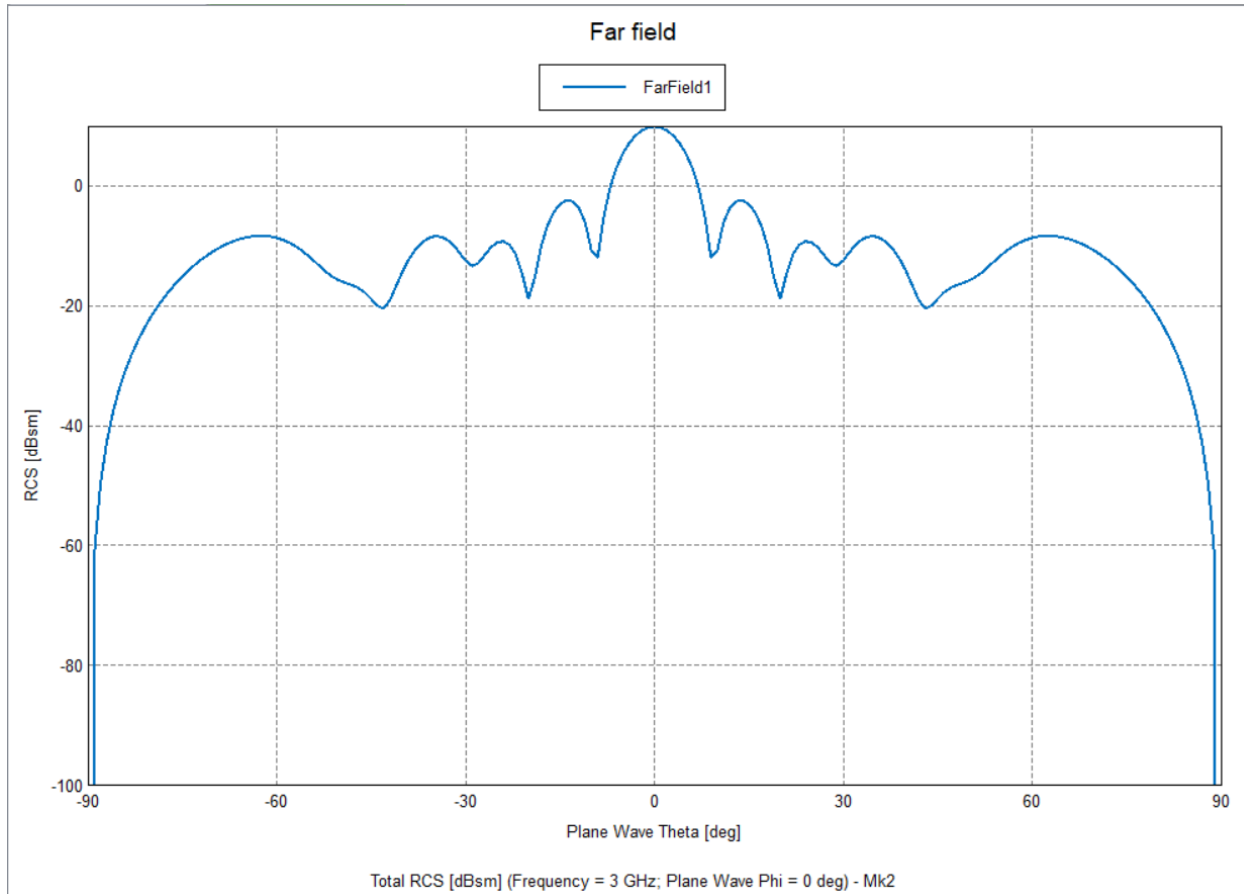


Figure 1: MoM Far Field RCS

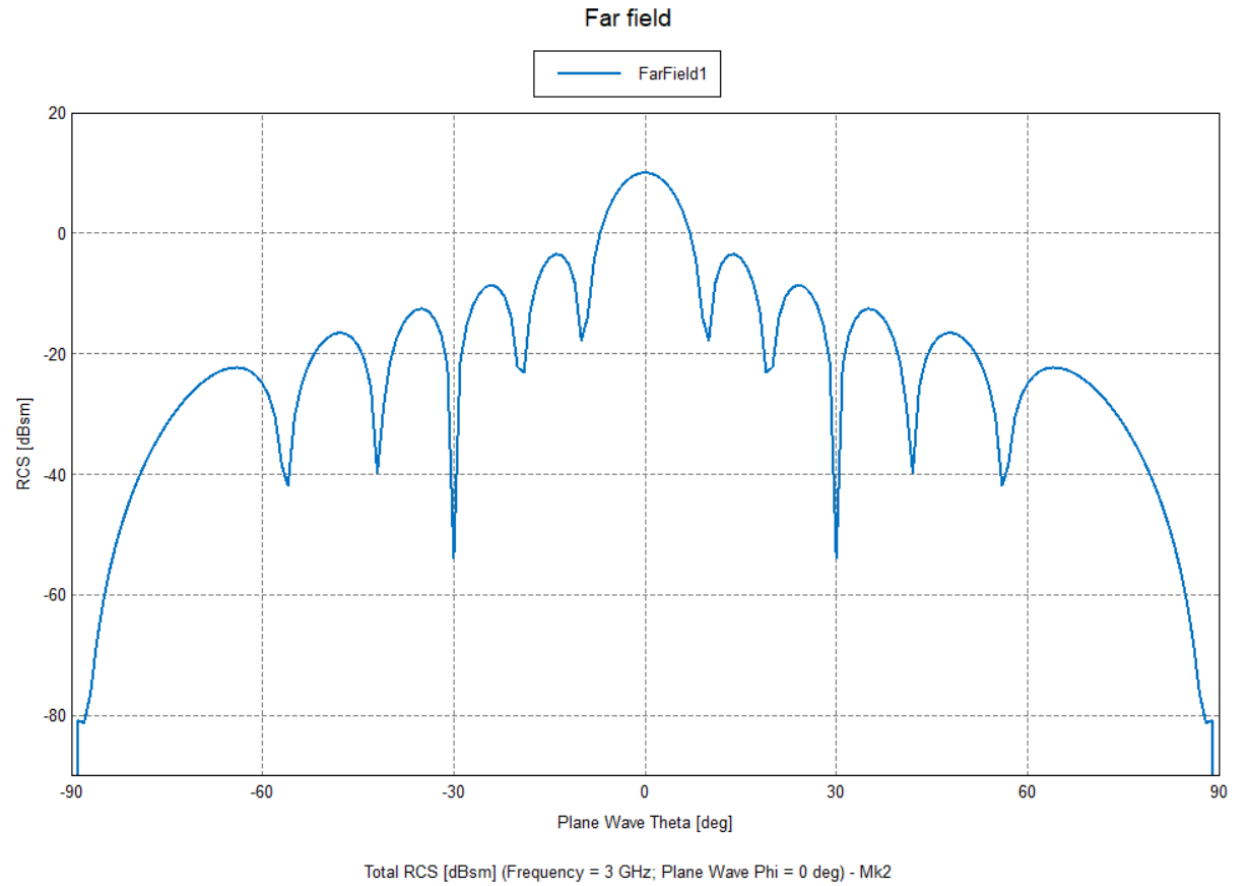


Figure 2: Physical Optics Far Field RCS.

- c. Compare the results with PO solution and discuss the difference.
 - a. The PO solution and the MoM solutions are quite similar, however, the PO solution does not lose shape continuity after 60/-60 degrees. Also, the PO solution is computed much faster.