The following table is from Particle Data Group.¹

μ MEAN LIFE τ

Measurements with an error $> 0.001 \times 10^{-6}$ s have been omitted.

<i>VALUE</i> (10 ⁻⁶ s)	DOCUMENT ID		TECN	CHG	COMMENT				
2.1969811±0.0000022 OUR AVERAGE									
$2.1969803 \pm 0.0000021 \pm 0.0000007$	^L TISHCHENKO								
$2.197083 \pm 0.000032 \pm 0.000015$	BARCZYK				Muons from π^+ decay at rest				
$2.197013 \pm 0.000021 \pm 0.000011$	CHITWOOD	07	CNTR	+	Surface μ^+ at PSI				
2.197078 ± 0.000073	BARDIN	84	CNTR	+					
2.197025 ± 0.000155	BARDIN	84	CNTR	_					
2.19695 ± 0.00006	GIOVANETTI	84	CNTR	+					
2.19711 ± 0.00008	BALANDIN	74	CNTR	+					
2.1973 ± 0.0003	DUCLOS	73	CNTR	+					
ullet $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$									
$2.1969803\!\pm\!0.0000022$	WEBBER	11	CNTR	+	Surface μ^+ at PSI				
1 TISHCHENKO 13 uses $1.6 imes 10^{12}~\mu^+$ events and supersedes WEBBER 11.									

From "V minus A" theory we have the following formula for muon lifetime τ .

$$\tau = \frac{96\pi^2 h}{G_F^2 \left(m_\mu c^2\right)^5}$$

Symbol G_F is Fermi coupling constant, m_{μ} is muon mass.

From NIST² we have

$$G_F = 1.1663787 \times 10^{-5} \text{ GeV}^{-2}$$

 $m_{\mu} = 1.883531627 \times 10^{-28} \text{ kilogram}$
 $h = 6.62607015 \times 10^{-34} \text{ joule second (exact)}$
 $c = 299792458 \text{ meter second}^{-1} \text{ (exact)}$
 $1 \text{ eV} = 1.602176634 \times 10^{-19} \text{ joule (exact)}$

Hence

$$\tau = 2.18735 \times 10^{-6} \, \text{second}$$

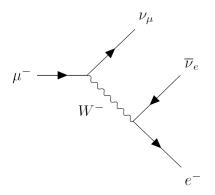
The result is a bit smaller than the PDG value.

$$\frac{\tau}{2.1969811 \times 10^{-6} \; \mathrm{second}} = 0.9956$$

A muon decays into a muon neutrino, an electron anti-neutrino, and an electron.

https://pdg.lbl.gov/2020/listings/rpp2020-list-muon.pdf

²https://physics.nist.gov/cuu/Constants/index.html



Particle	Symbol	Momentum	Spinor (up)	Spinor (down)
Muon	μ^-	p_1	u_{11}	u_{12}
Muon neutrino	$ u_{\mu}$	p_2	u_{21}	u_{22}
Electron anti-neutrino	$ar{ u}_e$	p_3	v_{31}	v_{32}
Electron	e^{-}	p_4	u_{41}	u_{42}

We will use the following momentum vectors.

$$p_{1} = \begin{pmatrix} E_{1} \\ p_{1x} \\ p_{1y} \\ p_{1z} \end{pmatrix} \qquad p_{2} = \begin{pmatrix} E_{2} \\ p_{2x} \\ p_{2y} \\ p_{2z} \end{pmatrix} \qquad p_{3} = \begin{pmatrix} E_{3} \\ p_{3x} \\ p_{3y} \\ p_{3z} \end{pmatrix} \qquad p_{4} = \begin{pmatrix} E_{4} \\ p_{4x} \\ p_{4y} \\ p_{4z} \end{pmatrix}$$

$$\mu^{-} \qquad \nu_{\mu} \qquad \bar{\nu}_{e} \qquad e^{-}$$

And we will also use the following Dirac spinors.

