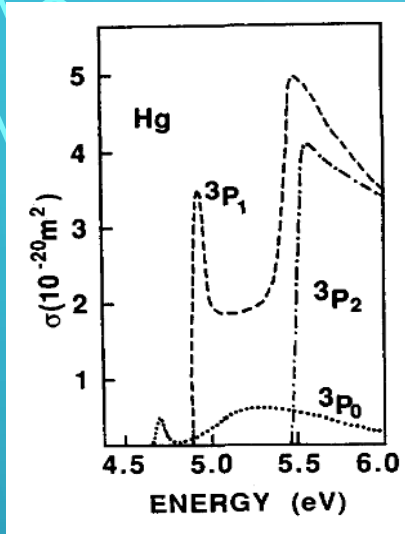


A decorative graphic on the left side of the slide, consisting of a network of white lines and small circles on a blue gradient background, resembling a circuit board or a stylized tree structure.

# THE FRANCK-HERTZ EXPERIMENT

KITTY HARRIS AND JOSH ELSARBOUKH

# ELECTRON ENERGY AND SCATTERING CROSS-SECTION



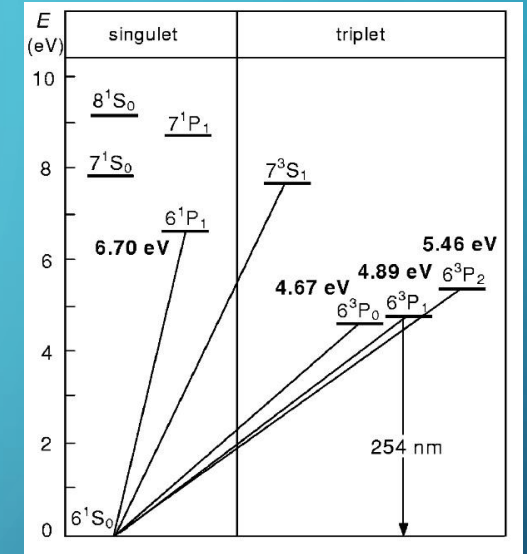
Cross-section of inelastic collisions vs. electron energy. Uncertainty of 30% due to measurement imprecisions. Haken & Wolf, p 698

The Franck-Hertz experiment, presented in 1914 by James Franck and Gustav Hertz, involves accelerating electrons through mercury (also done with neon) vapor via a potential grid, towards a beam current collecting anode. As the accelerated electrons gain sufficient energy for excitation, the cross-section of inelastic collisions increases. This can be seen in the diagram to the left, which shows roughly the cross section of inelastic electron-mercury collisions as a function of electron energy. The points of highest cross-section correspond approximately to the three lowest energy states at 4.67, 4.89, and 5.46 volts, respectively.

The cross-section can best be understood as an effective 'area' inside of which a scattering event can occur. It is not only related to accelerating electron energy, but also the mean free path of the electrons, as well as temperature.

$$\sigma = \frac{k_B T}{P \lambda}$$

(Where  $P$  is the pressure, which is a function of the temperature  $T$ , and  $\lambda$  is the mean free path, which is dependent on electron energy.)



Lowest energy states of mercury. Haken & Wolf, p 697

# CURRENT MINIMA AT ENERGY INTERVALS

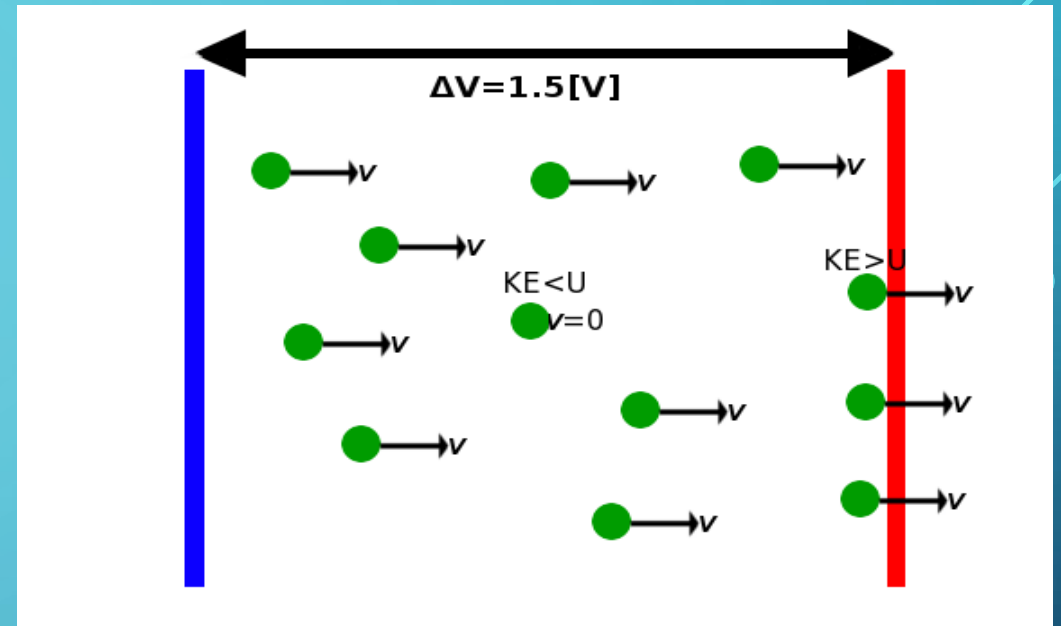
- To reach the collection plate, electrons must overcome a potential difference.

