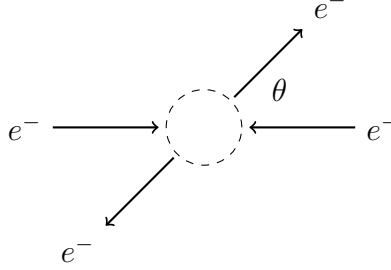


# Moller scattering

Moller scattering is the interaction  $e^- + e^- \rightarrow e^- + e^-$ .



Define the following momentum vectors and spinors. Symbol  $p$  is incident momentum. Symbol  $E$  is total energy  $E = \sqrt{p^2 + m^2}$  where  $m$  is electron mass. Polar angle  $\theta$  is the observed scattering angle. Azimuth angle  $\phi$  cancels out in scattering calculations.

$$\begin{array}{cccc}
 p_1 = \begin{pmatrix} E \\ 0 \\ 0 \\ p \end{pmatrix} & p_2 = \begin{pmatrix} E \\ 0 \\ 0 \\ -p \end{pmatrix} & p_3 = \begin{pmatrix} E \\ p \sin \theta \cos \phi \\ p \sin \theta \sin \phi \\ p \cos \theta \end{pmatrix} & p_4 = \begin{pmatrix} E \\ -p \sin \theta \cos \phi \\ -p \sin \theta \sin \phi \\ -p \cos \theta \end{pmatrix} \\
 \text{inbound} & \text{inbound} & \text{outbound} & \text{outbound} \\
 \text{electron} & \text{electron} & \text{electron} & \text{electron} \\
 \\
 u_{11} = \begin{pmatrix} E + m \\ 0 \\ p \\ 0 \end{pmatrix} & u_{21} = \begin{pmatrix} E + m \\ 0 \\ -p \\ 0 \end{pmatrix} & u_{31} = \begin{pmatrix} E + m \\ 0 \\ p_{3z} \\ p_{3x} + ip_{3y} \end{pmatrix} & u_{41} = \begin{pmatrix} E + m \\ 0 \\ p_{4z} \\ p_{4x} + ip_{4y} \end{pmatrix} \\
 \text{spin up} & \text{spin up} & \text{spin up} & \text{spin up} \\
 \\
 u_{12} = \begin{pmatrix} 0 \\ E + m \\ 0 \\ -p \end{pmatrix} & u_{22} = \begin{pmatrix} 0 \\ E + m \\ 0 \\ p \end{pmatrix} & u_{32} = \begin{pmatrix} 0 \\ E + m \\ p_{3x} - ip_{3y} \\ -p_{3z} \end{pmatrix} & u_{42} = \begin{pmatrix} 0 \\ E + m \\ p_{4x} - ip_{4y} \\ -p_{4z} \end{pmatrix} \\
 \text{spin down} & \text{spin down} & \text{spin down} & \text{spin down}
 \end{array}$$

The spinors are not individually normalized. Instead, a combined spinor normalization constant  $N = (E + m)^4$  will be used.

This is the probability density for spin state  $abcd$ . The formula is derived from Feynman diagrams for Moller scattering.

$$|\mathcal{M}_{abcd}|^2 = \frac{e^4}{N} \left| \frac{1}{t} (\bar{u}_{3c} \gamma^\mu u_{1a}) (\bar{u}_{4d} \gamma_\mu u_{2b}) - \frac{1}{u} (\bar{u}_{4d} \gamma^\nu u_{1a}) (\bar{u}_{3c} \gamma_\nu u_{2b}) \right|^2$$

no electron interchange                      electron interchange

Symbol  $e$  is electron charge and

$$\begin{aligned}
 t &= (p_1 - p_3)^2 = (p_1 - p_3)^\mu g_{\mu\nu} (p_1 - p_3)^\nu \\
 u &= (p_1 - p_4)^2 = (p_1 - p_4)^\mu g_{\mu\nu} (p_1 - p_4)^\nu
 \end{aligned}$$

Let

$$a_1 = (\bar{u}_{3c}\gamma^\mu u_{1a})(\bar{u}_{4d}\gamma_\mu u_{2b}), \quad a_2 = (\bar{u}_{4d}\gamma^\nu u_{1a})(\bar{u}_{3c}\gamma_\nu u_{2b})$$

