

Let $|\Psi\rangle$ be a coherent state where \bar{n} is the expected number of photons.

$$|\Psi\rangle = \sum_{n=0}^{\infty} \sqrt{\frac{\bar{n}^n \exp(-\bar{n})}{n!}} \exp\left(-i\left(n + \frac{1}{2}\right)\omega t\right) |n\rangle$$

It can be shown that

$$\hat{a}|\Psi\rangle = \sqrt{\bar{n}} \exp(-i\omega t)|\Psi\rangle$$

It follows that

$$\langle\Psi|\hat{a}^\dagger = (\hat{a}|\Psi\rangle)^\dagger = \sqrt{\bar{n}} \exp(i\omega t)\langle\Psi|$$

Let \hat{E} be the electric field operator

$$\hat{E} = i\sqrt{\frac{\hbar\omega}{2\epsilon_0}}(\hat{a} - \hat{a}^\dagger)$$

The expected electric field is

$$\langle\hat{E}\rangle = \langle\Psi|\hat{E}|\Psi\rangle = i\sqrt{\frac{\hbar\omega}{2\epsilon_0}}\langle\Psi|(\hat{a} - \hat{a}^\dagger)|\Psi\rangle$$

By distributive law

$$\langle\hat{E}\rangle = i\sqrt{\frac{\hbar\omega}{2\epsilon_0}}(\langle\Psi|\hat{a}|\Psi\rangle - \langle\Psi|\hat{a}^\dagger|\Psi\rangle)$$

Substitute eigenvalues for operators.

$$\langle\hat{E}\rangle = i\sqrt{\frac{\hbar\omega}{2\epsilon_0}}(\sqrt{\bar{n}} \exp(-i\omega t)\langle\Psi|\Psi\rangle - \sqrt{\bar{n}} \exp(i\omega t)\langle\Psi|\Psi\rangle)$$

By $\langle\Psi|\Psi\rangle = 1$ we have

$$\langle\hat{E}\rangle = i\sqrt{\frac{\hbar\omega}{2\epsilon_0}}(\sqrt{\bar{n}} \exp(-i\omega t) - \sqrt{\bar{n}} \exp(i\omega t))$$

Recalling that

$$2\sin(\omega t) = i\exp(-i\omega t) - i\exp(i\omega t)$$

we have

$$\langle\hat{E}\rangle = \sqrt{\frac{2\bar{n}\hbar\omega}{\epsilon_0}}\sin(\omega t)$$

Let \hat{B} be the magnetic field operator

$$\hat{B} = \sqrt{\frac{\hbar\omega\mu_0}{2}}(\hat{a} + \hat{a}^\dagger)$$

Then by deduction similar to that for $\langle \hat{E} \rangle$ we obtain

$$\langle \hat{B} \rangle = \sqrt{2\bar{n}\hbar\omega\mu_0} \cos(\omega t)$$

The energy of an electromagnetic wave is

$$U = \frac{\epsilon_0}{2} |\mathbf{E}|^2 + \frac{1}{2\mu_0} |\mathbf{B}|^2$$

For linear polarization and a suitable rotation matrix R we have

$$R\mathbf{E} = \begin{pmatrix} E \\ 0 \\ 0 \end{pmatrix}, \quad R\mathbf{B} = \begin{pmatrix} 0 \\ B \\ 0 \end{pmatrix}$$

Hence in the rotated frame

$$U = \frac{\epsilon_0}{2} E^2 + \frac{1}{2\mu_0} B^2$$

For a quantum field we have

$$U = \frac{\epsilon_0}{2} \langle \hat{E}^2 \rangle + \frac{1}{2\mu_0} \langle \hat{B}^2 \rangle$$

where

$$\begin{aligned} \langle \hat{E}^2 \rangle &= \langle \Psi | \hat{E} \hat{E} | \Psi \rangle = -\frac{\hbar\omega}{2\epsilon_0} \langle \Psi | (\hat{a} - \hat{a}^\dagger)(\hat{a} - \hat{a}^\dagger) | \Psi \rangle \\ \langle \hat{B}^2 \rangle &= \langle \Psi | \hat{B} \hat{B} | \Psi \rangle = \frac{\hbar\omega\mu_0}{2} \langle \Psi | (\hat{a} + \hat{a}^\dagger)(\hat{a} + \hat{a}^\dagger) | \Psi \rangle \end{aligned}$$

For the coherent state

$$\begin{aligned} \langle \Psi | \hat{a} \hat{a} | \Psi \rangle &= (\sqrt{\bar{n}} \exp(-i\omega t))^2 &= \bar{n} \exp(-2i\omega t) \\ \langle \Psi | \hat{a} \hat{a}^\dagger | \Psi \rangle &= \langle \Psi | (\hat{a}^\dagger \hat{a} + 1) | \Psi \rangle &= \bar{n} + 1 \\ \langle \Psi | \hat{a}^\dagger \hat{a} | \Psi \rangle &= (\sqrt{\bar{n}} \exp(i\omega t)) (\sqrt{\bar{n}} \exp(-i\omega t)) &= \bar{n} \\ \langle \Psi | \hat{a}^\dagger \hat{a}^\dagger | \Psi \rangle &= (\sqrt{\bar{n}} \exp(i\omega t))^2 &= \bar{n} \exp(2i\omega t) \end{aligned}$$

