

CMS-EXPERIMENT
JET REKONSTRUKTION
DATEN SIMULATION ÜEBEREINSTIMMUNG
DOMAIN ADAPTATION

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(Dr. Matthias Ulrich Mozer)

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Karlsruhe, den 03. September 2017

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Deutsche Zusammenfassung

CMS EXPERIMENT
JET RECONSTRUCTION
DOMAIN ADAPTATION

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MASTER THESIS

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Introduction

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1 Theory

2 The Compact Muon Solenoid Experiment at the Large Hadron Collider

Particle accelerator experiments have a great story of success in the investigation of physical laws. With these machines it was able to prove the existence of different elementary particles of the standard model. Particle accelerators use the electric charge of particles to accelerate them near to the speed of light. One kind of accelerator forces particles which move in opposite direction to collide at some point. The collision releases a high amount of energy and because of the laws of physics, heavy particles are produced which are highly unstable and decay after a short amount of time. To measure these particles, directly or indirectly, detectors are used.

This chapter gives a brief overview of the accelerator and detector system that has produced the events used in this work.

2.1 The Large Hadron Collider

The Large Hadron Collider (LHC) [32] [?] is a circular particle accelerator and located at the European Organization for Nuclear Research (CERN) in the region of Geneva. With a circumference of 27 km it is the world largest and most powerful particle accelerator and lies in a tunnel about 100 meters beneath the earth. The LHC is used to accelerate protons, xenon ions and lead ions. These particles undergo several steps of creation and pre-acceleration before they are injected in the LHC. An illustration of the system is shown in figure 2.1.

The LHC was constructed with two beam pipes, this allows the acceleration of protons in opposite direction. In this setup, for head on head collisions with equal beam energy, the center of mass energy is just twice the beam energy:

$$\sqrt{s} = 2 \cdot E \quad (2.1)$$

In 2010, the first collisions with a beam energy of 3.5 TeV and therefore $\sqrt{s} = 7$ TeV started. On this energy, the LHC was run until 2012 the center of mass energy was raised to $\sqrt{s} = 8$ TeV. From 2013 to 2015 the LHC was shut down for maintenance. After this the run 2 started in 2015 with $\sqrt{s} = 13$ TeV. After a shutdown in 2018, the LHC is planned to run on $\sqrt{s} = 14$ TeV in 2021 for run 3 (as of 2018).

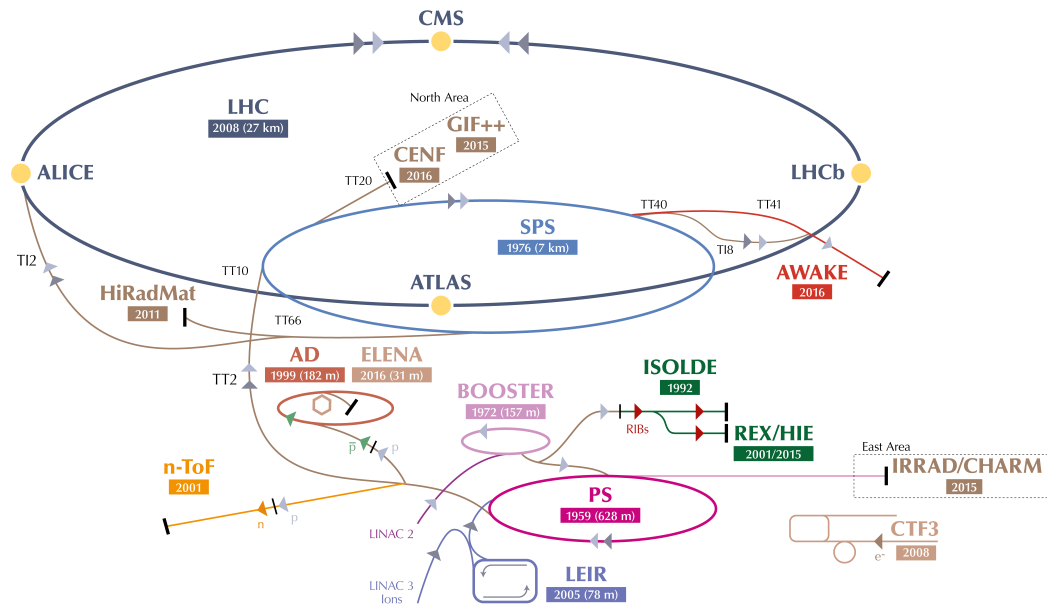


Figure 2.1: The CERN accelerator complex. The figure shows several accelerator facilities at CERN. The protons which end up in the LHC are obtained from hydrogen atoms and first accelerated with the LINAC2 to 50 MeV. Next they are injected into the Booster and accelerated to 1.4 GeV. It follows the acceleration of the Proton Synchrotron (PS) to 25 GeV and the Super Proton Synchrotron (SPS) to 450 GeV. Finally the protons are injected into the LHC in both directions. From there they reach a final energy of 6.5 TeV within 20 minutes. At the LHC the four main detectors are shown. Source: [?]

For run 2, the protons in the beam are packed in 2808 bunches, each bunch containing about $1.2 \cdot 10^{11}$ of them. For each beam, 8 superconducting cavities, cooled down to 4.5 K, generate an electromagnetic potential with a frequency of 400 MHz to accelerate and tighten the bunches. Each of the cavities is delivering 2 MV of energy. To avoid collisions with gas molecules during operation, the beam pipes have a ultrahigh vacuum pressure of 10^{-13} atm. 1232 dipole magnets with up to 7.74 T each, force the particle beam to stay on the circular path. Another 392 quadrupole magnets are used to focus the beam. To keep the magnets in an superconducting state, superfluid helium is used to cool them down to 1.9 K.