Then

$$\begin{aligned} |\mathcal{M}_{abcd}|^2 &= \frac{e^4}{N} \left| \frac{a_1}{t} - \frac{a_2}{u} \right|^2 \\ &= \frac{e^4}{N} \left( \frac{a_1}{t} - \frac{a_2}{u} \right) \left( \frac{a_1}{t} - \frac{a_2}{u} \right)^* \\ &= \frac{e^4}{N} \left( \frac{a_1 a_1^*}{t^2} - \frac{a_1 a_2^*}{tu} - \frac{a_1^* a_2}{tu} + \frac{a_2 a_2^*}{u^2} \right) \end{aligned}$$

The expected probability density  $\langle |\mathcal{M}|^2 \rangle$  is computed by summing  $|\mathcal{M}_{abcd}|^2$  over all spin states and dividing by the number of inbound states. There are four inbound states.

$$\begin{aligned} \langle |\mathcal{M}|^2 \rangle &= \frac{1}{4} \sum_{a=1}^2 \sum_{b=1}^2 \sum_{c=1}^2 \sum_{d=1}^2 |\mathcal{M}_{abcd}|^2 \\ &= \frac{e^4}{4N} \sum_{a=1}^2 \sum_{b=1}^2 \sum_{c=1}^2 \sum_{d=1}^2 \left( \frac{a_1 a_1^*}{t^2} - \frac{a_1 a_2^*}{tu} - \frac{a_1^* a_2}{tu} + \frac{a_2 a_2^*}{u^2} \right) \end{aligned}$$

The Casimir trick uses matrix arithmetic to compute sums.

$$\begin{aligned} f_{11} &= \frac{1}{N} \sum_{abcd} a_1 a_1^* = \text{Tr} \left( (\not{p}_3 + m) \gamma^\mu (\not{p}_1 + m) \gamma^\nu \right) \text{Tr} \left( (\not{p}_4 + m) \gamma_\mu (\not{p}_2 + m) \gamma_\nu \right) \\ f_{12} &= \frac{1}{N} \sum_{abcd} a_1 a_2^* = \text{Tr} \left( (\not{p}_3 + m) \gamma^\mu (\not{p}_1 + m) \gamma^\nu (\not{p}_4 + m) \gamma_\mu (\not{p}_2 + m) \gamma_\nu \right) \\ f_{22} &= \frac{1}{N} \sum_{abcd} a_2 a_2^* = \text{Tr} \left( (\not{p}_4 + m) \gamma^\mu (\not{p}_1 + m) \gamma^\nu \right) \text{Tr} \left( (\not{p}_3 + m) \gamma_\mu (\not{p}_2 + m) \gamma_\nu \right) \end{aligned}$$

Hence

$$\langle |\mathcal{M}|^2 \rangle = \frac{e^4}{4} \left( \frac{f_{11}}{t^2} - \frac{f_{12}}{tu} - \frac{f_{12}^*}{tu} + \frac{f_{22}}{u^2} \right)$$

The following formulas are equivalent to the Casimir trick. (Recall that  $a \cdot b = a^\mu g_{\mu\nu} b^\nu$ )

$$\begin{aligned} f_{11} &= 32(p_1 \cdot p_2)^2 + 32(p_1 \cdot p_4)^2 - 64m^2(p_1 \cdot p_3) + 64m^4 \\ f_{12} &= -32(p_1 \cdot p_2)^2 + 32m^2(p_1 \cdot p_2) + 32m^2(p_1 \cdot p_3) + 32m^2(p_1 \cdot p_4) - 32m^4 \\ f_{22} &= 32(p_1 \cdot p_2)^2 + 32(p_1 \cdot p_3)^2 - 64m^2(p_1 \cdot p_4) + 64m^4 \end{aligned}$$

For Mandelstam variables

$$\begin{aligned} s &= (p_1 + p_2)^2 = 4E^2 \\ t &= (p_1 - p_3)^2 \\ u &= (p_1 - p_4)^2 \end{aligned}$$

the formulas are

$$\begin{aligned} f_{11} &= 8s^2 + 8u^2 - 64sm^2 - 64um^2 + 192m^4 \\ f_{12} &= -8s^2 + 64sm^2 - 96m^4 \\ f_{22} &= 8s^2 + 8t^2 - 64sm^2 - 64tm^2 + 192m^4 \end{aligned}$$

