Spin

Spin state $|s\rangle$ is a normalized vector in \mathbb{C}^2 .

$$|s\rangle = \begin{pmatrix} c_+ \\ c_- \end{pmatrix}, \quad |c_+|^2 + |c_-|^2 = 1$$

Spin measurement probabilities are the transition probabilities from $|s\rangle$ to an eigenstate.

For spin measurements in the z direction we have

$$\Pr\left(S_z = +\frac{\hbar}{2}\right) = |\langle z_+|s\rangle|^2$$

$$\Pr\left(S_z = -\frac{\hbar}{2}\right) = |\langle z_-|s\rangle|^2$$

Define the z eigenstates as

$$|z_{+}\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |z_{-}\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

By definition of expectation value we have

$$\langle S_z \rangle = \frac{\hbar}{2} \Pr \left(S_z = +\frac{\hbar}{2} \right) - \frac{\hbar}{2} \Pr \left(S_z = -\frac{\hbar}{2} \right)$$

Rewrite as

$$\langle S_z \rangle = \frac{\hbar}{2} |\langle z_+ | s \rangle|^2 - \frac{\hbar}{2} |\langle z_- | s \rangle|^2$$

Rewrite again as

$$\langle S_z \rangle = \frac{\hbar}{2} \langle s|z_+ \rangle \langle z_+|s \rangle - \frac{\hbar}{2} \langle s|z_- \rangle \langle z_-|s \rangle$$

Then by

$$\langle S_z \rangle = \langle s | S_z | s \rangle$$

we have

$$S_z = \frac{\hbar}{2} |z_+\rangle \langle z_+| - \frac{\hbar}{2} |z_-\rangle \langle z_-| = \frac{\hbar}{2} \begin{pmatrix} 1 & 0\\ 0 & -1 \end{pmatrix}$$

From the commutator

$$S_{+}S_{-} - S_{-}S_{+} = 2\hbar S_{z}$$

we have

$$S_{+}S_{-} - S_{-}S_{+} = \hbar^{2}|z_{+}\rangle\langle z_{+}| - \hbar^{2}|z_{-}\rangle\langle z_{-}|$$

Rewrite as

$$S_+S_- - S_-S_+ = \hbar^2 |z_+\rangle \langle z_-|z_-\rangle \langle z_+| - \hbar^2 |z_-\rangle \langle z_+|z_+\rangle \langle z_-|$$

Hence

$$S_{+} = \hbar |z_{+}\rangle\langle z_{-}| = \hbar \begin{pmatrix} 0 & 1\\ 0 & 0 \end{pmatrix}$$
$$S_{-} = \hbar |z_{-}\rangle\langle z_{+}| = \hbar \begin{pmatrix} 0 & 0\\ 1 & 0 \end{pmatrix}$$

Then by

$$S_{+} = S_x + iS_y$$
$$S_{-} = S_x - iS_y$$

we obtain

$$S_x = \frac{S_+ + S_-}{2} = \frac{\hbar}{2} \begin{pmatrix} 0 & 1\\ 1 & 0 \end{pmatrix}$$
$$S_y = \frac{S_+ - S_-}{2i} = \frac{\hbar}{2} \begin{pmatrix} 0 & -i\\ i & 0 \end{pmatrix}$$

By solving for the eigenstates in

$$S_x|x_{\pm}\rangle = \pm \frac{\hbar}{2}|x_{\pm}\rangle$$

 $S_y|y_{\pm}\rangle = \pm \frac{\hbar}{2}|y_{\pm}\rangle$

we obtain

$$|x_{+}\rangle = \frac{|z_{+}\rangle + |z_{-}\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\1 \end{pmatrix}$$
$$|x_{-}\rangle = \frac{|z_{+}\rangle - |z_{-}\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\-1 \end{pmatrix}$$

and

$$|y_{+}\rangle = \frac{|z_{+}\rangle + i|z_{-}\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\i \end{pmatrix}$$
$$|y_{-}\rangle = \frac{|z_{+}\rangle - i|z_{-}\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\-i \end{pmatrix}$$

