

# LXe scintillation model

January 27, 2020

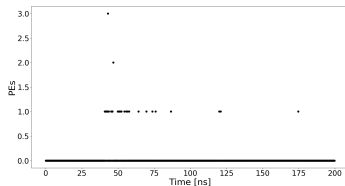
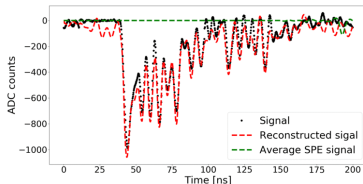
# Objective

We want to use the data collected by the Cs and Co runs to estimate the parameters for the scintillation model.

In the following we show that there is a  $\delta(t)$  component in the model with high statistical significance and propose how to go on to claim a discovery.

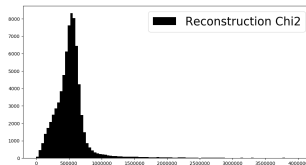
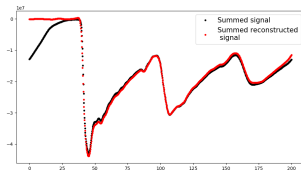
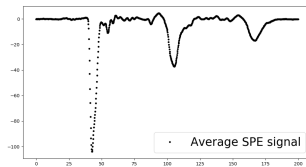
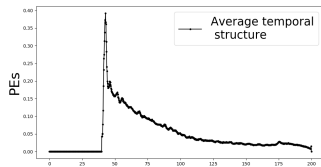
# Signal Reconstruction

To study the temporal structure of the photon emission a processing algorithm uses a template of the average SPE signal to reconstruct the temporal PE pattern in each event for each PMT separately. The output of this process is the number of PEs created in the PMT in each digitization point.



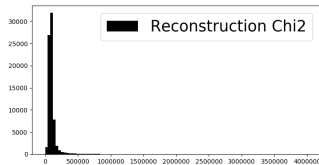
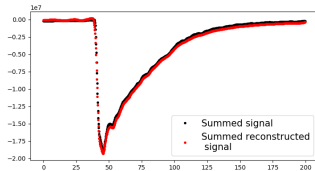
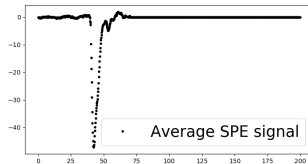
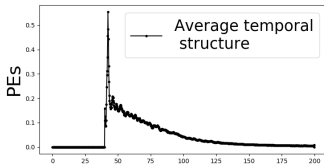
# Reconstruction results

Co - PMT 0 - 76000 events



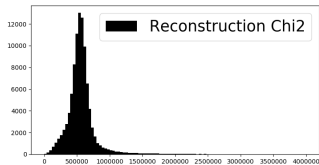
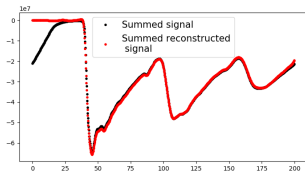
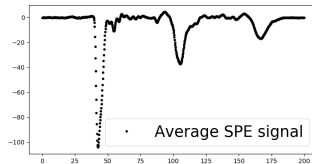
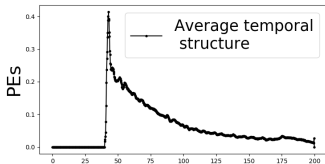
# Reconstruction results

Co57 - PMT 1 - 75000 events

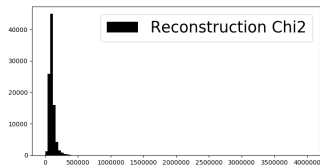
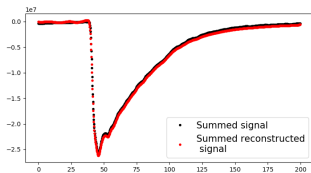
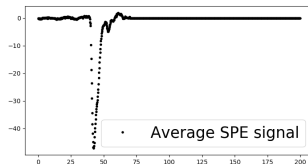
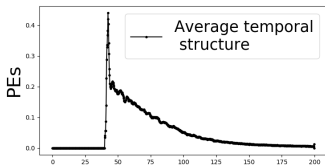


