Improvement of the single top-quark detection in the s-channel at the CMS Experiment

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BACHELOR THESIS

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I hereby certify that the enclosed thesis is my own work, that I have not sought or used inadmissible help of third parties to produce this work and that I have clearly referenced all sources used in the text.

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Abstract

Introduction

Contents

1 Theoretical Foundation

1.1 The Standard Model of Particle Physics

Fermions	*	Mass	Electric charge (e)
Weak isospin			

The standard model of particle physics (SM) describes all known elementary particles and three of the four known fundamental forces. These forces are the electromagnetic interaction, the weak interaction and the strong interaction. Gravity is not included yet, but is negligible on a microscopic scale. The SM predicted many particles that were found later in experiments, for example the top quark in 1995 [?], the tau neutrino in 2000 [?] or the Higgs boson in 2012 [? ?]. It also provides precise predictions about the properties of elementary particles like the Landé g-factor [?]. Despite the great success of the SM, hints exist that the particles and their interactions are described by a more fundamental mechanism. For example, many free parameters, like the masses of particles, cannot be predicted by the SM and have to be determined by measurements. Furthermore, dark matter and dark energy are not described by the SM. Therefore, to find evidences for new physics, it is important to measure the particle properties predicted by the SM with best possible precision.

This chapter gives a brief summary of the SM with focus on the electroweak interaction and quantum chromodynamics. A more detailed description of the SM can be found in references [? ? ? ?]. Natural units with $\hbar = c = 1$ are used to simplify the formulas quoted in this chapter.

1.2 Overview

The SM is a quantum field theory that describes the particles as fields and that includes quantum mechanics as well as special relativity. It is a gauge theory in which underlying interactions are described by the $U(1)_Y \otimes SU(2)_L \otimes SU(3)_C$ symmetry group.

Special relativity requires a Lorentz invariant description of the SM. The spin of a particle defines its representation of the Lorentz group and has a direct influence on its behavior under Lorentz transformation. With respect to the spin quantum number, the particles in the SM can be divided into two groups, fermions with half integer spin and bosons with integer spin. Fermions are described as spinor fields and are the

Table 1.1: The fermions of the standard model. All quarks and leptons of the SM and their elec-
troweak charges are listed. The fermions occur with negative chirality (left handed) in weak isospin
doublets or with positive chirality (right handed) in weak isospin singlets. Under the assumtion
that the neutrino mass is zero, there is no right handed neutrino in the SM. For each fermion, an
antifermion exists with opposite charge and chirality.

| Fermions | Generation | | | Weak hyper-charge Y_W | 3rd comp. of isospin I_3 | Electric charge Q |
	1	2	3			
Quarks	$\begin{pmatrix} u \\ d \end{pmatrix}_L$	$\begin{pmatrix} c \\ s \end{pmatrix}_L$	$\begin{pmatrix} t \\ b \end{pmatrix}_L$	$+1/3$ $+1/3$	$+1/2$ $-1/2$	$+2/3$ $-1/3$
	u_R	c_R	t_R	$+4/3$	0	$+2/3$
	d_R	s_R	b_R	$-2/3$	0	$-1/3$
Leptons	$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$	$\begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L$	$\begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L$	-1 -1	$+1/2$ $-1/2$	0 -1
	e_R	μ_R	τ_R	-2	0	-1

constituents of all microscopic matter, for example atoms. They exist in three gener-
ations, where only the first generation of particles is stable and does not decay over
time. Fermions can be further distinguished into quarks and leptons. Quarks are the
constituents of protons and neutrons, they interact through all forces and can be dis-
tinguished by their electric charge. Quarks with an electric charge of $+2/3\,e$ are called
up-type quarks, while down-type quarks carry an electric charge of $-1/3\,e$. The charged
leptons, the electron, the muon and the tauon can interact via the electromagnetic and
weak force. The neutral leptons, the neutrinos, can only interact via the weak force. A
summary of the fermions is given in table **??**. The second group of the SM consists
of bosons. Bosons with spin 1, described by vector fields, are the mediator particles
for the three fundamental forces of the SM: the photon (γ) for the electromagnetic
interaction, the $W^\pm$ and Z bosons for the weak interaction and the gluons (g) for the
strong interaction. The Higgs boson is the only particle of the SM with spin 0 and is
described by a scalar field. It interacts with all massive particles. A summary of the
bosons is shown in table **??**.

1.3 Electroweak Interaction

The electromagnetic interaction and the weak interaction are described in a unified
electroweak theory by the $U(1)_Y \otimes SU(2)_L$ group. By demanding the theory to be locally
invariant under $U(1)_Y$ transformations, there must exist a massless vector boson (B),
and a conserved quantity, the weak hypercharge Y_W. To obtain a local $SU(2)_L$ symmetry,

Table 1.2: Summary of the bosons of the SM. The vector bosons with spin 1, the photon, the $W^{\pm}$ and Z bosons and the gluons are the mediator particles of the three fundamental forces that are described by the SM. The Higgs boson with spin 0 is not mediating a force, but it couples to the mass of the particles, masses are taken from [?]

