3_1



Introduction

Johannes Kepler: Motion of planets governed by three laws:

- 1. Each planet moves in an elliptical orbit, with the Sun at one focus of the ellipse. ("Astronomia Nova", 1609)
- 2. A line from the Sun to a given planet sweeps out equal areas in equal times. ("Astronomia Nova", 1609)
- 3. The square of the orbital periods of the planets is proportional to the cube of the major axes. ("Harmonice Mundi", 1619)

Isaac Newton ("Principia", 1687): Kepler's laws are consequence of gravitational interaction between planets and the Sun, and the gravitational force is

$$F_1 = -\frac{Gm_1m_2}{r_{12}^2} \frac{r_{21}}{r_{12}} \tag{3.1}$$

where F_1 is the gravitational force exerted on object 1, m_1 , m_2 are the masses of the interacting objects, r their distance, and r_{21}/r_{12} the unit vector joining the objects, $r_{21} = r_2 - r_1$, $r_{12} = -r_{21}$ and $r_{12} = |r_{12}| = |r_{21}|$.

Planetary Dynamics

universität innsbruck

Kepler's 1st Law

Planet

*
Focus 2 Sun

Kepler's 1st Law: The orbits of the planets are ellipses and the Sun is at one focus of the ellipse.

For the planets of the solar system, the ellipses are almost circular, for comets they can be very eccentric.

Kepler's Laws



3-4

Minor Axis: 2b

Kepler's 1st Law

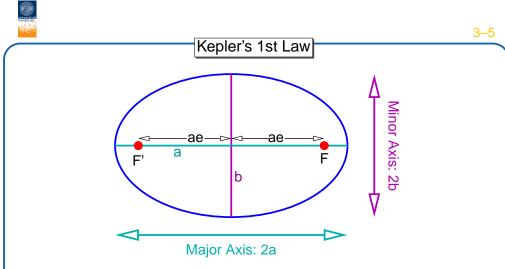
Major Axis: 2a

Definition: Ellipse = Sum of distances r, r' from any point on ellipse to two fixed points (foci, singular: focus), F, F', is constant:

$$r + r' = 2a \tag{3.2}$$

where \boldsymbol{a} is called the semi-major axis of the ellipse.

Kepler's Laws 2 Kepler's Laws 3



Definition: Eccentricity e: ratio between distance from centre of ellipse to focal point and semi-major axis.

So circles have e = 0.

F' 2ae π-θ F

Kepler's 1st Law

Major Axis: 2a

Law of cosines: $r'^2 = r^2 + (2ae)^2 - 2 \cdot r \cdot 2ae \cdot \cos(\pi - \theta)$

use r + r' = 2a and solve for r to find the polar coordinate form of the ellipse:

$$r = \frac{a(1 - e^2)}{1 + e\cos\theta}$$

(3.3)

5

3-8

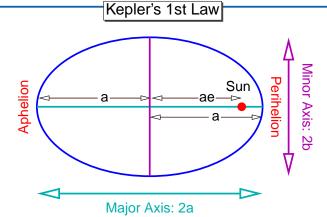
3-6

Check this for yourself! θ is called the *true anomaly*.

Kepler's Laws



3–7



Finally, we need the closest and farthest point from a focus:

closest point :
$$d_{\text{perihelion}} = a - ae = a(1 - e)$$

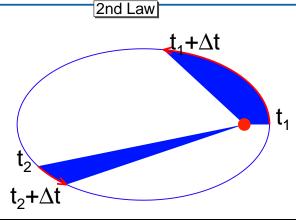
farthest point : $d_{\text{aphelion}} = a + ae = a(1 + e)$ (3.4)

for stars: periastron and apastron,

for satellites circling the Earth: perigee and apogee.

Kepler's Laws





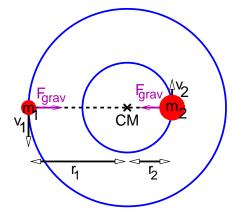
Kepler's 2nd Law: The radius vector to a planet sweeps out equal areas in equal intervals of time.

- 1. Kepler's 2nd Law is also called the law of areas.
- 2. perihelion: planet nearest to Sun ⇒ planet is fastest
- 3. aphelion: planet farthest from Sun ⇒ planet is slowest

3rd Law

tional to the cubes of the semimajor axes, a, of their orbits: $P^2 \propto a^3$.

Kepler's 3rd Law: The squares of the periods of the planets, P, are propor-



Calculating the motion of two bodies of mass m_1 and m_2 gives Newton's form of Kepler's third law:

$$P^2 = \frac{4\pi^2}{G(m_1 + m_2)} R^3 \tag{3.5}$$

where $r_1 + r_2 = R$ (for elliptical orbits: R is the semi-major axis).

3–9

3rd Law

Newton's form of Kepler's 3rd law is the most general form of the law.