Besides the center of mass energy, the luminosity is another important quantity for particle accelerators. The luminosity L determines the event rate and therefore, a high **Luminosity** is desired:

$$\frac{dN}{dt} = \sigma L \quad (2.2)$$

Where σ is the cross section and is given by the physical laws. The luminosity is dependent of the properties of the particle accelerator:

$$L = \frac{nN_1N_2f}{A} \quad (2.3)$$

Where, n is the total number of bunches in the collider, N_1 and N_2 are the number of particles in the colliding bunches, f is the circulating frequency for one single bunch and A is the cross sectional area. On this formula you can see the importance of having a good focusing, that will shrink the A . To get the total number of events, one can integrate formula 2.1 over time to

$$N = \sigma \int L dt = \sigma L_{int} \quad (2.4)$$

Where L_{int} is the integrated luminosity. In 2017 a peak luminosity of $2.06 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ was reached with an integrated luminosity of 50 fb^{-1} ($1 \text{ b} = 10^{-24} \text{ cm}^2$) for the total year [?]. After run 3 at the LHC, a high luminosity upgrade is planned to increase the luminosity to $5 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$.

The particle beams at the LHC are crossing at four different points, where at each point one detector is located. One is for the ALICE experiment which studies collisions of lead ions. Another one is for the LHCb experiment which is designed for the studies of b hadron physics. The ALICE and the CMS experiments are general purpose experiments for the measurement of the standard model particles and the search for new elementary particles.

2.2 The Compact Muon Solenoid Experiment

The Compact Muon Solenoid (CMS) detector was built up with the main motivation in the search for the higgs boson and the search for physics beyond the standard model. As

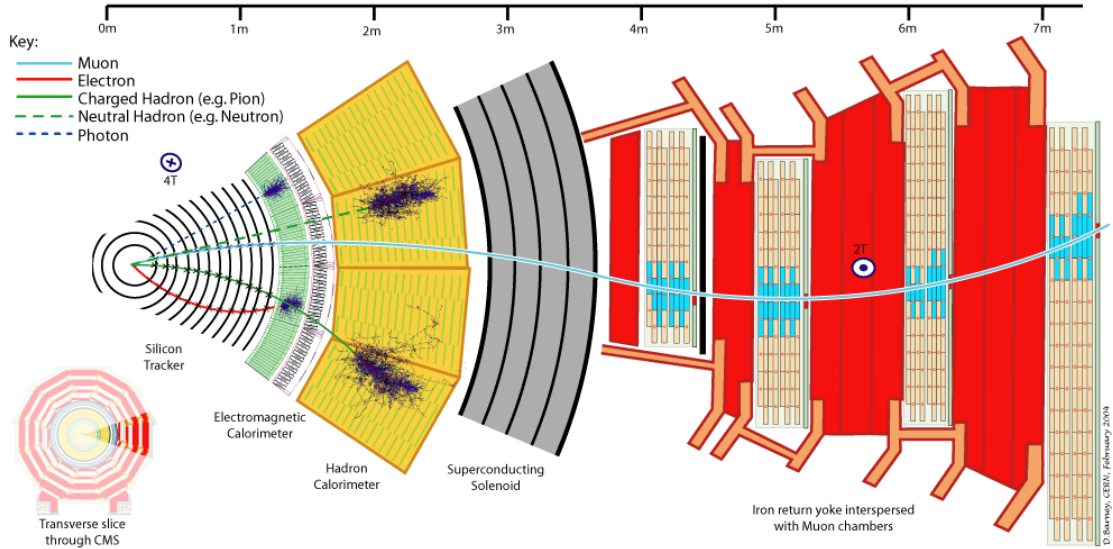


Figure 2.2: Slice of the CMS detector The scheme shows the different detector systems as well as some particle tracks and their interactions with the detector. Source: [40]

we know, the former has already been found, but we still search for alternative theories like supersymmetry or extra dimensions. Besides of those two big tasks, the CMS detector is used for a lot of other studies. For example the measurement of couplings of different particles [?]. In order to achieve these goals, the detector is required to measure the involved particles with best possible precision in position and energy. A sketch of a slice of the CMS detector with it's different layers is shown in figure ??.

The detector also needs a filter system, the so called trigger system, to store only the interesting events and reduce the amount of data to an amount that can be stored for further analysis. In this chapter, these systems are described shortly, in the following the typically used coordinate system is outlined.

coordinate system

The origin of the coordinate system is chosen to be the intended collision point. The z -axis points in the beam direction where the x lies in the plane of the collider ring and points to its center. The y axis points vertically upwards. The radial distance is $r = \sqrt{x^2 + y^2}$. ϕ is used as the azimuthal angle in the x - y plane measured from the x -axis. The polar angle θ is measured from the z -axis. Another important quantity is the pseudorapidity η . It is used to denote the angle between a vector and the z -axis:

$$\eta = -\ln \left[\tan \left(\frac{\theta}{2} \right) \right] \quad (2.5)$$

A value of $\eta = 0$ is equal to $\theta = 90^\circ$, an angle of $\theta = 0^\circ$ equals $\eta = \infty$. It is preferred over θ since the number of produced particles per η interval is constant. Because the momentum and energy of the initial colliding particles is not determined in a hadron collider, most of the time only the transversal quantities to the beam axis are of interest. These quantities are invariant under lorentz boosts in z -direction and therefore be used. They are given as $p_T = \sqrt{p_x^2 + p_y^2}$ and $E_T = \sqrt{E_x^2 + E_y^2}$ respectively. All transverse momentum and energy should sum up to one, in the measurement. The missing amount is denoted as p_T^{miss} and E_T^{miss} .

Silicon Tracker

The most interesting things in particle physics happen immediately after the collision. Heavy objects decay in less than nanoseconds and with sufficient effort some can be reconstructed. The innermost part of the detector is therefore of highest importance and has several requirements to fulfill. First, a good measurement of trajectories of charged particles is necessary to identify primary and secondary vertices. This also allows a momentum measurement as the full tracker is inside the CMS solenoid. This solenoid produces a homogeneous magnetic field of 4 T. Next, at each bunch crossing, in average about 1000 particles are created and travel through the detector. To distinguish the particles from each other and to assign the trajectories to the corresponding vertices, a high granularity is needed. Further, the time between two bunch crossings is about 25 ns therefore the response time and the cooling down period has to be short. Last but not least, the detector components have to be resilient against this high amount of radiation.

