

Since S_1 does not depend on \mathbf{A} or ϕ , we only need S_2 and S_3 .

$$S_2 = - \sum_i e_i \int \left(\phi(\mathbf{x}_i(t), t) - \frac{1}{c} \dot{\mathbf{x}}_i(t) \cdot \mathbf{A}(\mathbf{x}_i(t), t) \right) dt \quad (9.25)$$

$$S_3 = \frac{1}{8\pi} \int \int \left(\left| -\nabla\phi - \frac{1}{c} \frac{\partial \mathbf{A}}{\partial t} \right|^2 - |\nabla \times \mathbf{A}|^2 \right) d^3\mathbf{r} dt \quad (9.26)$$

Consider equation (2.7), the classical Lagrangian equation of motion.

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = \frac{\partial L}{\partial x} \quad (2.7)$$

Extend (2.7) to three dimensions.

$$\frac{d}{dt} \dot{\nabla} L = \nabla L \quad (1)$$

where

$$\dot{\nabla} = \mathbf{i} \frac{\partial}{\partial \dot{x}} + \mathbf{j} \frac{\partial}{\partial \dot{y}} + \mathbf{k} \frac{\partial}{\partial \dot{z}} \quad \nabla = \mathbf{i} \frac{\partial}{\partial x} + \mathbf{j} \frac{\partial}{\partial y} + \mathbf{k} \frac{\partial}{\partial z}$$

and

$$\mathbf{i} = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad \mathbf{j} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad \mathbf{k} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

From equation (9.25) for a single particle, let

$$L = \phi - \frac{1}{c} (\dot{x}A_x + \dot{y}A_y + \dot{z}A_z)$$

Then

$$\begin{aligned} \frac{d}{dt} \dot{\nabla} L &= -\frac{1}{c} \frac{d}{dt} (A_x \mathbf{i} + A_y \mathbf{j} + A_z \mathbf{k}) \\ &= -\frac{1}{c} \frac{d}{dt} \mathbf{A} \end{aligned} \quad (2)$$

and

$$\begin{aligned} \nabla L &= \nabla \phi - \frac{1}{c} \left(\dot{x} \frac{\partial A_x}{\partial x} \mathbf{i} + \dot{y} \frac{\partial A_y}{\partial y} \mathbf{j} + \dot{z} \frac{\partial A_z}{\partial z} \mathbf{k} \right) \\ &= \nabla \phi - \frac{1}{c} \nabla (\dot{\mathbf{x}} \cdot \mathbf{A}) \end{aligned} \quad (3)$$

Hence by equations (1), (2), and (3)

$$-\frac{1}{c} \frac{d}{dt} \mathbf{A} = \nabla \phi - \frac{1}{c} \nabla (\dot{\mathbf{x}} \cdot \mathbf{A})$$

FIXME