# Reconstruction results

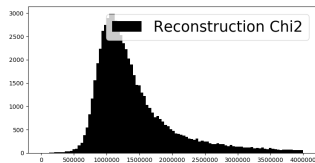
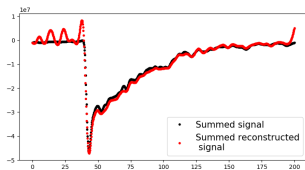
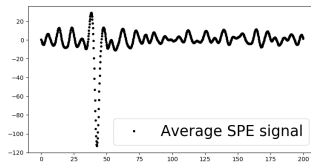
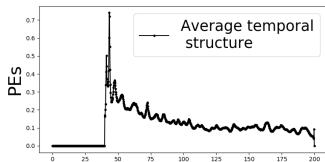
Cs137 - PMT 0 - 97000 events



# Cs137 - PMT 1 - 96000 events

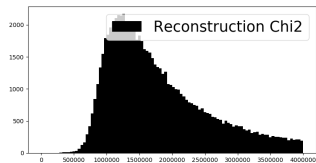
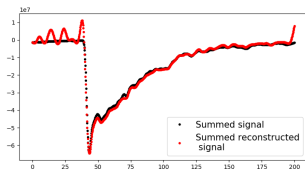
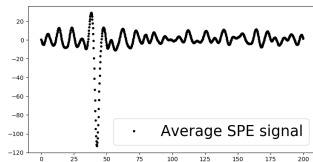
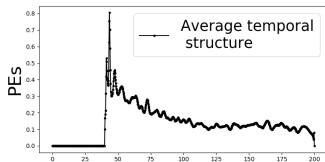