To fulfill these requirements, silicon technology is used for the whole tracker system. When charged particles travel through this materials, electron-hole pairs are produced, which in turn generates an electric current in the read out electronics. In the innermost part from $r = 4.4$ cm to 10.2 cm, three barrel layers of pixel detectors are surrounding the beam axis. Each pixel detector has a size of $100 \times 150 \mu\text{m}^2$. With this design, the occupancy for each pixel and bunch crossing is of the order of 10^{-4} . The pixel detectors are surrounded by 10 layers of silicon micro-strip detectors up to a radius of 1.1 m. As the particle flow decreases with the radius, larger detector elements can be used. This also reduces the amount of read out electronics. The inner micro strip detectors have a size of $10 \text{ cm} \times 80 \mu\text{m}$ where the size of the outer one is $25 \text{ cm} \times 180 \mu\text{m}$. At the endcaps, 2 disks of pixel detectors and 3 smaller and 9 larger discs of strip detectors are installed. This allows a reconstruction of charged particle tracks up to $|\eta| < 2.5$. In total the silicon tracker has a length of 5.8 m with a diameter of 2.5 m. At each bunch crossing, the occupancy of the detectors is at about 2-3 % for the inner micro strip detectors and lower for the other one. For high energetic muons (100 GeV), the resolution for the impact parameter, which is the distance of a reconstructed track from a vertex, is about $10 \mu\text{m}$.

Electromagnetic Calorimeter

The electromagnetic calorimeter (ECAL) surrounds the silicon tracker system hermetically and measures the energy of electrons, positrons and photons. Other **electromagnetic** interacting particles leave tracks but are not fully absorbed in the ECAL. The ECAL consists of homogeneous lead tungstate (PbWO_4) crystals which is the absorber and scintillator material at the same time. When electromagnetically interacting particles propagate through it, they produce photons through bremsstrahlung, which in turn create electron-positron pairs through pair production. An electromagnetic shower emerges. When the energy of the electrons drops to a critical amount, they mainly excite the atoms which emit photons with characteristic wavelength (scintillation). These photons are then measured **with photomultiplier**. The number of photons is proportional to the deposited energy of the particle. The radiation length of PbWO_4 is $X_0 = 0.89$ cm. This means that after this distance, the energy of an electron or positron **reduced** by the factor of $1/e$. The thickness of the ECAL with 23 cm is equivalent to $25.8 X_0$, therefore all electrons **positrons** and photons are absorbed completely. The ECAL at the CMS **was** build with the main motivation for a **good measurement** of the $H \rightarrow \gamma\gamma$ channel. Therefore it has a good energy resolution. To reduce the background in this channel, a preshower detector is installed just before the ECAL. This element helps to identify neutral pions that decay into two photons, **it** also improves the position determination of electrons, positrons and photons.

Hadronic Calorimeter

Hadrons propagate through the ECAL **with** losing only **small parts** of their energy, therefore an hadronic calorimeter (HCAL) is surrounding the ECAL which absorbs all hadrons and measures their energy. Therefore **it's** important for measuring hadron jets as well as the total energy of the system and the resulting missing transverse energy. Several elements allow to measure forward jets up to $|\eta| = 5.2$.

The HCAL at CMS is a sampling calorimeter which exists of alternating absorbing layers and scintillating layers. The absorbing layer, **the so called brass**, consists in main parts of copper and tin. **The** brass has a nuclear interaction length of $\lambda_I = 16.42$ cm. This is the mean distance for a hadronic particle after which it undergoes an inelastic nuclear interaction. Similar to the **ecal**, a chain reaction starts and an hadronic shower emerges. The total absorber material ECAL and HCAL makes up almost $7 \lambda_I$ and is therefore sufficient to absorb **all hadrons**. For the scintillating layers, plastic is used. When the hadronic shower hits these layers, the scintillator material gets excited. Photons with characteristic wavelength are emitted. They are guided through optical cables and measured by photodiodes.

Muon Detector

The muon detector is the part furthest from the collision point. The only particles that reach this part of the detector are neutrinos, which can not be detected with the CMS detector, and muons. Muons have a high mass and often high energy. Because of this, they lose only few energy by propagating through the inner parts and are usually not totally absorbed in the muon detector either. Muon identification and position determination is done with gaseous detectors. As they pass through these detectors, muons are ionizing the gas atoms, the resulting electrons and ions can then be measured as an electric current. The muon system consists of a iron return yoke to guide the magnetic field lines of the solenoid through this part of the detector. This creates a magnetic field of 2 T. As a consequence, muons travel a circular path. The momentum can then be determined by measuring the radius.

Trigger System

The LHC has a pp interaction rate of more than 1 GHz at each crossing point. This results in too much data to store every event. Anyway only in small parts, interesting physic processes occur. To reduce the event rate to a storable amount by selecting only interesting events, a trigger system consisting of two steps is used at the CMS [?]. The first selection is called the Level-1 (L1) trigger and is realized by hardware electronics. It has to decide in 4 μ s if the event is worth to keep. In this short time only parts of information from the calorimeters and the muon chamber can be used. The L1 trigger decides if the information is consistent with a physics objects like an electron, muon or others. At the end the event rate is reduced to less than 100 MHz.

The second trigger selection step is done by the High-Level Trigger (HLT). The HLT is realized with software tools where complex calculations and a more precise reconstruction can be done. Further is the information of the tracking system now included.

This selection reduces the event rate further to an average of 400 Hz. These events are then stored offline for further analysis.



