In his 1917 paper, "On the Quantum Theory of Radiation," Einstein uses the following argument to derive Planck's law. The argument requires induced emission, a process which was a theoretical discovery by Einstein. Prior to Einstein, no one was aware that induced emission existed.

Consider a gas at temperature T. Let N be the number of molecules in the gas and let N_n be the number of molecules with internal energy ε_n . By the Maxwell-Boltzmann distribution we have

$$\frac{N_n}{N} = p_n \exp\left(-\frac{\varepsilon_n}{kT}\right)$$

where k is Boltzmann's constant. The coefficient p_n is a statistical weighting factor that does not depend on T.

Let us now consider the processes by which a molecule transitions between energy levels. The processes are absorption, induced emission, and spontaneous emission. Let ε_m be an energy level such that $\varepsilon_m > \varepsilon_n$. Let B_{nm} , B_{mn} , and A_{mn} be coefficients of transition rates such that

$$\frac{dN_n}{dt} = B_{nm}\rho(\nu)N_n \qquad \frac{dN_m}{dt} = B_{mn}\rho(\nu)N_m \qquad \frac{dN_m}{dt} = A_{mn}N_m$$
absorption
$$\underset{\varepsilon_m \to \varepsilon_m}{\text{absorption}} \qquad \underset{\varepsilon_m \to \varepsilon_n}{\text{induced emission}} \qquad \underset{\varepsilon_m \to \varepsilon_n}{\text{spontaneous emission}}$$

Absorption and induced emission are proportional to $\rho(\nu)$ which is the radiant energy density of the gas as a function of radiant frequency ν . The A and B coefficients are presumed to not depend on temperature T.

At equilibrium, transition rates between ε_m and ε_n are equal.

Divide through by N to obtain

$$B_{nm}\rho(\nu)p_n \exp\left(-\frac{\varepsilon_n}{kT}\right) = B_{mn}\rho(\nu)p_m \exp\left(-\frac{\varepsilon_m}{kT}\right) + A_{mn}p_m \exp\left(-\frac{\varepsilon_m}{kT}\right)$$
absorption induced emission spontaneous emission
$$\varepsilon_m \to \varepsilon_m$$
 spontaneous emission
$$\varepsilon_m \to \varepsilon_n$$

Multiply both sides by $\exp(\varepsilon_m/kT)$.

$$B_{nm}\rho(\nu)p_n \exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) = B_{mn}\rho(\nu)p_m + A_{mn}p_m$$
absorption
$$\sum_{\varepsilon \to \varepsilon_n} a_{mn} = B_{mn}\rho(\nu)p_m + A_{mn}p_m$$
induced emission
$$\sum_{\varepsilon \to \varepsilon_n} a_{mn} = B_{mn}\rho(\nu)p_m + A_{mn}p_m$$

Note that for increasing T we have

$$\lim_{T \to \infty} \exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) = 1$$

It follows that for $T \to \infty$ the equilibrium formula is

$$B_{nm}\rho(\nu)p_n = B_{mn}\rho(\nu)p_m + A_{mn}p_m$$

Divide through by $\rho(\nu)$.

$$B_{nm}p_n = B_{mn}p_m + \frac{A_{mn}p_m}{\rho(\nu)}$$

Energy density $\rho(\nu)$ increases with temperature T hence $A_{mn}p_m/\rho(\nu)$ vanishes for $T\to\infty$ leaving

$$B_{nm}p_n = B_{mn}p_m \tag{1}$$

Einstein reasoned that equation (1) is true in general based on the assumption that the factors involved do not depend on T. By substitution in the absorption term we can now eliminate $B_{nm}p_n$ and obtain

$$B_{mn}\rho(\nu)p_m \exp\left(\frac{\varepsilon_m-\varepsilon_n}{kT}\right) = B_{mn}\rho(\nu)p_m + A_{mn}p_m$$
induced emission spontaneous emission
$$\varepsilon_m \to \varepsilon_n \qquad \text{spontaneous emission}$$

Divide both sides by $B_{mn}p_m$.

$$\rho(\nu) \exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) = \rho(\nu) + \frac{A_{mn}}{B_{mn}}$$
absorption
$$\varepsilon_m \to \varepsilon_n \qquad \text{spontaneous emission}$$

$$\varepsilon_m \to \varepsilon_n \qquad \text{spontaneous emission}$$

$$\varepsilon_m \to \varepsilon_n$$

Rearrange terms.

$$\rho(\nu) \exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) - \rho(\nu) = \frac{A_{mn}}{B_{mn}}$$
absorption
$$\varepsilon_m \to \varepsilon_n \text{ spontaneous emission}$$

$$\varepsilon_m \to \varepsilon_n$$
spontaneous emission
$$\varepsilon_m \to \varepsilon_n$$

Factor out $\rho(\nu)$.

$$\rho(\nu) \left(\exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) - 1 \right) = \frac{A_{mn}}{B_{mn}}$$

Solve for energy density $\rho(\nu)$.

$$\rho(\nu) = \frac{A_{mn}}{B_{mn}} \frac{1}{\exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) - 1}$$

We now consider the limit of $\rho(\nu)$ as $\varepsilon_m - \varepsilon_n \to \infty$.

$$\lim_{\varepsilon_m - \varepsilon_n \to \infty} \rho(\nu) = \frac{A_{mn}}{B_{mn}} \exp\left(-\frac{\varepsilon_m - \varepsilon_n}{kT}\right)$$

Then by equivalence with Wien's law (which is accurate for large ν)

$$\rho_{\text{wien}}(\nu) = \frac{2h\nu^3}{c^2} \exp\left(-\frac{h\nu}{kT}\right)$$

we have

$$\frac{A_{mn}}{B_{mn}} = \frac{2h\nu^3}{c^2} \tag{2}$$

and

$$\varepsilon_m - \varepsilon_n = h\nu$$

Then by substitution we obtain Planck's law.

$$\rho(\nu) = \frac{A_{mn}}{B_{mn}} \frac{1}{\exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) - 1}$$
$$= \frac{2h\nu^3}{c^2} \frac{1}{\exp\left(\frac{h\nu}{kT}\right) - 1}$$

Coefficient A_{mn} can be computed from first principles of quantum mechanics. For example, for hydrogen we have

$$A_{21} = \frac{e^{10} m_e}{26244 \pi^5 \varepsilon_0^5 \hbar^6 c^3} = 6.27 \times 10^8 \,\text{second}^{-1}$$

The coefficient for induced emission can be obtained from equation (2).

$$B_{mn} = \frac{c^2}{2h\nu^3} A_{mn}$$

The coefficient for absorption can be computed from equation (1).

$$B_{nm} = \frac{p_m}{p_n} B_{mn}$$

The ratio p_m/p_n is equal to g_m/g_n where g_m is the multiplicity (number of degenerate states) associated with energy level m. Hence p_m/p_n is determined by the atomic species.