

(14.1) Fill in the algebra leading to eqn 14.32.

This is equation (14.32).

$$\hat{H} = \int d^3p \sum_{\lambda=1}^2 E_{\mathbf{p}} \hat{a}_{\mathbf{p}\lambda}^\dagger \hat{a}_{\mathbf{p}\lambda} \quad (14.32)$$

Consider equation (14.31).

$$\hat{A}^\mu(x) = \int \frac{d^3p}{(2\pi)^{\frac{3}{2}}} \frac{1}{(2E_p)^{\frac{1}{2}}} \sum_{\lambda=1}^2 (\epsilon_\lambda^\mu(p) \hat{a}_{p\lambda} \exp(-ip \cdot x) + \epsilon_\lambda^{\mu*}(p) \hat{a}_{p\lambda}^\dagger \exp(ip \cdot x)) \quad (14.31)$$

Let

$$\hat{A}^\mu = \begin{pmatrix} \hat{A}_0 \\ \hat{A}_1 \\ \hat{A}_2 \\ \hat{A}_3 \end{pmatrix}, \quad \epsilon_\lambda^\mu(p) = \begin{pmatrix} \epsilon_{0,p\lambda} \\ \epsilon_{1,p\lambda} \\ \epsilon_{2,p\lambda} \\ \epsilon_{3,p\lambda} \end{pmatrix}$$

Then in vector form we have

$$\begin{pmatrix} \hat{A}_0 \\ \hat{A}_1 \\ \hat{A}_2 \\ \hat{A}_3 \end{pmatrix} = \int \frac{d^3p}{(2\pi)^{\frac{3}{2}}} \frac{\hat{a}}{(2E_p)^{\frac{1}{2}}}$$

where

$$\begin{aligned} \hat{a} = \begin{pmatrix} \epsilon_{0,p1} \\ \epsilon_{1,p1} \\ \epsilon_{2,p1} \\ \epsilon_{3,p1} \end{pmatrix} \hat{a}_{p1} \exp(-ip \cdot x) + \begin{pmatrix} \epsilon_{0,p1}^* \\ \epsilon_{1,p1}^* \\ \epsilon_{2,p1}^* \\ \epsilon_{3,p1}^* \end{pmatrix} \hat{a}_{p1}^\dagger \exp(ip \cdot x) \\ + \begin{pmatrix} \epsilon_{0,p2} \\ \epsilon_{1,p2} \\ \epsilon_{2,p2} \\ \epsilon_{3,p2} \end{pmatrix} \hat{a}_{p2} \exp(-ip \cdot x) + \begin{pmatrix} \epsilon_{0,p2}^* \\ \epsilon_{1,p2}^* \\ \epsilon_{2,p2}^* \\ \epsilon_{3,p2}^* \end{pmatrix} \hat{a}_{p2}^\dagger \exp(ip \cdot x) \end{aligned}$$

Hence

$$\begin{aligned}\hat{A}_\nu &= \int \frac{d^3p}{(2\pi)^{\frac{3}{2}}} \frac{1}{(2E_p)^{\frac{1}{2}}} \\ &\times \left[(\epsilon_{\nu,p1} \hat{a}_{p1} + \epsilon_{\nu,p2} \hat{a}_{p2}) \exp(-ip \cdot x) + (\epsilon_{\nu,p1}^* \hat{a}_{p1}^\dagger + \epsilon_{\nu,p2}^* \hat{a}_{p2}^\dagger) \exp(ip \cdot x) \right]\end{aligned}$$

Recall that

$$p \cdot x = E_p t - p_1 x_1 - p_2 x_2 - p_3 x_3$$

Hence the partial derivative of \hat{A}_ν is

$$\begin{aligned}\partial_\mu \hat{A}_\nu &= \int \frac{d^3p}{(2\pi)^{\frac{3}{2}}} \frac{1}{(2E_p)^{\frac{1}{2}}} \\ &\times \left[-ip_\mu (\epsilon_{\nu,p1} \hat{a}_{p1} + \epsilon_{\nu,p2} \hat{a}_{p2}) \exp(-ip \cdot x) + ip_\mu (\epsilon_{\nu,p1}^* \hat{a}_{p1}^\dagger + \epsilon_{\nu,p2}^* \hat{a}_{p2}^\dagger) \exp(ip \cdot x) \right]\end{aligned}$$

