

# Lecture 12: Stars on the Main Sequence

Lamers & Levesque Ch. 13



H-R diagram from *Gaia* DR2; <https://sci.esa.int/web/gaia/-/60240-the-hertzsprung-russell-diagram>

# **Midterm: in class Tuesday March 8**

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**will cover course material through March 1st Lecture 13**

**things to review (roughly in order of importance):**

**problem sets 1–3 (& solutions)**

**lecture slides**

**Phillips Ch. 1–5 (esp. end of chapter summaries)**

**Lamers & Levesque Ch. 1–11, 13, 14, 16–19**

**exam will be 80 minutes; 441/541 have same exam**

**mix of conceptual questions (short answer), derivation,  
and some short calculations (bring a calculator)**

**you are allowed to bring a formula sheet:**

**one side only of a 8.5" x 11" (letter size) sheet of paper**

# Stellar Timescales

## 1) Dynamical:

free-fall: what if the pressure in a star suddenly vanished?

$$\left. \begin{aligned} \Delta r \sim g t^2 &\rightarrow t^2 \sim \Delta r / g \\ \Delta r \sim R \\ g \sim GM/R^2 \end{aligned} \right\} t_{\text{ff}} \simeq \sqrt{\frac{R}{GM/R^2}} \sim \sqrt{\frac{R^3}{GM}} \rightarrow t_{\text{ff}} \simeq \frac{1}{\sqrt{G\rho}}$$

sound speed crossing time: what if the *gravity* of a star vanished?

$$t_{\text{sound}} \sim R/c_s \quad \left. \right\} t_{\text{sound}} \sim \frac{R}{\sqrt{GM/R}} \sim \sqrt{\frac{R^3}{GM}} \sim \sqrt{\frac{1}{g\rho}} \rightarrow t_{\text{sound}} \simeq \frac{1}{\sqrt{G\rho}}$$

( $t_{\text{ff}} \simeq t_{\text{sound}}$ ! why?)

This is called the  
**dynamical timescale**:

$$\tau_{\text{dyn}} = \frac{1}{\sqrt{G\rho}}$$

# Stellar Timescales

1) Dynamical:

$$\tau_{\text{dyn}} = \frac{1}{\sqrt{G\rho}}$$

Sun:  $\tau_{\text{dyn}} \sim \text{hours}$

2) Thermal/Kelvin-Helmholtz:

How long can a star live if nuclear fusion shuts off?  
(i.e., thermal energy is the only energy source left)

$$\left. \begin{array}{l} \tau_{KH} \simeq E_{th}/L \\ \text{Virial } E_{th} = -\frac{1}{2}E_{\text{pot}} \end{array} \right\} \tau_{KH} \simeq -\frac{E_{\text{pot}}}{L} \simeq \frac{GM^2}{R}/L$$

$$\boxed{\tau_{KH} = \frac{GM^2}{RL}}$$

# Stellar Timescales

1) Dynamical:

$$\tau_{\text{dyn}} = \frac{1}{\sqrt{G\rho}}$$

Sun:  $\tau_{\text{dyn}} \sim \text{hours}$

2) Thermal/Kelvin-Helmholtz:

$$\tau_{KH} = \frac{GM^2}{RL}$$

3) Nuclear:

$$\tau_{\text{nucl}} \sim \frac{E_{\text{nucl}}}{L} \sim f_M M c^2 \varepsilon_n / L$$

Sun:  
 $\tau_{\text{nucl}} \sim 10^{10} \text{ years}$

$\varepsilon_n$ =efficiency=fraction of mass converted into energy ( $\varepsilon_n = \frac{\Delta m}{m}$ )  
(H=0.007, He = 0.0007)

$f_M$ =fraction of stellar mass that takes part in nuclear fusion  
(Sun:  $f_M \sim 0.1$ ; for MS stars,  $f_M \cdot \varepsilon_H \simeq 10^{-3}$ )

So  $\tau_{\text{nucl}} \sim 10^{-3} M c^2 / L$  (H-fusion),  $\tau_{\text{nucl}} \sim 10^{-4} M c^2 / L$  (He-fusion)

# Stellar Timescales

1) Dynamical:

$$\tau_{\text{dyn}} = \frac{1}{\sqrt{G\rho}}$$

Sun:  $\tau_{\text{dyn}} \sim$  hours

2) Thermal/Kelvin-Helmholtz:

$$\tau_{KH} = \frac{GM^2}{RL}$$

3) Nuclear:

$$\tau_{\text{nuc}} \sim \frac{E_{\text{nuc}}}{L} \sim f_M M c^2 \varepsilon_n / L$$

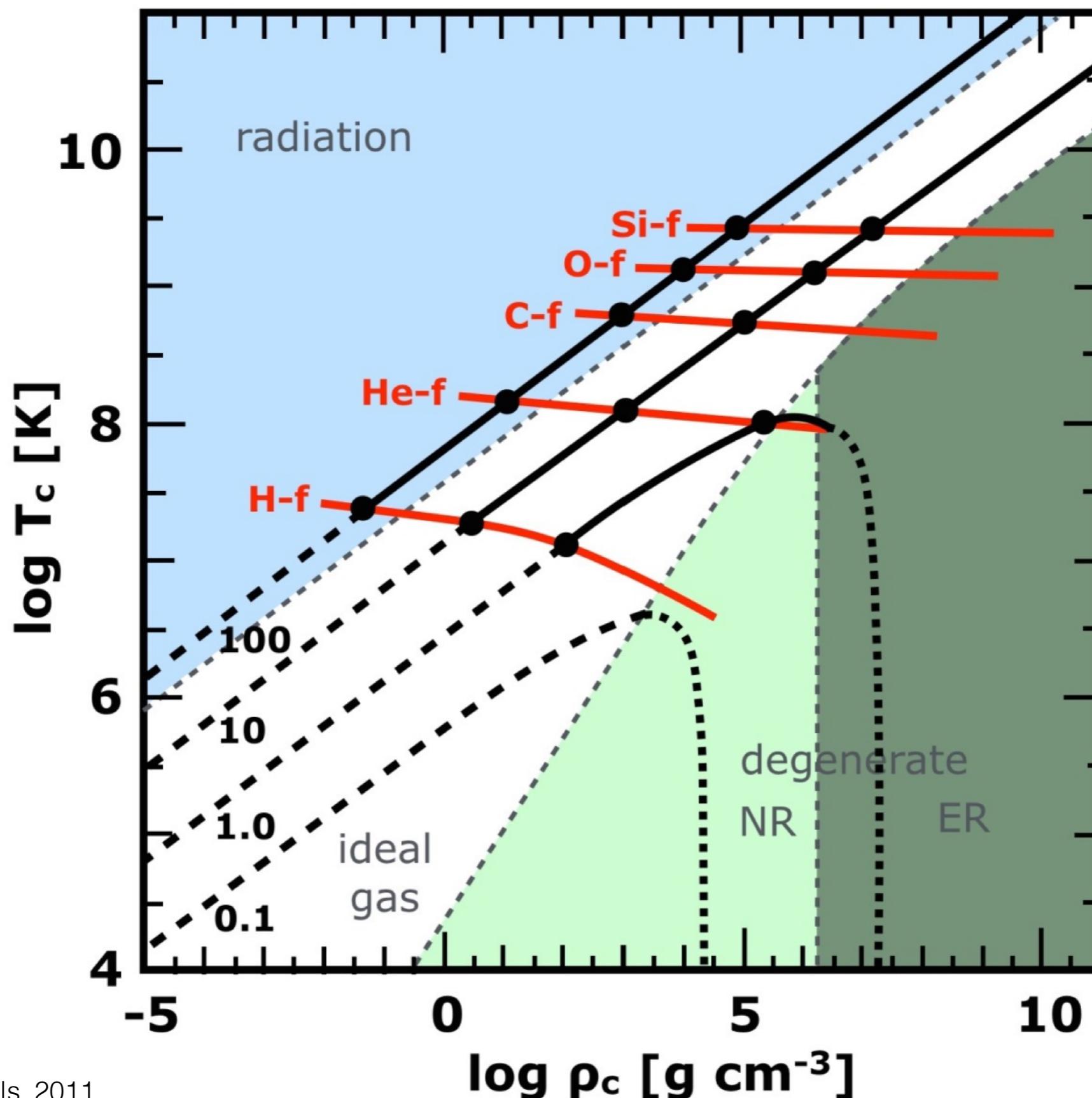
Sun:  
 $\tau_{\text{nuc}} \sim 10^{10}$  years