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Bibliography

- [1] G. Aad *et al.*, “Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC,” *Physics Letters B*, vol. 716, pp. 1–29, 2012.
- [2] S. Chatrchyan *et al.*, “Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC,” *Physics Letters B*, vol. 716, pp. 30–61, 2012.
- [3] The CMS Collaboration, “Search for $H \rightarrow b\bar{b}$ in association with a single top quark as a test of Higgs boson couplings at $\sqrt{s} = 13$ TeV,” CMS-PAS-HIG-16-019, CERN, Geneva, 2016.
- [4] The CMS Collaboration, “Search for $H \rightarrow b\bar{b}$ in association with single top quarks as a test of Higgs couplings,” CMS Physics Analysis Summary CMS-PAS-HIG-14-015, CERN, Geneva, 2014.
- [5] F. Englert and R. Brout, “Broken Symmetry and the Mass of Gauge Vector Mesons,” *Physical Review Letters*, vol. 13, pp. 321–323, 1964.
- [6] P. W. Higgs, “Broken Symmetries and the Masses of Gauge Bosons,” *Physical Review Letters*, vol. 13, pp. 508–509, 1964.
- [7] G. S. Guralnik, C. R. Hagen, and T. W. B. Kibble, “Global conservation laws and massless particles,” *Physical Review Letters*, vol. 13, pp. 585–587, 1964.
- [8] A. Einstein, “Die Grundlage der allgemeinen Relativitätstheorie,” *Annalen der Physik*, vol. 49, pp. 769–822, 1916.
- [9] K. A. Olive, “Review of Particle Physics,” *Chinese Physics C*, vol. 40, no. 10, p. 100001, 2016.
- [10] M. Gell-Mann and Y. Neeman, “The eightfold way,” WA Benjamin, Inc., New York, 1964.
- [11] H. Fritzsch, M. Gell-Mann, and H. Leutwyler, “Advantages of the color octet gluon picture,” *Physics Letters B*, vol. 47, no. 4, pp. 365 – 368, 1973.
- [12] S. Tomonaga, “On a relativistically invariant formulation of the quantum theory of wave fields,” *Progress of Theoretical Physics*, vol. 1, no. 2, pp. 27–42, 1946.

- [13] J. Schwinger, “Quantum electrodynamics. i. a covariant formulation,” *Physical Review*, vol. 74, pp. 1439–1461, 1948.
- [14] R. P. Feynman, “Space-time approach to quantum electrodynamics,” *Physical Review*, vol. 76, pp. 769–789, 1949.
- [15] S. Weinberg, “A model of leptons,” *Physical Review Letters*, vol. 19, pp. 1264–1266, 1967.
- [16] A. Salam, “Proceedings of the 8th Nobel Symposium, Stockholm, 1968,” *Almqvist and Wiksells*, 1968.
- [17] S. L. Glashow, J. Iliopoulos, and L. Maiani, “Weak interactions with lepton-hadron symmetry,” *Physical Review D*, vol. 2, pp. 1285–1292, 1970.
- [18] Y. Fukuda *et al.*, “Evidence for oscillation of atmospheric neutrinos,” *Physical Review Letters*, vol. 81, pp. 1562–1567, 1998.
- [19] J. Ellis, “Outstanding questions: physics beyond the standard model,” *Philosophical Transactions of the Royal Society A*, vol. 370, no. 1961, pp. 818–830, 2012.
- [20] J. D. Lykken, “Beyond the Standard Model,” in *CERN Yellow Report CERN-2010-002*, 101-109, 2010.
- [21] M. Peskin and D. Schroeder, “An introduction to quantum field theory,” Westview Press, Boulder, CO, 1995.
- [22] S. Dawson, “Introduction to electroweak symmetry breaking,” in *Proceedings, Summer School in High-energy physics and cosmology: Trieste, Italy, June 29-July 17, 1998*, pp. 1–83, 1998.
- [23] A. Zee, “Quantum field theory in a nutshell,” Princeton university press, Princeton, NJ, 2010.
- [24] L. Álvarez Gaumé and J. Ellis, “Eyes on a prize particle,” *Nature Physics*, vol. 7, no. 1, pp. 2–3, 2011.
- [25] M. Farina, C. Grojean, F. Maltoni, E. Salvioni, and A. Thamm, “Lifting degeneracies in Higgs couplings using single top production in association with a Higgs boson,” *Journal of High Energy Physics*, vol. 05, p. 022, 2013.
- [26] F. Demartin, F. Maltoni, K. Mawatari, and M. Zaro, “Higgs production in association with a single top quark at the LHC,” *The European Physical Journal C*, vol. 75, no. 6, p. 267, 2015.
- [27] F. Demartin, B. Maier, F. Maltoni, K. Mawatari, and M. Zaro, “tWH associated production at the LHC,” *arXiv:1607.05862*, 2016.

-
- [28] S. Dittmaier *et al.*, “Handbook of LHC Higgs Cross Sections: 1. Inclusive Observables,” *arXiv:1101.0593*, 2011.
- [29] F. Abe *et al.*, “Observation of top quark production in $\bar{p}p$ collisions,” *Physical Review Letters*, vol. 74, pp. 2626–2631, 1995.
- [30] S. Abachi *et al.*, “Observation of the top quark,” *Physical Review Letters*, vol. 74, pp. 2632–2637, 1995.
- [31] The ATLAS, CDF, CMS and DØ Collaborations, “First combination of Tevatron and LHC measurements of the top-quark mass,” *arXiv:1403.4427*, 2014.
- [32] C. Lefevre, “LHC: the guide (English version). Guide du LHC (version anglaise),” 2009.
- [33] CERN, “The Large Hadron Collider,” 2014. <http://home.cern/topics/large-hadron-collider>, last accessed on 2016-10-23.
- [34] C. De Melis, “The CERN accelerator complex. Complexe des accélérateurs du CERN,” 2016. <https://cds.cern.ch/record/2197559>, last accessed on 2016-10-18.
- [35] CERN, “LHC experiments back in business at record energy,” 2015. <http://home.cern/about/updates/2015/06/lhc-experiments-back-business-record-energy>, last accessed on 2016-10-23.
- [36] CERN, “LHC performance reaches new highs,” 2016. <https://home.cern/about/updates/2016/07/lhc-performance-reaches-new-highs>, last accessed on 2016-10-23.
- [37] The CMS Collaboration, “The CMS experiment at the CERN LHC,” *Journal of Instrumentation*, vol. 3, no. 08, p. S08004, 2008.
- [38] CERN, “CMS: The Compact Muon Solenoid,” 2012. <https://cds.cern.ch/record/1997263>, last accessed on 2016-11-22.
- [39] The CMS Collaboration, “Precise mapping of the magnetic field in the CMS barrel yoke using cosmic rays,” *Journal of Instrumentation*, vol. 5, no. 03, p. T03021, 2010.
- [40] S. R. Davis, “Interactive Slice of the CMS detector,” 2016. <https://cds.cern.ch/record/2205172>, last accessed on 2016-10-18.
- [41] V. Karimäki, M. Mannelli, P. Siegrist, H. Breuker, A. Caner, R. Castaldi, K. Freudenreich, G. Hall, R. Horisberger, M. Huhtinen, and A. Cattai, “The CMS tracker

