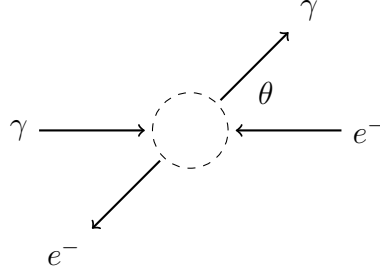


Compton scattering

Compton scattering is the interaction $e^- + \gamma \rightarrow e^- + \gamma$.



The following momentum vectors are for the center-of-mass frame with $E = \sqrt{\omega^2 + m^2}$.

$$p_1 = \begin{pmatrix} \omega \\ 0 \\ 0 \\ \omega \end{pmatrix} \quad p_2 = \begin{pmatrix} E \\ 0 \\ 0 \\ -\omega \end{pmatrix} \quad p_3 = \begin{pmatrix} \omega \\ \omega \sin \theta \cos \phi \\ \omega \sin \theta \sin \phi \\ \omega \cos \theta \end{pmatrix} \quad p_4 = \begin{pmatrix} E \\ -\omega \sin \theta \cos \phi \\ -\omega \sin \theta \sin \phi \\ -\omega \cos \theta \end{pmatrix}$$

$\gamma \rightarrow$ $\leftarrow e^-$ $\nearrow \gamma$ $\nwarrow e^-$

Spinors for p_2 .

$$u_{21} = \frac{1}{\sqrt{E+m}} \begin{pmatrix} E+m \\ 0 \\ -\omega \\ 0 \end{pmatrix} \quad u_{22} = \frac{1}{\sqrt{E+m}} \begin{pmatrix} 0 \\ E+m \\ 0 \\ \omega \end{pmatrix}$$

spin up spin down

Spinors for p_4 .

$$u_{41} = \frac{1}{\sqrt{E+m}} \begin{pmatrix} E+m \\ 0 \\ p_{4z} \\ p_{4x} + ip_{4y} \end{pmatrix} \quad u_{42} = \frac{1}{\sqrt{E+m}} \begin{pmatrix} 0 \\ E+m \\ p_{4x} - ip_{4y} \\ -p_{4z} \end{pmatrix}$$

spin up spin down

The scattering amplitude $\mathcal{M}_{ab}{}^{\mu\nu}$ for spin ab and polarization $\mu\nu$ is

$$\mathcal{M}_{ba}{}^{\mu\nu} = \mathcal{M}_{1ab}{}^{\mu\nu} + \mathcal{M}_{2ab}{}^{\nu\mu}$$

where

$$\mathcal{M}_{1ab}{}^{\mu\nu} = \frac{\bar{u}_{4b}(-ie\gamma^\mu)(\not{p}_1 + m)(-ie\gamma^\nu)u_{2a}}{s - m^2}$$

$$\mathcal{M}_{2ab}{}^{\nu\mu} = \frac{\bar{u}_{4b}(-ie\gamma^\nu)(\not{p}_2 + m)(-ie\gamma^\mu)u_{2a}}{u - m^2}$$

Matrices \not{q}_1 and \not{q}_2 represent momentum transfer.

$$\begin{aligned}\not{q}_1 &= (p_1 + p_2)^\alpha g_{\alpha\beta} \gamma^\beta \\ \not{q}_2 &= (p_4 - p_1)^\alpha g_{\alpha\beta} \gamma^\beta\end{aligned}$$

Scalars s and u are Mandelstam variables.

$$\begin{aligned}s &= (p_1 + p_2)^2 \\ u &= (p_1 - p_4)^2\end{aligned}$$

In component form (note that indices μ and ν are interchanged for \mathcal{M}_{2ab})

$$\begin{aligned}\mathcal{M}_{1ab}{}^{\mu\nu} &= \frac{(\bar{u}_{4b})_\alpha (-ie\gamma^{\mu\alpha}{}_\beta)(\not{q}_1 + m)^\beta{}_\rho (-ie\gamma^{\nu\rho}{}_\sigma)(u_{2a})^\sigma}{s - m^2} \\ \mathcal{M}_{2ab}{}^{\nu\nu} &= \frac{(\bar{u}_{4b})_\alpha (-ie\gamma^{\nu\alpha}{}_\beta)(\not{q}_2 + m)^\beta{}_\rho (-ie\gamma^{\mu\rho}{}_\sigma)(u_{2a})^\sigma}{u - m^2}\end{aligned}$$