- Only electrons with sufficient energy can overcome this difference:

$$\frac{m_e}{2} v^2 = KE \geq U = q\Delta V$$

So electrons which have lost too much energy in collisions can't make it across the 1.5[V] potential.

- Current reflects this:  $I = neAv$  increases as the number of electrons (that reach the collection plate, where the current is being read) increases.
- We see a smooth curve because velocity (and therefore energy) is a probabilistic distribution, not the same for each electron, and because not all electrons will collide.



# SETUP

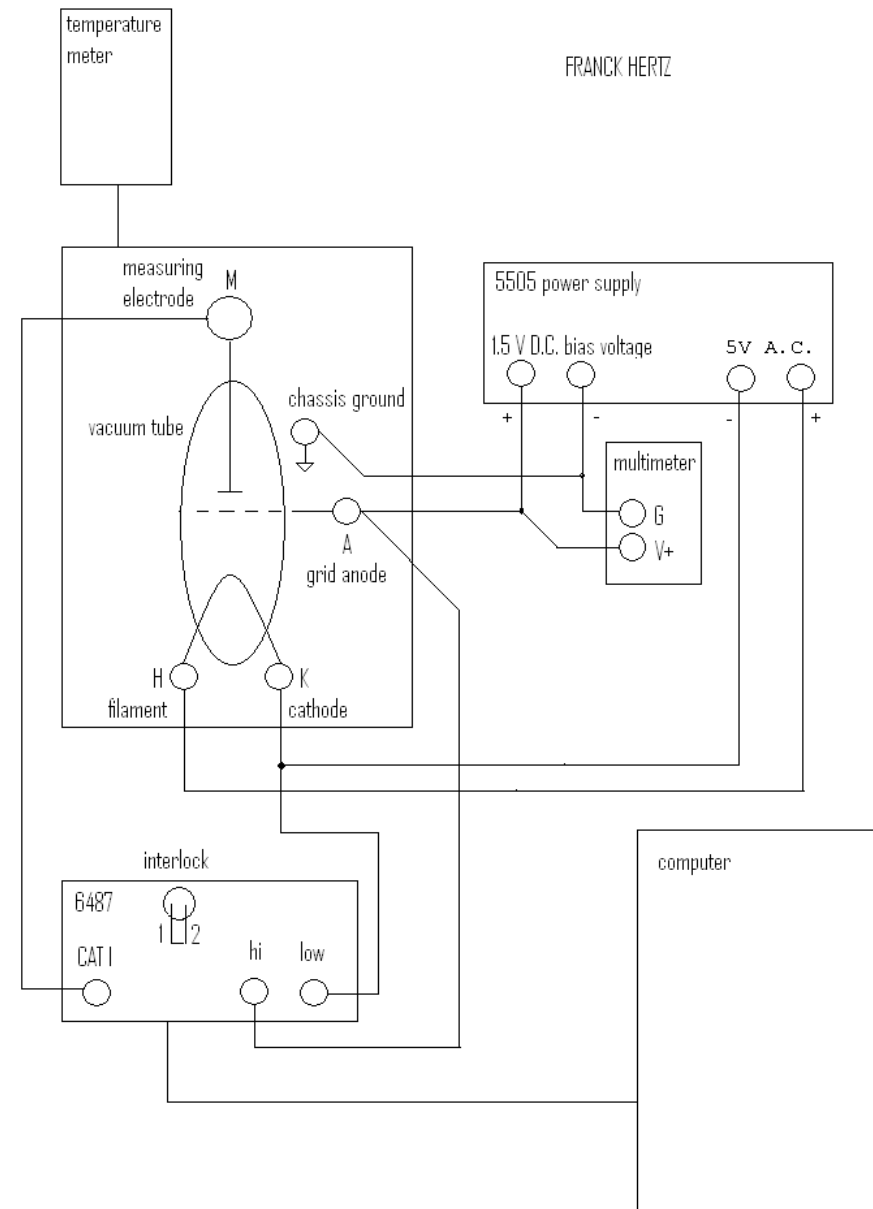
We will be focusing on the Keithley 6487 and the portion inside the quartz tube.