## High energy approximation

For high energy experiments  $E \gg m$  a useful approximation is to set  $m = 0$  and obtain

$$\begin{aligned} f_{11} &= 8s^2 + 8u^2 \\ f_{12} &= -8s^2 \\ f_{22} &= 8s^2 + 8t^2 \end{aligned}$$

Hence

$$\begin{aligned} \langle |\mathcal{M}|^2 \rangle &= \frac{e^4}{4} \left( \frac{f_{11}}{t^2} - \frac{f_{12}}{tu} - \frac{f_{12}^*}{tu} + \frac{f_{22}}{u^2} \right) \\ &= \frac{e^4}{4} \left( \frac{8s^2 + 8u^2}{t^2} - \frac{-8s^2}{tu} - \frac{-8s^2}{tu} + \frac{8s^2 + 8t^2}{u^2} \right) \\ &= 2e^4 \left( \frac{s^2 + u^2}{t^2} + \frac{2s^2}{tu} + \frac{s^2 + t^2}{u^2} \right) \end{aligned}$$

Combine terms so  $\langle |\mathcal{M}|^2 \rangle$  has a common denominator.

$$\langle |\mathcal{M}|^2 \rangle = 2e^4 \left( \frac{u^2(s^2 + u^2) + 2s^2tu + t^2(s^2 + t^2)}{t^2u^2} \right)$$

For  $m = 0$  the Mandelstam variables are

$$\begin{aligned} s &= 4E^2 \\ t &= 2E^2(\cos \theta - 1) \\ u &= -2E^2(\cos \theta + 1) \end{aligned}$$

Hence

$$\begin{aligned} \langle |\mathcal{M}|^2 \rangle &= 2e^4 \left( \frac{32E^8 \cos^4 \theta + 192E^8 \cos^2 \theta + 288E^8}{16E^8(\cos \theta - 1)^2(\cos \theta + 1)^2} \right) \\ &= 4e^4 \frac{(\cos^2 \theta + 3)^2}{(\cos \theta - 1)^2(\cos \theta + 1)^2} \\ &= 4e^4 \frac{(\cos^2 \theta + 3)^2}{\sin^4 \theta} \end{aligned}$$

The following equivalent formula can also be used.

$$\begin{aligned} \langle |\mathcal{M}|^2 \rangle &= 2e^4 \left( \frac{s^2 + u^2}{t^2} + \frac{2s^2}{tu} + \frac{s^2 + t^2}{u^2} \right) \\ &= 2e^4 \left( \underbrace{\frac{1 + \cos^4(\theta/2)}{\sin^4(\theta/2)}}_{\text{no electron interchange}} + \underbrace{\frac{2}{\sin^2(\theta/2) \cos^2(\theta/2)}}_{\text{interaction term}} + \underbrace{\frac{1 + \sin^4(\theta/2)}{\cos^4(\theta/2)}}_{\text{electron interchange}} \right) \end{aligned}$$

## Cross section

The differential cross section is

$$\frac{d\sigma}{d\Omega} = \frac{\langle |\mathcal{M}|^2 \rangle}{4(4\pi\epsilon_0)^2 s}, \quad s = (p_1 + p_2)^2 = 4E^2$$

For high energy experiments we have

$$\langle |\mathcal{M}|^2 \rangle = 4e^4 \frac{(\cos^2 \theta + 3)^2}{\sin^4 \theta}$$

Substitute for  $\langle |\mathcal{M}|^2 \rangle$ .

$$\frac{d\sigma}{d\Omega} = \frac{e^4}{(4\pi\epsilon_0)^2 s} \frac{(\cos^2 \theta + 3)^2}{\sin^4 \theta}$$

Noting that

$$e^2 = 4\pi\epsilon_0\alpha\hbar c$$

we can also write

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2(\hbar c)^2}{s} \frac{(\cos^2 \theta + 3)^2}{\sin^4 \theta}$$

We can integrate  $d\sigma$  to obtain a cumulative distribution function. Let  $I(\theta)$  be the following integral of  $d\sigma$ . (The  $\sin \theta$  is from  $d\Omega = \sin \theta d\theta d\phi$ .)

$$I(\theta) = \int \frac{(\cos^2 \theta + 3)^2}{\sin^4 \theta} \sin \theta d\theta$$

The result is

$$I(\theta) = -\frac{8 \cos \theta}{\sin^2 \theta} - \cos \theta$$

The cumulative distribution function is

$$F(\theta) = \frac{I(\theta) - I(a)}{I(\pi - a) - I(a)}, \quad a \leq \theta \leq \pi - a$$

Angular support is reduced by an arbitrary angle  $a > 0$  because  $I(0)$  and  $I(\pi)$  are undefined.

