## Rotating wave approximation

Let  $\Psi(\mathbf{r},t)$  be the following wave function for a two state system.

$$\Psi(\mathbf{r},t) = \psi_a(\mathbf{r})c_a(t)\exp(-\frac{i}{\hbar}E_at) + \psi_b(\mathbf{r})c_b(t)\exp(-\frac{i}{\hbar}E_bt)$$

Let  $\hat{H}(\mathbf{r},t)$  be the Hamiltonian

$$\hat{H}(\mathbf{r},t) = \hat{H}_0(\mathbf{r}) + \hat{H}_1(\mathbf{r},t)$$

where

$$\hat{H}_0 \psi_a = E_a \psi_a, \quad \hat{H}_0 \psi_b = E_b \psi_b, \quad \hat{H}_0 \Psi = (E_a + E_b) \Psi$$

It was shown that if  $\Psi$  is a solution to the Schrödinger equation then

$$\frac{d}{dt}c_a(t) = -\frac{i}{\hbar}\langle\psi_a|\hat{H}_1|\psi_b\rangle\exp(-i\omega_0t)c_b(t), \quad \frac{d}{dt}c_b(t) = -\frac{i}{\hbar}\langle\psi_b|\hat{H}_1|\psi_a\rangle\exp(i\omega_0t)c_a(t) \quad (1)$$

where

$$\omega_0 = \frac{E_b - E_a}{\hbar}$$

Let  $\hat{H}_1(\mathbf{r},t)$  be the perturbation

$$\hat{H}_1(\mathbf{r},t) = \hat{V}(\mathbf{r})\cos(\omega t)$$

Then

$$\langle \psi_a | \hat{H}_1 | \psi_b \rangle = \langle \psi_a | \hat{V} | \psi_b \rangle \left[ \frac{1}{2} \exp(i\omega t) + \frac{1}{2} \exp(-i\omega t) \right]$$

The rotating wave approximation discards the second term and asserts

$$\langle \psi_a | \hat{H}_1 | \psi_b \rangle = \frac{1}{2} \langle \psi_a | \hat{V} | \psi_b \rangle \exp(i\omega t)$$
 (2)

Substitute equation (2) into (1) to obtain

$$\frac{d}{dt}c_a(t) = -\frac{i}{2\hbar}\langle \psi_a | \hat{V} | \psi_b \rangle \exp(i(\omega - \omega_0)t)c_b(t)$$
(3)

and

$$\frac{d}{dt}c_b(t) = -\frac{i}{2\hbar} \langle \psi_b | \hat{V} | \psi_a \rangle \exp(i(\omega_0 - \omega)t) c_a(t)$$
(4)

Use Laplace transforms to solve for  $c_b(t)$  with initial conditions  $c_a(0) = 1$  and  $c_b(0) = 0$ .

$$c_b(t) = -\frac{i}{2\hbar\omega_r} \langle \psi_b | \hat{V} | \psi_a \rangle \sin(\omega_r t) \exp\left(\frac{i}{2}(\omega_0 - \omega)t\right)$$

Symbol  $\omega_r$  is the Rabi flopping frequency

$$\omega_r = \frac{1}{2} \sqrt{(\omega_0 - \omega)^2 + \left| \langle \psi_a | \hat{V} | \psi_b \rangle \right|^2 / \hbar^2}$$

Use the latter part of equation (1) and the solution for  $c_b(t)$  to solve for  $c_a(t)$ .

$$c_a(t) = \left[\cos(\omega_r t) + i\left(\frac{\omega_0 - \omega}{2\omega_r}\right)\sin(\omega_r t)\right] \exp\left(-\frac{i}{2}(\omega_0 - \omega)t\right)$$

Rewrite  $\omega_r$  as

$$\omega_r = \frac{1}{2\hbar} \sqrt{\hbar^2 (\omega_0 - \omega)^2 + \left| \langle \psi_a | \hat{V} | \psi_b \rangle \right|^2}$$

and note that for

$$\hbar^2(\omega_0 - \omega)^2 \gg \left| \langle \psi_a | \hat{V} | \psi_b \rangle \right|^2$$

we have

$$\omega_r \approx \frac{1}{2} |\omega_0 - \omega|$$

Using this approximation, the transition probability  $P_{a\to b}(t) = |c_b(t)|^2$  becomes identical to the first order perturbation result.