However, often shortcuts are possible.

Assume one central body dominates, $m_1 = M \gg m_2$:

$$\frac{P^2}{a^3} = \frac{4\pi^2}{GM} = \text{const.} = k$$
 (3.9)

So, if we know P and a for one body moving around m_1 , can calculate k.

For the Solar System, use Earth:

- P_{\oplus} = 1 year (by definition!)
- $a_{\oplus} =$ 1 AU (Astronomical Unit, 1 AU = 149.6 \times 10⁶ km)

$$\Longrightarrow k= 1\,\mathrm{yr^2\,AU^{-3}}$$

Kepler's Laws

Jupiter: $a_{1} = 5.2 \,\text{AU}$. What is its period?

Answer: $P_{7\!\!+}^2=$ 1 yr² AU $^{-3}\cdot$ 5.2³ AU $^3\sim$ 140 yr², or $P_{7\!\!+}\sim$ 12 years

(with pocket calculator: P_{7} = 11.86 years)

Kepler's Laws



3-11

3-12

5

3-10

Tidal forces: The Earth–Moon system

Gravitational acceleration for center of Earth (0) and for the point closest to a gravitating body (1):

$$a_0 = \frac{GM}{r^2}$$
 and $a_1 = \frac{GM}{(r - R_{\oplus})^2}$ (3.10)

such that difference in acceleration:

$$\Delta a = a_1 - a_0 = \frac{GM}{(r - R_{\oplus})^2} - \frac{GM}{r^2} = \frac{GM}{r^2} \left(\frac{1}{(1 - \frac{R_{\oplus}}{r})^2} - 1 \right)$$

$$\sim \frac{GM}{r^2} \left(1 + 2\frac{R_{\oplus}}{r} - 1 \right) = \frac{2GMR_{\oplus}}{r^3}$$
 (3.11)

Therefore the tides due to the Moon are

$$\Delta a_{\mathbb{C}} = \frac{2GM_{\mathbb{C}}R_{\oplus}}{r_{\mathbb{C}}^{3}} \tag{3.12}$$

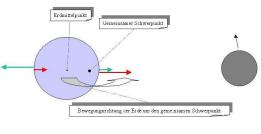
and the tides due to the Sun are

$$\Delta a_{\odot} = \frac{2GM_{\odot}R_{\oplus}}{(1.\text{AU})^3} \tag{3.13}$$

But since the mass of the Moon is $M_{\rm C} \sim M_{\oplus}/{\rm 81}$ and the $r_{\rm C} \sim {\rm 60} \times R_{\oplus}$

- $\Longrightarrow \Delta a_{\odot}/\Delta a_{\odot} = 0.46$
- ⇒ Moon twice as important as Sun; two sets of tidal bulges, spring-tides

Tidal forces: The Earth–Moon system







(Einsteins-Erben)

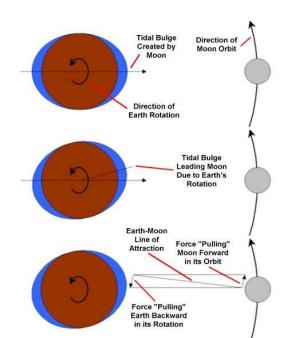
gravitational force of the Moon onto Earth

centrifugal force of Earth around centre of gravity

Tidal forces

Tidal forces

2





Stability of Satellites

Tidal forces also important for the stability of small bodies (e.g., moons) moving around a central body at a distance r.

Notation: mass $M_{\rm D,S}$, radius $R_{\rm D,S}$, density $\rho_{\rm D,S}$ where p: planet, s: satellite.

Satellite is bound by its own gravitational force. Satellite is not ripped apart, if binding force $F_{\rm G}$ > tidal force $F_{\rm T}$:

Binding force is \sim mutual attraction of two halfs of the satellite separated by R_s :

$$F_{\rm G} = \frac{GM_{\rm s}M_{\rm s}}{4R^2} > \frac{GM_{\rm p}M_{\rm s}R_{\rm s}}{r^3} = F_{\rm T}$$
 (3.14)

where binding force \sim gravitational attractions of two hemispheres of the satellite onto each other.

$$\implies \frac{GM_{\rm s}M_{\rm s}}{4R_{\rm s}^2} > \frac{GM_{\rm p}M_{\rm s}R_{\rm s}}{r^3} \implies \frac{M_{\rm s}}{4R_{\rm s}^3} > \frac{M_{\rm p}}{r^3} \quad \text{or} \quad \frac{\rho_{\rm s}}{4} > \frac{R^3\rho_{\rm p}}{r^3} \tag{3.15}$$

This means there is a critical distance (Roche, 1850):

$$\frac{r_{\rm c}}{R} = \sqrt[3]{\frac{4\rho_{\rm p}}{\rho_{\rm s}}} \qquad \frac{r_{\rm c}}{R} = 2.44\sqrt[3]{\frac{\rho_{\rm p}}{\rho_{\rm s}}}$$
 (3.16)

⇒ If of same density, the moon has to be more distant than 2.44 planetary radii from the planet in order to avoid disruption. This is important, e.g., for the formation of rings.



