Mott problem

Consider the emission of an α particle in a cloud chamber. The quantum mechanical model of an α particle is a spherical wave emanating from the origin. A spherical wave should ionize atoms throughout the cloud chamber. However, only straight tracks are observed. Nevill Mott showed that straight tracks are consistent with the Schrodinger equation.

Let **R** be the position of the α particle, let \mathbf{a}_1 and \mathbf{a}_2 be the positions of two atoms ionized by the α particle, and let \mathbf{r}_1 and \mathbf{r}_2 be the positions of the free electrons. The Hamiltonian for the system is

$$\hat{H} = \hat{K}_{\alpha} + \hat{K}_{1} + \hat{K}_{2} + U_{1} + U_{2} + V_{1} + V_{2}$$

where

$$\hat{K}_{\alpha} = -\frac{\hbar^2}{2M} \nabla_{\alpha}^2 \qquad \text{kinetic energy of } \alpha \text{ particle}$$

$$\hat{K}_1 = -\frac{\hbar^2}{2m} \nabla_1^2 \qquad \text{kinetic energy of 1st electron}$$

$$\hat{K}_2 = -\frac{\hbar^2}{2m} \nabla_2^2 \qquad \text{kinetic energy of 2nd electron}$$

$$U_1 = -\frac{e^2}{|\mathbf{r}_1 - \mathbf{a}_1|} \qquad \text{potential energy of 1st electron}$$

$$U_2 = -\frac{e^2}{|\mathbf{r}_2 - \mathbf{a}_2|} \qquad \text{potential energy of 2nd electron}$$

$$V_1 = -\frac{2e^2}{|\mathbf{R} - \mathbf{r}_1|} \qquad \text{potential energy of } \alpha \text{ and 1st electron}$$

$$V_2 = -\frac{2e^2}{|\mathbf{R} - \mathbf{r}_2|} \qquad \text{potential energy of } \alpha \text{ and 2nd electron}$$

Let ψ_1 and ψ_2 be atomic wavefunctions such that

$$\hat{H}\psi_1 = E_1\psi_1, \quad \hat{H}\psi_2 = E_2\psi_2$$

We want to find a wavefunction $F(\mathbf{R}, \mathbf{r}_1, \mathbf{r}_2)$ such that

$$\hat{H}F=EF$$

Let

$$F = F_0 + F_1 + F_2 + \cdots$$

and let

$$\hat{H}_0 = \hat{K}_\alpha + E_1 + E_2$$

Start by finding an F_0 such that

$$\hat{H}_0 F_0 = E F_0$$

The solution is

$$F_0(\mathbf{R}, \mathbf{r}_1, \mathbf{r}_2) = f_0(\mathbf{R})\psi_1(\mathbf{r}_1 - \mathbf{a}_1)\psi_2(\mathbf{r}_2 - \mathbf{a}_2)$$
(1)

where

$$f_0(\mathbf{R}) = \frac{1}{|\mathbf{R}|} \exp\left(\frac{ik|\mathbf{R}|}{\hbar}\right), \quad k = \sqrt{2M(E - E_1 - E_2)}$$

It follows that for the full Hamiltonian

$$\hat{H}F_0 = EF_0 + (V_1 + V_2)F_0$$

To cancel $(V_1 + V_2)F_0$ from the full Hamiltonian, find an F_1 such that

$$\hat{H}_0 F_1 = EF_1 - (V_1 + V_2)F_0$$

Rewrite as

$$(\hat{H}_0 - E) F_1 = -(V_1 + V_2) F_0$$

Expand F_1 and F_0 .

$$(\hat{H}_0 - E) f_1(\mathbf{R}) \psi_1(\mathbf{r}_1 - \mathbf{a}_1) \psi_2(\mathbf{r}_2 - \mathbf{a}_2) = -(V_1 + V_2) f_0(\mathbf{R}) \psi_1(\mathbf{r}_1 - \mathbf{a}_1) \psi_2(\mathbf{r}_2 - \mathbf{a}_2)$$

To solve for $f_1(\mathbf{R})$ multiply both sides by

$$\psi_1^*(\mathbf{r}_1 - \mathbf{a}_1)\psi_2^*(\mathbf{r}_2 - \mathbf{a}_2)$$

and integrate over \mathbf{r}_1 and \mathbf{r}_2 to obtain

$$\left(\hat{H}_0 - E\right) f_1(\mathbf{R}) = V_1(\mathbf{R}) f_0(\mathbf{R}) + V_2(\mathbf{R}) f_0(\mathbf{R}) \tag{2}$$

where

$$V_1(\mathbf{R}) = 2e^2 \int \frac{|\psi_1(\mathbf{r})|^2}{|\mathbf{R} - \mathbf{a}_1 - \mathbf{r}|} d\mathbf{r}, \quad V_2(\mathbf{R}) = 2e^2 \int \frac{|\psi_2(\mathbf{r})|^2}{|\mathbf{R} - \mathbf{a}_2 - \mathbf{r}|} d\mathbf{r}$$

