

(2.2) For the Hamiltonian  $\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2 + \lambda\hat{x}^4$ , where  $\lambda$  is small, show by writing the Hamiltonian in terms of creation and annihilation operators and using perturbation theory, that the energy eigenvalues of all the levels are given by

$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega + \frac{3\lambda}{4} \left(\frac{\hbar}{m\omega}\right)^2 (2n^2 + 2n + 1) \quad (2.67)$$

For the above Hamiltonian, let  $\hat{H} = \hat{H}_0 + \lambda\hat{V}$  with

$$\hat{H}_0 = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2 \quad \hat{V} = \hat{x}^4$$

From perturbation theory we have

$$\begin{aligned} (\hat{H}_0 + \lambda\hat{V})(\psi_0 + \lambda\psi_1 + \lambda^2\psi_2 + \dots) \\ = (E_0 + \lambda E_1 + \lambda^2 E_2 + \dots)(\psi_0 + \lambda\psi_1 + \lambda^2\psi_2 + \dots) \end{aligned}$$

Because  $\lambda$  is small, discard powers of  $\lambda$ .

$$\hat{H}_0(\psi_0 + \lambda\psi_1) + \lambda\hat{V}\psi_0 = E_0(\psi_0 + \lambda\psi_1) + \lambda E_1\psi_0$$

Cancel  $\hat{H}_0\psi_0$  with  $E_0\psi_0$ .

$$\lambda\hat{H}_0\psi_1 + \lambda\hat{V}\psi_0 = \lambda E_0\psi_1 + \lambda E_1\psi_0$$

Cancel  $\lambda$ .

$$\hat{H}_0\psi_1 + \hat{V}\psi_0 = E_0\psi_1 + E_1\psi_0$$

It follows that

$$\langle\psi_0|\hat{H}_0|\psi_1\rangle + \langle\psi_0|\hat{V}|\psi_0\rangle = E_0\langle\psi_0|\psi_1\rangle + E_1\langle\psi_0|\psi_0\rangle \quad (1)$$

Because  $\hat{H}_0$  is Hermitian we have

$$\langle\psi_0|\hat{H}_0|\psi_1\rangle = \left(\langle\psi_1|\hat{H}_0|\psi_0\rangle\right)^\dagger = (E_0\langle\psi_1|\psi_0\rangle)^\dagger = E_0\langle\psi_0|\psi_1\rangle$$

Hence the first and third terms in equation (1) cancel leaving

$$\langle\psi_0|\hat{V}|\psi_0\rangle = E_1\langle\psi_0|\psi_0\rangle$$

It follows that

$$E_1 = \frac{\langle \psi_0 | \hat{V} | \psi_0 \rangle}{\langle \psi_0 | \psi_0 \rangle}$$

Then for  $|\psi_0\rangle = |n\rangle$  we have

$$E_1 = \frac{\langle n | \hat{V} | n \rangle}{\langle n | n \rangle} = \langle n | \hat{V} | n \rangle$$

Consider equation (2.12).

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a} + \hat{a}^\dagger) \quad (2.12)$$

From  $\hat{V} = \hat{x}^4$  and from equation (2.12) we have

$$E_1 = \langle n | \hat{V} | n \rangle = \langle n | \hat{x}^4 | n \rangle = \left( \frac{\hbar}{2m\omega} \right)^2 \langle n | (\hat{a} + \hat{a}^\dagger)^4 | n \rangle$$

The following expectation values are from the expansion of  $(\hat{a} + \hat{a}^\dagger)^4$ . All other terms in the expansion vanish.

$$\begin{aligned} \langle n | \hat{a} \hat{a} \hat{a}^\dagger \hat{a}^\dagger | n \rangle &= \sqrt{n+1} \sqrt{n+2} \sqrt{n+2} \sqrt{n+1} &= n^2 + 3n + 2 \\ \langle n | \hat{a} \hat{a}^\dagger \hat{a} \hat{a}^\dagger | n \rangle &= \sqrt{n+1} \sqrt{n+1} \sqrt{n+1} \sqrt{n+1} &= n^2 + 2n + 1 \\ \langle n | \hat{a} \hat{a}^\dagger \hat{a}^\dagger \hat{a} | n \rangle &= \sqrt{n+1} \sqrt{n+1} \sqrt{n} \sqrt{n} &= n^2 + n \\ \langle n | \hat{a}^\dagger \hat{a} \hat{a} \hat{a}^\dagger | n \rangle &= \sqrt{n} \sqrt{n} \sqrt{n+1} \sqrt{n+1} &= n^2 + n \\ \langle n | \hat{a}^\dagger \hat{a} \hat{a}^\dagger \hat{a} | n \rangle &= \sqrt{n} \sqrt{n} \sqrt{n} \sqrt{n} &= n^2 \\ \langle n | \hat{a}^\dagger \hat{a}^\dagger \hat{a} \hat{a} | n \rangle &= \sqrt{n} \sqrt{n-1} \sqrt{n-1} \sqrt{n} &= n^2 - n \end{aligned}$$

Hence

$$E_1 = \left( \frac{\hbar}{2m\omega} \right)^2 (6n^2 + 6n + 3)$$

Therefore

$$E_n = E_0 + \lambda E_1 = \hbar\omega \left( n + \frac{1}{2} \right) + \lambda \left( \frac{\hbar}{2m\omega} \right)^2 (6n^2 + 6n + 3)$$