# Co57 - PMT 3 - 71000 events



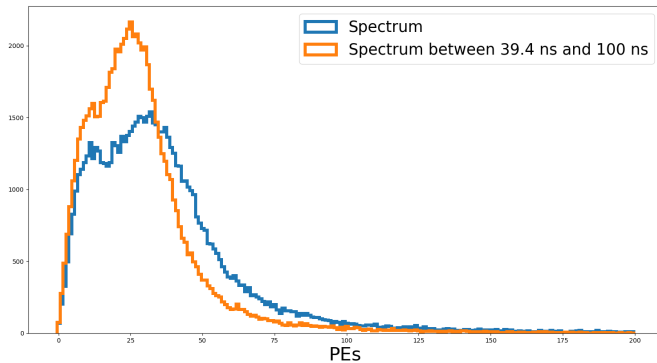


# Cs137 - PMT 3 - 86000 events

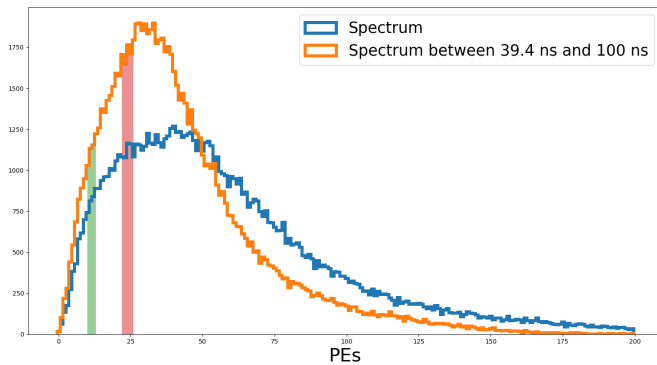


# 1 PMT Spectrum

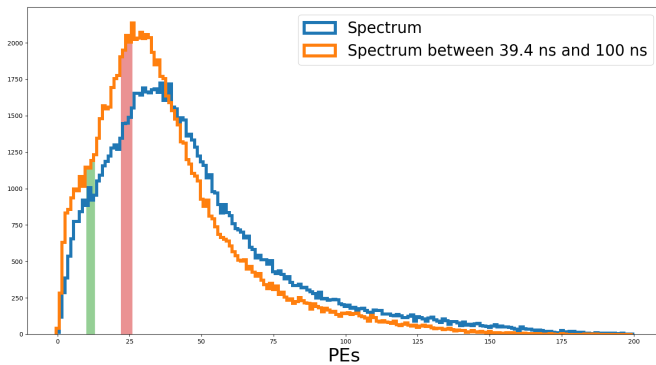
Co57 - PMT 0 - 76000 events



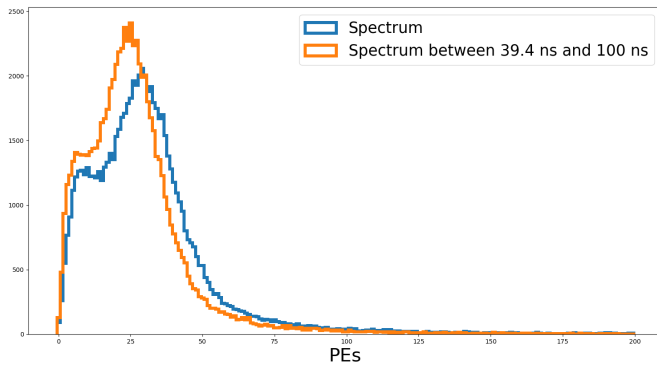
Cs137 - PMT 0 - 97000 events



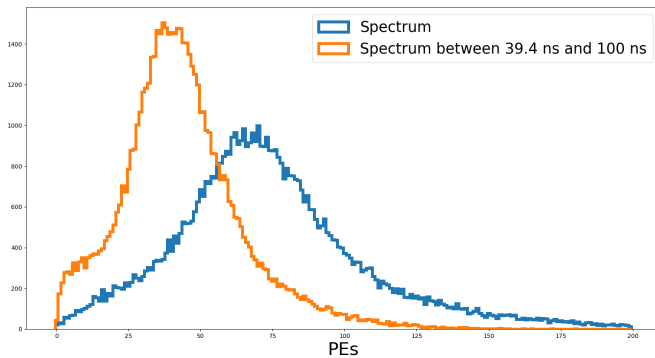
Cs137 - PMT 1 - 96000 events



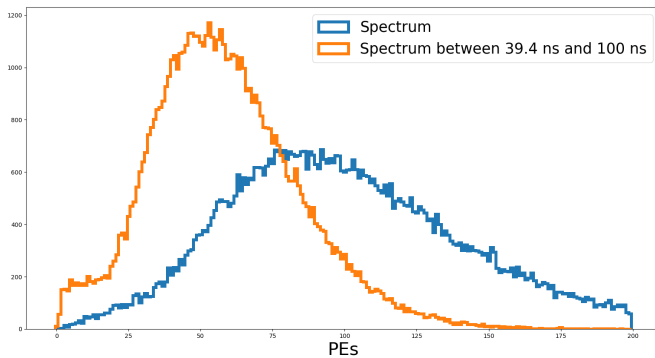
Co57 - PMT 1 - 75000 events



Co57 - PMT 3 - 71000 events

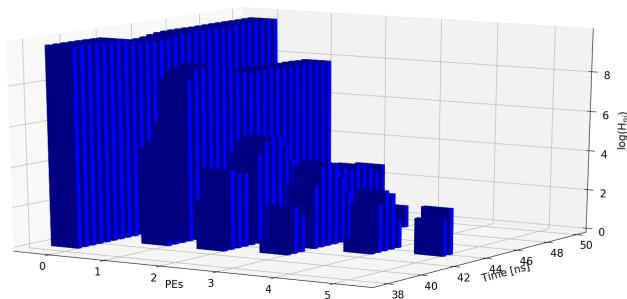


Cs137 - PMT 3 - 86000 events



# PE Histogram

The basis for analysis is a 2D histogram (for each PMT and each source)  $H_{ni}$  which holds the number of events in which  $n$  PEs were resolved at time  $i$ ,



The average temporal structure which is shown in the Reconstruction result slide is the average on each column.



# Model for $H_{ni}$

The number of photons emitted at time  $i$ ,

$$n_i^{ph} \sim \text{Poisson}(Y_i N), \quad (1)$$

where  $N$  is the average number of photon emitted in the events, and  $Y_i$  is the scintillation model (the probability to emit a photon at time  $i$ ).

The number of PEs created in the PMT at time  $i$ ,

$$n_i^{pe} \sim \text{Binom}(Q, \text{Poisson}(Y_i N)), \quad (2)$$

where  $Q$  is the photon detection efficiency (quantum efficiency, collection efficiency and double PE probability).

Each PMT has its temporal uncertainty combined with the code's temporal uncertainty. The number of PEs that need to be resolved at time  $i$ ,

$$\tilde{n}_i \sim \sum_j m_j^i, \quad (3)$$

where  $m_j^i$  is a random variable that represents the number of PEs that were created at time  $j$  but resolved at time  $i$ .

$$m_j^i \sim n_j^{pe} \text{Norm}(i - j - T | \sigma_t) \quad (4)$$

where  $\sigma_t$  is the temporal resolution of the PMT and the code and  $T$  is the average delay of the PMT.