Per Mott the solution to (2) is

$$f_1(\mathbf{R}) = \frac{M}{2\pi\hbar^2} \int \frac{V_1(\mathbf{r}) f_0(\mathbf{r})}{|\mathbf{R} - \mathbf{r}|} \exp\left(\frac{ik|\mathbf{R} - \mathbf{r}|}{\hbar}\right) d\mathbf{r} + \frac{M}{2\pi\hbar^2} \int \frac{V_2(\mathbf{r}) f_0(\mathbf{r})}{|\mathbf{R} - \mathbf{r}|} \exp\left(\frac{ik|\mathbf{R} - \mathbf{r}|}{\hbar}\right) d\mathbf{r}$$

Let I_1 be the first integral.

$$I_1(\mathbf{R}) = \frac{M}{2\pi\hbar^2} \int \frac{V_1(\mathbf{r}) f_0(\mathbf{r})}{|\mathbf{R} - \mathbf{r}|} \exp\left(\frac{ik|\mathbf{R} - \mathbf{r}|}{\hbar}\right) d\mathbf{r}$$

Substitute for $f_0(\mathbf{r})$ to obtain

$$I_1(\mathbf{R}) = \frac{M}{2\pi\hbar^2} \int \frac{V_1(\mathbf{r})}{|\mathbf{R} - \mathbf{r}||\mathbf{r}|} \exp\left(\frac{ik|\mathbf{R} - \mathbf{r}|}{\hbar} + \frac{ik|\mathbf{r}|}{\hbar}\right) d\mathbf{r}$$

Change of variable $\mathbf{r} \to \mathbf{r} + \mathbf{a}_1$

$$I_1(\mathbf{R}) = \frac{M}{2\pi\hbar^2} \int \frac{V_1(\mathbf{r} + \mathbf{a}_1)}{|\mathbf{R} - \mathbf{r} - \mathbf{a}_1||\mathbf{r} + \mathbf{a}_1|} \exp\left(\frac{ik|\mathbf{R} - \mathbf{r} - \mathbf{a}_1|}{\hbar} + \frac{ik|\mathbf{r} + \mathbf{a}_1|}{\hbar}\right) d\mathbf{r}$$

Per Mott (see also Figari and Teta)

$$I_{1}(\mathbf{R}) \approx \frac{1}{|\mathbf{R} - \mathbf{a}_{1}|} \exp\left(\frac{ik|\mathbf{R} - \mathbf{a}_{1}|}{\hbar}\right) \times \frac{M}{2\pi\hbar^{2}} \int \frac{V_{1}(\mathbf{r} + \mathbf{a}_{1})}{|\mathbf{r} + \mathbf{a}_{1}|} \exp\left(-\frac{ik\mathbf{u}_{1}(\mathbf{R}) \cdot \mathbf{r}}{\hbar} + \frac{ik|\mathbf{r} + \mathbf{a}_{1}|}{\hbar}\right) d\mathbf{r}$$

where

$$\mathbf{u}_1(\mathbf{R}) = rac{\mathbf{R} - \mathbf{a}_1}{|\mathbf{R} - \mathbf{a}_1|}$$

The condition for stationary phase is

$$g' = \frac{d}{d\mathbf{r}} \left(-\mathbf{u}_1(\mathbf{R}) \cdot \mathbf{r} + |\mathbf{r} + \mathbf{a}_1| \right) = -\mathbf{u}_1(\mathbf{R}) + \frac{\mathbf{r} + \mathbf{a}_1}{|\mathbf{r} + \mathbf{a}_1|} = 0$$

Note that $V_1(\mathbf{r} + \mathbf{a}_1)$ is small except for $\mathbf{r} \approx 0$ so we only require stationarity at the origin. Hence for $\mathbf{r} = 0$ the integral is stationary (g' = 0) when \mathbf{R} satisfies the condition

$$\mathbf{u}_1(\mathbf{R}) = rac{\mathbf{R} - \mathbf{a}_1}{|\mathbf{R} - \mathbf{a}_1|} = rac{\mathbf{a}_1}{|\mathbf{a}_1|}$$

By symmetry of the integrals, I_2 is stationary when **R** satisfies the condition

$$\mathbf{u}_2(\mathbf{R}) = \frac{\mathbf{R} - \mathbf{a}_2}{|\mathbf{R} - \mathbf{a}_2|} = \frac{\mathbf{a}_2}{|\mathbf{a}_2|}$$

Gentle reminder that

$$f_1(\mathbf{R}) = I_1(\mathbf{R}) + I_2(\mathbf{R})$$

Since both atoms are ionized, it follows that both V_1 and V_2 should contribute to f_1 . Hence both integrals should be stationary for the same values of \mathbf{R} . (If not stationary then the integral vanishes.) Simultaneous stationarity requires that

$$\mathbf{u}_1 = \mathbf{u}_2$$

Hence

$$\frac{\mathbf{R} - \mathbf{a}_1}{|\mathbf{R} - \mathbf{a}_1|} = \frac{\mathbf{R} - \mathbf{a}_2}{|\mathbf{R} - \mathbf{a}_2|}$$

and finally

$$\frac{\mathbf{a}_1}{|\mathbf{a}_1|} = \frac{\mathbf{a}_2}{|\mathbf{a}_2|}$$

This condition means that \mathbf{a}_1 and \mathbf{a}_2 are on the same ray emanating from the origin. If we keep \mathbf{a}_1 and choose a different \mathbf{a}_2 we get the same ray. It follows that all of the ionized atoms are on the same ray. Hence straight tracks are consistent with the Schrodinger equation.