$$u_{11} = \begin{pmatrix} E_1 + m_1 \\ 0 \\ p_{1z} \\ p_{1x} + ip_{1y} \end{pmatrix} \qquad u_{21} = \begin{pmatrix} E_2 + m_2 \\ 0 \\ p_{2z} \\ p_{2x} + ip_{2y} \end{pmatrix} \qquad v_{31} = \begin{pmatrix} p_{3z} \\ p_{3x} + ip_{3y} \\ E_3 + m_3 \\ 0 \end{pmatrix} \qquad u_{41} = \begin{pmatrix} E_4 + m_4 \\ 0 \\ p_{4z} \\ p_{4x} + ip_{4y} \end{pmatrix}$$

$$u_{12} = \begin{pmatrix} 0 \\ E_1 + m_1 \\ p_{1x} - ip_{1y} \\ -p_{1z} \end{pmatrix} \qquad u_{22} = \begin{pmatrix} 0 \\ E_2 + m_2 \\ p_{2x} - ip_{2y} \\ -p_{2z} \end{pmatrix} \qquad v_{32} = \begin{pmatrix} p_{3x} - ip_{3y} \\ -p_{3z} \\ 0 \\ E_3 + m_3 \end{pmatrix} \qquad u_{42} = \begin{pmatrix} 0 \\ E_4 + m_4 \\ p_{4x} - ip_{4y} \\ -p_{4z} \end{pmatrix}$$

The energy terms are total energy.

$$E_1 = \sqrt{(p_{1x})^2 + (p_{1y})^2 + (p_{1z})^2 + m_1^2}$$

$$E_2 = \sqrt{(p_{2x})^2 + (p_{2y})^2 + (p_{2z})^2 + m_2^2}$$

$$E_3 = \sqrt{(p_{3x})^2 + (p_{3y})^2 + (p_{3z})^2 + m_3^2}$$

$$E_4 = \sqrt{(p_{4x})^2 + (p_{4y})^2 + (p_{4z})^2 + m_4^2}$$

From the Feynman diagram above we have the following amplitude \mathcal{M}_{abcd} where each letter in abcd can be either 1 (spin up) or 2 (spin down).

$$\mathcal{M}_{abcd} = \frac{G_F}{\sqrt{N}} \left(\bar{u}_{4d} \gamma^{\mu} (1 - \gamma^5) v_{3c} \right) \left(\bar{u}_{2b} \gamma_{\mu} (1 - \gamma^5) u_{1a} \right)$$

Symbol N is the following normalization constant.

$$N = (E_1 + m_1)(E_2 + m_2)(E_3 + m_3)(E_4 + m_4)$$

Recall that the magnitude squared of an amplitude is a probability density and also an observable.

$$|\mathcal{M}_{abcd}|^2 = \mathcal{M}^*_{abcd} \mathcal{M}_{abcd}$$

In a typical muon decay experiment the spins are not observed. Consequently, the experimental result is an average of spin states. The average is computed by summing over all spin states and dividing by the number of inbound spin states, in this case four. (In the Feynman diagram, two particles have arrows pointing into vertices, these are the inbound particles. There are two spin states for μ^- and two spin states for $\bar{\nu}_e$. Hence there are $2 \times 2 = 4$ inbound spin states.)

$$\langle |\mathcal{M}|^2 \rangle = \frac{1}{4} \sum_{a=1}^2 \sum_{b=1}^2 \sum_{c=1}^2 \sum_{d=1}^2 |\mathcal{M}_{abcd}|^2$$

The result is a simple formula.

$$\langle |\mathcal{M}|^2 \rangle = 64G_F^2(p_1 \cdot p_3)(p_2 \cdot p_4) \tag{1}$$

In component notation we have

$$\langle |\mathcal{M}|^2 \rangle = 64G_F^2 \bigg((p_1)^{\alpha} g_{\alpha\beta} (p_3)^{\beta} \bigg) \bigg((p_2)^{\gamma} g_{\gamma\delta} (p_4)^{\delta} \bigg)$$

where

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

Muon decay rate Γ is an average over all possible momenta. (The simplified measure dp represents a lengthy product of terms.)

$$\Gamma = \int_{12 \text{ integrals}} \langle |\mathcal{M}|^2 \rangle \, dp$$

It can be shown that

$$\Gamma = \frac{G_F^2 m_\mu^5}{192\pi^3}$$

where $m_{\mu} = m_1$. Muon lifetime τ is the inverse of decay rate.

$$\tau = \frac{192\pi^3}{G_F^2 m_{_H}^5}$$

In physical units for c and h we have

$$\tau = \frac{96\pi^2 h}{G_F^2 \left(m_\mu c^2\right)^5}$$