# Computational Quantum Mechanics Prof. Gerson J. Ferreira INFIS/UFU 2020/1

Matrix representations for H

## General expansion

 $\varphi_1(x), \varphi_2(x), \varphi_3(x), \dots$ 

## General expansion

Consider a basis (plane-waves, polynomials, Hermite, ...):

$$\begin{array}{ll} \varphi_1(x), \varphi_2(x), \varphi_3(x), \dots & \text{Assuming orthonormal:} \\ |\varphi_1\rangle\,, |\varphi_2\rangle\,, |\varphi_3\rangle\,, \dots & \langle \varphi_i |\varphi_j\rangle = \delta_{i,j} \end{array}$$

Let's write the matrix elements as  $H_{j,i} = \langle arphi_j | H | arphi_i 
angle$ 

We want to solve  $H \ket{\psi} = E \ket{\psi}$  expanding on the basis  $\ket{\psi} = \sum_i c_i \ket{\varphi_i} = c_1 \ket{\varphi_1} + c_2 \ket{\varphi_2} + \dots$ 

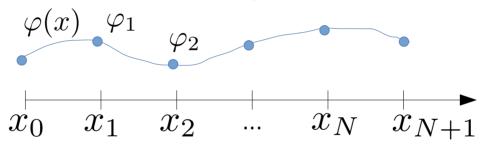
$$\sum_{i} c_{i} H |\varphi_{i}\rangle = \sum_{i} c_{i} E |\varphi_{i}\rangle$$

$$\sum_{i} H_{j,i} c_{i} = E c_{i}$$

$$\sum_{i} c_{i} \langle \varphi_{j} | H |\varphi_{i}\rangle = \sum_{i} c_{i} E \langle \varphi_{j} | \varphi_{i}\rangle$$

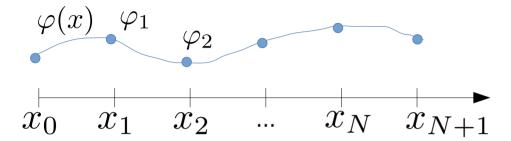
$$\bar{H}\vec{c} = E\vec{c}$$

# Finite differences



#### Finite differences: the second derivative

 $\rightarrow$  equivalent to a basis of discrete sites:  $\varphi_1 = \varphi(x_1), \varphi_2 = \varphi(x_2), ...$ 



We have seen that the second derivative can be written as

$$\varphi_n'' = \frac{\varphi_{n-1} - 2\varphi_n + \varphi_{n+1}}{\Delta x^2}$$

Apply rule for n=1 .. N (inner points)

It takes a matrix form:

$$\varphi_1'' = \frac{\varphi_0 - 2\varphi_1 + \varphi_2}{\Delta x^2}$$

$$\varphi_2'' = \frac{\varphi_1 - 2\varphi_2 + \varphi_3}{\Delta x^2}$$

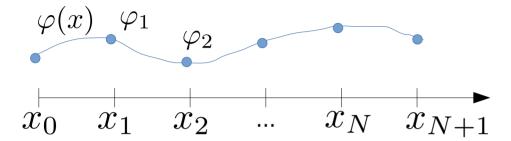
$$\varphi_3'' = \frac{\varphi_2 - 2\varphi_3 + \varphi_4}{\Delta x^2}$$

$$\frac{\partial^2}{\partial x^2} \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \vdots \\ \varphi_N \end{pmatrix} = \frac{1}{\Delta x^2} \begin{pmatrix} -2 & 1 & 0 & 0 \\ 1 & -2 & 1 & 0 \\ 0 & \ddots & \ddots & \ddots \\ 0 & 0 & 1 & -2 \end{pmatrix} \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \vdots \\ \varphi_N \end{pmatrix} + \frac{1}{\Delta x^2} \begin{pmatrix} \varphi_0 \\ 0 \\ \vdots \\ \varphi_{N+1} \end{pmatrix}$$

vanishing boundary conditions:  $\varphi_0 = \varphi_{N+1} = 0$ 

## Finite differences: the potential term

 $\rightarrow$  equivalent to a basis of discrete sites:  $\varphi_1 = \varphi(x_1), \varphi_2 = \varphi(x_2), \dots$ 



