

## Green's function

In this section we will find the Green's function  $G(\mathbf{x})$  such that

$$(\nabla^2 + k^2)G(\mathbf{x}) = \delta^3(\mathbf{x}) \quad (1)$$

Let  $g(\mathbf{y})$  be the Fourier transform of  $G(\mathbf{x})$  such that

$$G(\mathbf{x}) = \frac{1}{(2\pi)^{\frac{3}{2}}} \int \exp(i\mathbf{x} \cdot \mathbf{y}) g(\mathbf{y}) d\mathbf{y} \quad (2)$$

Substitute (2) into (1) to obtain

$$(\nabla^2 + k^2) \left[ \frac{1}{(2\pi)^{\frac{3}{2}}} \int \exp(i\mathbf{x} \cdot \mathbf{y}) g(\mathbf{y}) d\mathbf{y} \right] = \delta(\mathbf{x})$$

By linearity of differentiation the  $(\nabla^2 + k^2)$  can be moved inside the integral.

$$\frac{1}{(2\pi)^{\frac{3}{2}}} \int (\nabla^2 + k^2) \exp(i\mathbf{x} \cdot \mathbf{y}) g(\mathbf{y}) d\mathbf{y} = \delta(\mathbf{x})$$

Noting that

$$\nabla^2 \exp(i\mathbf{x} \cdot \mathbf{y}) = -y^2 \exp(i\mathbf{x} \cdot \mathbf{y})$$

and

$$\delta(\mathbf{x}) = \frac{1}{(2\pi)^3} \int \exp(i\mathbf{x} \cdot \mathbf{y}) d\mathbf{y}$$

we have

$$\frac{1}{(2\pi)^{\frac{3}{2}}} \int (-y^2 + k^2) \exp(i\mathbf{x} \cdot \mathbf{y}) g(\mathbf{y}) d\mathbf{y} = \frac{1}{(2\pi)^3} \int \exp(i\mathbf{x} \cdot \mathbf{y}) d\mathbf{y}$$

Hence

$$g(\mathbf{y}) = \frac{1}{(2\pi)^{\frac{3}{2}}(k^2 - y^2)}$$

Substitute for  $g(\mathbf{y})$  in (2) to obtain

$$G(\mathbf{x}) = \frac{1}{(2\pi)^3} \int \frac{\exp(i\mathbf{x} \cdot \mathbf{y})}{k^2 - y^2} d\mathbf{y}$$

Change to polar coordinates where  $x = |\mathbf{x}|$ ,  $y = |\mathbf{y}|$ , and  $\theta$  and  $\phi$  are the angular distance from  $\mathbf{x}$  to  $\mathbf{y}$ .

$$G(\mathbf{x}) = \frac{1}{(2\pi)^3} \int_0^\infty \int_0^\pi \int_0^{2\pi} \frac{\exp(ixy \cos \theta)}{k^2 - y^2} y^2 \sin \theta dy d\theta d\phi$$

For the integrals over  $\theta$  and  $\phi$  we have

$$\int_0^\pi \int_0^{2\pi} \exp(ixy \cos \theta) \sin \theta d\theta d\phi = \frac{4\pi \sin(xy)}{xy}$$

Hence

$$G(\mathbf{x}) = \frac{1}{2\pi^2 x} \int_0^\infty \frac{y \sin(xy)}{k^2 - y^2} dy$$

Noting that  $y \sin(xy)$  is an even function of  $y$  we can change the integral limits as follows.

$$G(\mathbf{x}) = \frac{1}{4\pi^2 x} \int_{-\infty}^\infty \frac{y \sin(xy)}{k^2 - y^2} dy$$

Negate the denominator.

$$G(\mathbf{x}) = \frac{1}{4\pi^2 x} \int_{-\infty}^\infty -\frac{y \sin(xy)}{y^2 - k^2} dy$$

Change the sine function to exponential form and factor the denominator.

$$G(\mathbf{x}) = \frac{i}{8\pi^2 x} \left[ \int_{-\infty}^\infty \frac{y \exp(ixy)}{(y - k)(y + k)} dy - \int_{-\infty}^\infty \frac{y \exp(-ixy)}{(y - k)(y + k)} dy \right]$$