4) Convective: ~week (Sun) to ~month (AGB star)

$$\tau_{\text{dyn}} \ll \tau_{\text{conv}} \ll \tau_{KH} \ll \tau_{\text{nuc}}$$

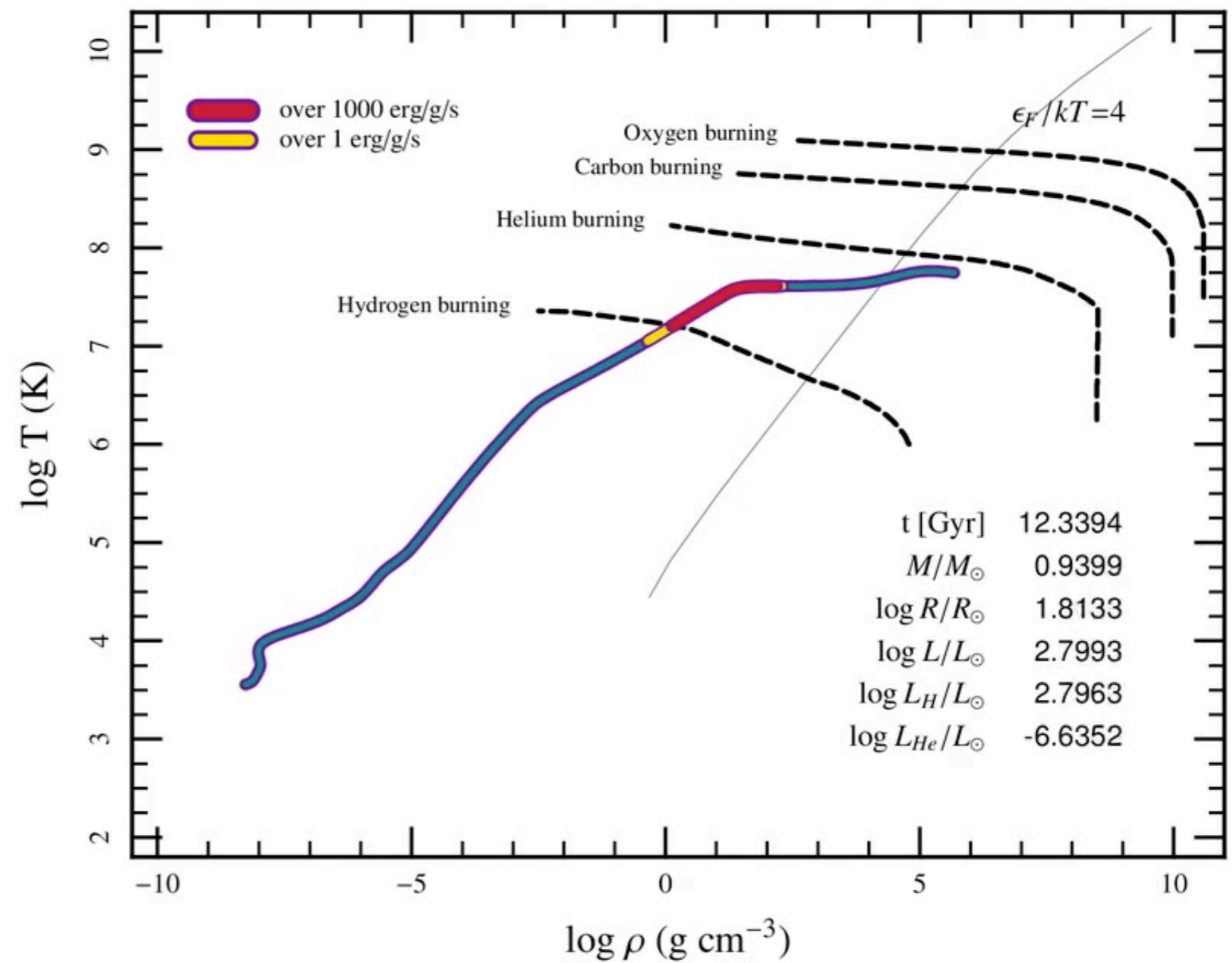
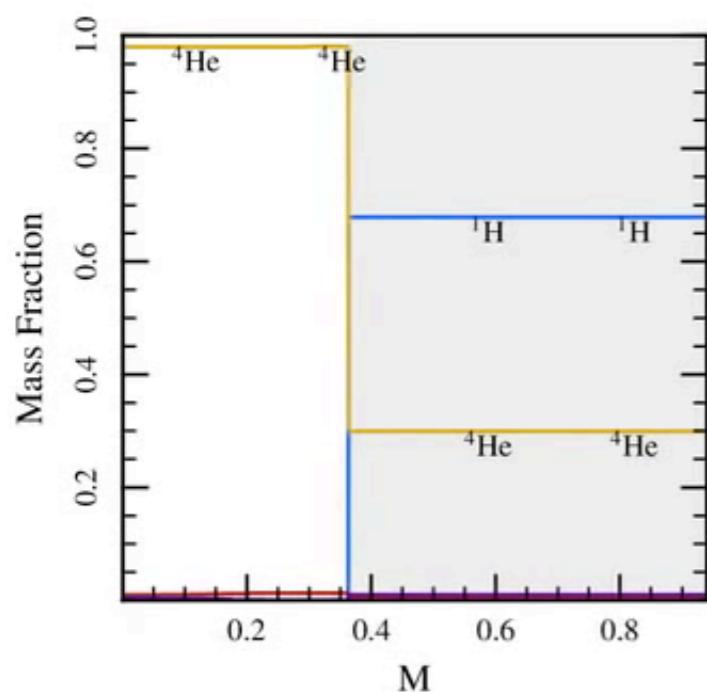
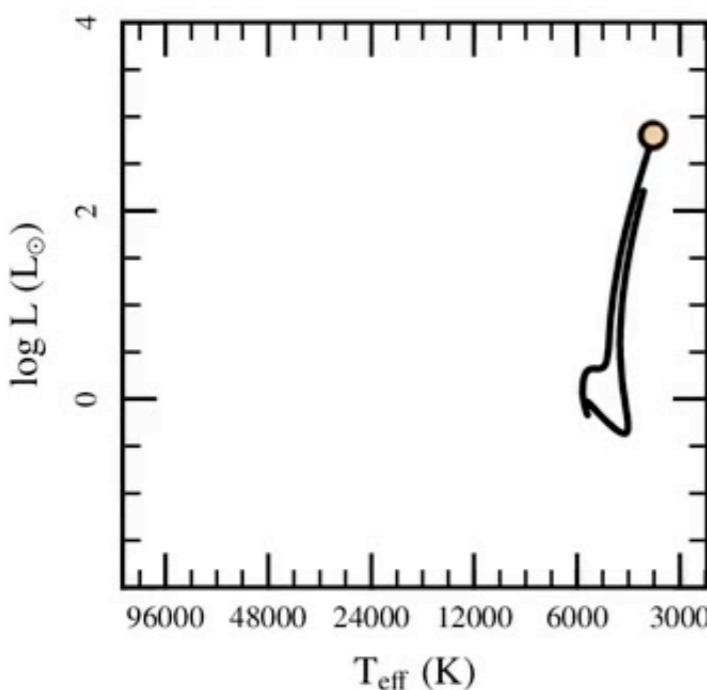
∴ ...except for explosive phases, stars are always in quasi-H.E.  
...contraction phases last about 1% as long as nuclear phases