Boson	Force	Couples to	Mass (GeV)
photon (γ)	electromagnetic	electric charge	0
$W^{\pm}$			80.385
	weak	weak charge	
Z			91.188
8 gluons (g)	strong	color charge	0
Higgs (H)	-	mass	125.09

three additional massless vector bosons (W_0, W_1 and W_2) are needed; the conserved quantity of this symmetry is the third component of the weak isospin I_3. The electric charge

$$Q = \frac{Y_W}{2} + I_3 \quad ,$$ (1.1)

is related with these conserved quantities and is hence conserved. The Higgs boson originates from a spontaneous symmetry breaking of the local $U(1)_Y \otimes SU(2)_L$ symmetry and leads to a non-diagonal mass matrix for the four vector bosons. This mass matrix can be diagonalized to find the mass eigenstates

$$\gamma = \cos(\theta_W)B + \sin(\theta_W)W_0 \quad ,$$
$$Z = -\sin(\theta_W)B + \cos(\theta_W)W_0 \quad ,$$
$$W^{\pm} = \frac{1}{\sqrt{2}}(W_1 \mp iW_2) \quad ,$$ (1.2)

with the Weinberg angle θ_W. In these mass eigenstates, the photon remains massless while the Z, W^+ and W^- bosons receive a mass. The Z boson and the photon carry no electric charge, while the $W^{\pm}$ bosons carry an electric charge of $\pm 1\,e$. The large mass of the Z and $W^{\pm}$ bosons is responsible for the short range of the weak interaction, as they have a lifetime of less than $10^{-24}\,$s [?]. The effective strength can be estimated with the Fermi coupling constant of $G_F \approx 10^{-5}\,$GeV^{-2} of the weak interaction. The electromagnetic coupling constant has a value in the order of 10^{-2}.

With respect to $SU(2)_L$ transformations, the quarks and leptons of one generation are grouped into weak isospin doublets, consisting of two left-handed spinor fields with $I_3 = \pm 1/2$. Weak isospin singlets with $I_3 = 0$ are represented by right-handed spinor fields, which cannot couple to charged currents. By interacting through $W^{\pm}$ bosons, flavor[1] changing charged currents are possible. They describe transitions of fermions

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[1]The flavor indicates the type of a particle. For example, there are six quark flavors (u, d, s, c, b, t) and

within a weak isospin doublet, for example $u \rightarrow d + W^+$. For quarks, this current is given by

$$J_{CC}^\mu = \begin{pmatrix} \bar{u} & \bar{c} & \bar{t} \end{pmatrix} \gamma^\mu \frac{1-\gamma^5}{2} \begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} \tag{1.3}$$

$$= \bar{u}\, \gamma^\mu \frac{1-\gamma^5}{2}\, d' + \bar{c}\, \gamma^\mu \frac{1-\gamma^5}{2}\, s' + \bar{t}\, \gamma^\mu \frac{1-\gamma^5}{2}\, b' \quad ,$$

with the Dirac matrices γ^μ and γ^5 and $\bar{q} = q^\dagger \gamma_0$ denoting the Dirac adjunct spinor of q. The quarks described in table **??**, are in their electroweak eigenstate and are distinguished from the observable mass eigenstates. Here, q' denotes a quark in its electroweak eigenstate and q in its mass eigenstate. The mass eigenstates of the quarks are therefore a mixture of their electroweak eigenstates and vice versa:

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = V \begin{pmatrix} d \\ s \\ b \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} \quad . \tag{1.4}$$

The 3×3 unitary matrix V is the Cabibbo-Kobayashi-Maskawa matrix (CKM-matrix) that allows transitions of quarks between two generations. The probability of these transitions is low due to small off-diagonal values of V [**?**]:

$$V = \begin{pmatrix} 0.97417 \pm 0.00021 & 0.2248 \pm 0.0006 & 0.00409 \pm 0.00039 \\ 0.220 \pm 0.005 & 0.995 \pm 0.016 & 0.0405 \pm 0.0015 \\ 0.0082 \pm 0.0006 & 0.0400 \pm 0.0027 & 1.009 \pm 0.031 \end{pmatrix} \quad . \tag{1.5}$$

The matrix V has four free parameters that cannot be predicted by the SM and have to be determined by measurements. By inserting equation **??** into equation **??**, one can, for instance, obtain the current for a transition of a b quark to a c quark . The relevant part is given by

$$J_{CC,cb}^\mu = \bar{c}\, \gamma^\mu \frac{1-\gamma^5}{2}\, V_{cb}\, b \quad , \tag{1.6}$$

which is suppressed by the factor of $V_{cb} = 0.0405$.