Consider equation (14.28).

$$\mathcal{H} = \frac{1}{2}(\mathbf{E}^2 + \mathbf{B}^2) \quad (14.28)$$

We need to integrate Hamiltonian density \mathcal{H} to obtain \hat{H} , that is,

$$\hat{H} = \int d^3x \mathcal{H} \quad (1)$$

Substitute (14.28) into (1).

$$\hat{H} = \frac{1}{2} \int d^3x (E_1^2 + E_2^2 + E_3^2 + B_1^2 + B_2^2 + B_3^2)$$

From equations (5.28) and (5.29) we have

$$\mathbf{E} = \begin{pmatrix} \partial_0 A_1 - \partial_1 A_0 \\ \partial_0 A_2 - \partial_2 A_0 \\ \partial_0 A_3 - \partial_3 A_0 \end{pmatrix}, \quad \mathbf{B} = \begin{pmatrix} \partial_3 A_2 - \partial_2 A_3 \\ \partial_1 A_3 - \partial_3 A_1 \\ \partial_2 A_1 - \partial_1 A_2 \end{pmatrix}$$

Start with E_1^2 .

$$E_1^2 = (\partial_0 \hat{A}_1)^2 - (\partial_0 \hat{A}_1)(\partial_1 \hat{A}_0) - (\partial_1 \hat{A}_0)(\partial_0 \hat{A}_1) + (\partial_1 \hat{A}_0)^2$$

Let

$$I_1 = \frac{1}{2} \int d^3x (\partial_0 \hat{A}_1)^2$$

Hence

$$\begin{aligned} I_1 &= \frac{1}{2} \int \frac{d^3x d^3p d^3q}{(2\pi)^3} \frac{1}{(2E_p)^{\frac{1}{2}} (2E_q)^{\frac{1}{2}}} \\ &\quad \times \sum_{\lambda=1}^2 \left(-iE_p \epsilon_{\lambda}^1(p) \hat{a}_{p\lambda} \exp(-ip \cdot x) + iE_p \epsilon_{\lambda}^{1*}(p) \hat{a}_{p\lambda}^{\dagger} \exp(ip \cdot x) \right) \\ &\quad \times \sum_{\lambda=1}^2 \left(-iE_q \epsilon_{\lambda}^1(q) \hat{a}_{q\lambda} \exp(-iq \cdot x) + iE_q \epsilon_{\lambda}^{1*}(q) \hat{a}_{q\lambda}^{\dagger} \exp(iq \cdot x) \right) \end{aligned}$$

Using the integral

$$\int d^3x \exp(ip \cdot x) = (2\pi)^3 \exp(iE_p t) \delta^{(3)}(\mathbf{p})$$

we have

$$\begin{aligned} I_1 &= \frac{1}{2} \int \frac{d^3p d^3q}{(2E_p)^{\frac{1}{2}} (2E_q)^{\frac{1}{2}}} \\ &\quad \times \delta^{(3)}(\mathbf{p}) \sum_{\lambda=1}^2 \left(-iE_p \epsilon_{\lambda}^1(p) \hat{a}_{p\lambda} \exp(-iE_p t) + iE_p \epsilon_{\lambda}^{1*}(p) \hat{a}_{p\lambda}^{\dagger} \exp(iE_p t) \right) \\ &\quad \times \delta^{(3)}(\mathbf{q}) \sum_{\lambda=1}^2 \left(-iE_q \epsilon_{\lambda}^1(q) \hat{a}_{q\lambda} \exp(-iE_q t) + iE_q \epsilon_{\lambda}^{1*}(q) \hat{a}_{q\lambda}^{\dagger} \exp(iE_q t) \right) \end{aligned}$$