### Nonlinear systems of equations

We saw earlier how to solve large systems of linear equations: collect them into a single matrix equation, and use an algorithm like Gaussian elimination to construct and solve a factorization.

We also saw how to make a linear approximation to a nonlinear function: if  $f \in \mathbb{R}^n \to \mathbb{R}^n$ , then we can get a first-order Taylor approximation by calculating the differential,

$$df(x) = f'(x)dx$$

These two tools can work together: suppose that we want to solve a nonlinear system of equations

$$f(x) = 0$$

If we start from a guess  $x_1$  at a solution, we can construct a first-order Taylor expansion

$$df=f^{\prime}(x_{1})dx$$

Holding  $x_1$  fixed, this is a *linear* equation for df in terms of dx. So we can ask to find dx that makes  $f(x_1) + df = 0$  — that is, we can solve a linear approximation to the original nonlinear equations. As before, Gaussian elimination or other factorizations can solve this linear system quickly and reliably.

If there are multiple solutions — i.e., if f' is singular — we need to pick one; methods for doing so are beyond the scope of these notes.

With the solution dx in hand, we can construct a new guess

$$x_2 = x_1 + dx$$

We can then make a new Taylor expansion around  $x_2$ , leading to a new linear approximation

$$f(x_2)+df=0 \qquad df=f'(x_2)dx$$

Repeating the process lets us construct  $x_3$ ,  $x_4$ , and so forth. Hopefully each successive  $x_t$  comes closer to satisfying  $f(x_t) = 0$ .

This process is called *Newton's method*, and it often converges rapidly to a solution of the nonlinear system f(x) = 0. In fact, the stationary points of Newton's method are

strongly related to the solutions: if dx=0 then the second equation implies df=0, so the first equation implies f(x)=0. In the other direction, if f(x)=0 then the first equation constrains df=0. Then if f'(x) is not singular, we have to have dx=0. (If f'(x) is singular, we might be at a stationary point without solving f(x)=0.)

If Newton's method fails, sometimes we can rescue it by *damping*, i.e., decreasing our step size: that is, we set  $x_{t+1} = x_t + \alpha_t dx$  for some  $\alpha_t \in (0,1)$ . But tuning the step size (and other methods beyond damped Newton) are beyond the scope of this set of notes.

#### **Example**

Let  $f(x) = e^x - 1$ , so that  $df = e^x dx$ . The solution to f(x) = 0 is x = 0, but let's see if we can find this by Newton's method, starting from somewhere else.

x	f	df	Equation	dx
1	e-1	e	e dx = 1 - e	$\frac{1-e}{e}$
-0.632	-0.468	0.532	0.532  dx = 0.468	0.880
0.248	0.281	1.281	1.281  dx = -0.281	-0.219

Quite rapidly we have reached x=0.029, very close to the true solution.

# **Unconstrained optimization**

Solving optimization problems is strongly related to solving systems of equations. In an unconstrained optimization problem

$$\min_{\theta} L(\theta)$$

we can try to find the solution by looking for a critical point: a place where, locally, changes to  $\theta$  do not change  $L(\theta)$ .

Critical points can be minima or maxima, and they can be either local or global. In addition, they can be neither: they can be places where the function flattens out temporarily, or places where it looks like a saddle. For now, we won't be concerned with checking which is which.

To find a critical point, we can look at the first order Taylor expansion of L:

$$dL = L'( heta) \, d heta$$

At a critical point, all possible changes  $d\theta$  should leave dL=0. That means we must have

$$L'(\theta) = 0$$

These equations are the *first-order optimality conditions* for  $L(\theta)$ . Geometrically, the Taylor expansion is flat (constant):

Of course, the system of equations  $L'(\theta)=0$  could be nonlinear. So, we can apply Newton's method — that is, we can set a first-order Taylor approximation of L' to zero and solve for  $d\theta$ :

$$L'(\theta) + dL' = 0$$
  $dL' = L''(\theta)d\theta$ 

We can find L' and L'' by differentiating L twice.

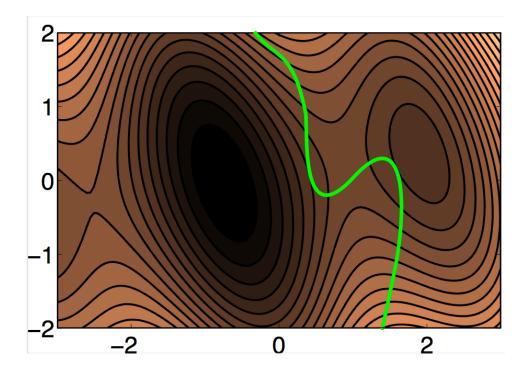
This is such a common application of Newton's method that it shares the same name. If necessary, we can distinguish by calling the two algorithms "Newton's method for solving a system of equations" and "Newton's method for optimizing a function".