# Fusion and Stellar Evolution



# MESA: open stellar modeling code

<http://mesa.sourceforge.net>



from Josiah Schwab, MESA model of a 1 solar mass star  
<https://www.youtube.com/watch?v=oZY3TtA63sE>

# MESA: open stellar modeling code

## MESA Stellar Evolution Movies

In the past, I've taught the Astro 310 course (Introduction to Stellar Astrophysics) with the help of [movies](#) that animate the evolution of stars of various masses. These movies have been very successful in conveying key aspects of main-sequence and RGB stellar evolution; however, they are based on models created using [EZ-Web](#), and are therefore increasingly unreliable beyond the onset of helium burning. (The helium flash is fudged in low-mass stars, and in high-mass stars the evolution has to be stopped during carbon burning).

Given that [MESA](#) is now my preferred way to calculate stellar models, *and* given that I've finally weaned myself off [IDL](#) and onto [Python](#), I thought it time to update the movies.

### Download Movies

The table below contains links to the "best-so-far" set of MESA-based movies I've created. Blanks in the table are partly deliberate (e.g., massive stars don't undergo helium flashes) and partly due to the fact that I'm still generating movies (which take a surprisingly long time to render). Please note that the already-extant movies are themselves a work-in-progress, and will change as I update movies with additional info, snazzier graphics, better axis scaling, etc.

Mass ( $M_{\odot}$ )	Evolutionary Phase				
	Main Sequence	Red-Giant Branch	Helium Flash	Helium Burning	Asympmtotic Giant Branch
0.3	<a href="#">EOS Composition Energy</a>				
1.0	<a href="#">EOS Composition Energy</a>				
3.0	<a href="#">EOS Composition Energy</a>	<a href="#">EOS Composition Energy</a>		<a href="#">EOS Composition Energy</a>	
10.0	<a href="#">EOS Composition Energy</a>	<a href="#">EOS Composition Energy</a>		<a href="#">EOS Composition Energy</a>	

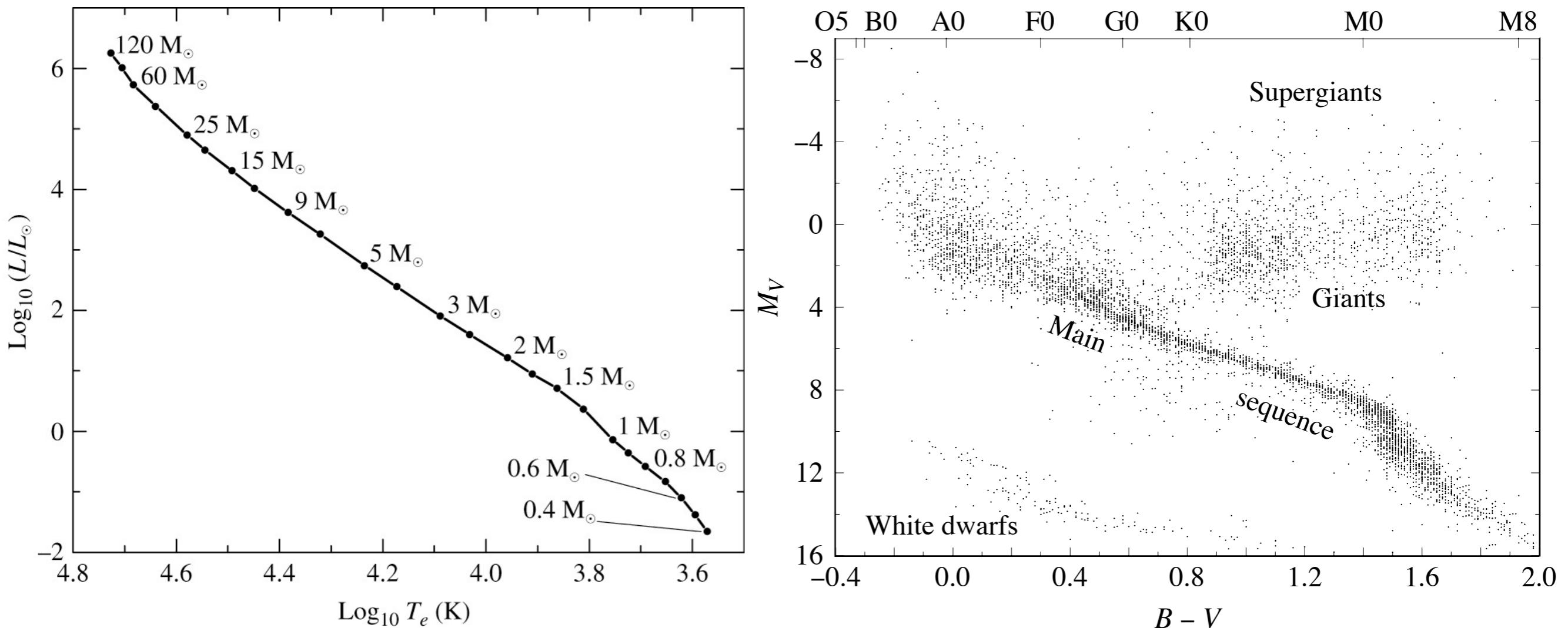
The 'EOS' links are to movies showing the equation-of-state variables ( $P$ ,  $\rho$ ,  $T$  and  $\mu$ ); the 'Composition' links are to movies showing the composition ( $X$ ,  $Y$  and CNO mass fractions); and the 'Energy' links are to movies showing the energy generation (PP, CNO, triple  $\alpha$ ).