The expected spin direction vector is

$$\langle \mathbf{S} \rangle = \begin{pmatrix} \langle s | S_x | s \rangle \\ \langle s | S_y | s \rangle \\ \langle s | S_z | s \rangle \end{pmatrix}, \quad |\langle \mathbf{S} \rangle| = \frac{\hbar}{2}$$

To convert a direction vector to a spin state, let θ and ϕ be the polar and azimuth angles such that

$$\langle \mathbf{S} \rangle = \frac{\hbar}{2} \begin{pmatrix} \sin \theta \cos \phi \\ \sin \theta \sin \phi \\ \cos \theta \end{pmatrix}$$

Then

$$|s\rangle = \begin{pmatrix} \cos(\theta/2) \\ \sin(\theta/2) \exp(i\phi) \end{pmatrix}$$

By the identities

$$\cos^{2}(\theta/2) = \frac{\cos \theta + 1}{2}, \quad \sin^{2}(\theta/2) = \frac{1 - \cos \theta}{2}$$

and noting that $0 \le \theta \le \pi$ we have

$$\cos(\theta/2) = \sqrt{\frac{\langle z \rangle + 1}{2}}, \quad \sin(\theta/2) \exp(i\phi) = \sqrt{\frac{1 - \langle z \rangle}{2}} \frac{\langle x \rangle + i \langle y \rangle}{\sqrt{\langle x \rangle^2 + \langle y \rangle^2}}$$

where

$$\langle x \rangle = \frac{2}{\hbar} \langle S_x \rangle = \sin \theta \cos \phi$$
$$\langle y \rangle = \frac{2}{\hbar} \langle S_y \rangle = \sin \theta \sin \phi$$
$$\langle z \rangle = \frac{2}{\hbar} \langle S_z \rangle = \cos \theta$$

The following commutator was used to derive S_x and S_y .

$$S_{+}S_{-} - S_{-}S_{+} = 2\hbar S_{z}$$

The commutator follows from

$$\hat{\mathbf{S}}\psi = (\mathbf{r} \times \hat{\mathbf{p}})\psi, \quad \mathbf{r} = \begin{pmatrix} x \\ y \\ z \end{pmatrix}, \quad \hat{\mathbf{p}} = -i\hbar \begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{pmatrix}$$

Rewrite in component form.

$$\hat{S}_x \psi = -i\hbar \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \psi$$

$$\hat{S}_y \psi = -i\hbar \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \psi$$

$$\hat{S}_z \psi = -i\hbar \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) \psi$$

By computer algebra we have

$$(\hat{S}_y \hat{S}_z - \hat{S}_z \hat{S}_y) \psi = i\hbar \hat{S}_x \psi$$
$$(\hat{S}_z \hat{S}_x - \hat{S}_x \hat{S}_z) \psi = i\hbar \hat{S}_y \psi$$
$$(\hat{S}_x \hat{S}_y - \hat{S}_y \hat{S}_x) \psi = i\hbar \hat{S}_z \psi$$

Let

$$\hat{S}_{+} = \hat{S}_x + i\hat{S}_y$$

$$\hat{S}_{-} = \hat{S}_x - i\hat{S}_y$$

Then by computer algebra

$$(\hat{S}_{+}\hat{S}_{-} - \hat{S}_{-}\hat{S}_{+})\psi = 2\hbar\hat{S}_{z}\psi$$

