

## Atomic transitions 3

In the previous section we obtained

$$c_b(t) = \frac{ieE_0}{m\hbar\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \frac{\sin(\frac{1}{2}(\omega_0 - \omega)t)}{\omega_0 - \omega} \exp(\frac{i}{2}(\omega_0 - \omega)t) \quad (1)$$

Use the dipole approximation

$$\exp(i\mathbf{k} \cdot \mathbf{r}) = 1 + i\mathbf{k} \cdot \mathbf{r} + \dots \approx 1$$

to obtain

$$\langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \approx \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle$$

By the identity

$$\mathbf{p} = \frac{im}{\hbar} [H_0, \mathbf{r}] \quad (2)$$

we have

$$\begin{aligned} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} | \psi_a \rangle &= \frac{im}{\hbar} \langle \psi_b | \boldsymbol{\epsilon} \cdot [H_0, \mathbf{r}] | \psi_a \rangle \\ &= \frac{im}{\hbar} \langle \psi_b | \boldsymbol{\epsilon} \cdot H_0 \mathbf{r} | \psi_a \rangle - \frac{im}{\hbar} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} H_0 | \psi_a \rangle \\ &= \frac{im}{\hbar} E_b \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle - \frac{im}{\hbar} E_a \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \\ &= \frac{im}{\hbar} (E_b - E_a) \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \\ &= im\omega_0 \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \end{aligned}$$

Hence

$$\langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{p} \exp(i\mathbf{k} \cdot \mathbf{r}) | \psi_a \rangle \approx im\omega_0 \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle$$

Substitute into (1) to obtain

$$c_b(t) = -\frac{eE_0}{\hbar} \frac{\omega_0}{\omega} \langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle \frac{\sin(\frac{1}{2}(\omega_0 - \omega)t)}{\omega_0 - \omega} \exp(\frac{i}{2}(\omega_0 - \omega)t)$$

Verify dimensions.

$$c_b(t) \propto \frac{\frac{e}{\text{coulomb}} \frac{E_0}{\text{joule second}}}{\hbar} \times \frac{\frac{\omega_0}{\text{second}^{-1}}}{\omega} \times \frac{\langle \psi_b | \boldsymbol{\epsilon} \cdot \mathbf{r} | \psi_a \rangle}{\frac{\omega_0 - \omega}{\text{second}^{-1}}} = 1$$

Note that for an experiment with  $\boldsymbol{\epsilon} \cdot \mathbf{r} = x$  we have

$$c_b(t) = -\frac{eE_0}{\hbar} \frac{\omega_0}{\omega} \langle \psi_b | x | \psi_a \rangle \frac{\sin(\frac{1}{2}(\omega_0 - \omega)t)}{\omega_0 - \omega} \exp(\frac{i}{2}(\omega_0 - \omega)t)$$