check out the movies at

<http://www.astro.wisc.edu/~townsend/static.php?ref=mesa-movies>

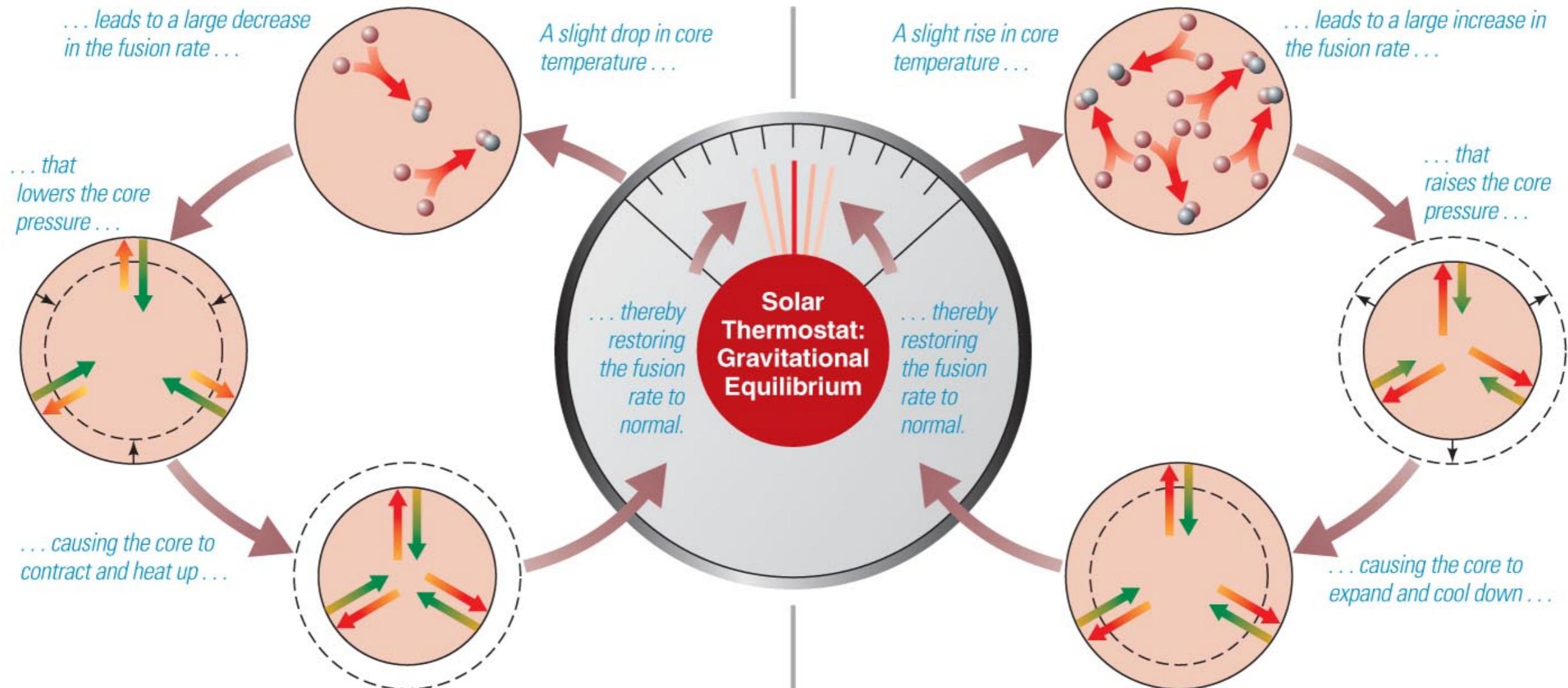
<http://www.astro.wisc.edu/~townsend/static.php?ref=ez-web-movies>

# HR diagram: theory and observations



the main sequence is a *mass* sequence; Carroll & Ostlie Figures 10.13 and 8.13

# Stellar Thermostat



Decline in core temperature causes fusion rate to drop, so core contracts and heats up.

Rise in core temperature causes fusion rate to rise, so core expands and cools down.

**essential point: temperature is regulated (stable) *when there is fusion in the core (center of star)***

# Main Sequence Evolution

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## Changes during H fusion

T<sub>c</sub>: very little; **why?**

# Main Sequence Evolution

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$T_c$ : very little; **why?** (fusion rate sensitivity)

# Main Sequence Evolution

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$\mu$ : increases ( $H \rightarrow He$ )

(from  $P_c/\rho_c \sim T_c/\mu_c$ ,  $\rho_c$  increases more than  $P_c$ )

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R: increases (denser core → envelope expands; virial eq.)

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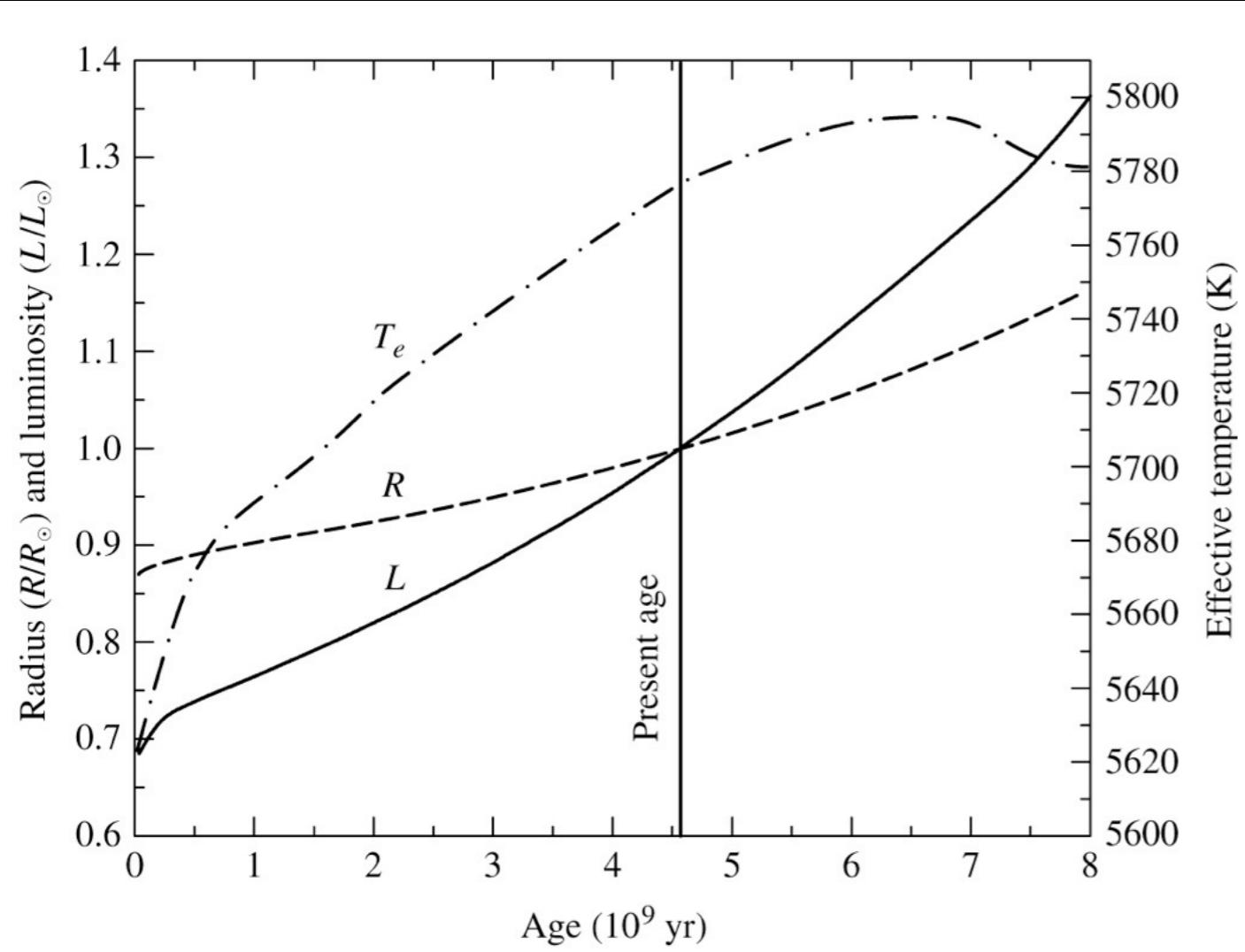
R: increases (denser core → envelope expands; virial eq.)