1. Verify that

$$S_x = \frac{\hbar}{2}(|x_+\rangle\langle x_+| - |x_-\rangle\langle x_-|)$$

$$S_y = \frac{\hbar}{2}(|y_+\rangle\langle y_+| - |y_-\rangle\langle y_-|)$$

$$S_z = \frac{\hbar}{2}(|z_+\rangle\langle z_+| - |z_-\rangle\langle z_-|)$$

```
zp = (1,0)
zm = (0,1)

xp = (zp + zm) / sqrt(2)
xm = (zp - zm) / sqrt(2)

yp = (zp + i zm) / sqrt(2)

ym = (zp - i zm) / sqrt(2)

Sx = hbar / 2 ((0,1),(1,0))
Sy = hbar / 2 ((0,-i),(i,0))
Sz = hbar / 2 ((1,0),(0,-1))

check(Sx == hbar / 2 (outer(xp,conj(xp)) - outer(xm,conj(xm))))
check(Sy == hbar / 2 (outer(yp,conj(yp)) - outer(ym,conj(ym))))
check(Sz == hbar / 2 (outer(zp,conj(zp)) - outer(zm,conj(zm))))
```

2. Let $|s\rangle$ be the following spin state.

$$|s\rangle = \begin{pmatrix} \frac{1}{3} - \frac{2}{3}i\\ \frac{2}{3} \end{pmatrix}$$

Verify that $|s\rangle$ is normalized and that

$$\langle \mathbf{S} \rangle = \langle s | \mathbf{S} | s \rangle = \frac{\hbar}{2} \begin{pmatrix} \frac{4}{9} \\ \frac{8}{9} \\ \frac{1}{9} \end{pmatrix}$$

where

$$\mathbf{S} = \begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix}$$

Note: In component form we have

$$\langle s|\mathbf{S}|s\rangle = s_{\beta}^* S^{\alpha\beta}{}_{\gamma} s^{\gamma}$$

Eigenmath requires a transpose so that the β indices are adjacent.

$$\langle s|\mathbf{S}|s\rangle = s_{\beta}^* S^{\beta\alpha}{}_{\gamma} s^{\gamma}$$

```
s = (1/3 - 2/3 i, 2/3)
check(dot(conj(s),s) == 1)

Sx = hbar / 2 ((0,1),(1,0))
Sy = hbar / 2 ((0,-i),(i,0))
Sz = hbar / 2 ((1,0),(0,-1))

S = (Sx,Sy,Sz)
check(dot(conj(s),transpose(S),s) == hbar / 2 (4/9, 8/9, 1/9))
```

3. Let $|s\rangle$ be the following spin state.

$$|s\rangle = \begin{pmatrix} \frac{1}{3} - \frac{2}{3}i\\ \frac{2}{3} \end{pmatrix}$$

Verify the following measurement probabilities for $|s\rangle$.

$$\Pr\left(S_x = +\frac{\hbar}{2}\right) = |\langle x_+|s\rangle|^2 = \frac{13}{18}$$

$$\Pr\left(S_x = -\frac{\hbar}{2}\right) = |\langle x_-|s\rangle|^2 = \frac{5}{18}$$

$$\Pr (S_y = +\frac{\hbar}{2}) = |\langle y_+ | s \rangle|^2 = \frac{17}{18}$$

$$\Pr (S_y = -\frac{\hbar}{2}) = |\langle y_- | s \rangle|^2 = \frac{1}{18}$$

$$\Pr \left(S_z = + \frac{\hbar}{2} \right) = |\langle z_+ | s \rangle|^2 = \frac{5}{9}$$

$$\Pr \left(S_z = -\frac{\hbar}{2} \right) = |\langle z_- | s \rangle|^2 = \frac{4}{9}$$

$$s = (1/3 - 2/3 i, 2/3)$$

 $zp = (1,0)$

zm = (0,1)

$$xp = (zp + zm) / sqrt(2)$$

$$xm = (zp - zm) / sqrt(2)$$

$$yp = (zp + i zm) / sqrt(2)$$

 $ym = (zp - i zm) / sqrt(2)$

$$check(Pr(xp,s) == 13/18)$$

$$check(Pr(xm,s) == 5/18)$$

$$check(Pr(yp,s) == 17/18)$$

$$check(Pr(ym,s) == 1/18)$$

$$check(Pr(zp,s) == 5/9)$$

$$check(Pr(zm,s) == 4/9)$$

4. Let $|s\rangle$ be the following spin state.

$$|s\rangle = \begin{pmatrix} \frac{1}{3} - \frac{2}{3}i\\ \frac{2}{3} \end{pmatrix}$$

