

Atomic transitions 2

For electric field

$$\mathbf{E} = E_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t) \boldsymbol{\epsilon}$$

the Hamiltonian is

$$H_1(\mathbf{r}, t) = -\frac{eE_0}{m\omega} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t) \boldsymbol{\epsilon} \cdot \mathbf{p}$$

In exponential form

$$H_1(\mathbf{r}, t) = -\frac{eE_0}{m\omega} \left(\frac{1}{2} \exp(i\mathbf{k} \cdot \mathbf{r} - i\omega t) + \frac{1}{2} \exp(-i\mathbf{k} \cdot \mathbf{r} + i\omega t) \right) \boldsymbol{\epsilon} \cdot \mathbf{p}$$

Given the initial condition $c_b(0) = 0$ the first-order approximation for $c_b(t)$ is

$$c_b(t) = -\frac{i}{\hbar} \int_0^t \langle \psi_b | H_1(\mathbf{r}, t') | \psi_a \rangle \exp(i\omega_0 t') dt', \quad \omega_0 = \frac{E_b - E_a}{\hbar}$$

Factor the integrand.

$$\begin{aligned} c_b(t) &= \frac{ieE_0}{2m\hbar\omega} \langle \psi_b | \exp(i\mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \int_0^t \exp(-i\omega t') \exp(i\omega_0 t') dt' \\ &\quad + \frac{ieE_0}{2m\hbar\omega} \langle \psi_b | \exp(-i\mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \int_0^t \exp(i\omega t') \exp(i\omega_0 t') dt' \end{aligned}$$

Solve the time integrals to obtain

$$\begin{aligned} c_b(t) &= \frac{eE_0}{2m\hbar\omega} \langle \psi_b | \exp(i\mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \frac{\exp(i(\omega_0 - \omega)t) - 1}{\omega_0 - \omega} \\ &\quad + \frac{eE_0}{2m\hbar\omega} \langle \psi_b | \exp(-i\mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \frac{\exp(i(\omega_0 + \omega)t) - 1}{\omega_0 + \omega} \quad (1) \end{aligned}$$

As an approximation, discard the second term since the first term dominates for $\omega \approx \omega_0$.

$$c_b(t) = \frac{eE_0}{2m\hbar\omega} \langle \psi_b | \exp(i\mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \frac{\exp(i(\omega_0 - \omega)t) - 1}{\omega_0 - \omega}$$

Rewrite $c_b(t)$ in the form of a sine function.

$$c_b(t) = \frac{ieE_0}{m\hbar\omega} \langle \psi_b | \exp(i\mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle \frac{\sin\left(\frac{1}{2}(\omega_0 - \omega)t\right)}{\omega_0 - \omega} \exp\left(\frac{i}{2}(\omega_0 - \omega)t\right) \quad (2)$$

Verify dimensions.

$$H_1(\mathbf{r}, t) \propto \frac{e}{C} \frac{E_0}{N C^{-1}} \times \frac{\boldsymbol{\epsilon} \cdot \mathbf{p}}{\frac{m}{kg} \frac{\omega}{s^{-1}}} = N m = J$$

$$c_b(t) \propto \frac{e}{C} \frac{E_0}{N C^{-1}} \times \frac{\langle \psi_b | \exp(i\mathbf{k} \cdot \mathbf{r}) \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle}{\frac{kg m s^{-1}}{\omega_0 - \omega}} = \frac{N m}{J} = 1$$

Wave functions ψ_a and ψ_b have dimension meter $^{-1/2}$ hence they cancel with $dx \propto$ meter in the integral leaving units of momentum due to \mathbf{p} .