- system project: Technical Design Report,” Technical Design Report CMS, CERN, Geneva, 1997.
- [42] The CMS Collaboration, “The CMS tracker: addendum to the Technical Design Report,” Technical Design Report CMS, CERN, Geneva, 2000.
 - [43] The CMS Collaboration, “Tracker detector,” 2011. <http://cms.web.cern.ch/news/tracker-detector>, last accessed on 2016-10-24.
 - [44] The CMS Collaboration, “The CMS electromagnetic calorimeter project: Technical Design Report,” Technical Design Report CMS, CERN, Geneva, 1997.
 - [45] P. Bloch, R. Brown, P. Lecoq, and H. Rykaczewski, “Changes to CMS ECAL electronics: addendum to the Technical Design Report,” Technical Design Report CMS, CERN, Geneva, 2002.
 - [46] The CMS Collaboration, “Electromagnetic Calorimeter,” 2011. <http://cms.web.cern.ch/news/electromagnetic-calorimeter>, last accessed on 2016-10-24.
 - [47] The CMS Collaboration, “The CMS hadron calorimeter project: Technical Design Report,” Technical Design Report CMS, CERN, Geneva, 1997.
 - [48] The CMS Collaboration, “Hadron Calorimeter,” 2011. <http://cms.web.cern.ch/news/hadron-calorimeter>, last accessed on 2016-10-24.
 - [49] The CMS Collaboration, “The CMS muon project: Technical Design Report,” Technical Design Report CMS, CERN, Geneva, 1997.
 - [50] The CMS Collaboration, “Muon Detectors,” 2011. <http://cms.web.cern.ch/news/muon-detectors>, last accessed on 2016-10-25.
 - [51] V. Khachatryan *et al.*, “The CMS trigger system,” *Submitted to: Journal of Instrumentation*, *arXiv:1609.02366*, 2016.
 - [52] D. Acosta, “CMS Trigger Improvements Towards Run II,” *Nuclear and Particle Physics Proceedings*, vol. 273–275, pp. 1008 – 1013, 2016. 37th International Conference on High Energy Physics (ICHEP).
 - [53] C. Eck *et al.*, “LHC computing Grid: Technical Design Report. Version 1.06 (20 Jun 2005),” Technical Design Report LCG, CERN, Geneva, 2005.
 - [54] I. Bird *et al.*, “Update of the Computing Models of the WLCG and the LHC Experiments,” CERN-LHCC-2014-014. LCG-TDR-002, 2014.
 - [55] WLCG Office, “Worldwide LHC Computing Grid (WLCG) – Tier centres,” 2016. <http://wlcg-public.web.cern.ch/tier-centres>, last accessed on 2016-10-25.

-
- [56] The CMS Collaboration, “Computing Grid,” 2011. <http://cms.web.cern.ch/news/computing-grid>, last accessed on 2016-10-25.
- [57] WLCG Project Office, “Documents & Reference - Tiers & Grid Images (2014),” 2014. <http://wlcg.web.cern.ch/documents-reference>, last accessed on 2016-10-18.
- [58] S. Höche, “Introduction to parton-shower event generators,” in *Theoretical Advanced Study Institute in Elementary Particle Physics: Journeys Through the Precision Frontier: Amplitudes for Colliders (TASI 2014) Boulder, Colorado, June 2-27, 2014*, 2014.
- [59] V. Gribov and L. Lipatov, “Deep inelastic $e p$ scattering in perturbation theory,” *Soviet Journal of Nuclear Physics*, vol. 15, pp. 438–450, 1972.
- [60] G. Altarelli and G. Parisi, “Asymptotic Freedom in Parton Language,” *Nuclear Physics B*, vol. 126, p. 298, 1977.
- [61] Y. Dokshitzer, “Calculation of the structure functions for deep inelastic scattering and e^+e^- annihilation by perturbation theory in quantum chromodynamics,” *Zh. Eksp. Teor. Fiz*, vol. 73, p. 1216, 1977.
- [62] J. S. Gainer, J. Lykken, K. T. Matchev, S. Mrenna, and M. Park, “The Matrix Element Method: Past, Present, and Future,” in *Proceedings, Community Summer Study 2013: Snowmass on the Mississippi (CSS2013): Minneapolis, MN, USA, July 29-August 6, 2013*, 2013.
- [63] The NNPDF Collaboration, “Neural Network Parton Distribution Functions (NNPDF) – Images,” 2016. <https://nnpdf.hepforge.org/images/>, last accessed on 2016-10-26.
- [64] B. Andersson, G. Gustafson, G. Ingelman, and T. Sjöstrand, “Parton Fragmentation and String Dynamics,” *Physics Reports*, vol. 97, pp. 31–145, 1983.
- [65] N. Metropolis and S. Ulam, “The Monte Carlo Method,” *Journal of the American Statistical Association*, vol. 44, no. 247, pp. 335–341, 1949.
- [66] J. Alwall, M. Herquet, F. Maltoni, O. Mattelaer, and T. Stelzer, “MadGraph 5 : Going Beyond,” *Journal of High Energy Physics*, vol. 06, p. 128, 2011.
- [67] F. Maltoni and T. Stelzer, “MadEvent: Automatic event generation with MadGraph,” *Journal of High Energy Physics*, vol. 02, p. 027, 2003.
- [68] J. Alwall, R. Frederix, S. Frixione, V. Hirschi, F. Maltoni, *et al.*, “The automated computation of tree-level and next-to-leading order differential cross sections, and their matching to parton shower simulations,” *Journal of High Energy Physics*, vol. 07, p. 079, 2014.