Expected probability density $\langle |\mathcal{M}|^2 \rangle$ is the sum of squared amplitudes divided by the number of inbound states.

$$\langle |\mathcal{M}|^2 \rangle = \frac{1}{4} \sum_{ab} \sum_{\mu\nu} |\mathcal{M}_{ab}{}^{\mu\nu}|^2$$

Summing over $\mu\nu$ requires $g_{\mu\nu}$ to lower indices.

$$\langle |\mathcal{M}|^2 \rangle = \frac{1}{4} \sum_{ab} \mathcal{M}_{ab}{}^{\mu\nu} (g_{\mu\alpha} \mathcal{M}_{ab}{}^{\alpha\beta} g_{\beta\nu})^*$$

Expand the summand and label the terms. By hermiticity $\boxed{2} = \boxed{3}$.

$$\begin{aligned}\langle |\mathcal{M}|^2 \rangle &= \frac{1}{4} \sum_{ab} \left[\underbrace{\mathcal{M}_{1ab}{}^{\mu\nu} (g_{\mu\alpha} \mathcal{M}_{1ab}{}^{\alpha\beta} g_{\beta\nu})^*}_{\boxed{1}} + \underbrace{\mathcal{M}_{1ab}{}^{\mu\nu} (g_{\nu\alpha} \mathcal{M}_{2ab}{}^{\alpha\beta} g_{\beta\mu})^*}_{\boxed{2}} \right. \\ &\quad \left. + \underbrace{\mathcal{M}_{2ab}{}^{\nu\mu} (g_{\mu\alpha} \mathcal{M}_{1ab}{}^{\alpha\beta} g_{\beta\nu})^*}_{\boxed{3}} + \underbrace{\mathcal{M}_{2ab}{}^{\nu\mu} (g_{\nu\alpha} \mathcal{M}_{2ab}{}^{\alpha\beta} g_{\beta\mu})^*}_{\boxed{4}} \right]\end{aligned}$$

The Casimir trick uses matrix arithmetic to sum over spin and polarization states.

$$\begin{aligned}\sum_{ab} \boxed{1} &= \frac{e^4}{(s - m^2)^2} \text{Tr} \left[(\not{p}_2 + m) \gamma^\mu (\not{q}_1 + m) \gamma^\nu (\not{p}_4 + m) \gamma_\nu (\not{q}_1 + m) \gamma_\mu \right] \\ \sum_{ab} \boxed{2} &= \frac{e^4}{(s - m^2)(u - m^2)} \text{Tr} \left[(\not{p}_2 + m) \gamma^\mu (\not{q}_2 + m) \gamma^\nu (\not{p}_4 + m) \gamma_\mu (\not{q}_1 + m) \gamma_\nu \right] \\ \sum_{ab} \boxed{4} &= \frac{e^4}{(u - m^2)^2} \text{Tr} \left[(\not{p}_2 + m) \gamma^\mu (\not{q}_2 + m) \gamma^\nu (\not{p}_4 + m) \gamma_\nu (\not{q}_2 + m) \gamma_\mu \right]\end{aligned}$$

Let

$$\begin{aligned} f_{11} &= \text{Tr} \left[(\not{p}_2 + m) \gamma^\mu (\not{q}_1 + m) \gamma^\nu (\not{p}_4 + m) \gamma_\nu (\not{q}_1 + m) \gamma_\mu \right] \\ f_{12} &= \text{Tr} \left[(\not{p}_2 + m) \gamma^\mu (\not{q}_2 + m) \gamma^\nu (\not{p}_4 + m) \gamma_\mu (\not{q}_1 + m) \gamma_\nu \right] \\ f_{22} &= \text{Tr} \left[(\not{p}_2 + m) \gamma^\mu (\not{q}_2 + m) \gamma^\nu (\not{p}_4 + m) \gamma_\nu (\not{q}_2 + m) \gamma_\mu \right] \end{aligned}$$

so that

$$\langle |\mathcal{M}|^2 \rangle = \frac{e^4}{4} \left[\frac{f_{11}}{(s - m^2)^2} + \frac{2f_{12}}{(s - m^2)(u - m^2)} + \frac{f_{22}}{(u - m^2)^2} \right]$$