Now let's consider the potential term as

$$V(x)\varphi(x) \to V_i\varphi_i$$

Apply rule for n=1 .. N (inner points)

$$\begin{aligned} V(x)\varphi(x)|_{x=x_1} &= V_1\varphi_1 \\ V(x)\varphi(x)|_{x=x_2} &= V_2\varphi_2 \\ V(x)\varphi(x)|_{x=x_3} &= V_3\varphi_3 \end{aligned} \qquad V(x)\varphi(x) \rightarrow \begin{pmatrix} V_1 & 0 & 0 & 0 & 0 \\ 0 & V_2 & 0 & 0 & 0 \\ 0 & 0 & \ddots & 0 & 0 \\ 0 & 0 & 0 & V_{N-1} & 0 \\ 0 & 0 & 0 & 0 & V_N \end{pmatrix} \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \vdots \\ \varphi_{N-1} \\ \varphi_N \end{pmatrix}$$

It takes a matrix form:

$$\begin{pmatrix} V_1 & 0 & 0 & 0 & 0 \\ 0 & V_2 & 0 & 0 & 0 \\ 0 & 0 & \ddots & 0 & 0 \\ 0 & 0 & 0 & V_{N-1} & 0 \\ 0 & 0 & 0 & 0 & V_N \end{pmatrix} \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \vdots \\ \varphi_{N-1} \\ \varphi_N \end{pmatrix}$$

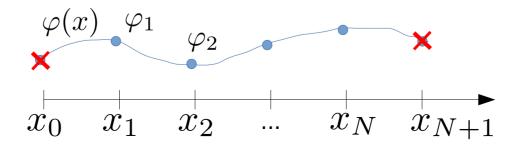
#### Finite differences

Therefore,

$$H\varphi(x) = E\varphi(x)$$
 
$$H = -\frac{1}{2}\frac{\partial^2}{\partial x^2} + V(x)$$

on a discrete lattice becomes

$$\bar{\bar{H}} = -\frac{1}{2} \frac{1}{\Delta x^2} \begin{pmatrix} -2 & 1 & 0 & 0 \\ 1 & -2 & 1 & 0 \\ 0 & \ddots & \ddots & \ddots \\ 0 & 0 & 1 & -2 \end{pmatrix} + \begin{pmatrix} V_1 & 0 & 0 & 0 \\ 0 & V_2 & 0 & 0 \\ 0 & 0 & \ddots & 0 \\ 0 & 0 & 0 & V_N \end{pmatrix}$$



## Plane waves

 $\exp(ikx)$ 

#### Plane waves

In this case, the basis is  $\varphi_k(x) = e^{ikx}$ 

**How to choose the values of k?** There's a relation between k and x resolutions, which is due to the periodicity imposed by the plane-waves

It's easier if we start in x-space  $\rightarrow$  let's assume that V(x+L) = V(x) is periodic

Consequently, we want

$$\varphi_k(x+L) = \varphi_k(x)$$

which gives

$$e^{ikL} = 1 \to k_j = \frac{2\pi}{L}j$$

and j=0..N, but how to choose N?



The reciprocal k-space is also periodic

$$\varphi_{j+N}(x) = \varphi_j(x)$$

$$e^{i\frac{2\pi}{L}Nx_n} = 1 \to x_n = \frac{L}{N}n$$

Plane waves  $\varphi_j(x_n) = e^{ik_jx_n}$ 

#### **Summarizing:**

Periodicity in x-space implies 
$$k_j = j\Delta k$$
  $\Delta k = 2\pi/L$   $j = 0...N$ 

Periodicity in k-space implies 
$$\ \ \vec{x}_n = n\Delta x \quad \ \Delta x = L/N \quad \ n = 0...N$$

Notice that the resolution in k and x spaces are related: 
$$\Delta x \Delta k = \frac{2\pi}{N}$$

For a fix N, if you increase the resolution in one space, you loose in the other → very **similar** to Heisenberg's uncertainty principle

#### Implementation?

- → go back to the General expansion slide
- → calculate the matrix elements of each k either analytically or numerically if needed.
- $\rightarrow$  for many introductory problems it is possible to calculate  $H_{ij}$  analytically, and use the analytical expression to build the matrices in Python and diagonalize.