- [69] P. Nason, “A New method for combining NLO QCD with shower Monte Carlo algorithms,” *Journal of High Energy Physics*, vol. 11, p. 040, 2004.
- [70] S. Frixione, P. Nason, and C. Oleari, “Matching NLO QCD computations with Parton Shower simulations: the POWHEG method,” *Journal of High Energy Physics*, vol. 11, p. 070, 2007.
- [71] S. Alioli, P. Nason, C. Oleari, and E. Re, “A general framework for implementing NLO calculations in shower Monte Carlo programs: the POWHEG BOX,” *Journal of High Energy Physics*, vol. 06, p. 043, 2010.
- [72] T. Sjöstrand, S. Mrenna, and P. Z. Skands, “PYTHIA 6.4 Physics and Manual,” *Journal of High Energy Physics*, vol. 05, p. 026, 2006.
- [73] T. Sjöstrand, S. Mrenna, and P. Z. Skands, “A Brief Introduction to PYTHIA 8.1,” *Computer Physics Communications*, vol. 178, pp. 852–867, 2008.
- [74] S. Agostinelli *et al.*, “GEANT4: A Simulation toolkit,” *Nuclear Instruments and Methods in Physics*, vol. A506, pp. 250–303, 2003.
- [75] J. Allison *et al.*, “Geant4 developments and applications,” *IEEE Transactions on Nuclear Science*, vol. 53, no. 1, pp. 270–278, 2006.
- [76] S. Chatrchyan *et al.*, “Description and performance of track and primary-vertex reconstruction with the CMS tracker,” *Journal of Instrumentation*, vol. 9, no. 10, p. P10009, 2014.
- [77] P. Billoir, “Progressive track recognition with a Kalman-like fitting procedure,” *Computer Physics Communications*, vol. 57, no. 1–3, pp. 390 – 394, 1989.
- [78] P. Billoir and S. Qian, “Simultaneous pattern recognition and track fitting by the Kalman filtering method,” *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 294, no. 1–2, pp. 219 – 228, 1990.
- [79] R. Frühwirth, “Application of Kalman filtering to track and vertex fitting,” *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 262, no. 2–3, pp. 444 – 450, 1987.
- [80] R. Frühwirth, W. Waltenberger, and P. Vanlaer, “Adaptive Vertex Fitting,” CMS-NOTE-2007-008, CERN, Geneva, 2007.
- [81] The CMS Collaboration, “Particle-Flow Event Reconstruction in CMS and Performance for Jets, Taus, and MET,” CMS Physics Analysis Summary CMS-PAS-PFT-09-001, CERN, Geneva, 2009.

-
- [82] N. Bartosik, "PF concept in CMS," 2016. http://bartosik.pp.ua/hep_sketches/cms_particle_flow, last accessed on 2016-11-01.
- [83] W. Adam, R. Frühwirth, A. Strandlie, and T. Todorov, "Reconstruction of electrons with the Gaussian-sum filter in the CMS tracker at the LHC," *Journal of Physics G Nuclear Physics*, vol. 31, p. 9, Sept. 2005.
- [84] V. Khachatryan *et al.*, "Performance of Electron Reconstruction and Selection with the CMS Detector in Proton-Proton Collisions at $\sqrt{s} = 8$ TeV," *Journal of Instrumentation*, vol. 10, no. 06, p. P06005, 2015.
- [85] S. Baffioni, C. Charlot, F. Ferri, D. Futyan, P. Meridiani, *et al.*, "Electron reconstruction in CMS," *Eur. Phys. J. C*, vol. 49, pp. 1099–1116, 2007.
- [86] S. Chatrchyan *et al.*, "Performance of CMS muon reconstruction in pp collision events at $\sqrt{s} = 7$ TeV," *Journal of Instrumentation*, vol. 7, p. P10002, 2012.
- [87] M. Cacciari, G. P. Salam, and G. Soyez, "The Anti-k(t) jet clustering algorithm," *Journal of High Energy Physics*, vol. 04, p. 063, 2008.
- [88] V. Khachatryan *et al.*, "Jet energy scale and resolution in the CMS experiment in pp collisions at 8 TeV," *Submitted to: Journal of Instrumentation*, *arXiv:1607.03663*, 2016.
- [89] The CMS Collaboration: Jet Energy Resolution and Corrections (JERC) Subgroup, "Introduction to Jet Energy Corrections at CMS," 2016. <https://twiki.cern.ch/twiki/bin/view/CMS/IntroToJEC>, last accessed on 2016-11-03.
- [90] The DØ Collaboration, "Observation of Single Top Quark Production – B-Jet Identification," 2009. https://www-d0.fnal.gov/Run2Physics/top/singletop_observation/, last accessed on 2016-10-27.
- [91] The CMS collaboration, "Identification of b-quark jets with the CMS experiment," *Journal of Instrumentation*, vol. 8, no. 04, p. P04013, 2013.
- [92] The CMS Collaboration, "Identification of b quark jets at the CMS Experiment in the LHC Run 2," CMS-PAS-BTV-15-001, CERN, Geneva, 2016.
- [93] V. Khachatryan *et al.*, "Measurement of $B\bar{B}$ Angular Correlations based on Secondary Vertex Reconstruction at $\sqrt{s} = 7$ TeV," *Journal of High Energy Physics*, vol. 03, p. 136, 2011.
- [94] V. Blobel and E. Lohrmann, "Statistische und numerische Methoden der Datenanalyse," Teubner, Stuttgart Leipzig, 1998.

