Fermi's golden rule

Find the transition rate $\Gamma_{1\rightarrow 2}$ by deriving Fermi's golden rule.

Let the perturbing Hamiltonian be

$$H_1(x,t) = 2V(x)\cos(\omega t)$$

Let $\Psi(x,t)$ be the following linear combination of the two eigenstates.

$$\Psi(x,t) = c_1(t)\psi_1(x)\exp(-i\omega_1 t) + c_2(t)\psi_2(x)\exp(-i\omega_2 t)$$

We need to solve for $c_2(t)$ to find the transition rate.

From the time-dependent Schrodinger equation we have

$$i\hbar\frac{\partial\Psi(x,t)}{\partial t} = H_0(x)\Psi(x,t) + H_1(x,t)\Psi(x,t)$$

The left-hand side works out to be

$$LHS = \left(i\hbar\frac{\partial c_1(t)}{\partial t} + \hbar\omega_1 c_1(t)\right)\psi_1(x)\exp(-i\omega_1 t)$$

$$+ \left(i\hbar\frac{\partial c_2(t)}{\partial t} + \hbar\omega_2 c_2(t)\right)\psi_2(x)\exp(-i\omega_2 t)$$

Using the time-independent Schrodinger equation the right-hand side is

$$RHS = c_1(t)E_1\psi_1(x)\exp(-i\omega_1 t) + c_2(t)E_2\psi_2(x)\exp(-i\omega_2 t) + H_1(x,t)\Psi(x,t)$$

The energy terms cancel by the substitutions $E_1 = \hbar\omega_1$ and $E_2 = \hbar\omega_2$ leaving

$$LHS = i\hbar\frac{\partial c_1(t)}{\partial t}\psi_1(x)\exp(-i\omega_1 t) + i\hbar\frac{\partial c_2(t)}{\partial t}\psi_2(x)\exp(-i\omega_2 t)$$

and

$$RHS = H_1(x,t)\Psi(x,t)$$

Hence by $LHS = RHS$ we have

$$i\hbar\frac{\partial c_1(t)}{\partial t}\psi_1(x)\exp(-i\omega_1 t) + i\hbar\frac{\partial c_2(t)}{\partial t}\psi_2(x)\exp(-i\omega_2 t) = H_1(x,t)\Psi(x,t)$$

Take the inner product of $\psi_2^*(x)$ with the above equation to obtain

$$i\hbar\frac{\partial c_2(t)}{\partial t}\exp(-i\omega_2 t) = 2c_1(t)M_{21}\cos(\omega t)\exp(-i\omega_1 t) + 2c_2(t)M_{22}\cos(\omega t)\exp(-i\omega_2 t)$$

where M_{21} and M_{22} are the matrix elements

$$M_{21} = \int \psi_2^*(x)V(x)\psi_1(x)\,dx$$

$$M_{22} = \int \psi_2^*(x)V(x)\psi_2(x)\,dx$$

Cancel exponentials.

$$i\hbar\frac{\partial c_2(t)}{\partial t} = 2c_1(t)M_{21}\cos(\omega t)\exp\big(i(\omega_2 - \omega_1)t\big) + 2c_2(t)M_{22}\cos(\omega t)$$

The initial state is $\Psi(x,0) = \psi_1(x)$ hence the initial conditions are

$$c_1(0) = 1, \quad c_2(0) = 0$$

Then for time t near the origin we can use the approximation

$$i\hbar\frac{\partial c_2(t)}{\partial t} = 2M_{21}\cos(\omega t)\exp\big(i(\omega_2 - \omega_1)t\big)$$

Solve for $c_2(t)$ by integrating.

$$c_2(t) = \frac{2M_{21}}{i\hbar}\int_0^t \cos(\omega t')\exp\big(i(\omega_2 - \omega_1)t'\big)\,dt'$$

The solution is

$$c_2(t) = -\frac{M_{21}}{\hbar}\left(\frac{\exp\big(i(\omega_2 - \omega_1 - \omega)t\big) - 1}{\omega_2 - \omega_1 - \omega} - \frac{\exp\big(i(\omega_2 - \omega_1 + \omega)t\big) - 1}{\omega_2 - \omega_1 + \omega}\right) \tag{1}$$

For ω such that $\omega \approx \omega_2 - \omega_1$ the first term dominates so discard the second term and write

$$c_2(t) = -\frac{M_{21}}{\hbar}\left(\frac{\exp\big(i(\omega_2 - \omega_1 - \omega)t\big) - 1}{\omega_2 - \omega_1 - \omega}\right)$$

Rewrite as

$$c_2(t) = -\frac{it}{\hbar}M_{21}\exp\left(i\frac{\omega_2 - \omega_1 - \omega}{2}t\right)\mathrm{sinc}\left(\frac{\omega_2 - \omega_1 - \omega}{2}t\right) \tag{2}$$

Hence

$$|c_2(t)|^2 = \frac{t^2}{\hbar^2}|M_{21}|^2\,\mathrm{sinc}^2\left(\frac{\omega_2 - \omega_1 - \omega}{2}t\right) \tag{3}$$

Let

$$E = \hbar(\omega_2 - \omega_1)$$

The total probability is

$$P_{tot} = \frac{t^2}{\hbar^2}|M_{21}|^2\int_{E-\hbar\omega/2}^{E+\hbar\omega/2}\mathrm{sinc}^2\left(\frac{E'/\hbar - \omega}{2}t\right)g(E')\,dE'$$

where $g(E')$ is the density of photon states for energy E'.

Use the approximation $g(E') \approx g(\hbar\omega)$ to obtain

$$P_{tot} = \frac{t^2}{\hbar^2}|M_{21}|^2 g(\hbar\omega)\int_{E-\hbar\omega/2}^{E+\hbar\omega/2}\mathrm{sinc}^2\left(\frac{E'/\hbar - \omega}{2}t\right)dE'$$

Let

$$y = \frac{E'/\hbar - \omega}{2} t$$

It follows that

$$E' = \frac{2\hbar y}{t} + \hbar\omega$$

and

$$dE' = \frac{2\hbar}{t} dy$$

The integration limits transform as

$$E \pm \hbar\omega/2 \rightarrow \frac{(E \pm \hbar\omega/2)/\hbar - \omega}{2} t = \frac{Et}{2\hbar} - \frac{\omega t}{2} \pm \frac{\omega t}{4} \approx \pm \frac{\omega t}{4}$$

Hence

$$P_{tot} = \frac{2t}{\hbar} |M_{21}|^2 g(\hbar\omega) \int_{-\omega t/4}^{+\omega t/4} \text{sinc}^2 y \, dy$$

Use the approximation

$$\int_{-\omega t/4}^{+\omega t/4} \text{sinc}^2 y \, dy \approx \int_{-\infty}^{\infty} \text{sinc}^2 y \, dy = \pi$$

to obtain

$$P_{tot} = \frac{2\pi t}{\hbar} |M_{21}|^2 g(\hbar\omega)$$

The transition rate is

$$\Gamma = \frac{d}{dt} P_{tot} = \frac{2\pi}{\hbar} |M_{21}|^2 g(\hbar\omega)$$

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