(36.1) (a) Show that the Dirac equation can be recast in the form

$$i\frac{\partial\psi}{\partial t} = \hat{H}_D\psi \tag{36.33}$$

where $\hat{H}_D = \boldsymbol{\alpha} \cdot \hat{\mathbf{p}} + \beta m$ and find $\boldsymbol{\alpha}$ and β in terms of the γ matrices.

(b) Evaluate \hat{H}_D^2 and show that for a Klein-Gordon dispersion to result we must have:

(i) that the α^i and β objects all anticommute with each other; and

(ii)
$$(\alpha^i)^2 = (\beta)^2 = 1$$
.

(c) Prove the following commutation relations

(i)
$$[\hat{H}, \hat{L}^i] = i(\hat{\mathbf{p}} \times \boldsymbol{\alpha})^i$$
 where $\hat{\mathbf{L}} = \hat{\mathbf{x}} \times \hat{\mathbf{p}}$.

(ii)
$$[\hat{H}, \hat{S}^i] = -i(\hat{\mathbf{p}} \times \boldsymbol{\alpha})^i$$
 where $\hat{\mathbf{S}} = \frac{1}{2}\boldsymbol{\Sigma}$ and we define $\boldsymbol{\Sigma} = \frac{i}{2}\boldsymbol{\gamma} \times \boldsymbol{\gamma}$.

(a) Consider the following form of the Dirac equation.

$$i\left(\gamma^0 \frac{\partial}{\partial t} + \gamma^1 \frac{\partial}{\partial x} + \gamma^2 \frac{\partial}{\partial y} + \gamma^3 \frac{\partial}{\partial z}\right) \psi = m\psi$$

Rewrite as

$$i\gamma^0 \frac{\partial}{\partial t} \psi = -i \left(\gamma^1 \frac{\partial}{\partial x} + \gamma^2 \frac{\partial}{\partial y} + \gamma^3 \frac{\partial}{\partial z} \right) \psi + m \psi$$

Noting that $\gamma^0 \gamma^0 = I$, multiply both sides by γ^0 to obtain

$$i\frac{\partial}{\partial t}\psi = -i\gamma^0 \left(\gamma^1 \frac{\partial}{\partial x} + \gamma^2 \frac{\partial}{\partial y} + \gamma^3 \frac{\partial}{\partial z}\right)\psi + m\gamma^0 \psi$$

Hence for $\hat{\mathbf{p}} = -i\nabla$ we have

$$\boldsymbol{lpha} = \gamma^0 \begin{pmatrix} \gamma^1 \\ \gamma^2 \\ \gamma^3 \end{pmatrix}, \quad \beta = \gamma^0$$

(b) The dispersion relation is

$$\hat{H}_D^2 = \hat{\mathbf{p}}^2 + m^2$$

Squaring \hat{H}_D we have

$$\hat{H}_D^2 = (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}} + \beta m)(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}} + \beta m)$$

= $(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}}) + (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\beta m + \beta m(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}}) + \beta^2 m^2$

The middle terms must cancel, that is

$$(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\beta m + \beta m(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}}) = 0$$

Hence

$$\alpha^i \beta = -\beta \alpha^i$$

Cross terms must cancel, that is

$$\left(-i\alpha^1\frac{\partial}{\partial x} - i\alpha^2\frac{\partial}{\partial y} - i\alpha^3\frac{\partial}{\partial z}\right)^2 = -(\alpha^1)^2\frac{\partial^2}{\partial x^2} - (\alpha^2)^2\frac{\partial^2}{\partial y^2} - (\alpha^3)^2\frac{\partial^2}{\partial z^2}$$

Hence

$$\alpha^i \alpha^j = -\alpha^j \alpha^i, \quad i \neq j$$

With the above anticommutation relations we now have

$$\hat{H}_D^2 = -(\alpha^1)^2 \frac{\partial^2}{\partial x^2} - (\alpha^2)^2 \frac{\partial^2}{\partial y^2} - (\alpha^3)^2 \frac{\partial^2}{\partial z^2} + \beta^2 m = \hat{\mathbf{p}}^2 + m^2$$

Hence

$$(\alpha^i)^2 = I$$
 and $\beta^2 = I$

(c) We have

$$\hat{\mathbf{L}}\psi = \hat{\mathbf{x}} \times \hat{\mathbf{p}}\psi$$

If follows that

$$[\hat{H}_D, \hat{\mathbf{L}}]\psi = \hat{H}_D \hat{\mathbf{L}}\psi - \hat{\mathbf{L}}\hat{H}_D \psi$$

= $(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}} + \beta m)(\hat{\mathbf{x}} \times \hat{\mathbf{p}}\psi) - \hat{\mathbf{x}} \times \hat{\mathbf{p}}(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}} + \beta m)\psi$

The βm terms cancel by the following identity.

$$\beta m \hat{\mathbf{x}} \times \hat{\mathbf{p}} \psi = \hat{\mathbf{x}} \times \beta m \hat{\mathbf{p}} \psi = \hat{\mathbf{x}} \times \hat{\mathbf{p}} (\beta m \psi)$$

Hence

$$[\hat{H}_D, \hat{\mathbf{L}}]\psi = (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})(\hat{\mathbf{x}} \times \hat{\mathbf{p}}\psi) - \hat{\mathbf{x}} \times \hat{\mathbf{p}}(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\psi$$

By the product rule for differentiation

$$(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})(\hat{\mathbf{x}} \times \hat{\mathbf{p}}\psi) = (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\hat{\mathbf{x}} \times \hat{\mathbf{p}}\psi + \hat{\mathbf{x}} \times (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\hat{\mathbf{p}}\psi$$

Hence

$$[\hat{H}_D, \hat{\mathbf{L}}]\psi = (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\hat{\mathbf{x}} \times \hat{\mathbf{p}}\psi + \hat{\mathbf{x}} \times (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\hat{\mathbf{p}}\psi - \hat{\mathbf{x}} \times \hat{\mathbf{p}}(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\psi$$
$$= (\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\hat{\mathbf{x}} \times \hat{\mathbf{p}}\psi$$

Note that

$$(\boldsymbol{\alpha} \cdot \hat{\mathbf{p}})\hat{\mathbf{x}} = \boldsymbol{\alpha} \cdot (-i\nabla \hat{\mathbf{x}}) = -i\boldsymbol{\alpha}$$

Hence

$$[\hat{H}_D, \hat{\mathbf{L}}]\psi = -i\boldsymbol{\alpha} \times \hat{\mathbf{p}}\psi = i(\hat{\mathbf{p}} \times \boldsymbol{\alpha})\psi$$