The following formulas are equivalent to the Casimir trick. (Recall that $a \cdot b = a^\mu g_{\mu\nu} b^\nu$)

$$\begin{aligned} f_{11} &= 32(p_1 \cdot p_2)(p_1 \cdot p_4) + 32(p_1 \cdot p_2)m^2 + 32m^4 \\ f_{12} &= 16(p_1 \cdot p_2)m^2 - 16(p_1 \cdot p_4)m^2 + 32m^4 \\ f_{22} &= 32(p_1 \cdot p_2)(p_1 \cdot p_4) - 32(p_1 \cdot p_4)m^2 + 32m^4 \end{aligned}$$

In Mandelstam variables

$$\begin{aligned} f_{11} &= -8su + 24sm^2 + 8um^2 + 8m^4 \\ f_{12} &= 8sm^2 + 8um^2 + 16m^4 \\ f_{22} &= -8su + 8sm^2 + 24um^2 + 8m^4 \end{aligned} \tag{2}$$

Compton scattering experiments are typically done in the lab frame where the electron is at rest. Define Lorentz boost Λ for transforming momentum vectors to the lab frame.

$$\Lambda = \begin{pmatrix} E/m & 0 & 0 & \omega/m \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \omega/m & 0 & 0 & E/m \end{pmatrix}$$

The electron is at rest in the lab frame.

$$\Lambda p_2 = \begin{pmatrix} m \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

Mandelstam variables are invariant under a boost.

$$\begin{aligned} s &= (p_1 + p_2)^2 = (\Lambda p_1 + \Lambda p_2)^2 \\ t &= (p_1 - p_3)^2 = (\Lambda p_1 - \Lambda p_3)^2 \\ u &= (p_1 - p_4)^2 = (\Lambda p_1 - \Lambda p_4)^2 \end{aligned}$$

In the lab frame, let ω_L be the angular frequency of the incident photon and let ω'_L be the angular frequency of the scattered photon.

$$\omega_L = \Lambda p_1 \cdot \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = \frac{\omega^2}{m} + \frac{\omega E}{m}$$

$$\omega'_L = \Lambda p_3 \cdot \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = \frac{\omega^2 \cos \theta}{m} + \frac{\omega E}{m}$$

It can be shown that

$$\begin{aligned} s &= m^2 + 2m\omega_L \\ t &= 2m(\omega'_L - \omega_L) \\ u &= m^2 - 2m\omega'_L \end{aligned} \tag{3}$$

Then by (1), (2), and (3) we have

$$\langle |\mathcal{M}|^2 \rangle = 2e^4 \left[\frac{\omega_L}{\omega'_L} + \frac{\omega'_L}{\omega_L} + \left(\frac{m}{\omega_L} - \frac{m}{\omega'_L} + 1 \right)^2 - 1 \right]$$

Lab scattering angle θ_L is given by the Compton equation

$$\cos \theta_L = \frac{m}{\omega_L} - \frac{m}{\omega'_L} + 1$$

Hence

$$\begin{aligned} \langle |\mathcal{M}|^2 \rangle &= 2e^4 \left(\frac{\omega_L}{\omega'_L} + \frac{\omega'_L}{\omega_L} + \cos^2 \theta_L - 1 \right) \\ &= 2e^4 \left(\frac{\omega_L}{\omega'_L} + \frac{\omega'_L}{\omega_L} - \sin^2 \theta_L \right) \end{aligned}$$

Cross section

Now that we have derived $\langle |\mathcal{M}|^2 \rangle$ we can investigate the angular distribution of scattered photons. For simplicity let us drop the L subscript from lab variables. From now on the symbols ω , ω' , and θ will be lab frame variables.

The differential cross section is

$$\frac{d\sigma}{d\Omega} = \frac{1}{4(4\pi\epsilon_0)^2 s} \left(\frac{\omega'}{\omega} \right)^2 \langle |\mathcal{M}|^2 \rangle$$

where

$$s = m^2 + 2m\omega = (mc^2)^2 + 2(mc^2)(\hbar\omega)$$

and ω' is given by the Compton equation

$$\omega' = \frac{\omega}{1 + \frac{\hbar\omega}{mc^2}(1 - \cos\theta)}$$

For the lab frame we have

$$\langle |\mathcal{M}|^2 \rangle = 2e^4 \left(\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - \sin^2\theta \right)$$

Hence in the lab frame

$$\frac{d\sigma}{d\Omega} = \frac{e^4}{2(4\pi\epsilon_0)^2 s} \left(\frac{\omega'}{\omega} \right)^2 \left(\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - \sin^2\theta \right)$$

Noting that

$$e^2 = 4\pi\epsilon_0\alpha\hbar c$$

we have

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2(\hbar c)^2}{2s} \left(\frac{\omega'}{\omega} \right)^2 \left(\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - \sin^2\theta \right)$$

