

9-8. For the state for which there is just one photon present in level 1, \mathbf{k} , all of the factors in the wave function are ϕ_0 except one, which is ϕ_1 . But for an oscillator $\phi_1(x) = \sqrt{2}x\phi_0(x)$. The wave function representing an excited running wave is a linear superposition of the state with the cosine mode excited and i times the state with the sine wave excited, so show that the unnormalized wave function for just one photon present in 1, \mathbf{k} is $\bar{a}_{1,\mathbf{k}}^*\Phi_0$. The normalization is $\int \Phi_0^* \bar{a}_{1,\mathbf{k}} \bar{a}_{1,\mathbf{k}}^* \Phi_0 d\bar{a}$, or the expectation of $\bar{a}_{1,\mathbf{k}} \bar{a}_{1,\mathbf{k}}^*$ for the vacuum, which we have seen in the preceding problem is $\hbar/2kc$. Hence the normalized one-photon state is $\sqrt{2kc/\hbar} \bar{a}_{1,\mathbf{k}}^* \Phi_0$.

From problem 9-6, let

$$\Phi_0 = \exp \left(-\frac{kc}{4\hbar} (\bar{a}_{1,\mathbf{k}}^c)^2 - \frac{kc}{4\hbar} (\bar{a}_{1,\mathbf{k}}^s)^2 \right)$$

It follows that

$$\Phi_0^* \Phi_0 = \exp \left(-\frac{kc}{2\hbar} (\bar{a}_{1,\mathbf{k}}^c)^2 - \frac{kc}{2\hbar} (\bar{a}_{1,\mathbf{k}}^s)^2 \right)$$

The expectation of Φ_0 is

$$\langle \Phi_0 \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Phi_0^* \Phi_0 d\bar{a}_{1,\mathbf{k}}^c d\bar{a}_{1,\mathbf{k}}^s = \frac{2\pi\hbar}{kc} \quad (1)$$

Let

$$\Phi_1 = \bar{a}_{1,\mathbf{k}}^* \Phi_0$$

Then

$$\Phi_1^* \Phi_1 = \Phi_0^* \frac{(\bar{a}_{1,\mathbf{k}}^c)^2 + (\bar{a}_{1,\mathbf{k}}^s)^2}{2} \Phi_0$$

The expectation of Φ_1 is

$$\langle \Phi_1 \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Phi_1^* \Phi_1 d\bar{a}_{1,\mathbf{k}}^c d\bar{a}_{1,\mathbf{k}}^s = \frac{\hbar}{kc} \frac{2\pi\hbar}{kc} \quad (2)$$

Hence the expectation of Φ_1 is \hbar/kc times the vacuum state.

The energy of a single photon is $\hbar\omega = \hbar/kc$.