- [95] The CMS Collaboration, “Documentation of the RooStats-based statistics tools for Higgs PAG,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/SWGuideHiggsAnalysisCombinedLimit>, last accessed on 2016-11-10.
- [96] The CMS Collaboration, “Welcome to the RooStats Wiki,” 2016. <https://twiki.cern.ch/twiki/bin/view/RooStats/WebHome>, last accessed on 2016-11-10.
- [97] T. Junk, “Confidence level computation for combining searches with small statistics,” *Nuclear Instruments and Methods in Physics A*, vol. 434, pp. 435–443, 1999.
- [98] A. L. Read, “Modified frequentist analysis of search results (the CL_s method),” 2000. CERN-OPEN-2000-205.
- [99] A. L. Read, “Presentation of search results: the CL_s technique,” *Journal of Physics G: Nuclear and Particle Physics*, vol. 28, no. 10, p. 2693, 2002.
- [100] J. Neyman and E. S. Pearson, “On the Problem of the Most Efficient Tests of Statistical Hypotheses,” *Philosophical Transactions of the Royal Society of London. Series A, Containing Papers of a Mathematical or Physical Character*, vol. 231, pp. 289–337, 1933.
- [101] G. Cowan, K. Cranmer, E. Gross, and O. Vitells, “Asymptotic formulae for likelihood-based tests of new physics,” *European Physical Journal C*, vol. 71, p. 1554, 2011.
- [102] S. S. Wilks, “The Large-Sample Distribution of the Likelihood Ratio for Testing Composite Hypotheses,” *Annals of Mathematical Statistics*, vol. 9, no. 1, pp. 60–62, 1938.
- [103] L. Breiman, J. H. Friedman, R. A. Olshen, and C. J. Stone, “Classification and regression trees,” CRC Press, New York, 1984.
- [104] T. J. Hastie, R. J. Tibshirani, and J. H. Friedman, “The elements of statistical learning: data mining, inference, and prediction,” Springer, Berlin, 2009.
- [105] G. Bohm and G. Zech, “Introduction to statistics and data analysis for physicists,” Verl. Dt. Elektronen-Synchrotron, Hamburg, 2010.
- [106] A. Hoecker, P. Speckmayer, J. Stelzer, J. Therhaag, E. von Toerne, and H. Voss, “TMVA: Toolkit for Multivariate Data Analysis,” *Proceeding of Science*, vol. ACAT, p. 040, 2007.
- [107] Y. Freund and R. E. Schapire, “A decision-theoretic generalization of on-line learning and an application to boosting,” *Journal of Computer and System Sciences*, vol. 55, no. 1, pp. 119 – 139, 1997.

-
- [108] The CMS Collaboration, “Investigating the top-Yukawa coupling with the production of a Higgs boson in association with a single top quark in the $H \rightarrow b\bar{b}$ decay channel,” CMS Physics Analysis Note, CMS AN-16-065, CERN, 2016.
- [109] CERN, “SM Higgs Branching Ratios and Total Decay Widths (update in CERN Report4 2016),” 2016. https://twiki.cern.ch/twiki/bin/view/LHCPhysics/CERNYellowReportPageBR#Higgs_2_fermions, last accessed on 2016-11-29.
- [110] The ATLAS and CMS Collaborations, “Measurements of the Higgs boson production and decay rates and constraints on its couplings from a combined ATLAS and CMS analysis of the LHC pp collision data at $\sqrt{s} = 7$ and 8 TeV,” *Journal of High Energy Physics*, vol. 08, p. 045, 2016.
- [111] The CMS Collaboration, “Measurement of the cross section ratio $\sigma_{t\bar{t}b\bar{b}}/\sigma_{t\bar{t}jj}$ in pp collisions at $\sqrt{s} = 8$ TeV,” *Physics Letters B*, vol. 746, pp. 132 – 153, 2015.
- [112] The ATLAS Collaboration, “Study of heavy-flavor quarks produced in association with top-quark pairs at $\sqrt{s} = 7$ TeV using the ATLAS detector,” *Physical Review D*, vol. 89, no. 7, p. 072012, 2014.
- [113] The CMS Collaboration, “Multivariate Electron Identification for Run2 – Triggering electron MVA details and working points,” 2016. https://twiki.cern.ch/twiki/bin/view/CMS/MultivariateElectronIdentificationRun2#Triggering_electron_MVA_details, last accessed on 2016-11-30.
- [114] The CMS Collaboration, “Baseline muon selections for Run-II – Muon Identification,” 2016. https://twiki.cern.ch/twiki/bin/viewauth/CMS/SWGuideMuonIdRun2#Tight_Muon, last accessed on 2016-12-01.
- [115] The CMS Collaboration, “Jet Identification – Recommendations for 13 TeV data analysis (74X, 76X, 80X),” 2016. https://twiki.cern.ch/twiki/bin/view/CMS/JetID#Recommendations_for_13_TeV_data, last accessed on 2016-12-01.
- [116] The CMS Collaboration, “Usage of b Tag Objects for 13 TeV Data with 25ns bunch spacing and 76X ReReco – Supported Algorithms and Operating Points,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/BtagRecommendation76X>, last accessed on 2016-12-01.
- [117] The CMS Collaboration, “MET Corrections and Uncertainties for Run-II,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/MissingETRRun2Corrections>, last accessed on 2016-12-02.

