Spin

Spin state $|s\rangle$ is a normalized vector in $\mathbb{C}^2$.

$$|s\rangle = \begin{pmatrix} c_+ \\ c_- \end{pmatrix}, \quad |c_+|^2 + |c_-|^2 = 1$$

Spin measurement probabilities are the transition probabilities from $|s\rangle$ to an eigenstate.

For spin measurements in the z direction we have

$$\Pr\left(S_z = +\tfrac{\hbar}{2}\right) = |\langle z_+|s\rangle|^2 = |c_+|^2$$
$$\Pr\left(S_z = -\tfrac{\hbar}{2}\right) = |\langle z_-|s\rangle|^2 = |c_-|^2$$

where the eigenstates are

$$|z_+\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |z_-\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

By definition of expectation value we have

$$\langle S_z \rangle = \tfrac{\hbar}{2}\Pr\left(S_z = +\tfrac{\hbar}{2}\right) - \tfrac{\hbar}{2}\Pr\left(S_z = -\tfrac{\hbar}{2}\right)$$

Rewrite as

$$\langle S_z \rangle = \tfrac{\hbar}{2}|\langle z_+|s\rangle|^2 - \tfrac{\hbar}{2}|\langle z_-|s\rangle|^2$$

Rewrite again as

$$\langle S_z \rangle = \tfrac{\hbar}{2}\langle s|z_+\rangle\langle z_+|s\rangle - \tfrac{\hbar}{2}\langle s|z_-\rangle\langle z_-|s\rangle$$

Then by

$$\langle S_z \rangle = \langle s|S_z|s\rangle$$

we have

$$S_z = \tfrac{\hbar}{2}|z_+\rangle\langle z_+| - \tfrac{\hbar}{2}|z_-\rangle\langle z_-| = \frac{\hbar}{2}\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

From the commutator

$$S_+ S_- - S_- S_+ = 2\hbar S_z$$

we have

$$S_+ S_- - S_- S_+ = \hbar^2|z_+\rangle\langle z_+| - \hbar^2|z_-\rangle\langle z_-|$$

Rewrite as

$$S_+ S_- - S_- S_+ = \hbar^2|z_+\rangle\langle z_-|z_-\rangle\langle z_+| - \hbar^2|z_-\rangle\langle z_+|z_+\rangle\langle z_-|$$

Hence

$$S_+ = \hbar|z_+\rangle\langle z_-| = \hbar\begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$S_- = \hbar|z_-\rangle\langle z_+| = \hbar\begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

Then by

$$S_+ = S_x + iS_y$$
$$S_- = S_x - iS_y$$

we obtain

$$S_x = \frac{S_+ + S_-}{2} = \frac{\hbar}{2}\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$S_y = \frac{S_+ - S_-}{2i} = \frac{\hbar}{2}\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

By solving for the eigenstates in

$$S_x|x_\pm\rangle = \pm\tfrac{\hbar}{2}|x_\pm\rangle$$
$$S_y|y_\pm\rangle = \pm\tfrac{\hbar}{2}|y_\pm\rangle$$

we obtain

$$|x_+\rangle = \frac{|z_+\rangle + |z_-\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}}\begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$|x_-\rangle = \frac{|z_+\rangle - |z_-\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}}\begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

and

$$|y_+\rangle = \frac{|z_+\rangle + i|z_-\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}}\begin{pmatrix} 1 \\ i \end{pmatrix}$$

$$|y_-\rangle = \frac{|z_+\rangle - i|z_-\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}}\begin{pmatrix} 1 \\ -i \end{pmatrix}$$

Expectation values for spin operators.

$$\langle S_x\rangle = \langle s|S_x|s\rangle = \frac{\hbar}{2}\left(c_+c_-^* + c_+^*c_-\right)$$

$$\langle S_y\rangle = \langle s|S_y|s\rangle = \frac{i\hbar}{2}\left(c_+c_-^* - c_+^*c_-\right)$$

$$\langle S_z\rangle = \langle s|S_z|s\rangle = \frac{\hbar}{2}\left(|c_+|^2 - |c_-|^2\right)$$

Expected spin vector.

$$\langle \mathbf{S}\rangle = \begin{pmatrix} \langle S_x\rangle \\ \langle S_y\rangle \\ \langle S_z\rangle \end{pmatrix}$$

Consider the case of having determined $\langle \mathbf{S}\rangle$ by experiment. To convert $\langle \mathbf{S}\rangle$ to a spin state $|s\rangle$, let

$$x = \frac{2}{\hbar}\langle S_x\rangle = \sin\theta\cos\phi$$

$$y = \frac{2}{\hbar}\langle S_y\rangle = \sin\theta\sin\phi$$

$$z = \frac{2}{\hbar}\langle S_z\rangle = \cos\theta$$

Then
$$|s\rangle = \begin{pmatrix} \cos(\theta/2) \\ \sin(\theta/2)\exp(i\phi) \end{pmatrix}$$

where
$$\cos(\theta/2) = \sqrt{\frac{\cos\theta + 1}{2}} = \sqrt{\frac{z+1}{2}}$$
$$\sin(\theta/2) = \sqrt{\frac{1 - \cos\theta}{2}} = \sqrt{\frac{1-z}{2}}$$

and
$$\exp(i\phi) = \cos\phi + i\sin\phi = \frac{x + iy}{\sqrt{x^2 + y^2}}$$

Spin angular momentum magnitude squared operator S^2.
$$S^2 = S_x^2 + S_y^2 + S_z^2 = \frac{3\hbar^2}{4}\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Expectation value for S^2.
$$\langle s|S^2|s\rangle = \frac{3\hbar^2}{4}\left(|c_+|^2 + |c_-|^2\right) = \frac{3\hbar^2}{4}$$