## **Constrained optimization**

In a constrained problem

$$\min_{\theta} L(\theta)$$
 s.t.  $g(\theta) = 0$ 

we don't need  $L'(\theta)$  to be zero: it's OK if there's a direction of decrease in  $L(\theta)$  as long the constraint prevents us from moving in this direction.



To encode this condition, we need to be a bit clever. First note that the solutions to the following problem

$$\min_{\theta} \left[ L(\theta) + \alpha g(\theta) \right] \quad \text{s.t.} \quad g(\theta) = 0$$

are the same as the solutions to our original problem, no matter what the value of  $\alpha$  is, since  $\alpha q(\theta) = 0$  for any feasible  $\theta$ .

Then note that, by choosing  $\alpha$  appropriately, we can rule out any direction of decrease in L that doesn't satisfy the constraint: if L would decrease on the side of the constraint where  $g(\theta)>0$ , then we choose  $\alpha$  to be very positive, so that any motion in this direction would cause  $L(\theta)+\alpha g(\theta)$  to increase instead of decreasing. Similarly, if L would decrease on the other side if the constraint, where  $g(\theta)<0$ , we choose  $\alpha$  to be very negative.

Given this new objective, we can use a Taylor expansion the same way as before. We ask for a critical point: a  $\theta$  where, to first order, the objective doesn't change as we change  $d\theta$ . That is,

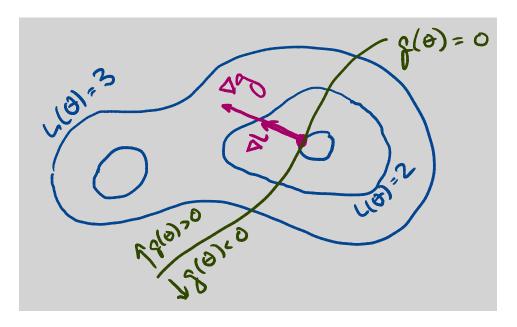
$$0 = d(L(\theta) + \alpha g(\theta)) = L'(\theta)d\theta + \alpha g'(\theta)d\theta$$

which implies

$$L'(\theta) + \alpha g'(\theta) = 0$$

Geometrically, this equation tells us that at a critical point we can only change  $L(\theta)$  by

changing  $\theta$  in a direction *orthogonal* to the constraint (parallel to  $g'(\theta)$ ): sliding in any direction along the constraint doesn't change L, at least to first order.



Interestingly, that means that we didn't have to choose  $\alpha$  a priori: any  $\theta$  and  $\alpha$  that satisfy

$$g(\theta) = 0$$
  $L'(\theta) + \alpha g'(\theta) = 0$ 

will represent a critical point. So, as before, we've turned our optimization problem into a possibly-nonlinear system of equations. We can solve this system with Newton's method or any other appropriate tool.

The new variable  $\alpha$  is called a *Lagrange multiplier* or *dual variable*. We can interpret  $-L'(\theta)$  as a force that wants to push our current point  $\theta$  downhill, toward a minimum of L. We can then think of  $-\alpha g'(\theta)$  as a force that pushes back, keeping  $\theta$  from violating the constraint. At the solution, the two forces balance exactly.

By introducing the dual variable, we've transformed our optimization problem into a system of simultaneous equations, where the objective and the constraints are treated the same way. This transformation was what let us apply Newton's method.

Practice: solve the following problem by introducing a Lagrange multiplier.

$$\min_{x,y} rac{1}{2} (x^2 + y^2) \quad ext{s.t.} \quad x + 2y = 1$$

## Multiple constraints

Suppose we have more than one constraint:

$$\min_{\theta} L(\theta) \quad \mathrm{s.t.} \quad g(\theta) = 0$$

where the output of  $g(\theta)$  is in  $\mathbb{R}^d$  instead of  $\mathbb{R}$ . The solution in this case is almost identical: we can still solve

$$g(\theta) = 0$$
  $L'(\theta) + \alpha g'(\theta) = 0$ 

But now, instead of  $\alpha \in \mathbb{R}$ , we need  $\alpha \in \mathbb{R}^{1 \times d}$ , so that  $\alpha g'(\theta)$  is the same shape as  $L'(\theta)$ .

Each coordinate  $\alpha_i$  is still called a Lagrange multiplier. The geometric interpretation is only slightly different from before: we think of each  $\alpha_i$  as controlling a separate force, in a direction that's normal to the corresponding constraint  $g_i(\theta) = 0$ . At equilibrium, all of the forces  $\alpha_i g_i'(\theta)$  combine to cancel out  $L'(\theta)$ .