L: increases (“μ-effect”; higher  $\rho$  increases fusion rate)

T<sub>eff</sub>: from  $T_{\text{eff}} \sim (L/R^2)^{1/4}$

( $R^2$  increases more than  $L$ ),  $T_{\text{eff}}$  *decreases*

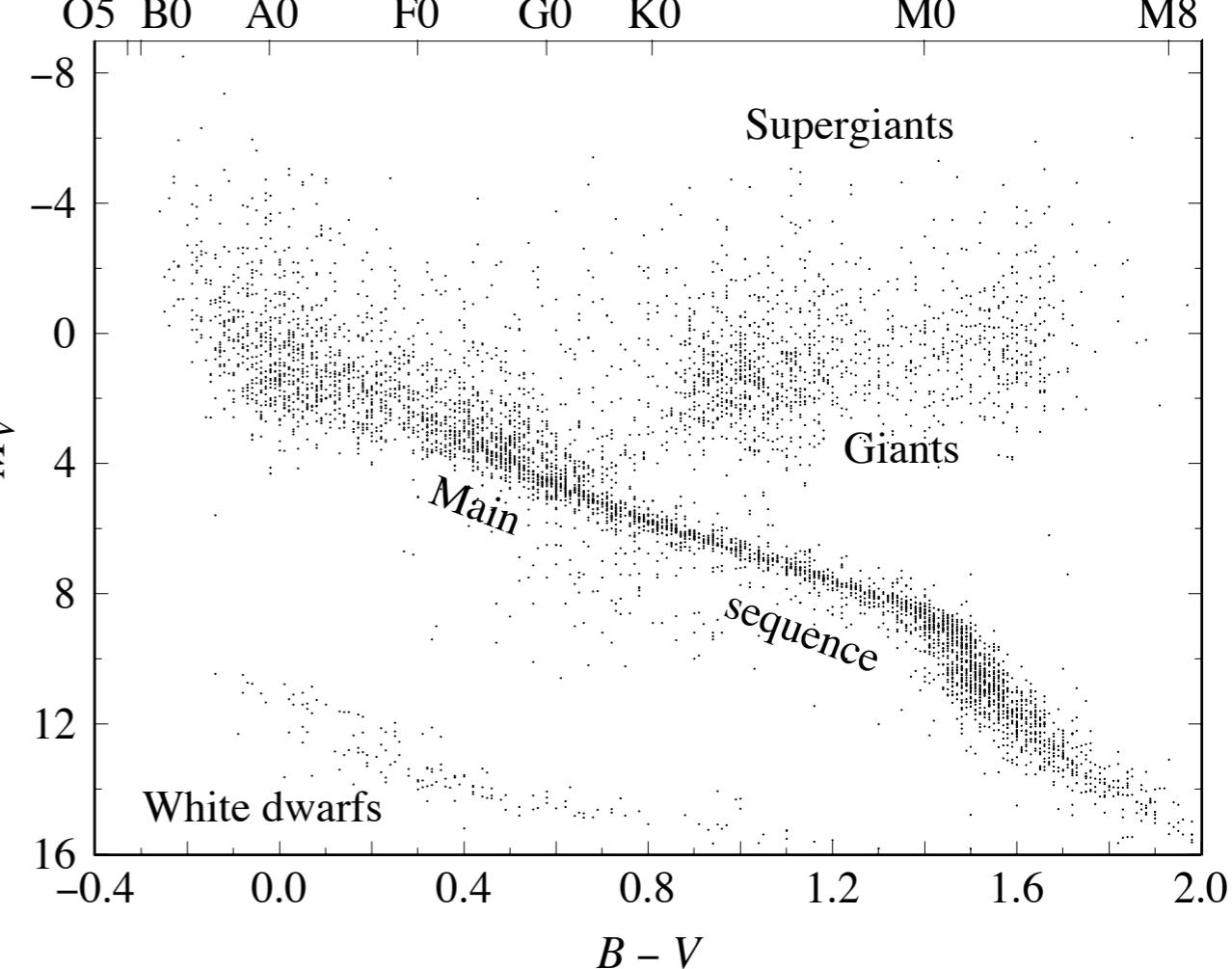
# width of main-sequence: evolution



**Table 13.1.** The Parameters of a  $1 M_\odot$  Model at Three Different MS Ages (Data are from Schaller et al. 1992)

Age Phase	$L/L_\odot$	$R/R_\odot$	$T_{\text{eff}}$	$T_c$	$P_c$	$\rho_c$	$X_c$	$Y_c$	$Z_c$	$\mu_c$
0 Gyr ZAMS	0.69	0.66	6540	13.6	1.24	78	0.680	0.301	0.02	0.71
4.5 Gyr Now	1.00	1.00	5820	15.9	2.36	157	0.302	0.677	0.02	0.88
9.4 Gyr TAMS	1.55	1.25	5820	18.7	6.73	572	0.000	0.980	0.02	1.32

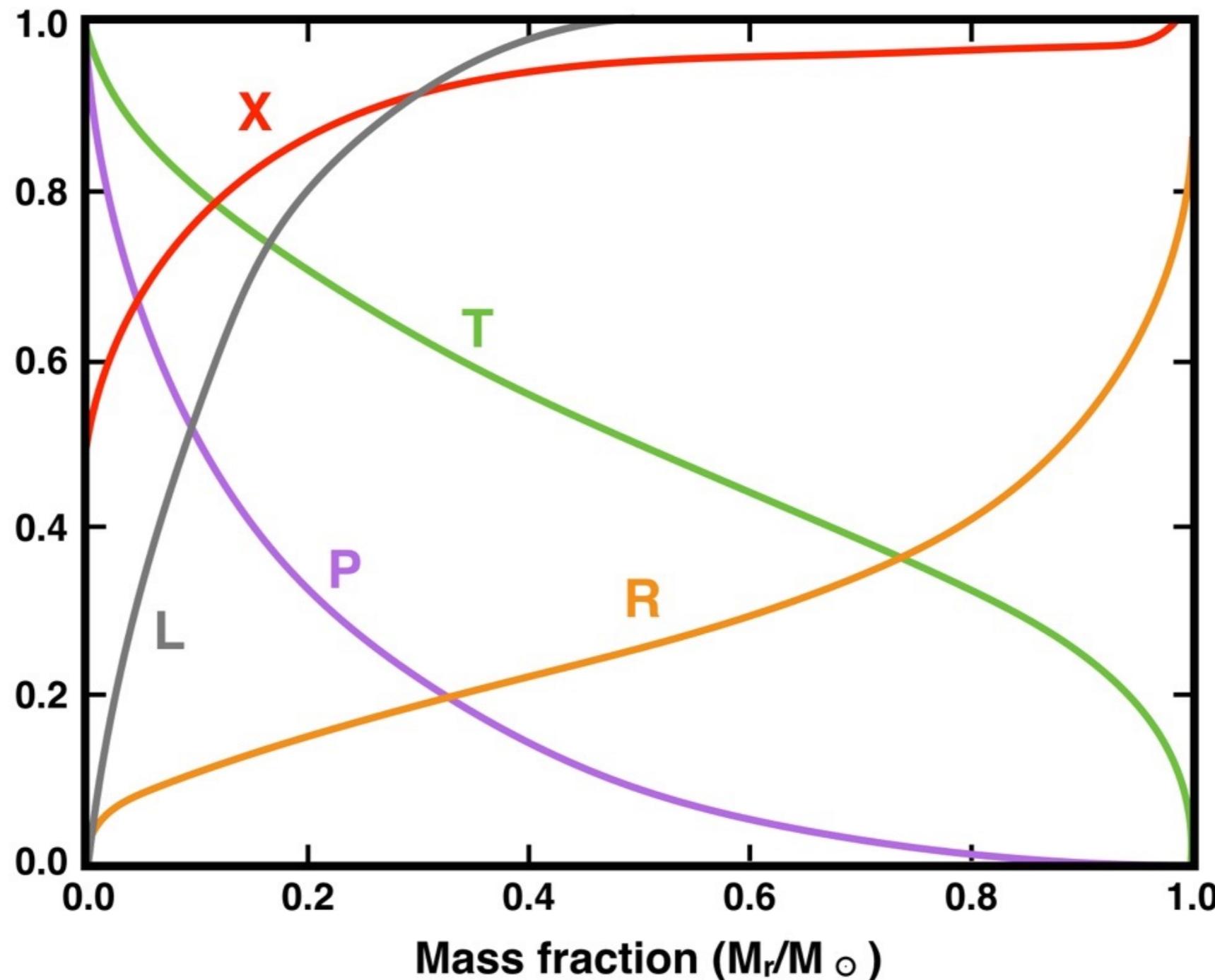
**Note.** Units of  $T_{\text{eff}}$  are in K,  $T_c$  are in MK,  $P_c$  are in  $10^{17}$  dyn cm $^{-2}$ , and  $\rho_c$  are in g cm $^{-3}$ .