Verify that the following spin state $|\chi\rangle$ is indistinguishable from $|s\rangle$.

$$|\chi\rangle = \begin{pmatrix} \cos(\theta/2) \\ \sin(\theta/2) \exp(i\phi) \end{pmatrix}$$

where

$$\cos(\theta/2) = \sqrt{\frac{\langle z \rangle + 1}{2}} = \frac{\sqrt{5}}{3}$$

and

$$\sin(\theta/2)\exp(i\phi) = \sqrt{\frac{1-\langle z\rangle}{2}} \frac{\langle x\rangle + i\langle y\rangle}{\sqrt{\langle x\rangle^2 + \langle y\rangle^2}} = \frac{2+4i}{3\sqrt{5}}$$

with

$$\langle x \rangle = \frac{2}{\hbar} \langle S_x \rangle$$
$$\langle y \rangle = \frac{2}{\hbar} \langle S_y \rangle$$
$$\langle z \rangle = \frac{2}{\hbar} \langle S_z \rangle$$

```
s = (1/3 - 2/3 i, 2/3)

Sx = hbar / 2 ((0,1),(1,0))
Sy = hbar / 2 ((0,-i),(i,0))
Sz = hbar / 2 ((1,0),(0,-1))

S = (Sx,Sy,Sz)

x = 2 / hbar dot(conj(s),Sx,s)
y = 2 / hbar dot(conj(s),Sy,s)
z = 2 / hbar dot(conj(s),Sz,s)

cp = sqrt((z + 1) / 2)
cm = sqrt((1 - z) / 2) (x + i y) / sqrt(x^2 + y^2)

check(cp == sqrt(5) / 3)
check(cm == (2 + 4 i) / (3 sqrt(5)))

chi = (cp,cm)

check(dot(conj(s),transpose(S),s) == dot(conj(chi),transpose(S),chi))
```

5. Verify the following spin commutation relations using $\mathbf{S}\psi = (\mathbf{r} \times \mathbf{p})\psi$.

$$\begin{split} [S_x, S_y] &= i\hbar S_z \\ [S_y, S_z] &= i\hbar S_x \\ [S_z, S_x] &= i\hbar S_y \\ [S^2, S_x] &= 0 \\ [S^2, S_y] &= 0 \\ [S^2, S_z] &= 0 \\ [S_+, S_-] &= 2\hbar S_z \end{split}$$

where

$$S^2 = S_x^2 + S_y^2 + S_z^2$$

and

$$S_{+} = S_x + iS_y$$
$$S_{-} = S_x - iS_y$$

```
Sx(psi) = -i hbar (y d(psi,z) - z d(psi,y))
Sy(psi) = -i hbar (z d(psi,x) - x d(psi,z))
Sz(psi) = -i hbar (x d(psi,y) - y d(psi,x))

psi = Psi()

check(Sy(Sz(psi)) - Sz(Sy(psi)) == i hbar Sx(psi))
check(Sz(Sx(psi)) - Sx(Sz(psi)) == i hbar Sy(psi))
check(Sx(Sy(psi)) - Sy(Sx(psi)) == i hbar Sz(psi))

S2(psi) = Sx(Sx(psi)) + Sy(Sy(psi)) + Sz(Sz(psi))

check(S2(Sx(psi)) - Sx(S2(psi)) == 0)
check(S2(Sy(psi)) - Sy(S2(psi)) == 0)
check(S2(Sz(psi)) - Sz(S2(psi)) == 0)

Sp(psi) = Sx(psi) + i Sy(psi)
Sm(psi) = Sx(psi) - i Sy(psi)

check(Sp(Sm(psi)) - Sm(Sp(psi)) == 2 hbar Sz(psi))
```