Since  $\tilde{n}_i$  is a sum of independent random variables its distribution can be approximated by a distribution with mean and variance  $\langle \tilde{n}_i \rangle = \sum_j \langle m \rangle_j^i$ ,  $\text{Var}_i = \sum_j \text{Var}_j^i$ , where,

$$\langle m \rangle_j^i = \text{Var}_j = Q N_0 Y_j \text{Norm}(i - j - T | \sigma_t) \quad (5)$$

Poisson has an equal mean and variance, so

$$\tilde{n}_i \sim \text{Poisson} \left( Q N \sum_j Y_j \text{Norm}(i - j - T | \sigma_t) \right) \quad (6)$$

# Scintillation model

$Y_j$  is the probability to emit a photon at time  $t_j\Delta t$ .

$$Y_j = \Delta t F \delta(t_j - T) + \Delta t \frac{1 - F}{\tau_s} e^{-t_j/\tau_s} \quad (7)$$

$$\begin{aligned} \langle \tilde{n}_i \rangle &= \frac{QN\Delta t}{\sqrt{2\pi}\sigma_t} \int_0^\infty Y(\tilde{t}) e^{-(\tilde{t}-t_i-T)^2/2\sigma_t^2} = \\ NQ\Delta t &\left[ \frac{e^{-(t_i-T)^2/2\sigma_t^2}}{\sqrt{2\pi}\sigma_t} + \frac{K}{\tau} e^{-t/\tau} \left( 1 - \operatorname{erf} \left( \frac{\sigma}{\sqrt{2}\tau} - \frac{t-T}{\sqrt{2}\sigma} \right) \right) \right] \end{aligned} \quad (8)$$

Where  $K$  is a normalization constant,

$$K = \left[ 1 - \operatorname{erf} \left( \frac{\sigma}{\sqrt{2}\tau} + \frac{T}{\sqrt{2}\sigma} \right) + e^{-\sigma^2/2\tau^2 - T/\tau} \left( 1 + \operatorname{erf} \left( \frac{T}{\sqrt{2}\sigma} \right) \right) \right]^{-1} \quad (9)$$

# SPE resolution

Finally, the code has its resolution for SPEs (its accuracy to resolve the correct number of PEs from the signal at a given time). The probability to resolve  $n$  PEs given  $\tilde{n}$  is the correct number,

$$n_i \sim P_n(\tilde{n}|\sigma_{pe}) = \sum_m \text{Norm}(m-1|\sigma_{pe}) P_{n-m}(\tilde{n}-1|\sigma_{pe}) \quad (10)$$

## Expectation for $H_{ni}$

The expected number of events in which  $n$  PEs were resolved at time  $i$  is,

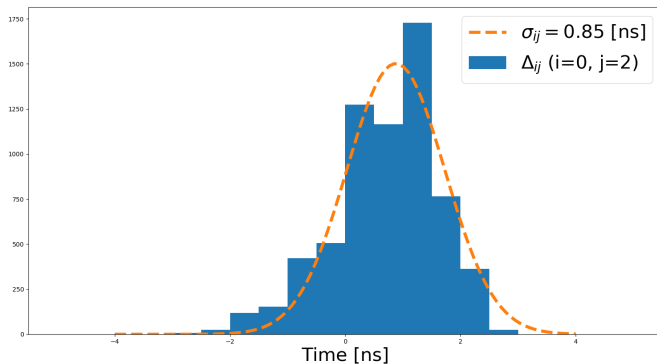
$$h_{ni} = N_{ev} \sum_m P_n(m|\sigma_{pe}) \text{Poisson}(m|\langle \tilde{n}_i \rangle) \quad (11)$$

This is a model with 7 parameters:  $N, Q, F, \tau_s, T, \sigma_t, \sigma_{PE}$ . Some of them can be constrained by calibration data.

