

Atomic transitions 4

From the previous section the transition rate is

$$R_{a \rightarrow b} = \frac{\pi e^2}{\epsilon_0 m^2 \hbar^2} \left| \langle \psi_b | \exp(i \mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \right|^2 \frac{\rho(\omega_0)}{\omega_0^2}$$

Substitute the dipole approximation

$$\exp(i \mathbf{k} \cdot \mathbf{r}) \approx 1$$

to obtain

$$R_{a \rightarrow b} = \frac{\pi e^2}{\epsilon_0 m^2 \hbar^2} \left| \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \right|^2 \frac{\rho(\omega_0)}{\omega_0^2}$$

By the identity

$$\mathbf{p} = \frac{im}{\hbar} [H_0, \mathbf{r}] \quad (1)$$

we have

$$\begin{aligned} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle &= \frac{im}{\hbar} \langle \psi_b | \boldsymbol{\epsilon} \cdot [H_0, \mathbf{r}] | \psi_a \rangle \\ &= \frac{im}{\hbar} \langle \psi_b | \boldsymbol{\epsilon} \cdot H_0 \mathbf{r} | \psi_a \rangle - \frac{im}{\hbar} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} H_0 | \psi_a \rangle \\ &= \frac{im}{\hbar} E_b \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle - \frac{im}{\hbar} E_a \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \\ &= \frac{im}{\hbar} (E_b - E_a) \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \\ &= im \omega_0 \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \end{aligned}$$

Hence

$$R_{a \rightarrow b} = \frac{\pi e^2}{\epsilon_0 \hbar^2} \left| \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \right|^2 \rho(\omega_0)$$

Verify dimensions.

$$R_{a \rightarrow b} \propto \frac{\frac{e^2}{\text{C}^2}}{\frac{\epsilon_0}{\text{C}^2 \text{ J}^{-1} \text{ m}^{-1}} \frac{\hbar^2}{\text{J}^2 \text{ s}^2}} \times \left| \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \right|^2 \times \frac{\rho(\omega_0)}{\text{m}^2 \text{ J m}^{-3} \text{ s}} = \text{s}^{-1}$$

Eigenmath code