The probability of observing scattering events in the interval  $\theta_1$  to  $\theta_2$  is

$$P(\theta_1 \leq \theta \leq \theta_2) = F(\theta_2) - F(\theta_1)$$

Let  $N$  be the total number of scattering events from an experiment. Then the number of scattering events in the interval  $\theta_1$  to  $\theta_2$  is predicted to be

$$N \times (F(\theta_2) - F(\theta_1))$$

The probability density function is

$$f(\theta) = \frac{dF(\theta)}{d\theta} = \frac{1}{I(\pi - a) - I(a)} \frac{(\cos^2 \theta + 3)^2}{\sin^4 \theta} \sin \theta$$

Note that if we had carried through the  $\alpha^2(\hbar c)^2/s$  in  $I(\theta)$ , it would have canceled out in  $F(\theta)$ .

## Notes

1. A. Zee page 134 has the cross section

$$\frac{d\sigma}{d\Omega} = \left(\frac{e^2}{4\pi}\right)^2 \frac{1}{8E^2} f(\theta)$$

where  $f(\theta)$  is the probability density function

$$f(\theta) = \frac{1 + \cos^4(\theta/2)}{\sin^4(\theta/2)} + \frac{2}{\sin^2(\theta/2) \cos^2(\theta/2)} + \frac{1 + \sin^4(\theta/2)}{\cos^4(\theta/2)}$$

The probability density function is equivalent to

$$f(\theta) = \frac{2(\cos^2 \theta + 3)^2}{\sin^4 \theta}$$

Hence for natural units  $\varepsilon_0 = \hbar = c = 1$  and  $e^2 = 4\pi\alpha$  the above cross section is equivalent to

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 (\hbar c)^2 (\cos^2 \theta + 3)^2}{4E^2 \sin^4 \theta}$$

2. In component notation, the trace operators of the Casimir trick become sums over a repeated index, in this case  $\alpha$ .

$$\begin{aligned} f_{11} &= \left( (\not{p}_3 + m)^\alpha {}_\beta \gamma^{\mu\beta} {}_\rho (\not{p}_1 + m)^\rho {}_\sigma \gamma^{\nu\sigma} {}_\alpha \right) \left( (\not{p}_4 + m)^\alpha {}_\beta \gamma_\mu{}^\beta {}_\rho (\not{p}_2 + m)^\rho {}_\sigma \gamma_\nu{}^\sigma {}_\alpha \right) \\ f_{12} &= (\not{p}_3 + m)^\alpha {}_\beta \gamma^{\mu\beta} {}_\rho (\not{p}_1 + m)^\rho {}_\sigma \gamma^{\nu\sigma} {}_\tau (\not{p}_4 + m)^\tau {}_\delta \gamma_\mu{}^\delta {}_\eta (\not{p}_2 + m)^\eta {}_\xi \gamma_\nu{}^\xi {}_\alpha \\ f_{22} &= \left( (\not{p}_4 + m)^\alpha {}_\beta \gamma^{\mu\beta} {}_\rho (\not{p}_1 + m)^\rho {}_\sigma \gamma^{\nu\sigma} {}_\alpha \right) \left( (\not{p}_3 + m)^\alpha {}_\beta \gamma_\mu{}^\beta {}_\rho (\not{p}_2 + m)^\rho {}_\sigma \gamma_\nu{}^\sigma {}_\alpha \right) \end{aligned}$$

To convert the above formulas to Eigenmath code, the  $\gamma$  tensors need to be transposed so that repeated indices are adjacent to each other. Also, multiply  $\gamma^\mu$  by the metric tensor to lower the index.

$$\begin{aligned} \gamma^{\beta\mu} {}_\rho &\rightarrow \text{gammaT} = \text{transpose}(\text{gamma}) \\ \gamma^\beta {}_{\mu\rho} &\rightarrow \text{gammaL} = \text{transpose}(\text{dot}(\text{gmunu}, \text{gamma})) \end{aligned}$$

