

This is the transition rate A_{nm} for a spontaneous emission process $\psi_n \rightarrow \psi_m$.

$$A_{nm} = \frac{e^2}{3\pi\epsilon_0\hbar c^3} \omega_{nm}^3 |\langle r_{nm} \rangle|^2$$

The transition frequency is

$$\omega_{nm} = \frac{1}{\hbar}(E_n - E_m)$$

For a hydrogen atom we have

$$E_n = -\frac{\mu}{2n^2} \left(\frac{e^2}{4\pi\epsilon_0\hbar} \right)^2$$

where μ is reduced electron mass.

The radial density is

$$|\langle r_{nm} \rangle|^2 = |\langle x_{nm} \rangle|^2 + |\langle y_{nm} \rangle|^2 + |\langle z_{nm} \rangle|^2$$

where

$$\begin{aligned} \langle x_{nm} \rangle &= \int \psi_m^* (r \sin \theta \cos \phi) \psi_n dV \\ \langle y_{nm} \rangle &= \int \psi_m^* (r \sin \theta \sin \phi) \psi_n dV \\ \langle z_{nm} \rangle &= \int \psi_m^* (r \cos \theta) \psi_n dV \end{aligned}$$

Let us compute A_{21} for a hydrogen atom. For $n = 2$ there are four possible states.

n	ℓ	m
2	1	1
2	1	-1
2	1	0
2	0	0

The following table shows the radial density for every possible transition.

	$\psi_{2,1,1} \rightarrow \psi_{1,0,0}$	$\psi_{2,1,-1} \rightarrow \psi_{1,0,0}$	$\psi_{2,1,0} \rightarrow \psi_{1,0,0}$	$\psi_{2,0,0} \rightarrow \psi_{1,0,0}$
$\langle x_{21} \rangle =$	$-\frac{128}{243} a_0$	$\frac{128}{243} a_0$	0	0
$\langle y_{21} \rangle =$	$-\frac{128}{243} i a_0$	$-\frac{128}{243} i a_0$	0	0
$\langle z_{21} \rangle =$	0	0	$\frac{128}{243} \sqrt{2} a_0$	0
$ \langle r_{21} \rangle ^2 =$	$\frac{32768}{59049} a_0^2$	$\frac{32768}{59049} a_0^2$	$\frac{32768}{59049} a_0^2$	0

Note that the transition rate of $\psi_{2,0,0} \rightarrow \psi_{1,0,0}$ is zero. For the allowed transitions, the radial density $|\langle r_{21} \rangle|^2$ is independent of ℓ and m .

Symbol a_0 is the Bohr radius

$$a_0 = \frac{4\pi\epsilon_0\hbar^2}{e^2 m_e}$$

For the transition frequency we have

$$\omega_{21} = \frac{1}{\hbar}(E_2 - E_1) = \frac{3e^4\mu}{128\pi^2\varepsilon_0^2\hbar^3}$$

Hence

$$A_{21} = \frac{e^2}{3\pi\varepsilon_0\hbar c^3} \underbrace{\left(\frac{3e^4\mu}{128\pi^2\varepsilon_0^2\hbar^3}\right)^3}_{\omega_{21}} \underbrace{\frac{32768}{59049} \left(\frac{4\pi\varepsilon_0\hbar^2}{e^2m_e}\right)^2}_{|\langle r_{21} \rangle|^2} = \frac{e^{10}\mu^3}{26244\pi^5\varepsilon_0^5\hbar^6c^3m_e^2} = 6.26 \times 10^8 \text{ second}^{-1}$$

Let us analyze the units involved. For the coefficient of A_{nm} we have

$$\frac{e^2}{3\pi\varepsilon_0\hbar c^3} \propto \frac{\text{ampere}^2 \text{second}^2}{\underbrace{e^2}_{\varepsilon_0} \underbrace{\left(\frac{\text{kilogram meter}^2}{\text{second}}\right)}_{\hbar} \underbrace{\left(\frac{\text{meter}^3}{\text{second}^3}\right)}_{c^3}} = \frac{\text{second}^2}{\text{meter}^2}$$

For the transition frequency we have

$$\omega_{21} = \frac{3e^4\mu}{128\pi^2\varepsilon_0^2\hbar^3} \propto \frac{\underbrace{(\text{ampere}^4 \text{second}^4)}_{e^4} \underbrace{\text{kilogram}}_{\mu}}{\underbrace{\left(\frac{\text{ampere}^4 \text{second}^8}{\text{kilogram}^2 \text{meter}^6}\right)}_{\varepsilon_0^2} \underbrace{\left(\frac{\text{kilogram}^3 \text{meter}^6}{\text{second}^3}\right)}_{\hbar^3}} = \text{second}^{-1}$$

For the Bohr radius we have

$$a_0 = \frac{4\pi\varepsilon_0\hbar^2}{e^2m_e} \propto \frac{\underbrace{\left(\frac{\text{ampere}^2 \text{second}^4}{\text{kilogram meter}^3}\right)}_{\varepsilon_0} \underbrace{\left(\frac{\text{kilogram}^2 \text{meter}^4}{\text{second}^2}\right)}_{\hbar^2}}{\underbrace{(\text{ampere}^2 \text{second}^2)}_{e^2} \underbrace{\text{kilogram}}_{m_e}} = \text{meter}$$

Hence

$$A_{nm} \propto \frac{\text{second}^2}{\underbrace{\text{meter}^2}_{\omega_{nm}^3}} \times \underbrace{\text{second}^{-3}}_{\omega_{nm}^3} \times \underbrace{\text{meter}^2}_{a_0^2} = \text{second}^{-1}$$