- [118] T. Chwalek, “Measurement of the W-Boson Helicity-Fractions in Top-Quark Decays with the CDF II Experiment and Prospects for an Early $t\bar{t}$ Cross-Section Measurement with the CMS Experiment,” Karlsruhe Institute of Technology, CERN-THESIS-2010-255, 2010.
- [119] The CMS Collaboration, “Pileup Reweighting Utilities,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/PileupMCReweightingUtilities>, last accessed on 2016-12-03.
- [120] The CMS Collaboration, “Muon T&P Instructions for Run-II,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/MuonTagAndProbeTreesRun2>, last accessed on 2016-12-04.
- [121] The CMS Collaboration, “Reference muon id, isolation and trigger efficiencies for Run-II,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/MuonReferenceEffsRun2>, last accessed on 2016-12-04.
- [122] The CMS Collaboration, “Electron and photon ID – Efficiencies and scale factors,” 2016. https://twiki.cern.ch/twiki/bin/view/CMS/EgammaIDRecipesRun2#Efficiencies_and_scale_factors, last accessed on 2016-12-04.
- [123] The CMS Collaboration, “Details of the Tag and Probe procedure for Egamma,” 2016. <https://twiki.cern.ch/twiki/bin/view/CMS/ElectronScaleFactorsRun2>, last accessed on 2016-12-04.
- [124] The CMS Collaboration, “Btag Shape Calibration – Event reweighting using scale factors calculated with a tag and probe method,” 2016. <https://twiki.cern.ch/twiki/bin/view/CMS/BTagShapeCalibration>, last accessed on 2016-12-04.
- [125] The CMS Collaboration, “Calibration of the Combined Secondary Vertex b-Tagging discriminant using dileptonic $t\bar{t}$ and Drell-Yan events,” CMS Physics Analysis Note, CMS AN-13-130, CERN, 2013.
- [126] E. Boos *et al.*, “Generic user process interface for event generators,” in *Physics at TeV colliders. Proceedings, Euro Summer School, Les Houches, France, May 21-June 1, 2001*, 2001.
- [127] J. Alwall, A. Ballestrero, P. Bartalini, S. Belov, E. Boos, *et al.*, “A Standard format for Les Houches event files,” *Computer Physics Communications*, vol. 176, pp. 300–304, 2007.
- [128] V. D. Barger, J. Ohnemus, and R. J. N. Phillips, “Event shape criteria for single lepton top signals,” *Physical Review D*, vol. 48, pp. 3953–3956, 1993.

-
- [129] G. C. Fox and S. Wolfram, “Observables for the analysis of event shapes in e^+e^- annihilation and other processes,” *Physical Review Letters*, vol. 41, no. 23, p. 1581, 1978.
- [130] “CMS Luminosity Measurement for the 2015 Data Taking Period,” CMS-PAS-LUM-15-001, CERN, Geneva, 2016.
- [131] The CMS Collaboration, “Estimating Systematic Errors Due to Pileup Modeling,” 2013. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/PileupSystematicErrors>, last accessed on 2016-12-13.
- [132] The CMS Collaboration, “Utilities for Accessing Pileup Information for Data – Pileup JSON Files For Run II,” 2016. https://twiki.cern.ch/twiki/bin/view/CMS/PileupJSONFileforData#Pileup_JSON_Files_For_Run_II, last accessed on 2016-12-13.
- [133] The CMS Collaboration, “Jet Energy Resolution – JER Scaling factors and Uncertainty for 13 TeV (2015 and 2016),” 2016. https://twiki.cern.ch/twiki/bin/viewauth/CMS/JetResolution#JER_Scaling_factors_and_Uncertai, last accessed on 2016-12-13.
- [134] The CMS Collaboration, “Recommended Jet Energy Corrections and Uncertainties For Data and MC – Jet Energy Corrections in Run2,” 2016. <https://twiki.cern.ch/twiki/bin/view/CMS/JECDataMC>, last accessed on 2016-12-13.
- [135] The CMS Collaboration, “Performance of the CMS missing transverse momentum reconstruction in pp data at $\sqrt{s} = 8$ TeV,” *Journal of Instrumentation*, vol. 10, no. 02, p. P02006, 2015.
- [136] R. Barlow and C. Beeston, “Fitting using finite Monte Carlo samples,” *Computer Physics Communications*, vol. 77, no. 2, pp. 219 – 228, 1993.
- [137] J. S. Conway, “Incorporating Nuisance Parameters in Likelihoods for Multisource Spectra,” in *Proceedings, PHYSTAT 2011 Workshop on Statistical Issues Related to Discovery Claims in Search Experiments and Unfolding*, CERN, Geneva, Switzerland 17-20 January 2011, pp. 115–120, 2011.
- [138] The CMS Collaboration, “SM Higgs production cross sections at $\sqrt{s} = 13$ TeV (update in CERN Report4 2016),” 2016. <https://twiki.cern.ch/twiki/bin/view/LHCPhysics/CERNYellowReportPageAt13TeV>, last accessed on 2016-12-14.
- [139] The CMS Collaboration, “NNLO+NNLL top-quark-pair cross sections,” 2016. <https://twiki.cern.ch/twiki/bin/view/LHCPhysics/TtbarNNLO>, last accessed on 2016-12-14.

- [140] F. Maltoni, D. Pagani, and I. Tsirikos, “Associated production of a top-quark pair with vector bosons at NLO in QCD: impact on $t\bar{t}H$ searches at the LHC,” *Journal of High Energy Physics*, vol. 02, p. 113, 2016.
- [141] The CMS Collaboration, “NLO single-top channel cross sections,” 2016. <https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SingleTopRefXsec>, last accessed on 2016-12-14.
- [142] The CMS Collaboration, “Documentation on c-tag MVA trainings,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/BTagCTagger>, last accessed on 2016-12-16.
- [143] S. Moortgat, “Development of new charm-tagging methods for the search for Flavour Changing top-quark dark matter interactions at the LHC,” Vrije Universiteit Brussel, Master thesis, 2015.
- [144] The ATLAS Collaboration, “Performance and Calibration of the JetFitterCharm Algorithm for c-Jet Identification,” ATL-PHYS-PUB-2015-001, CERN, Geneva, 2015.
- [145] The CMS Collaboration, “Modelling of the single top-quark production in association with the Higgs boson at 13 TeV,” 2016. <https://twiki.cern.ch/twiki/bin/viewauth/CMS/SingleTopHiggsGeneration13TeV>, last accessed on 2016-12-18.

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