

Feynman and Hibbs problem 4-2

For a particle of charge e in a magnetic field the Lagrangian is

$$L(\dot{\mathbf{x}}, \mathbf{x}) = \frac{m}{2} \dot{\mathbf{x}}^2 + \frac{e}{c} \dot{\mathbf{x}} \cdot \mathbf{A}(\mathbf{x}, t) - e\phi(\mathbf{x}, t)$$

where $\dot{\mathbf{x}}$ is the velocity vector, c is the velocity of light, and \mathbf{A} and ϕ are the vector and scalar potentials. Show that the corresponding Schrodinger equation is

$$\frac{\partial \psi}{\partial t} = -\frac{i}{\hbar} \left(\frac{1}{2m} \left(\frac{\hbar}{i} \nabla - \frac{e}{c} \mathbf{A} \right) \cdot \left(\frac{\hbar}{i} \nabla - \frac{e}{c} \mathbf{A} \right) \psi + e\phi \psi \right)$$

From equation (4.3)

$$\psi(\mathbf{x}, t + \epsilon) = \frac{1}{A} \int_{\mathbb{R}^3} \exp \left(\frac{i\epsilon}{\hbar} L \left(\frac{\mathbf{x} - \mathbf{y}}{\epsilon}, \frac{\mathbf{x} + \mathbf{y}}{2} \right) \right) \psi(\mathbf{y}, t) dy_1 dy_2 dy_3 \quad (1)$$

This is the Lagrangian with arguments from (1).

$$\begin{aligned} L \left(\frac{\mathbf{x} - \mathbf{y}}{\epsilon}, \frac{\mathbf{x} + \mathbf{y}}{2} \right) \\ = \frac{m}{2\epsilon^2} (\mathbf{x} - \mathbf{y})^2 + \frac{e}{c\epsilon} (\mathbf{x} - \mathbf{y}) \cdot \mathbf{A} \left(\frac{\mathbf{x} + \mathbf{y}}{2}, t \right) - e\phi \left(\frac{\mathbf{x} + \mathbf{y}}{2}, t \right) \end{aligned}$$

Hence

$$\begin{aligned} \psi(\mathbf{x}, t + \epsilon) &= \frac{1}{A} \int_{\mathbb{R}^3} \\ &\exp \left(\frac{im}{2\hbar\epsilon} (\mathbf{x} - \mathbf{y})^2 + \frac{ie}{\hbar c} (\mathbf{x} - \mathbf{y}) \cdot \mathbf{A} \left(\frac{\mathbf{x} + \mathbf{y}}{2}, t \right) - \frac{ie\epsilon}{\hbar} \phi \left(\frac{\mathbf{x} + \mathbf{y}}{2}, t \right) \right) \\ &\times \psi(\mathbf{y}, t) dy_1 dy_2 dy_3 \end{aligned}$$

Let

$$\mathbf{y} = \mathbf{x} + \boldsymbol{\eta}$$

Then

$$\mathbf{x} - \mathbf{y} = -\boldsymbol{\eta}, \quad \frac{\mathbf{x} + \mathbf{y}}{2} = \mathbf{x} + \frac{1}{2}\boldsymbol{\eta}, \quad dy_1 dy_2 dy_3 = d\eta_1 d\eta_2 d\eta_3$$

Hence

$$\begin{aligned} \psi(\mathbf{x}, t + \epsilon) = \frac{1}{A} \int_{\mathbb{R}^3} \exp \left(\frac{im}{2\hbar\epsilon} \boldsymbol{\eta}^2 + \frac{ie}{\hbar c} \boldsymbol{\eta} \cdot \mathbf{A} \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) - \frac{ie\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \\ \times \psi(\mathbf{x} + \boldsymbol{\eta}, t) d\eta_1 d\eta_2 d\eta_3 \end{aligned}$$

Factor the exponential.

$$\begin{aligned} \psi(\mathbf{x}, t + \epsilon) = \frac{1}{A} \int_{\mathbb{R}^3} \\ \exp \left(\frac{im}{2\hbar\epsilon} \boldsymbol{\eta}^2 \right) \exp \left(\frac{ie}{\hbar c} \boldsymbol{\eta} \cdot \mathbf{A} \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \exp \left(-\frac{ie\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \\ \times \psi(\mathbf{x} + \boldsymbol{\eta}, t) d\eta_1 d\eta_2 d\eta_3 \end{aligned} \quad (2)$$

From the identity $\exp(i\theta) = \cos(\theta) + i \sin(\theta)$ we have

$$\begin{aligned} \exp \left(-\frac{ie\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \\ = \cos \left(-\frac{e\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) + i \sin \left(-\frac{e\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \end{aligned}$$