3-14

Bound Rotation

- Tidal bulges preceed rotation by 2.5 h due to friction (ocean ground)
- Rotation of Earth slows down ⇒ length of day increases by

1.6 ms century⁻¹

confirmed by historic solar eclipse observations (e.g., Thales, 585 B.C., eclipse observations by Chinese astronomers)

- 370 Mio. years ago (Devon): 1 Year~400 days (from coral growth)
- Moon is accelerated, its distance increases by 4 cm yr⁻¹
- length of the month grows
- friction ceases when 1 day = 1 month (synodic) = 50 days (now)
- equilibrium, i.e., bound rotation, is reached in $\sim 10^{12}$ years
- Moon's rotation bound to orbital motion by tidal friction on the liquid interior of the Moon soon after formation
- Other moons are also synchronized: e.g. Galilean moons of Jupiter
- Pluto/Charon both synchronized

Tidal forces



3 - 15

3–16

Precession and Nutation

Earth is \sim rotational ellipsoid, orbits of Sun and Moon are $\it not$ in plane of equator

(Earth's axis has tilt of ~23.5°, Moon's orbit tilted by 7° against ecliptic)

 \Longrightarrow Sun and Moon excert torques onto Earth

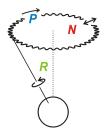
Earth's rotational axis is not stable in space.

Two major effects:

luni-solar precession: Earth's axis rotates around pole of ecliptic once every 25800 years (\sim 50" per year).

Already discovered by Hipparcos in \sim 200 BC!

nutation: "Wobble" with \sim 18 year periodicity caused by short-term perturbations caused by Moon and Sun.



Tidal forces 5 Precession and Nutation 1

N-Body Problem

3 - 173 - 18N-Body Problem

periodic perturbations: Terms containing time in \sin - and \cos -functions.

secular perturbations: Long term changes which depend on time (usually as a

Analytical approach is very important for understanding the underlying physics,

⇒ New high precision calculations are all based on numerical simulations, i.e.,

Today's standard: DE102, DE405, DE414 from Jet Propulsion Laboratory, Pasadena, and INPOP06 from

but mathematically very tedious. Series do not converge on long timescales

Perturbation theory yields two kinds of perturbations:

direct solution of equation of motion on computers.

Apart from Sun, motion of planets also influenced by forces between planets:

Total equation of motion for the *i*-th object:

$$m_i \ddot{r}_i = -\sum_{k=1}^{N} \frac{Gm_i m_k}{r_{i,k}^2} \frac{r_i - r_k}{r_{i,k}}$$
 (3.17)

 \Longrightarrow 3N differential equations of 2nd order, requiring 6N integrations for their solution.

Closed solution only possible for 10 of these (6: from motion of center of mass, 3: conservation of angular momentum, 1: conservation of energy).

Analytic solution: "Perturbation theory":

- 1. Assume two body motion around Sun for all planets
- 2. Evaluate force based on this motion.
- 3. Update positions with this "perturbation".
- 4. Iterate (i.e., goto step 2)

N-Body Problem



3 - 19

polynomial).

(1000's of years).

Laskar et al., IMCCE, Observatoire de Paris.

N-Body Problem

Long-Term Evolution of the Solar System

eccentricity 0. Time (Gyr) Numerical simulations allow to obtain good information about behavior of solar system for timescales of a few 10 million years around the present \Longrightarrow Important, e.g., for paleoclimatology.

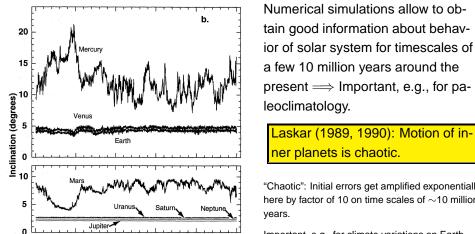
Laskar (1989, 1990): Motion of inner planets is chaotic.

"Chaotic": Initial errors get amplified exponentially, here by factor of 10 on time scales of \sim 10 million

Important, e.g., for climate variations on Earth ("Milankovitch cycles").

Also found with different methods by Wisdom and Suskind.

Long-Term Evolution of the Solar System



Time (Gyr)

Laskar, J. 1990, Icarus, 88, 266

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3 - 19

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(Laskar 1994)

3 N-Body Problem N-Body Problem

(Laskar 1994)

Laskar, J. 1989, Nature, 338, 237

Laskar, J. 1994, A&A, 287, L9