The following commutator was used to derive S_x and S_y.
$$S_+S_- - S_-S_+ = 2\hbar S_z$$

The commutator is a consequence of the following wave equation for spin.
$$\hat{\mathbf{S}}\psi = (\mathbf{r} \times \hat{\mathbf{p}})\psi, \quad \mathbf{r} = \begin{pmatrix} x \\ y \\ z \end{pmatrix}, \quad \hat{\mathbf{p}} = -i\hbar\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{pmatrix}$$

Rewrite in component form.
$$\hat{S}_x\psi = -i\hbar\left(y\frac{\partial}{\partial z} - z\frac{\partial}{\partial y}\right)\psi$$
$$\hat{S}_y\psi = -i\hbar\left(z\frac{\partial}{\partial x} - x\frac{\partial}{\partial z}\right)\psi$$
$$\hat{S}_z\psi = -i\hbar\left(x\frac{\partial}{\partial y} - y\frac{\partial}{\partial x}\right)\psi$$

By computer algebra we have
$$(\hat{S}_y\hat{S}_z - \hat{S}_z\hat{S}_y)\psi = i\hbar\hat{S}_x\psi$$
$$(\hat{S}_z\hat{S}_x - \hat{S}_x\hat{S}_z)\psi = i\hbar\hat{S}_y\psi$$
$$(\hat{S}_x\hat{S}_y - \hat{S}_y\hat{S}_x)\psi = i\hbar\hat{S}_z\psi$$

Let

$$\hat{S}_+ = \hat{S}_x + i\hat{S}_y$$
$$\hat{S}_- = \hat{S}_x - i\hat{S}_y$$

By computer algebra

$$(\hat{S}_+\hat{S}_- - \hat{S}_-\hat{S}_+)\psi = 2\hbar\hat{S}_z\psi$$

Exercises

1. Verify spin operators and expectation values.

2. Let $|s\rangle$ be the following spin state.

$$|s\rangle = \begin{pmatrix} \frac{1}{3} - \frac{2}{3}i \\ \frac{2}{3} \end{pmatrix}$$

Verify that $|s\rangle$ is normalized and that

$$\langle \mathbf{S} \rangle = \langle s|\mathbf{S}|s\rangle = \frac{\hbar}{2} \begin{pmatrix} \frac{4}{9} \\ \frac{8}{9} \\ \frac{1}{9} \end{pmatrix}$$

where

$$\mathbf{S} = \begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix}$$

Note: In component form we have

$$\langle s|\mathbf{S}|s\rangle = s_\beta^* S^{\alpha\beta}{}_\gamma s^\gamma$$

Eigenmath requires a transpose so that the β indices are adjacent.

$$\langle s|\mathbf{S}|s\rangle = s_\beta^* S^{\beta\alpha}{}_\gamma s^\gamma$$

3. Let $|s\rangle$ be the following spin state.

$$|s\rangle = \begin{pmatrix} \frac{1}{3} - \frac{2}{3}i \\ \frac{2}{3} \end{pmatrix}$$

Verify the following probabilities.

$$\mathrm{Pr}\left(S_x = +\tfrac{\hbar}{2}\right) = |\langle x_+|s\rangle|^2 = \tfrac{13}{18}$$
$$\mathrm{Pr}\left(S_x = -\tfrac{\hbar}{2}\right) = |\langle x_-|s\rangle|^2 = \tfrac{5}{18}$$

$$\mathrm{Pr}\left(S_y = +\tfrac{\hbar}{2}\right) = |\langle y_+|s\rangle|^2 = \tfrac{17}{18}$$
$$\mathrm{Pr}\left(S_y = -\tfrac{\hbar}{2}\right) = |\langle y_-|s\rangle|^2 = \tfrac{1}{18}$$

$$\mathrm{Pr}\left(S_z = +\tfrac{\hbar}{2}\right) = |\langle z_+|s\rangle|^2 = \tfrac{5}{9}$$
$$\mathrm{Pr}\left(S_z = -\tfrac{\hbar}{2}\right) = |\langle z_-|s\rangle|^2 = \tfrac{4}{9}$$

4. Let $|s\rangle$ be the following spin state.

$$|s\rangle = \begin{pmatrix} \frac{1}{3} - \frac{2}{3}i \\ \frac{2}{3} \end{pmatrix}$$

Verify that for θ and ϕ determined by $|s\rangle$, the following spin state $|\chi\rangle$ is indistinguishable from $|s\rangle$.

$$|\chi\rangle = \begin{pmatrix} \cos(\theta/2) \\ \sin(\theta/2)\exp(i\phi) \end{pmatrix} = \begin{pmatrix} \frac{\sqrt{5}}{3} \\ \frac{2+4i}{3\sqrt{5}} \end{pmatrix}$$

5. Verify the following commutators for the wave equation $\mathbf{S}\psi = (\mathbf{r} \times \mathbf{p})\psi$.

$$[S_y, S_z] = i\hbar S_x$$
$$[S_z, S_x] = i\hbar S_y$$
$$[S_x, S_y] = i\hbar S_z$$

$$[S_+, S_-] = 2\hbar S_z$$