# Calibration $\sigma_t$

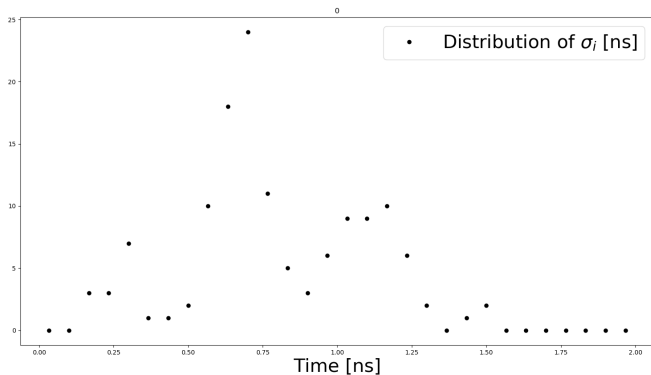
We can calibrate  $\sigma_t$  for each PMT by the pulser data. Consider  $\text{PMT}_i$  that we want to calibrate and a any other  $\text{PMT}_j$ . For each pulser event in which both PMTs saw SPE we measure  $\Delta_{ij}$ , the time difference between the two signals.

$$\Delta_{ij} \sim \text{Norm}(T_{ij}, \sigma_{ij}).$$



$$\sigma_i = \sqrt{\frac{1}{2}(\sigma_{ij}^2 + \sigma_{ik}^2 - \sigma_{jk}^2)} \text{ for any } j \neq k.$$

So a pair of PMTs ( $k, j$ ) is used to estimate the  $\sigma_i$ . Each PMT <sub>$i$</sub>  has  $\frac{1}{2}\binom{19}{2}$  estimates for  $\sigma_i$  and their mean is  $\sigma_t$  for that PMT.

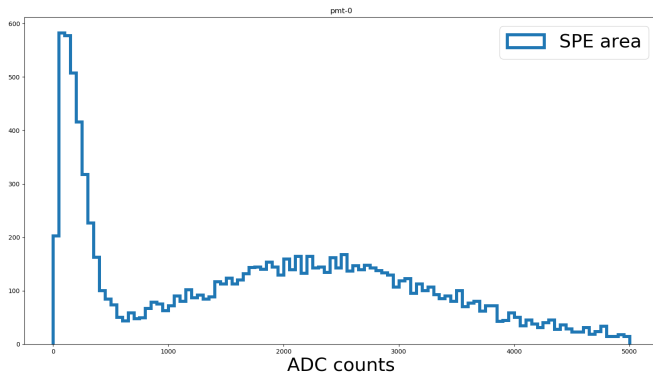




# Calibration of $\sigma_{PE}$

$\sigma_{PE}$  is the standard deviation of the distribution of the number of PEs that are resolved from a SPE created in the PMT. This is calibrated from the SPE area distribution.

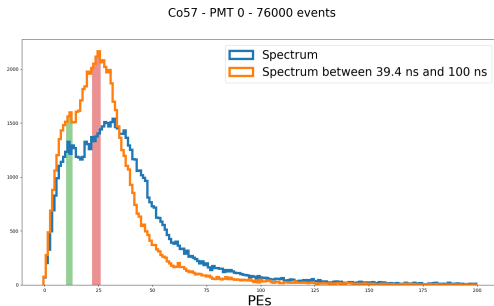
$$\text{Area}_{spe} \sim \text{Norm}(\bar{A}, \bar{A} \cdot \sigma_{PE}) \quad \rightarrow \quad n \sim \text{Norm}(1, \sigma_{PE}) \quad (12)$$



# Constrain on $NQ$

We choose the data to fit our model to from a narrow PE bands in the spectrum, where it can be estimated that the origin of the events in these bands is from the same energy and their variation is due to statistical fluctuation and no physical. Thus we expect that

$$N_{PE} \sim \text{Poisson}\left(\sum_{39.4ns}^{100ns} \langle h_{ni} \rangle_n\right) \quad (13)$$



# Fit

To find the parameters which give  $h_{ni}$  as close to  $H_{ni}$  we maximize

$$L(\text{data}, \text{model}(\theta)) = \sum_{\text{dataset}} \sum_{j \in J} \text{Poisson}(\text{data}_j | \text{model}_j) / \text{len}(J) \quad (14)$$

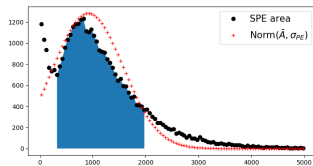
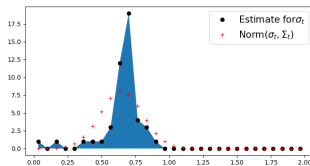
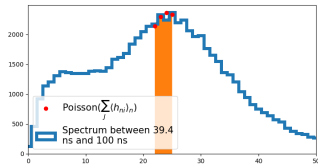
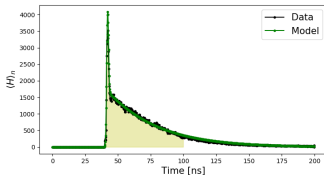
Where  $\theta$  is the parameter array and  $J$  is some sub-range of the dataset. The datasets in the summation:

- ▶ Dataset -  $H_{ni}$ , model -  $h_{ni}$ .
- ▶ Dataset - distribution of estimates for  $\sigma_t$ , model -  $\text{Norm}(\sigma_t, \Sigma_t)$ .
- ▶ Dataset - distribution of SPE areas, model -  $\text{Norm}(\bar{A}, \bar{A} \cdot \sigma_{PE})$ .
- ▶ Dataset - PE spectrum, model -  $\text{Poisson}(\sum_{39.4ns}^{100ns} \langle h_{ni} \rangle_n)$ .

This adds two dummy parameters,  $\Sigma_t$  is the standard deviation of the estimates of  $\sigma_t$  and  $\bar{A}$  is the mean SPE area.

# Results ( $^{57}\text{Co}$ peak 2)

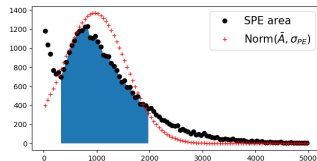
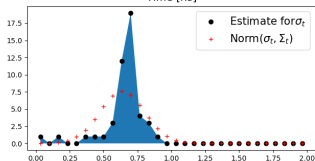
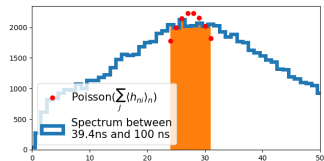
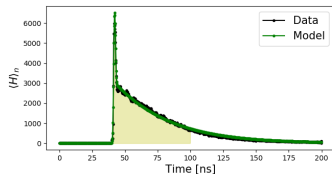
PMT 1



$NQ$	$F$	$\sigma_t[ns]$	$\sigma_t[ns]$	$\sigma_{PE}$
33	0.1	34	0.6	0.7

# Results ( $^{137}\text{Cs}$ )

PMT 1

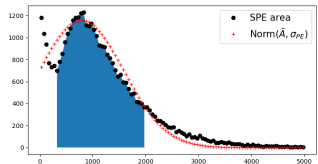
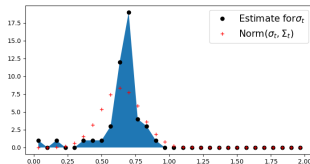
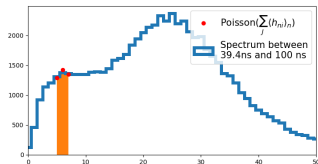
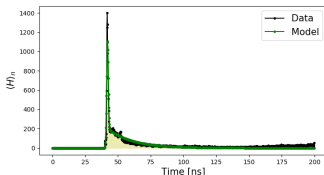


$NQ$	$F$	$\sigma_t[ns]$	$\sigma_t[ns]$	$\sigma_{PE}$
39	0.09	38	0.6	0.6

# Results ( $^{57}\text{Co}$ peak 1)

Notice that for this peak  $F$  is much greater and  $\tau_s$  is much smaller.

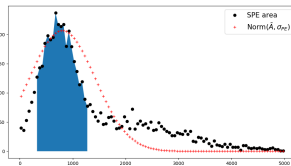
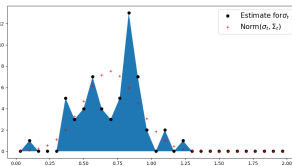
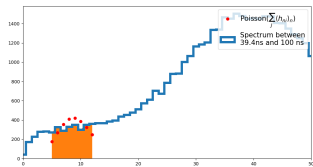
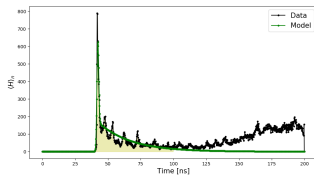
PMT 1