Then for small ϵ

$$\exp \left(-\frac{ie\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \approx 1 - \frac{ie\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right)$$

The authors write that the $\boldsymbol{\eta}$ term can be dropped “because the error is of higher order than ϵ .” Hence

$$\exp \left(-\frac{ie\epsilon}{\hbar} \phi \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \approx 1 - \frac{ie\epsilon}{\hbar} \phi(\mathbf{x}, t) \quad (3)$$

Substitute (3) into (2).

$$\begin{aligned} \psi(\mathbf{x}, t + \epsilon) = \frac{1}{A} \left(1 - \frac{ie\epsilon}{\hbar} \phi(\mathbf{x}, t) \right) \int_{\mathbb{R}^3} \\ \exp \left(\frac{im}{2\hbar\epsilon} \boldsymbol{\eta}^2 \right) \exp \left(\frac{ie}{\hbar c} \boldsymbol{\eta} \cdot \mathbf{A} \left(\mathbf{x} + \frac{1}{2} \boldsymbol{\eta}, t \right) \right) \psi(\mathbf{x} + \boldsymbol{\eta}, t) d\eta_1 d\eta_2 d\eta_3 \end{aligned}$$

Let

$$T = -\frac{ie}{\hbar c} (\eta_1 A_1 + \eta_2 A_2 + \eta_3 A_3)$$

Then use the approximation (need to justify this)

$$\exp\left(\frac{ie}{\hbar c}\boldsymbol{\eta} \cdot \mathbf{A}(\mathbf{x} + \tfrac{1}{2}\boldsymbol{\eta}, t)\right) \approx 1 + T + \tfrac{1}{2}T^2$$

Hence

$$\begin{aligned} \psi(\mathbf{x}, t + \epsilon) &= \frac{1}{A} \left(1 - \frac{ie\epsilon}{\hbar} \phi(\mathbf{x}, t)\right) \\ &\times \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) (1 + T + \tfrac{1}{2}T^2) \psi(\mathbf{x} + \boldsymbol{\eta}, t) d\eta_1 d\eta_2 d\eta_3 \end{aligned} \quad (4)$$

Next we will use the following Taylor series approximations.

$$\begin{aligned} \psi(\mathbf{x}, t + \epsilon) &\approx \psi(\mathbf{x}, t) + \epsilon \frac{\partial \psi}{\partial t} \\ \psi(\mathbf{x} + \boldsymbol{\eta}, t) &\approx \psi(\mathbf{x}, t) + \boldsymbol{\eta} \cdot \nabla \psi + \tfrac{1}{2}\boldsymbol{\eta} \cdot \nabla(\boldsymbol{\eta} \cdot \nabla \psi) \end{aligned} \quad (5)$$

Note: In component notation

$$\boldsymbol{\eta} \cdot \nabla \psi = \eta_1 \frac{\partial \psi}{\partial x_1} + \eta_2 \frac{\partial \psi}{\partial x_2} + \eta_3 \frac{\partial \psi}{\partial x_3}$$

and

$$\begin{aligned} \boldsymbol{\eta} \cdot \nabla(\boldsymbol{\eta} \cdot \nabla \psi) &= \eta_1^2 \frac{\partial^2 \psi}{\partial x_1^2} + \eta_2^2 \frac{\partial^2 \psi}{\partial x_2^2} + \eta_3^2 \frac{\partial^2 \psi}{\partial x_3^2} \\ &+ 2\eta_1\eta_2 \frac{\partial^2 \psi}{\partial x_1 \partial x_2} + 2\eta_1\eta_3 \frac{\partial^2 \psi}{\partial x_1 \partial x_3} + 2\eta_2\eta_3 \frac{\partial^2 \psi}{\partial x_2 \partial x_3} \end{aligned}$$

Substitute the approximations (5) into (4).

$$\begin{aligned} \psi(\mathbf{x}, t) + \epsilon \frac{\partial \psi}{\partial t} &= \frac{1}{A} \left(1 - \frac{ie\epsilon}{\hbar} \phi(\mathbf{x}, t)\right) \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \\ &\times (1 + T + \tfrac{1}{2}T^2) \left(\psi(\mathbf{x}, t) + \boldsymbol{\eta} \cdot \nabla \psi + \tfrac{1}{2}\boldsymbol{\eta} \cdot \nabla(\boldsymbol{\eta} \cdot \nabla \psi)\right) d\eta_1 d\eta_2 d\eta_3 \end{aligned} \quad (6)$$

