The following table of hydrogen transition data is from "Atomic Transition Probabilities," 1966.

| Transition | λ(Å) | $E_i(\text{cm}^{-1})$ | $E_k(\mathrm{cm}^{-1})$ | gi | gk | $A_{ki}(\sec^{-1})$ |
|---|---------|-----------------------|-------------------------|---------|-------------|--|
| 1. 0 | 1215.67 | 0 | 82259 | 9 | 6 | 6.065 × 108 |
| $ \begin{array}{l} 1s-2p\\1s-3p \end{array} $ | 1025.72 | 0 | 97492 | 2 | 6 6 | 6.265×10^{8} 1.672×10^{8} |
| 1s-3p 1s-4p | 972.537 | ŏ | 102824 | 2 | 6 | 6.818×10^{7} |
| 1s-4p 1s-5p | 949.743 | ŏ | 105292 | 2 | 6 | 3.437×10^7 |
| 1s-6p | 937.804 | 0 | 106632 | 2 2 2 2 | 6 | 1.973×10^{7} |
| 0- 2- | 6562.86 | 82259 | 97492 | 6 | 2 | 6 212 × 106 |
| 2p-3s | 4861.35 | 82259 | | 6 | 2 | 6.313×10^{6} |
| 2p-4s | | 82259 | 102824 | 6 | 2 2 2 | 2.578×10^{6} |
| 2p-5s | 4340.48 | 82259 | 105292 | 6 | 2 | 1.289×10^{6} |
| 2p-6s | 4101.75 | 62239 | 106632 | 0 | | 7.350×10^{5} |
| 2s-3p | 6562.74 | 82259 | 97492 | 2 | 6 | 2.245×10^{7} |
| 2s-4p | 4861.29 | 82259 | 102824 | 2 2 2 2 | 6 | 9.668×10^{6} |
| 2s-5p | 4340.44 | 82259 | 105292 | 2 | 6 | 4.948×10^{6} |
| 2s-6p | 4101.71 | 82259 | 106632 | 2 | 6 | 2.858×10^{6} |
| 2p-3d | 6562.81 | 82259 | 97492 | 6 | 10 | 6.465×10^{7} |
| 2p-4d | 4861.33 | 82259 | 102824 | 6 | 10 | 2.062×10^{7} |
| 2p-5d | 4340.47 | 82259 | 105292 | 6 | 10 | 9.425×10^{6} |
| 2p-6d | 4101.74 | 82259 | 106632 | 6 | 10 | 5.145×10^{6} |

The 2-3 transitions emit the bright red H- α line.

| Transition | λ (Å) | $A_{ki} (\operatorname{second}^{-1})$ |
|------------|---------|---------------------------------------|
| 2p-3s | 6562.86 | 6.313×10^{6} |
| 2s-3p | 6562.74 | 2.245×10^{7} |
| 2p-3d | 6562.81 | 6.465×10^{7} |

Let us compute the spontaneous emission coefficients A_{ki} for H- α and see if the results match the table.

The orbital names correspond to the following angular momenta.

| Letter | Angular momentum ℓ |
|--------|-------------------------|
| s | 0 |
| p | 1 |
| d | 2 |

Because of the magnetic quantum number m_{ℓ} there are multiple ways for each orbital transition to occur. $(m_{\ell} = 0, \pm 1, \dots, \pm \ell)$

There are three transitions for $3s \to 2p$.

$$\psi_{3,0,0} \to \psi_{2,1,1} \psi_{3,0,0} \to \psi_{2,1,0} \psi_{3,0,0} \to \psi_{2,1,-1}$$

There are three transitions for $3p \rightarrow 2s$.

$$\psi_{3,1,1} \to \psi_{2,0,0}$$

$$\psi_{3,1,0} \to \psi_{2,0,0}$$

$$\psi_{3,1,-1} \to \psi_{2,0,0}$$

Finally, there are fifteen transitions for $3d \rightarrow 2p$.

For each H- α line, an average A_{ki} is obtained by summing A_{ki} for individual transitions and dividing by the number of distinct initial states.

For example, $3d \rightarrow 2p$ has five distinct initial states, so the divisor is five.

 A_{ki} can be computed using the following Heisenberg formula.

$$A_{ki} = \frac{e^2}{3\pi\varepsilon_0\hbar c^3} \,\omega_{ki}^3 \,|r_{ki}|^2$$

The transition frequency ω_{ki} is given by Bohr's frequency condition.

$$\omega_{ki} = \frac{1}{\hbar} (E_k - E_i)$$

The transition probability (multiplied by a physical constant) is

$$|r_{ki}|^2 = |x_{ki}|^2 + |y_{ki}|^2 + |z_{ki}|^2$$

For wave functions ψ in spherical coordinates we have the following transition amplitudes.

$$x_{ki} = \int \psi_k^* (r \sin \theta \cos \phi) \, \psi_i \, dV$$
$$y_{ki} = \int \psi_k^* (r \sin \theta \sin \phi) \, \psi_i \, dV$$
$$z_{ki} = \int \psi_k^* (r \cos \theta) \, \psi_i \, dV$$

Using Eigenmath we obtain

$$A_{3s2p} = 6.31358 \times 10^6 \text{ second}^{-1}$$

 $A_{3p2s} = 2.24483 \times 10^7 \text{ second}^{-1}$
 $A_{3d2p} = 6.4651 \times 10^7 \text{ second}^{-1}$

which is very close to the values shown in the table.

Some of the $|r_{ki}|^2$ are zero, indicating forbidden transitions.

The following tables show $|r_{ki}|^2$ for each transition (multiply all by a_0^2).

Each row is an initial state ψ_i and each column is a final state ψ_k .

| | $\psi_{2,1,1}$ | $\psi_{2,1,0}$ | $\psi_{2,1,-1}$ |
|-----------------|-----------------|----------------|-----------------|
| $\psi_{3,0,0}$ | 0.293534 | 0.293534 | 0.293534 |
| | | | |
| | | $\psi_{2,0,0}$ | |
| | $\psi_{3,1,1}$ | 3.13103 | |
| | $\psi_{3,1,0}$ | 3.13103 | |
| | $\psi_{3,1,-1}$ | 3.13103 | |
| | , | , | , |
| | $\psi_{2,1,1}$ | $\psi_{2,1,0}$ | $\psi_{2,1,-1}$ |
| $\psi_{3,2,2}$ | 9.01737 | 0 | 0 |
| $\psi_{3,2,1}$ | 4.50868 | 4.50868 | 0 |
| $\psi_{3,2,0}$ | 1.50289 | 6.01158 | 1.50289 |
| $\psi_{3,2,-1}$ | 0 | 4.50868 | 4.50868 |
| $\psi_{3,2,-2}$ | $_{2}$ 0 | 0 | 9.01737 |