# Main Sequence Evolution

## Changes during H fusion

Example: Sun-like star...



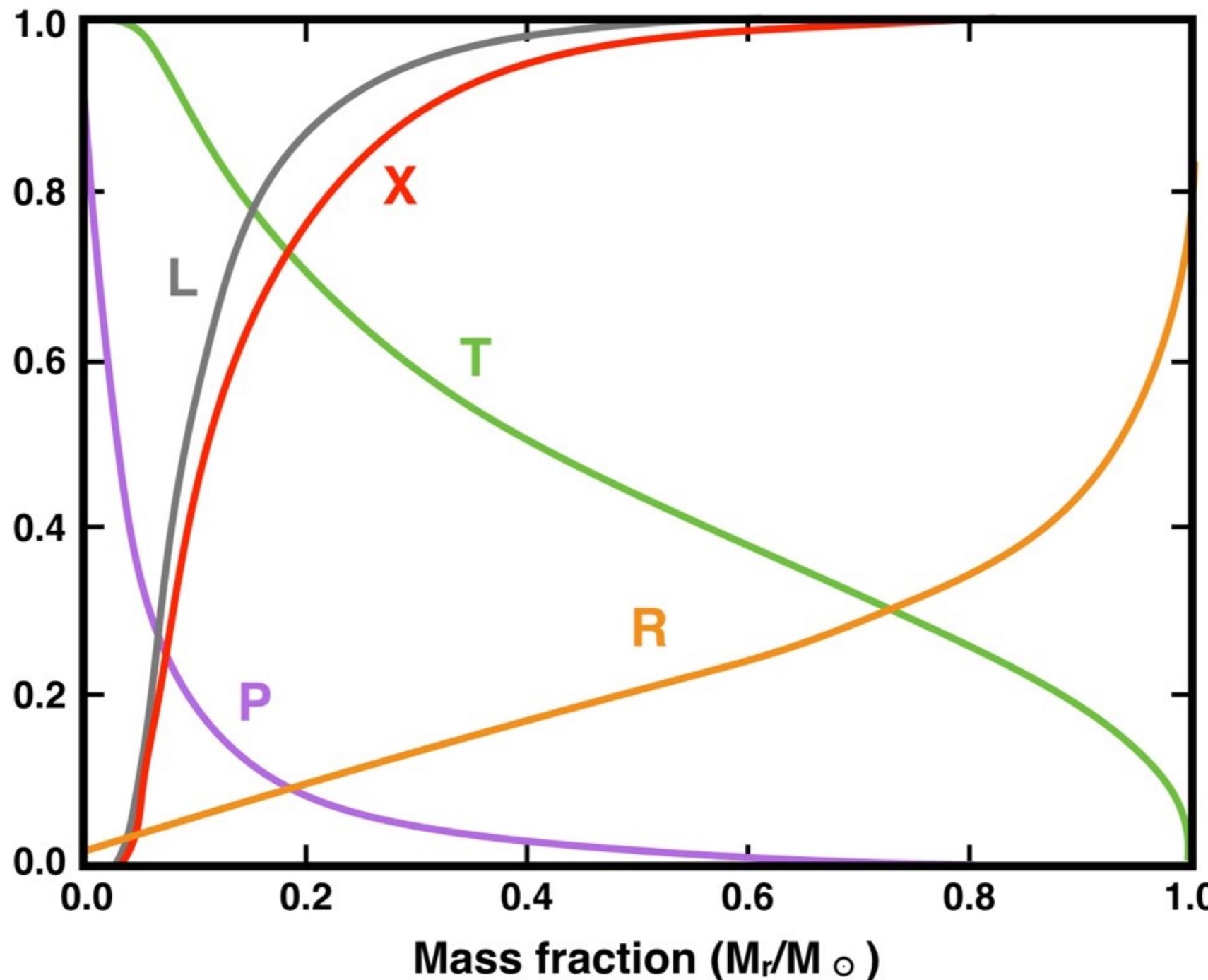
at 4.5 Gyr (sun now)

Reproduced from Iben J, 'Stellar Evolution.VI. Evolution from the Main Sequence to the Red-Giant Branch for Stars of Mass  $1 M_\odot$ ,  $1.25 M_\odot$ , and  $1.5 M_\odot$ ', Astrophysical Journal, vol. 147, p.624, 1967."

# Main Sequence Evolution

## Changes during H fusion

Example: Sun-like star...