To solve the above integral, we will use the following formulas provided by the authors.

$$I_k = \int_{-\infty}^{\infty} \exp\left(\frac{im\eta_k^2}{2\hbar\epsilon}\right) d\eta_k = \left(\frac{2\pi i\hbar\epsilon}{m}\right)^{1/2} \quad \text{FH (4.7)}$$

$$J_k = \int_{-\infty}^{\infty} \eta_k \exp\left(\frac{im\eta_k^2}{2\hbar\epsilon}\right) d\eta_k = 0 \quad \text{FH (4.9)}$$

$$K_k = \int_{-\infty}^{\infty} \eta_k^2 \exp\left(\frac{im\eta_k^2}{2\hbar\epsilon}\right) d\eta_k = \frac{i\hbar\epsilon}{m} \left(\frac{2\pi i\hbar\epsilon}{m}\right)^{1/2} \quad \text{FH (4.10)}$$

Hence

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \psi(\mathbf{x}, t) d\eta_1 d\eta_2 d\eta_3 \\
&= I_1 I_2 I_3 \psi(\mathbf{x}, t) \\
&= \left(\frac{2\pi i\hbar\epsilon}{m}\right)^{3/2} \psi(\mathbf{x}, t)
\end{aligned} \tag{7}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \boldsymbol{\eta} \cdot \nabla \psi(\mathbf{x}, t) d\eta_1 d\eta_2 d\eta_3 \\
&= J_1 I_2 I_3 \frac{\partial \psi}{\partial x_1} + I_1 J_2 I_3 \frac{\partial \psi}{\partial x_2} + I_1 I_2 J_3 \frac{\partial \psi}{\partial x_3} \\
&= 0
\end{aligned} \tag{8}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \frac{1}{2} \boldsymbol{\eta} \cdot \nabla (\boldsymbol{\eta} \cdot \nabla \psi) d\eta_1 d\eta_2 d\eta_3 \\
&= \frac{1}{2} K_1 I_2 I_3 \frac{\partial^2 \psi}{\partial x_1^2} + \frac{1}{2} I_1 K_2 I_3 \frac{\partial^2 \psi}{\partial x_2^2} + \frac{1}{2} I_1 I_2 K_3 \frac{\partial^2 \psi}{\partial x_3^2} \\
&+ J_1 J_2 I_1 \frac{\partial^2 \psi}{\partial x_1 x_2} + J_1 I_2 J_3 \frac{\partial^2 \psi}{\partial x_1 x_3} + I_1 J_2 J_3 \frac{\partial^2 \psi}{\partial x_2 x_3} \\
&= \frac{i\hbar\epsilon}{2m} \left(\frac{2\pi i\hbar\epsilon}{m}\right)^{3/2} \nabla^2 \psi
\end{aligned} \tag{9}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) T\psi(\mathbf{x}, t) d\eta_1 d\eta_2 d\eta_3 \\
&= J_1 I_2 I_3 \frac{-ie}{\hbar c} A_1 \psi + I_1 J_2 I_3 \frac{-ie}{\hbar c} A_2 \psi + I_1 I_2 J_3 \frac{-ie}{\hbar c} A_3 \psi \\
&= 0
\end{aligned} \tag{10}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) T(\boldsymbol{\eta} \cdot \nabla\psi) d\eta_1 d\eta_2 d\eta_3 \\
&= \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \\
&\times \frac{-ie}{\hbar c}(\eta_1 A_1 + \eta_2 A_2 + \eta_3 A_3) \left(\eta_1 \frac{\partial\psi}{\partial x_1} + \eta_2 \frac{\partial\psi}{\partial x_2} + \eta_3 \frac{\partial\psi}{\partial x_3} \right) d\eta_1 d\eta_2 d\eta_3 \\
&= K_1 I_2 I_3 \frac{-ieA_1}{\hbar c} \frac{\partial\psi}{\partial x_1} + J_1 J_2 I_3 \frac{-ieA_1}{\hbar c} \frac{\partial\psi}{\partial x_2} + J_1 I_2 J_3 \frac{-ieA_1}{\hbar c} \frac{\partial\psi}{\partial x_3} \\
&+ J_1 J_2 I_3 \frac{-ieA_2}{\hbar c} \frac{\partial\psi}{\partial x_1} + I_1 K_2 I_3 \frac{-ieA_2}{\hbar c} \frac{\partial\psi}{\partial x_2} + I_1 J_2 J_3 \frac{-ieA_2}{\hbar c} \frac{\partial\psi}{\partial x_3} \\
&+ J_1 I_2 J_3 \frac{-ieA_3}{\hbar c} \frac{\partial\psi}{\partial x_1} + I_1 J_2 J_3 \frac{-ieA_3}{\hbar c} \frac{\partial\psi}{\partial x_2} + I_1 I_2 K_3 \frac{-ieA_3}{\hbar c} \frac{\partial\psi}{\partial x_3} \\
&= \frac{i\hbar\epsilon}{m} \left(\frac{2\pi i\hbar\epsilon}{m} \right)^{3/2} \frac{-ie}{\hbar c} \left(A_1 \frac{\partial\psi}{\partial x_1} + A_2 \frac{\partial\psi}{\partial x_2} + A_3 \frac{\partial\psi}{\partial x_3} \right) \\
&= \frac{i\hbar\epsilon}{m} \left(\frac{2\pi i\hbar\epsilon}{m} \right)^{3/2} \frac{-ie}{\hbar c} \mathbf{A} \nabla\psi \tag{11}
\end{aligned}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) T\left(\frac{1}{2}\boldsymbol{\eta} \cdot (\boldsymbol{\eta} \cdot \nabla\psi)\right) d\eta_1 d\eta_2 d\eta_3 \\
&= \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \frac{-ie}{\hbar c}(\eta_1 A_1 + \eta_2 A_2 + \eta_3 A_3) \\
&\times \left(\frac{1}{2}\eta_1^2 \frac{\partial^2\psi}{\partial x_1^2} + \frac{1}{2}\eta_2^2 \frac{\partial^2\psi}{\partial x_2^2} + \frac{1}{2}\eta_3^2 \frac{\partial^2\psi}{\partial x_3^2} + \eta_1\eta_2 \frac{\partial^2\psi}{\partial x_1\partial x_2} + \eta_1\eta_3 \frac{\partial^2\psi}{\partial x_1\partial x_3} + \eta_2\eta_3 \frac{\partial^2\psi}{\partial x_2\partial x_3} \right) \\
&= 0 \tag{12}
\end{aligned}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \frac{1}{2}T^2\psi(\mathbf{x},t) d\eta_1 d\eta_2 d\eta_3 \\
&= \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \left(-\frac{e^2}{2\hbar^2c^2}\right) (\eta_1A_1 + \eta_2A_2 + \eta_3A_3)^2\psi(\mathbf{x},t) d\eta_1 d\eta_2 d\eta_3 \\
&= \frac{i\hbar\epsilon}{m} \left(\frac{2\pi i\hbar\epsilon}{m}\right)^{3/2} \left(-\frac{e^2}{2\hbar^2c^2}\right) (A_1^2 + A_2^2 + A_3^2)\psi(\mathbf{x},t) \\
&= \frac{i\hbar\epsilon}{m} \left(\frac{2\pi i\hbar\epsilon}{m}\right)^{3/2} \left(-\frac{e^2}{2\hbar^2c^2}\right) \mathbf{A}^2\psi
\end{aligned} \tag{13}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \frac{1}{2}T^2(\boldsymbol{\eta} \cdot \nabla\psi) d\eta_1 d\eta_2 d\eta_3 \\
&= \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \left(-\frac{e^2}{2\hbar^2c^2}\right) \\
&\quad \times (\eta_1A_1 + \eta_2A_2 + \eta_3A_3)^2 \left(\eta_1\frac{\partial\psi}{\partial x_1} + \eta_2\frac{\partial\psi}{\partial x_2} + \eta_3\frac{\partial\psi}{\partial x_3}\right) d\eta_1 d\eta_2 d\eta_3 \\
&= 0
\end{aligned} \tag{14}$$