Noting that

$$d\Omega = \sin\theta d\theta d\phi$$

we also have

$$d\sigma = \frac{\alpha^2(\hbar c)^2}{2s} \left(\frac{\omega'}{\omega} \right)^2 \left(\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - \sin^2\theta \right) \sin\theta d\theta d\phi$$

Let $S(\theta_1, \theta_2)$ be the following integral of $d\sigma$.

$$S(\theta_1, \theta_2) = \int_0^{2\pi} \int_{\theta_1}^{\theta_2} d\sigma$$

The solution is

$$S(\theta_1, \theta_2) = \frac{\pi\alpha^2(\hbar c)^2}{s} [I(\theta_2) - I(\theta_1)]$$

where

$$\begin{aligned} I(\theta) = & -\frac{\cos\theta}{R^2} + \log(1 + R(1 - \cos\theta)) \left(\frac{1}{R} - \frac{2}{R^2} - \frac{2}{R^3} \right) \\ & - \frac{1}{2R(1 + R(1 - \cos\theta))^2} + \frac{1}{1 + R(1 - \cos\theta)} \left(-\frac{2}{R^2} - \frac{1}{R^3} \right) \end{aligned}$$

and

$$R = \frac{\hbar\omega}{mc^2}$$

The cumulative distribution function is

$$F(\theta) = \frac{S(0, \theta)}{S(0, \pi)} = \frac{I(\theta) - I(0)}{I(\pi) - I(0)}, \quad 0 \leq \theta \leq \pi$$

The probability of observing scattering events in the interval θ_1 to θ_2 is

$$P(\theta_1 < \theta \leq \theta_2) = F(\theta_2) - F(\theta_1)$$

The probability density function is

$$f(\theta) = \frac{dF(\theta)}{d\theta} = \frac{1}{I(\pi) - I(0)} \left(\frac{\omega'}{\omega} \right)^2 \left(\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - \sin^2 \theta \right) \sin \theta$$

Thomson scattering

For $\hbar\omega \ll mc^2$ we have

$$\omega' = \frac{\omega}{1 + \frac{\hbar\omega}{mc^2} (1 - \cos \theta)} \approx \omega$$

Hence we can use the approximations

$$\omega = \omega' \quad \text{and} \quad s = (mc^2)^2$$

to obtain

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 \hbar^2}{2m^2 c^2} (1 + \cos^2 \theta)$$

which is the formula for Thomson scattering.

High energy approximation

For $\omega \gg m$ a useful approximation is to set $m = 0$ and obtain

$$\begin{aligned} f_{11} &= -8su \\ f_{12} &= 0 \\ f_{22} &= -8su \end{aligned}$$

Hence

$$\begin{aligned} \langle |\mathcal{M}|^2 \rangle &= \frac{e^4}{4} \left(\frac{-8su}{s^2} + \frac{-8su}{u^2} \right) \\ &= 2e^4 \left(-\frac{u}{s} - \frac{s}{u} \right) \end{aligned}$$

The Mandelstam variables for $m = 0$ are

$$\begin{aligned} s &= 4\omega^2 \\ u &= -2\omega^2(\cos \theta + 1) \end{aligned}$$

Hence

$$\langle |\mathcal{M}|^2 \rangle = 2e^4 \left(\frac{\cos \theta + 1}{2} + \frac{2}{\cos \theta + 1} \right)$$

In the center of mass frame

$$\frac{d\sigma}{d\Omega} = \frac{\langle |\mathcal{M}|^2 \rangle}{4s(4\pi\epsilon_o)^2} = \frac{e^4}{2s(4\pi\epsilon_o)^2} \left(\frac{\cos \theta + 1}{2} + \frac{2}{\cos \theta + 1} \right)$$

Substitute $e^4 = (4\pi\epsilon_0\alpha\hbar c)^2$ to obtain

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{2s} \left(\frac{\cos\theta + 1}{2} + \frac{2}{\cos\theta + 1} \right) \times (\hbar c)^2$$

It follows that

$$\frac{d\sigma}{d\cos\theta} = 2\pi \frac{d\sigma}{d\Omega} = \frac{\pi\alpha^2}{s} \left(\frac{\cos\theta + 1}{2} + \frac{2}{\cos\theta + 1} \right) \times (\hbar c)^2$$

Cf. equation (1) of arxiv.org/pdf/hep-ex/0504012.