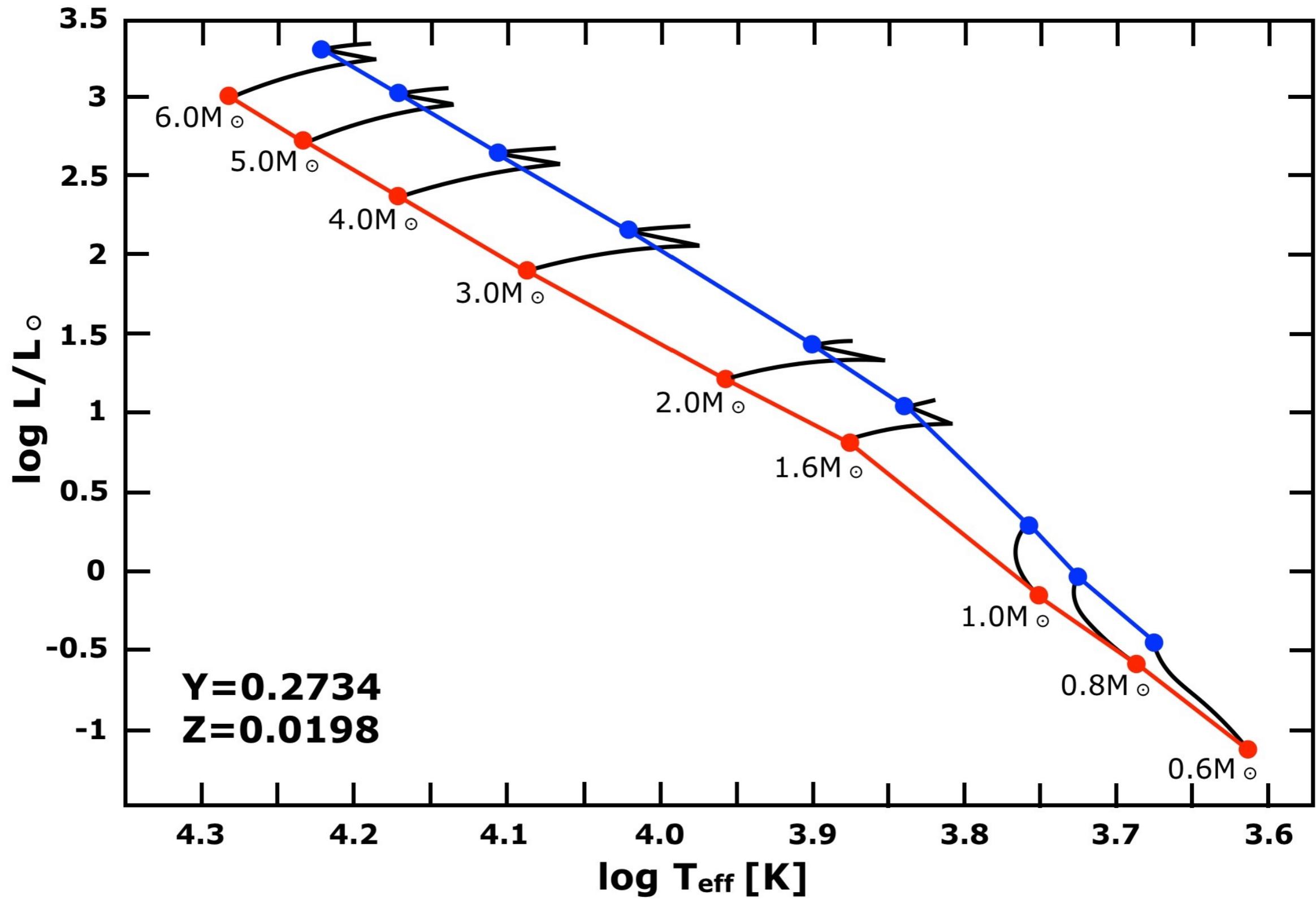


at 9.2 Gyr (TAMS)

Reproduced from Iben J, 'Stellar Evolution.VI. Evolution from the Main Sequence to the Red-Giant Branch for Stars of Mass  $1 M_\odot$ ,  $1.25 M_\odot$ , and  $1.5 M_\odot$ ', Astrophysical Journal, vol. 147, p.624, 1967."

# Main Sequence Evolution

## Changes during H fusion



# Main Sequence Evolution

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(from  $P_c/\rho_c \sim T_c/\mu_c$ ,  $\rho_c$  increases more than  $P_c$ )

R: increases (denser core → envelope expands; virial eq.)

L: increases (“μ-effect”; higher  $\rho$  increases fusion rate)

T<sub>eff</sub>: from  $T_{\text{eff}} \sim (L/R^2)^{1/4}$

( $R^2$  increases more than  $L$ ),  $T_{\text{eff}}$  *decreases*

# Main Sequence Evolution

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$\mu$ : increases ( $H \rightarrow He$ )

(from  $P_c/\rho_c \sim T_c/\mu_c$ ,  $\rho_c$  increases more than  $P_c$ )

$R$ : increases (denser core  $\rightarrow$  envelope expands; virial eq.)

$L$ : increases (" $\mu$ -effect"; higher  $\rho$  increases fusion rate)

$T_{\text{eff}}$ : from  $T_{\text{eff}} \sim (L/R^2)^{1/4}$

( $R^2$  increases more than  $L$ ),  $T_{\text{eff}}$  *decreases*

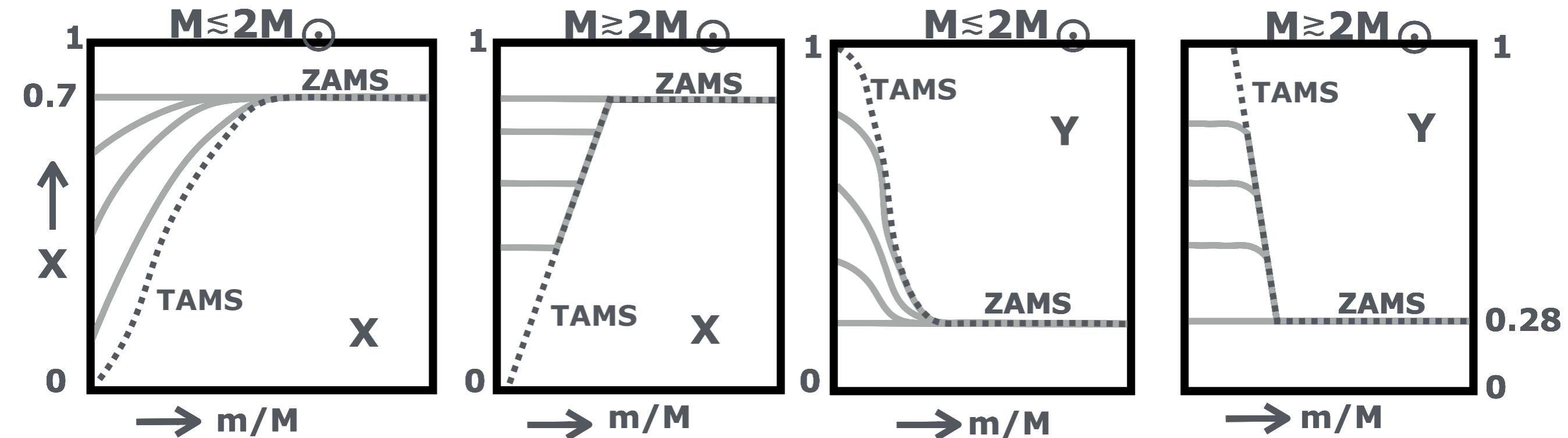
Chemical profile (for  $M \geq 1.2M_\odot$ ): H fusion happens via

T-sensitive CNO cycle and in a convective core...

# Main Sequence Evolution

## Changes during H fusion

- Chemical profile (for  $M \gtrsim 1.2M_{\odot}$ ): H fusion happens via T-sensitive CNO cycle and in a convective core...
- i. convection refuels core with fresh H, extends MS lifetime
  - ii. chemical profile is flat in the center (as opposed to peaked for lower-mass stars)
  - iii. as core H-fraction decreases, convection zone shrinks in mass