$$\begin{aligned}
& \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \frac{1}{2}T^2\left(\frac{1}{2}\boldsymbol{\eta} \cdot (\boldsymbol{\eta} \cdot \nabla\psi)\right) d\eta_1 d\eta_2 d\eta_3 \\
&= \int_{\mathbb{R}^3} \exp\left(\frac{im}{2\hbar\epsilon}\boldsymbol{\eta}^2\right) \left(-\frac{e^2}{2\hbar^2c^2}\right) (\eta_1A_1 + \eta_2A_2 + \eta_3A_3)^2 \\
&\quad \times \left(\frac{1}{2}\eta_1^2\frac{\partial^2\psi}{\partial x_1^2} + \frac{1}{2}\eta_2^2\frac{\partial^2\psi}{\partial x_2^2} + \frac{1}{2}\eta_3^2\frac{\partial^2\psi}{\partial x_3^2} + \eta_1\eta_2\frac{\partial^2\psi}{\partial x_1x_2} + \eta_1\eta_3\frac{\partial^2\psi}{\partial x_1x_3} + \eta_2\eta_3\frac{\partial^2\psi}{\partial x_2x_3}\right) \\
&= \frac{e^2\epsilon^2}{2m^2c^2} \left(\frac{2\pi i\hbar\epsilon}{m}\right)^2 \\
&\quad \times \left(\frac{1}{2}\mathbf{A}^2\nabla^2\psi + \mathbf{A}\nabla\nabla\psi\mathbf{A} - \frac{3}{2}\left(A_1^2\frac{\partial^2}{\partial x_1^2} + A_2^2\frac{\partial^2}{\partial x_2^2} + A_3^2\frac{\partial^2}{\partial x_3^2}\right)\right)
\end{aligned} \tag{15}$$