$NQ$	$F$	$\sigma_t[ns]$	$\sigma_t[ns]$	$\sigma_{PE}$
7	0.4	18	0.6	1

# Results ( $^{57}\text{Co}$ peak 1 - strange result)

PMT 3



# Parameter estimation

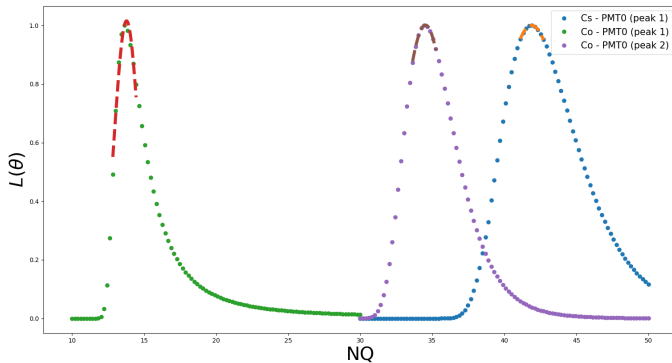
$\hat{\theta}$  is the parameter array that maximize  $L$ , so for any parameter  $\theta_i$ ,

$$\begin{aligned}\partial_{\theta_i} L = 0 \quad \rightarrow \quad L(\hat{\theta} + \Delta\theta) &= L_{max} \left( 1 + \frac{\partial_{\theta_i}^2 L}{2L_{max}} \Delta\theta^2 \right) \quad \rightarrow \\ L(\hat{\theta} + \Delta\theta) &\approx L_{max} e^{-\Delta\theta^2 / 2\sigma_{\theta}^2}\end{aligned}\tag{15}$$



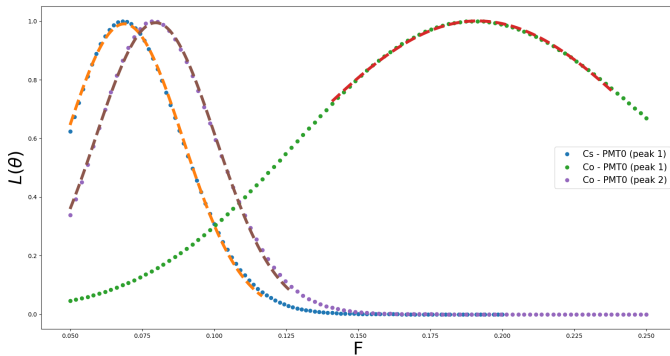
For each parameter  $L$  was maximized while holding the parameter fixed (for a range of parameters, for each PMT individually)

# NQ

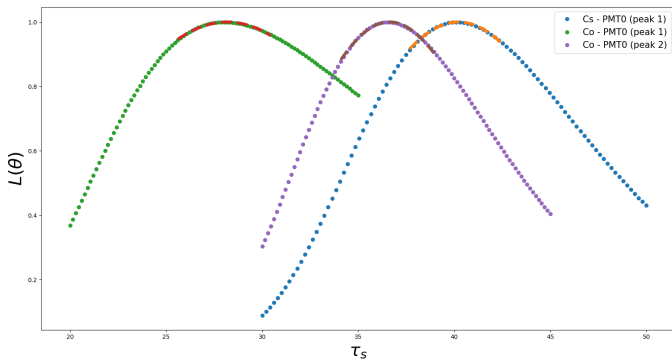


Source	$\hat{N}Q$	$\sigma_{NQ}$
Cs	42	2.5
Co (peak 1)	14	0.8
Co (peak 2)	34	1.7

F



Source	$\hat{F}$	$\sigma_F$	significance [ $\sigma$ ]
Cs	0.07	0.02	3
Co (peak 1)	0.2	0.06	3
Co (peak 2)	0.08	0.02	4

$\tau_s$ 

Source	$\hat{\tau}_s [ns]$	$\sigma_{\tau_s} [ns]$
Cs	40	6
Co (peak 1)	28	8
Co (peak 2)	37	5

# How to go on

- ▶ Preform this analysis for all PMTs individually for both sources and both orientations (source to the lab / to the door).
- ▶ Study the low energy component in the Co spectrum and if it has a different scintillation regime (NR?).
- ▶ Maximize  $L$  with the constrain that  $N, F, \tau_s$  is the same for all PMTs (but different between sources). This will brake the degeneracy of  $NQ$ .
- ▶ Maximize  $L$  with only  $N, \tau_s$  fixed for all PMTs. This will show the anisotropy in the fast component.
- ▶ Maximize  $L$  with  $N, \tau_s$  fixed and constrain a 90 deg between the symmetry axis of the anisotropy of  $F$  in the two orientations.
- ▶ Maximize  $L$  with the above with the constrain that  $F$  is isotropic in the BG data.