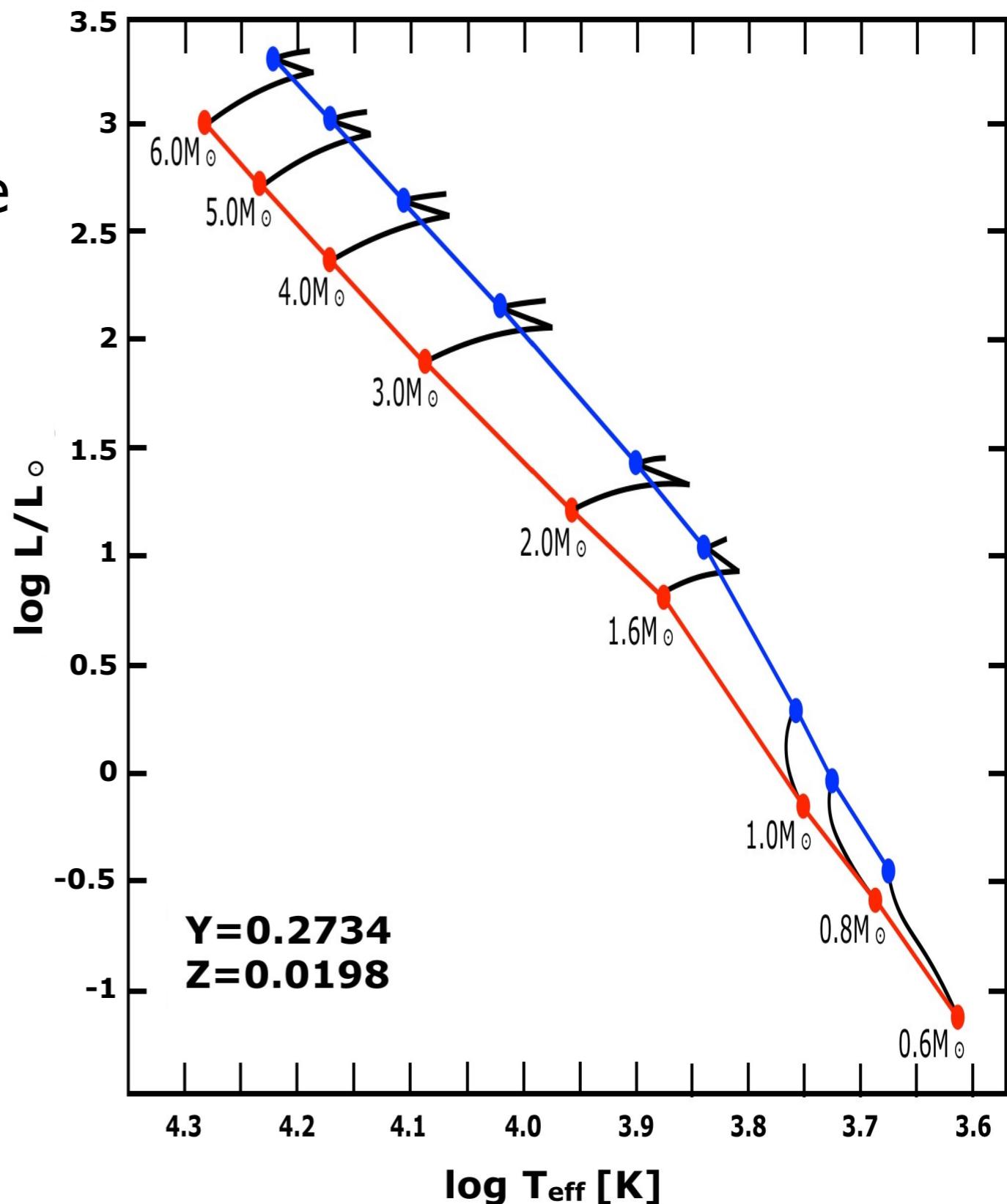
# Main Sequence Evolution

## End of the main sequence

When H fusion stops, the core contracts to compensate for radiative losses.

$M < 1.2M_{\odot}$ : core contracts gradually as H-fusion slowly peters out

$M \gtrsim 1.2M_{\odot}$ : contraction happens suddenly; H is exhausted all at once due to convective mixing



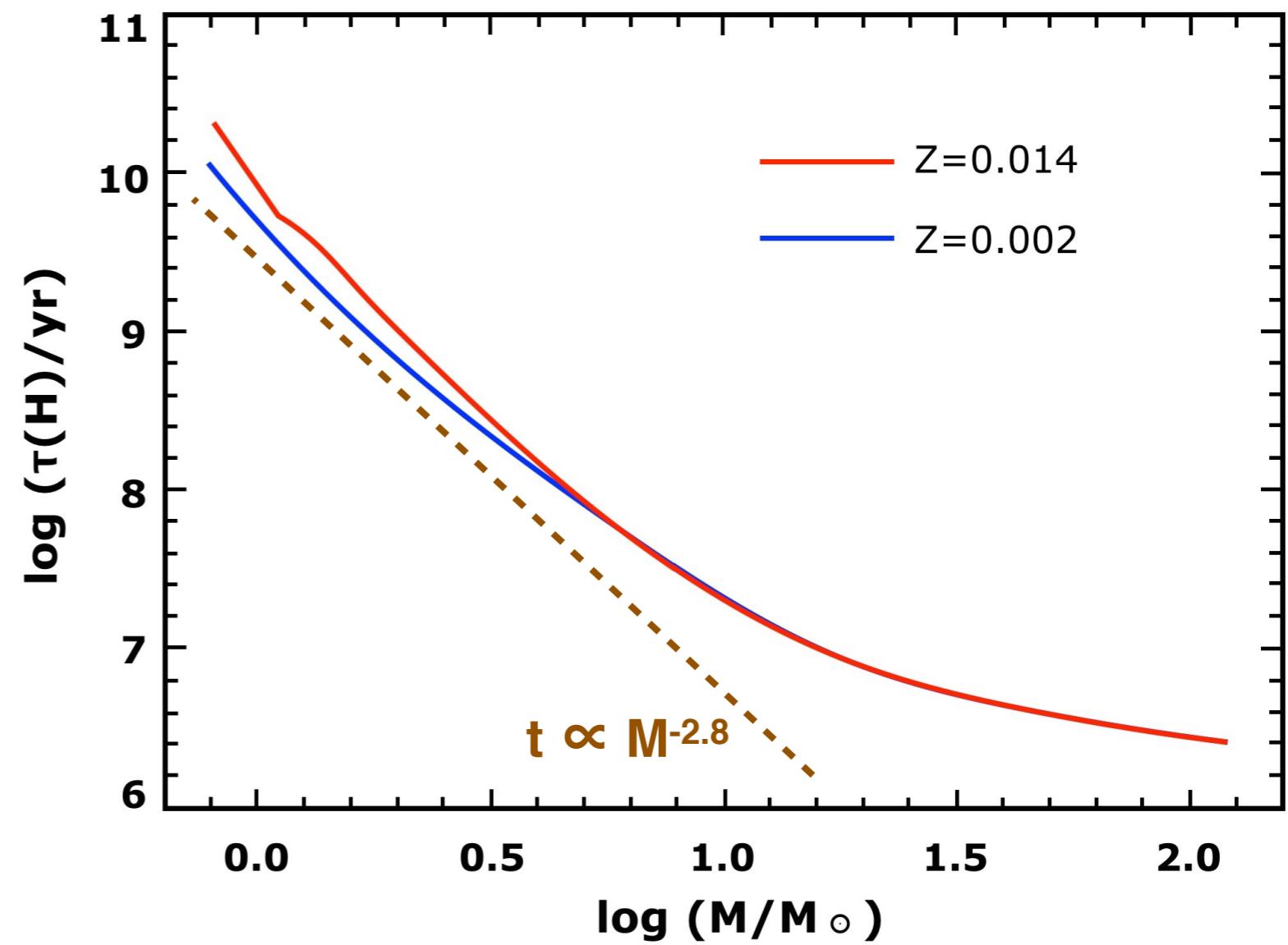
# Main Sequence Evolution

## The main sequence lifetime

Duration of the MS phase depends on M and L. MS lifetimes would be proportional to  $M/L$  (with  $L \sim M^{3.8}$  so  $t \sim M^{-2.8}$ ) if the same mass fraction of all stars was used for core H fusion.

It is not quite so steep because of two effects:

- 1)  $T_c$  increases with M, so H fusion is possible in a larger mass fraction.
- 2) If  $M_i > 1.2M_\odot$ , core convection brings material into the core fusion region.



For  $z=0.014$

$$\log \tau(H)/\text{yr} = 9.9551 - 3.3674 x + 0.4794 x^2 + 0.3676 x^3 - 0.1013 x^4, \quad x = \log(M/M_\odot)$$