

## Atomic transitions 2

Start with the perturbing Hamiltonian where  $E_0$  is electric field strength.

$$H_1(\mathbf{r}, t) = -\frac{eE_0\boldsymbol{\epsilon} \cdot \mathbf{p}}{m\omega} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t)$$

In exponential form

$$H_1(\mathbf{r}, t) = -\frac{eE_0\boldsymbol{\epsilon} \cdot \mathbf{p}}{m\omega} \left( \frac{1}{2} \exp(i\mathbf{k} \cdot \mathbf{r} - i\omega t) + \frac{1}{2} \exp(-i\mathbf{k} \cdot \mathbf{r} + i\omega t) \right)$$

Given the initial condition  $c_b(0) = 0$  the first-order approximation for  $c_b(t)$  is

$$c_b(t) = -\frac{i}{\hbar} \int_0^t \langle \psi_b | H_1(\mathbf{r}, t') | \psi_a \rangle \exp(i\omega_0 t') dt', \quad \omega_0 = \frac{E_b - E_a}{\hbar}$$

Factor the integrand.

$$\begin{aligned} c_b(t) = & \frac{ieE_0}{2m\hbar\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \int_0^t \exp(-i\omega t') \exp(i\omega_0 t') dt' \\ & + \frac{ieE_0}{2m\hbar\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(-i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \int_0^t \exp(i\omega t') \exp(i\omega_0 t') dt' \end{aligned}$$

Solve the integrals to obtain

$$\begin{aligned} c_b(t) = & \frac{eE_0}{2m\hbar\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \frac{\exp(i(\omega_0 - \omega)t) - 1}{\omega_0 - \omega} \\ & + \frac{eE_0}{2m\hbar\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(-i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \frac{\exp(i(\omega_0 + \omega)t) - 1}{\omega_0 + \omega} \quad (1) \end{aligned}$$

As an approximation, discard the second term since the first term dominates for  $\omega \approx \omega_0$ .

$$c_b(t) = \frac{eE_0}{2m\hbar\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \frac{\exp(i(\omega_0 - \omega)t) - 1}{\omega_0 - \omega}$$

Rewrite  $c_b(t)$  in the form of a sine function.

$$c_b(t) = \frac{ieE_0}{m\hbar\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \frac{\sin\left(\frac{1}{2}(\omega_0 - \omega)t\right)}{\omega_0 - \omega} \exp\left(\frac{i}{2}(\omega_0 - \omega)t\right) \quad (2)$$

Verify dimensions.

$$\begin{aligned} \frac{eE_0\boldsymbol{\epsilon} \cdot \mathbf{p}}{m\omega} & \propto \frac{\text{coulomb}}{m} \frac{e}{\text{kilogram}} \frac{E_0}{\text{joule second}} \frac{\boldsymbol{\epsilon} \cdot \mathbf{p}}{\omega} = \text{joule} \\ & \quad \text{second}^{-1} \\ c_b(t) & \propto \frac{\text{coulomb}}{m} \frac{e}{\text{kilogram}} \frac{E_0}{\text{joule second}} \frac{\langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle}{\omega} \\ & \quad \frac{\text{momentum}}{\text{second}^{-1}} = 1 \end{aligned}$$

Wave functions  $\psi_a$  and  $\psi_b$  have dimension meter<sup>-1/2</sup> hence they cancel with  $dx \propto$  meter in the integral leaving units of momentum due to  $\mathbf{p}$ .