Diagrams uploaded by Dr. Tagg to

<https://sites.google.com/site/experimentalphysicsdecathlon/home/04-fundamental-quantum-behavior/franck-hertz-experiment>.

Diagram to the right is titled “FRANCK HERTZ.bmp”.

Diagram in the next slide is pulled from the safety instructions in “device.pdf”.



# INSIDE THE QUARTZ TUBE

## Quartz Tube

- Electromagnetically Non-Conductive
- Transparent
- Common: Makes it relatively cheap compared to alternatives
- Average thermal conductivity prevents it from having a major contribution to temperature.
- Vacuum prevents vapor from escaping, prevents other gases from entering, and keeps pressure low.

## Mercury Vapor

- Boiling Point: 629.88[K] – much greater than the box.
  - However, pressure is very low ( $\sim 10^{-3}$ [atm]), so vapor is easy to maintain.
- Excited states are distinct.

## Tungsten Filament (Cathode)

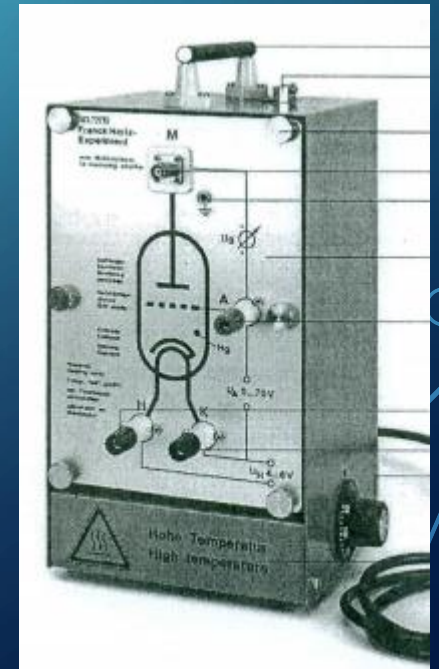
- Electrons are boiled off – high resistance > excess heat > high electron energy

## Grid (Anode)

- This is where we sweep potential relative to the cathode.

## Collection Plate

- Swept to stay at a 1.5[V] difference from the anode, reducing KE of incoming electrons.
- This is where current is measured.





# KEITHLEY 6487



- Dual picoammeter and variable voltage source
- Ideal for low-, dark-, and beam-current measurements
- Fine resolution, low noise
- I/O trigger modes for automated production

- Ammeter Resolution: 10[fA]
- Sweeping Voltage:
  - 200 microvolts to 500 volts
  - 200 steps per second
  - 0.2 millivolt resolution
- Burden Voltage: <200 microvolts

Range	5½ Digit Default Resolution	Accuracy (1 Year) <sup>1</sup> ±(% rdg. + offset) 18°–28°C, 0–70% RH	Typical RMS Noise <sup>2</sup>	Typical Analog Rise Time (10% to 90%) <sup>3</sup> Damping <sup>4</sup>	
				Off	On
2 nA	10 fA	0.3 % + 400 fA	20 fA	4 ms	80 ms
20 nA	100 fA	0.2 % + 1 pA	20 fA	4 ms	80 ms
200 nA	1 pA	0.15% + 10 pA	1 pA	300 μs	1 ms
2 μA	10 pA	0.15% + 100 pA	1 pA	300 μs	1 ms
20 μA	100 pA	0.1 % + 1 nA	100 pA	110 μs	110 μs
200 μA	1 nA	0.1 % + 10 nA	100 pA	110 μs	110 μs
2 mA	10 nA	0.1 % + 100 nA	10 nA	110 μs	110 μs
20 mA	100 nA	0.1 % + 1 μA	10 nA	110 μs	110 μs