Substitute the solved integrals into (6) to obtain

$$\psi(\mathbf{x},t) + \epsilon\frac{\partial\psi}{\partial t} = \frac{1}{A} \left(1 - \frac{ie\epsilon}{\hbar}\phi(\mathbf{x},t)\right) I$$

where

$$I = \left(\frac{2\pi i \hbar \epsilon}{m} \right)^{3/2} \psi(\mathbf{x}, t) \quad \text{from (7)}$$

$$+ \frac{i \hbar \epsilon}{2m} \left(\frac{2\pi i \hbar \epsilon}{m} \right)^{3/2} \nabla^2 \psi \quad \text{from (9)}$$

$$+ \frac{i \hbar \epsilon}{m} \left(\frac{2\pi i \hbar \epsilon}{m} \right)^{3/2} \frac{-ie}{\hbar c} \mathbf{A} \nabla \psi \quad \text{from (11)}$$

$$+ \frac{i \hbar \epsilon}{m} \left(\frac{2\pi i \hbar \epsilon}{m} \right)^{3/2} \left(-\frac{e^2}{2\hbar^2 c^2} \right) \mathbf{A}^2 \psi \quad \text{from (13)}$$

$$+ 0 \quad \text{from (15)}$$

The result from (15) is discarded because it is proportional to ϵ^2 .

Simplify I as follows.

$$I = \left(\frac{2\pi i \hbar \epsilon}{m} \right)^{3/2} \left(\psi(\mathbf{x}, t) + \frac{i \hbar \epsilon}{2m} \nabla^2 \psi + \frac{e \epsilon}{m c} \mathbf{A} \nabla \psi - \frac{i e^2 \epsilon}{2m \hbar c^2} \mathbf{A}^2 \psi \right) \quad (16)$$

In the limit as $\epsilon \rightarrow 0$ we have

$$\psi(\mathbf{x}, t) = \frac{1}{A} \left(\frac{2\pi i \hbar \epsilon}{m} \right)^{3/2} \psi(\mathbf{x}, t)$$

hence

$$A = \left(\frac{2\pi i \hbar \epsilon}{m} \right)^{3/2}$$

Cancel A with the coefficient in I to obtain

$$\begin{aligned} \psi(\mathbf{x}, t) + \epsilon \frac{\partial \psi}{\partial t} &= \left(1 - \frac{i e \epsilon}{\hbar} \phi(\mathbf{x}, t) \right) \\ &\times \left(\psi(\mathbf{x}, t) + \frac{i \hbar \epsilon}{2m} \nabla^2 \psi + \frac{e \epsilon}{m c} \mathbf{A} \nabla \psi - \frac{i e^2 \epsilon}{2m \hbar c^2} \mathbf{A}^2 \psi \right) \end{aligned}$$

Expand the product and discard terms of order greater than ϵ .

$$\begin{aligned} \psi(\mathbf{x}, t) + \epsilon \frac{\partial \psi}{\partial t} &= \psi(\mathbf{x}, t) + \frac{i \hbar \epsilon}{2m} \nabla^2 \psi + \frac{e \epsilon}{m c} \mathbf{A} \nabla \psi - \frac{i e^2 \epsilon}{2m \hbar c^2} \mathbf{A}^2 \psi - \frac{i e \epsilon}{\hbar} \phi(\mathbf{x}, t) \psi \end{aligned}$$

Cancel leading terms $\psi(\mathbf{x}, t)$ and divide through by ϵ .

$$\frac{\partial \psi}{\partial t} = \frac{i\hbar}{2m} \nabla^2 \psi + \frac{e}{mc} \mathbf{A} \nabla \psi - \frac{ie^2}{2m\hbar c^2} \mathbf{A}^2 \psi - \frac{ie}{\hbar} \phi(\mathbf{x}, t) \psi \quad (17)$$