1.4 Quantum Chromodynamics

The strong interaction is responsible for the interaction between quarks and gluons and is described by quantum chromodynamics (QCD) with the underlying summetry

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six lepton flavors (e, μ, τ , ν_e, ν_μ, ν_τ).

of the $SU(3)_C$ group. Eight gluons are needed to obtain a $SU(3)_C$ invariant description. The charge of this symmetry is called color, of which three different types exist: red, green and blue. For each color, there is also a corresponding anticolor. As this symmetry is not broken, the gluons remain massless and indistinguishable. They carry one color and one anticolor charge, this allows interactions of gluons among themselves. Quarks on the other hand carry one color charge while antiquarks are charged with an anticolor.

The nature of the strong force leads to a potential for a quark-antiquark pair of

$$V(r) = -a \cdot \frac{\alpha_s(r)}{r} + b \cdot r \quad, \tag{1.7}$$

with the positive constants a and b, the radius between the two particles r and the strong couplings constant $\alpha_s(r)$, which is dependent on the radius and is in the order of 1^2. For small radii, the first term is dominant and allows bound states. The lower the distance between two particles, the lower is the strong interaction between them, this effect is called asymptotic freedom. On the other hand, the potential increases linearly with the distance between the particles. If a quark in a bound state is pulled apart from the other quark(s) in this bound state, the energy of the system is steadily increasing. When the energy is high enough, a quark-antiquark pair is produced in order to form two color neutral objects with the separated particles. This effect is called confinement and is the reason that no color-charged isolated object exists. A color-neutral object that consists of quarks is called hadron.

Hadrons can be grouped into mesons and baryons. Mesons are quark-antiquark pairs, whereas baryons consist of three quarks. Besides the top quark, all quarks can be bound in hadrons, therefore a huge number of different hadrons exists. Hadrons are produced by the strong force and decay in flavor-changing charged currents as described in equation **??**. Therefore, the lifetime of hadrons [**?**] varies from 10^{-8} s for charged pions (u$\bar{\text{d}}$, $\bar{\text{u}}$d), 10^{-12} s for B mesons (u$\bar{\text{b}}$, $\bar{\text{u}}$b, $\bar{\text{b}}$d, b$\bar{\text{d}}$) to values up to 10^{-24} s for hadrons in excited states like the ρ meson (u$\bar{\text{d}}$, $\bar{\text{u}}$d). The only stable hadron is the proton (uud), as the neutron (udd) is only stable in atomic nuclei. The top quark is the only quark that cannot occur in a bound state because its lifetime of about $5 \cdot 10^{-25}$ s [**?**] is roughly 20 times smaller than the hadronization time.

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[2]For instance, at the energy of the mass of the Z boson, $M_Z = 91.188\,\text{GeV}$, $\alpha_s(M_Z^2) = 0.1181 \pm 0.0013$ [**?**]

2 The Compact Muon Solenoid Experiment

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Bibliography

[1] F. Abe *et al.*, "Observation of top quark production in $\bar{p}p$ collisions", *Physical Review Letters*, vol. 74, pp. 2626–2631, 1995.

[2] K. Kodama *et al.*, "Observation of tau neutrino interactions", *Physics Letters B*, vol. 504, pp. 218–224, 2001.

[3] G. Aad *et al.*, "Observation of a new particle in the search for the standard model higgs boson with the atlas detector at the lhc", *Physics Letters B*, vol. 716, pp. 1–29, 2012.

[4] S. Chatrchyan *et al.*, "Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC", *Physics Letters B*, vol. 716, pp. 30–61, 2012.

[5] D. Hanneke, S. F. Hoogerheide, and G. Gabrielse, "Cavity Control of a Single-Electron Quantum Cyclotron: Measuring the Electron Magnetic Moment", *Physical Review A*, vol. 83, p. 052122, 2011.

[6] F. Halzen and A. D. Martin, "Quarks and Leptons: An Introductory Course in Modern Particle Physics", 1984.

[7] G. L. Kane, "Modern elementary particle physics: the fundamental particles and forces?", updated ed., 4. pr. ed., 1998.

[8] J. Schwichtenberg, "Physics from symmetry", Undergraduate Lecture Notes in Physics, Springer, Cham, 2nd ed. 2018 ed., 2018.

[9] M. E. Peskin and D. V. Schroeder, "An Introduction to quantum field theory", Addison-Wesley, Reading, USA, 1995.

[10] C. Patrignani *et al.*, "Review of Particle Physics", *Chinese Physics C*, vol. 40, no. 10, p. 100001, 2016.

[11] A. Quadt, "Top quark physics at hadron colliders", *The European Physical Journal C - Particles and Fields*, vol. 48, pp. 835–1000, Dec 2006.

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