

Stern-Gerlach 2

From the previous section we have the following Schrödinger equations for a Stern-Gerlach experiment.

$$\begin{aligned} -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} \psi_1 + \mu_B(B_0 + \alpha z)\psi_1 &= i\hbar \frac{\partial}{\partial t} \psi_1 \\ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} \psi_2 - \mu_B(B_0 + \alpha z)\psi_2 &= i\hbar \frac{\partial}{\partial t} \psi_2 \end{aligned}$$

We now seek solutions for ψ_1 and ψ_2 .

An online paper¹ provides the following solutions.

$$\begin{aligned} \psi_1 &= \text{Ai}\left[\left(\frac{\mu_B \alpha m}{\hbar^2}\right)^{1/3} \left(\frac{\mu_B \alpha}{4m}\right) t^2 + \left(\frac{\mu_B \alpha m}{\hbar^2}\right)^{1/3} z\right] \exp\left(-\frac{i\alpha \mu_B z t}{2\hbar}\right) \exp\left(-\frac{\mu_B B_0 t}{\hbar}\right) \\ \psi_2 &= \text{Ai}\left[\left(\frac{\mu_B \alpha m}{\hbar^2}\right)^{1/3} \left(\frac{\mu_B \alpha}{4m}\right) t^2 - \left(\frac{\mu_B \alpha m}{\hbar^2}\right)^{1/3} z\right] \exp\left(+\frac{i\alpha \mu_B z t}{2\hbar}\right) \exp\left(+\frac{\mu_B B_0 t}{\hbar}\right) \end{aligned}$$

Let us try verifying ψ_1 and ψ_2 with a Taylor series expansion of the Airy function $\text{Ai}(x)$.

We have

$$\text{Ai}(x) \approx \sum_{k=0}^n a_k x^k$$

where

$$\begin{aligned} a_0 &= \frac{1}{3^{2/3} \Gamma(2/3)} \\ a_1 &= -\frac{1}{3^{1/3} \Gamma(1/3)} \\ a_2 &= 0 \\ a_{n+3} &= \frac{a_n}{(n+3)(n+2)} \end{aligned}$$

For $n = 9$ we obtain

$$\begin{aligned} \text{Ai}(x) &\approx 2.73941 \times 10^{-5} x^9 - 0.000513531 x^7 + 0.00197238 x^6 \\ &\quad - 0.0215683 x^4 + 0.0591713 x^3 - 0.258819 x + 0.355028 \end{aligned}$$

Calculate ψ_1 and ψ_2 using the $\text{Ai}(x)$ approximation.

To verify ψ_1 and ψ_2 , calculate departures from equality of the Schrödinger equations.

$$\begin{aligned} \epsilon_1 &= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} \psi_1 + \mu_B(B_0 + \alpha z)\psi_1 - i\hbar \frac{\partial}{\partial t} \psi_1 \\ \epsilon_2 &= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} \psi_2 - \mu_B(B_0 + \alpha z)\psi_2 - i\hbar \frac{\partial}{\partial t} \psi_2 \end{aligned}$$

¹ “Construction of Exact Solutions for the Stern-Gerlach Effect” by Bulnes and Oliveira.

After calculating ϵ_1 and ϵ_2 , cancel exponentials and zero out factors t^n and z^m where $n \geq 8$ and $m \geq 5$.

The results are

$$\begin{aligned}\epsilon_1 = & \boxed{-6.50521 \times 10^{-19}} \frac{\alpha^{16/3} \mu_B^{16/3} t^6 z^2}{\hbar^{8/3} m^{5/3}} + \boxed{5.42101 \times 10^{-20}} \frac{\alpha^5 \mu_B^5 t^6 z}{\hbar^2 m^2} \\ & - \boxed{1.73472 \times 10^{-18}} \frac{\alpha^{13/3} \mu_B^{13/3} t^4 z^3}{\hbar^{8/3} m^{2/3}} - \boxed{3.46945 \times 10^{-18}} \frac{\alpha^{10/3} m^{1/3} \mu_B^{10/3} t^2 z^4}{\hbar^{8/3}} \\ & - \boxed{2.1684 \times 10^{-19}} \frac{i \alpha^{16/3} \mu_B^{16/3} t^5 z^4}{\hbar^{11/3} m^{2/3}}\end{aligned}$$

and

$$\begin{aligned}\epsilon_2 = & \boxed{-6.50521 \times 10^{-19}} \frac{\alpha^{16/3} \mu_B^{16/3} t^6 z^2}{\hbar^{8/3} m^{5/3}} - \boxed{5.42101 \times 10^{-20}} \frac{\alpha^5 \mu_B^5 t^6 z}{\hbar^2 m^2} \\ & + \boxed{1.73472 \times 10^{-18}} \frac{\alpha^{13/3} \mu_B^{13/3} t^4 z^3}{\hbar^{8/3} m^{2/3}} - \boxed{3.46945 \times 10^{-18}} \frac{\alpha^{10/3} m^{1/3} \mu_B^{10/3} t^2 z^4}{\hbar^{8/3}} \\ & - \boxed{2.1684 \times 10^{-19}} \frac{i \alpha^{16/3} \mu_B^{16/3} t^5 z^4}{\hbar^{11/3} m^{2/3}}\end{aligned}$$

The numerical values are negligible hence ψ_1 and ψ_2 are confirmed as solutions.

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