By Cauchy's integral formula we have

$$\int_{-\infty}^\infty \frac{y \exp(ixy)}{y + k} \frac{1}{y - k} dy = i\pi \exp(ikx)$$

and

$$\int_{-\infty}^\infty \frac{y \exp(-ixy)}{y - k} \frac{1}{y + k} dy = -i\pi \exp(ikx)$$

Hence

$$\boxed{G(\mathbf{x}) = -\frac{\exp(ikx)}{4\pi x}} \quad (3)$$

where

$$x = |\mathbf{x}|$$

Verify that (3) satisfies (1).

We will need the following formula from Griffiths and Schroeter problem 10.8.

$$\nabla^2(1/r) = -4\pi\delta^3(\mathbf{r}) \quad (4)$$

Recall that  $\nabla^2 = \nabla \cdot \nabla$  and

$$\nabla \cdot (f\mathbf{A}) = \nabla f \cdot \mathbf{A} + f\nabla \cdot \mathbf{A}$$

Substituting  $\mathbf{r}$  for  $\mathbf{x}$  in (3) we have for the Laplacian of  $G(\mathbf{r})$

$$\begin{aligned} \nabla^2 G(\mathbf{r}) &= -\frac{1}{4\pi} \nabla \cdot \nabla \left( \frac{e^{ikr}}{r} \right) \\ &= -\frac{1}{4\pi} \nabla \cdot \left( \frac{1}{r} \nabla e^{ikr} + e^{ikr} \nabla \frac{1}{r} \right) \\ &= -\frac{1}{4\pi} \left( \underbrace{\nabla \frac{1}{r} \cdot \nabla e^{ikr}}_{\text{subst. (6)}} + \underbrace{\frac{1}{r} \nabla^2 e^{ikr}}_{\text{subst. (7)}} + \underbrace{\nabla e^{ikr} \cdot \nabla \frac{1}{r}}_{\text{subst. (6)}} + \underbrace{e^{ikr} \nabla^2 \frac{1}{r}}_{\text{subst. (4)}} \right) \end{aligned} \quad (5)$$

In spherical coordinates

$$\nabla \frac{1}{r} \cdot \nabla e^{ikr} = \left( -\frac{1}{r^2} \mathbf{e}_r + 0 \mathbf{e}_\theta + 0 \mathbf{e}_\phi \right) \cdot (ike^{ikr} \mathbf{e}_r + 0 \mathbf{e}_\theta + 0 \mathbf{e}_\phi) = -\frac{ike^{ikr}}{r^2} \quad (6)$$

and

$$\begin{aligned} \frac{1}{r} \nabla^2 e^{ikr} &= \frac{1}{r^2} \frac{\partial^2}{\partial r^2} (r e^{ikr}) \\ &= \frac{1}{r^2} \frac{\partial}{\partial r} (e^{ikr} + ikr e^{ikr}) \\ &= \frac{1}{r^2} (2ike^{ikr} - k^2 r e^{ikr}) \\ &= \frac{2ike^{ikr}}{r^2} - \frac{k^2 e^{ikr}}{r} \end{aligned} \quad (7)$$

Substitute (4), (6), and (7) into (5) to obtain

$$\nabla^2 G(\mathbf{r}) = \frac{k^2 e^{ikr}}{4\pi r} + \delta^3(\mathbf{r}) e^{ikr}$$

Then by equation (3)

$$\nabla^2 G(\mathbf{r}) = -k^2 G(\mathbf{r}) + \delta^3(\mathbf{r}) e^{ikr}$$

Noting that  $e^{ikr} = 1$  at  $r = 0$ , the  $e^{ikr}$  term can be discarded leaving

$$\nabla^2 G(\mathbf{r}) = -k^2 G(\mathbf{r}) + \delta^3(\mathbf{r}) \quad (8)$$

Hence

$$(\nabla^2 + k^2)G(\mathbf{r}) = \delta^3(\mathbf{r})$$

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