The expectation $\bar{n} + 1$ for $\hat{a}\hat{a}^\dagger$ is from the commutator

$$\hat{a}\hat{a}^\dagger - \hat{a}^\dagger\hat{a} = 1$$

By substitution

$$\begin{aligned}\langle \hat{E}^2 \rangle &= -\frac{\hbar\omega}{2\epsilon_0} (\bar{n} \exp(-2i\omega t) + \bar{n} \exp(2i\omega t) - 2\bar{n} - 1) \\ \langle \hat{B}^2 \rangle &= \frac{\hbar\omega\mu_0}{2} (\bar{n} \exp(-2i\omega t) + \bar{n} \exp(2i\omega t) + 2\bar{n} + 1)\end{aligned}$$

Noting that

$$\begin{aligned}4 \sin(\omega t)^2 &= -\exp(-2i\omega t) - \exp(2i\omega t) + 2 \\ 4 \cos(\omega t)^2 &= \exp(-2i\omega t) + \exp(2i\omega t) + 2\end{aligned}$$

we have

$$\begin{aligned}\langle \hat{E}^2 \rangle &= \frac{2\bar{n}\hbar\omega}{\epsilon_0} \sin(\omega t)^2 + \frac{\hbar\omega}{2\epsilon_0} \\ \langle \hat{B}^2 \rangle &= 2\bar{n}\hbar\omega\mu_0 \cos(\omega t)^2 + \frac{\hbar\omega\mu_0}{2}\end{aligned}$$

Rewrite as

$$\begin{aligned}\frac{\epsilon_0}{2} \langle \hat{E}^2 \rangle &= \bar{n}\hbar\omega \sin(\omega t)^2 + \frac{\hbar\omega}{4} \\ \frac{1}{2\mu_0} \langle \hat{B}^2 \rangle &= \bar{n}\hbar\omega \cos(\omega t)^2 + \frac{\hbar\omega}{4}\end{aligned}$$

Hence the total energy per unit volume is

$$U = \frac{\epsilon_0}{2} \langle \hat{E}^2 \rangle + \frac{1}{2\mu_0} \langle \hat{B}^2 \rangle = \left(\bar{n} + \frac{1}{2}\right) \hbar\omega$$

Checking physical dimensions we have

$$\hbar\omega = h\nu \propto \text{joule second} \times \frac{1}{\text{second}} = \text{joule}$$

We will now show that

$$\hat{a}|\Psi\rangle = \sqrt{\bar{n}} \exp(-i\omega t)|\Psi\rangle$$

Apply operator \hat{a} to coherent state $|\Psi\rangle$ to obtain

$$\hat{a}|\Psi\rangle = \sum_{n=0}^{\infty} \sqrt{\frac{\bar{n}^n \exp(-\bar{n})}{n!}} \exp\left(-i\left(n + \frac{1}{2}\right)\omega t\right) \sqrt{n}|n-1\rangle$$

The $n = 0$ term vanishes hence the sum can start from $n = 1$.

$$\hat{a}|\Psi\rangle = \sum_{n=1}^{\infty} \sqrt{\frac{\bar{n}^n \exp(-\bar{n})}{n!}} \exp\left(-i\left(n + \frac{1}{2}\right)\omega t\right) \sqrt{n}|n-1\rangle$$

The \sqrt{n} cancels with n factorial.

$$\hat{a}|\Psi\rangle = \sum_{n=1}^{\infty} \sqrt{\frac{\bar{n}^n \exp(-\bar{n})}{(n-1)!}} \exp\left(-i\left(n + \frac{1}{2}\right)\omega t\right) |n-1\rangle$$

Factor out $\sqrt{\bar{n}} \exp(-i\omega t)$.

$$\hat{a}|\Psi\rangle = \sqrt{\bar{n}} \exp(-i\omega t) \sum_{n=1}^{\infty} \sqrt{\frac{\bar{n}^{n-1} \exp(-\bar{n})}{(n-1)!}} \exp\left(-i\left(n - \frac{1}{2}\right)\omega t\right) |n-1\rangle$$

Substitute $n + 1$ for index n .

$$\hat{a}|\Psi\rangle = \sqrt{\bar{n}} \exp(-i\omega t) \sum_{n=0}^{\infty} \sqrt{\frac{\bar{n}^n \exp(-\bar{n})}{n!}} \exp\left(-i\left(n + \frac{1}{2}\right)\omega t\right) |n\rangle$$

Hence

$$\hat{a}|\Psi\rangle = \sqrt{\bar{n}} \exp(-i\omega t) |\Psi\rangle$$