

1 Introduction

We consider two-dimensional correlated Brownian motion with absorbing boundaries:

$$X(t) = x_0 + \mu_x t + \sigma_x W_x(t) \quad a_x < X(t) < b_x \quad (1)$$

$$Y(t) = y_0 + \mu_y t + \sigma_y W_y(t) \quad a_y < Y(t) < b_y \quad (2)$$

where W_i are standard Brownian motions with $\text{Cov}(W_1(t), W_2(t)) = \rho t$ for $0 < t' \leq t$. In particular, we find the joint transition density function for $(X(t), Y(t))$ under the boundary conditions:

$$p(X(t) = x, Y(t) = y | a_x < X(t') < b_x, a_y < Y(t') < b_y, 0 < t' \leq t, X(0) = x_0, Y(0) = y_0). \quad (3)$$

This density, which we shorten to $p(x, y, t)$ from now on, is the solution to the Fokker-Planck equation [?]:

$$\frac{\partial}{\partial t} p(x, y, t') = -\mu_x \frac{\partial}{\partial x} p(x, y, t') - \mu_y \frac{\partial}{\partial y} p(x, y, t') + \frac{1}{2} \sigma_x^2 \frac{\partial^2}{\partial x^2} p(x, y, t') + \rho \sigma_x \sigma_y \frac{\partial^2}{\partial x \partial y} p(x, y, t') + \frac{1}{2} \sigma_y^2 \frac{\partial^2}{\partial y^2} p(x, y, t'), \quad (4)$$

$$p(a_x, y, t') = p(b_x, y, t') = p(x, a_y, t') = p(x, b_y, t') = 0 \quad 0 < t' \leq t. \quad (5)$$

Differentiating $p(x, y, t)$ with respect to the boundaries produces the transition density of a particle beginning and ending at the points $(X_1(0), X_2(0))$ and $(X_1(t), X_2(t))$ respectively, while attaining the minima a_x/a_y and maxima b_x/b_y in each coordinate direction:

$$\frac{\partial^4}{\partial a_x \partial b_x \partial a_y \partial b_y} p(x, y, t) =$$

$$p\left(X(t) = x, Y(t) = y \mid \min_{t'} X(t') = a_x, \max_{t'} X(t') = b_x, \min_{t'} Y(t') = a_y, \max_{t'} Y(t') = b_y, 0 < t' \leq t, X(0) = x_0, Y(0) = y_0\right). \quad (6)$$

The transition density for the considered system has been used in computing first passage times [??], with application to structural models in credit risk and default correlations [??]. [?] use variants of the differentiated solutions with respect to some of the boundaries to price financial derivative instruments whose payoff depends on observed maxima/minima.

Closed-form solutions to (4) - (5) are available for some parameter regimes. When $\rho = 0$, the transition density of the process can be obtained with a Fourier expansion. When $a_1 = -\infty$ and $b_1 = \infty$, the method of images can be used to enforce the remaining boundaries. For either $a_1, a_2 = -\infty$ or $b_1, b_2 = \infty$, the Fokker-Planck equation is a Sturm-Liouville problem in radial coordinates. Both of these techniques are used by ?. However, to the best of our knowledge, there is no closed-form solution to the general problem in (4) - (5).

It is still possible to approach the general problem by proposing a Fourier expansions. However, a drawback of this out-of-the-box solution is that the system matrix for the corresponding eigenvalue problem is large and dense. An alternative is to use a finite difference scheme. However, discretization of the initial condition introduces a numerical bias in the estimation procedure.

In this paper, we propose a solution to the general problem (4) - (5) which is obtained by combining a small-time analytic solution with a finite-element method.

ADD OUT APPLICATION (ESTIMATION)

2 Numerical Method

2.1 Fourier Expansion

The formal Fourier expansion for the problem is

$$f(x, y | t) = \lim_{K \rightarrow \infty} \sum_{k=1}^K c_k(t) \sin(2\pi k x / L).$$

Plugging into the PDE yields an eigenvalue problem which is slow to solve.

2.2 Finite Difference

2.3 Finite Element

References