Define the following  $4 \times 4$  matrices.

$$\begin{aligned} (\not{p}_1 + m) &\rightarrow \text{X1} = \text{pslash1} + \text{m I} \\ (\not{p}_2 + m) &\rightarrow \text{X2} = \text{pslash2} + \text{m I} \\ (\not{p}_3 + m) &\rightarrow \text{X3} = \text{pslash3} + \text{m I} \\ (\not{p}_4 + m) &\rightarrow \text{X4} = \text{pslash4} + \text{m I} \end{aligned}$$

Then for  $f_{11}$  we have the following Eigenmath code. The contract function sums over  $\alpha$ .

$$\begin{aligned} (\not{p}_3 + m)^\alpha {}_\beta \gamma^{\mu\beta} {}_\rho (\not{p}_1 + m)^\rho {}_\sigma \gamma^{\nu\sigma} {}_\alpha &\rightarrow \text{T1} = \text{contract}(\text{dot}(\text{X3}, \text{gammaT}, \text{X1}, \text{gammaT}), 1, 4) \\ (\not{p}_4 + m)^\alpha {}_\beta \gamma_\mu{}^\beta {}_\rho (\not{p}_2 + m)^\rho {}_\sigma \gamma_\nu{}^\sigma {}_\alpha &\rightarrow \text{T2} = \text{contract}(\text{dot}(\text{X4}, \text{gammaL}, \text{X2}, \text{gammaL}), 1, 4) \end{aligned}$$

Next, multiply then sum over repeated indices. The dot function sums over  $\nu$  then the contract function sums over  $\mu$ . The transpose makes the  $\nu$  indices adjacent as required by the dot function.

$$f_{11} = \text{Tr}(\cdots \gamma^\mu \cdots \gamma^\nu) \text{Tr}(\cdots \gamma_\mu \cdots \gamma_\nu) \rightarrow \text{contract}(\text{dot}(\text{T1}, \text{transpose}(\text{T2})))$$

Follow suit for  $f_{22}$ .

$$\begin{aligned} (\not{p}_4 + m)^\alpha_\beta \gamma^{\mu\beta}_\rho (\not{p}_1 + m)^\rho_\sigma \gamma^{\nu\sigma}_\alpha &\rightarrow \text{T1} = \text{contract}(\text{dot}(\text{X4}, \text{gammaT}, \text{X1}, \text{gammaT}), 1, 4) \\ (\not{p}_3 + m)^\alpha_\beta \gamma_\mu^\beta_\rho (\not{p}_2 + m)^\rho_\sigma \gamma_\nu^\sigma_\alpha &\rightarrow \text{T2} = \text{contract}(\text{dot}(\text{X3}, \text{gammaL}, \text{X2}, \text{gammaL}), 1, 4) \end{aligned}$$

Then

$$f_{22} = \text{Tr}(\cdots \gamma^\mu \cdots \gamma^\nu) \text{Tr}(\cdots \gamma_\mu \cdots \gamma_\nu) \rightarrow \text{contract}(\text{dot}(\text{T1}, \text{transpose}(\text{T2})))$$

The calculation of  $f_{12}$  begins with

$$\begin{aligned} (\not{p}_3 + m)^\alpha_\beta \gamma^{\mu\beta}_\rho (\not{p}_1 + m)^\rho_\sigma \gamma^{\nu\sigma}_\tau (\not{p}_4 + m)^\tau_\delta \gamma_\mu^\delta_\eta (\not{p}_2 + m)^\eta_\xi \gamma_\nu^\xi_\alpha \\ \rightarrow \text{T} = \text{contract}(\text{dot}(\text{X3}, \text{gammaT}, \text{X1}, \text{gammaT}, \text{X4}, \text{gammaL}, \text{X2}, \text{gammaL}), 1, 6) \end{aligned}$$

Then sum over repeated indices  $\mu$  and  $\nu$ .

$$f_{12} = \text{Tr}(\cdots \gamma^\mu \cdots \gamma^\nu \cdots \gamma_\mu \cdots \gamma_\nu) \rightarrow \text{contract}(\text{contract}(\text{T}, 1, 3))$$