# TEMPERATURE-DEPENDENCE OF $\Delta E$

Temperature [K]	433	447	462	476
Energy ( $E_a$ ) [eV]	4.85	5.03	4.99	4.91

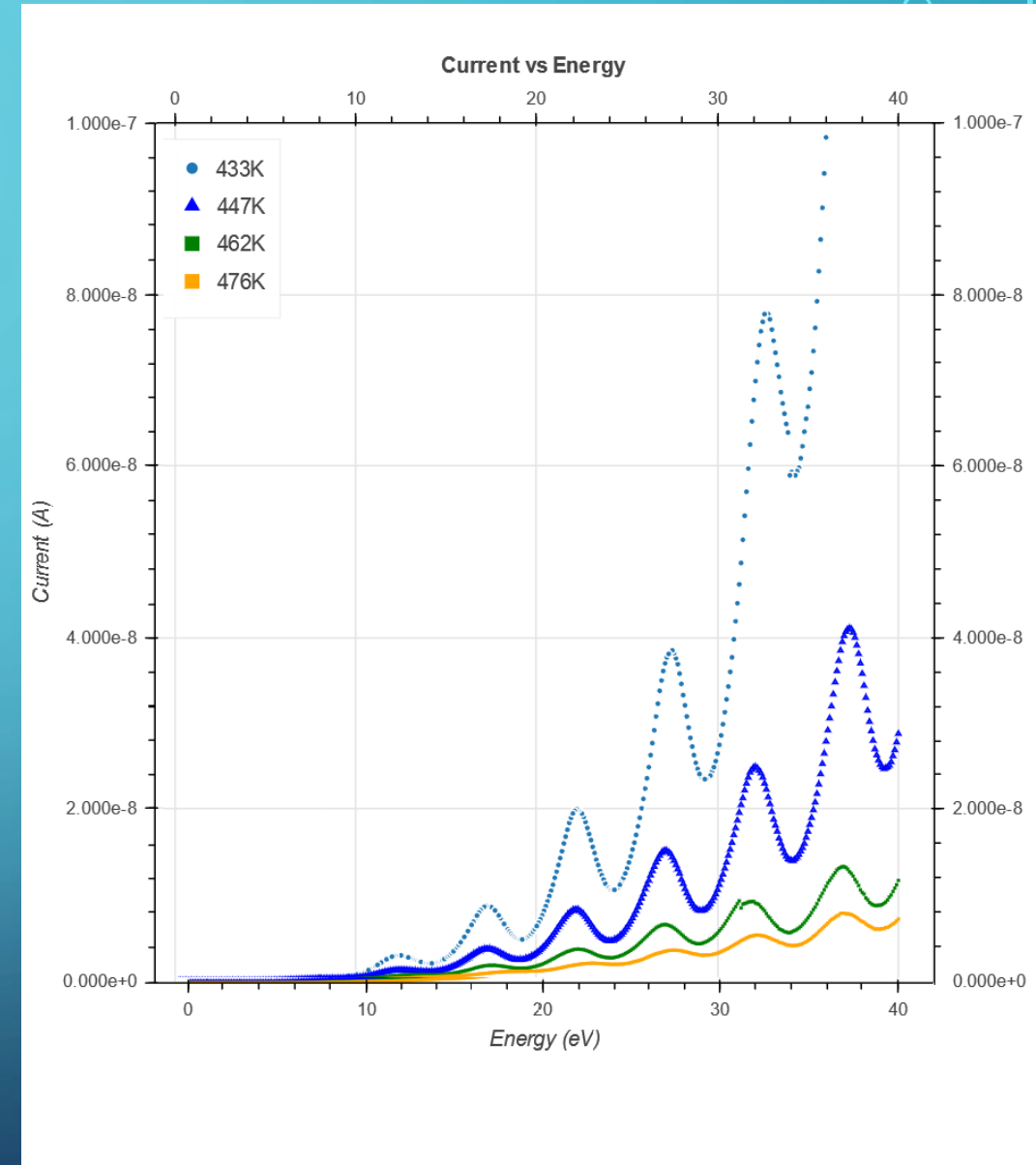
Uncertainties  
are  $\pm 0.1$  [eV]

Since our carriers are electrons, voltage [V] and energy [eV] measurements are identical.

Minima were determined by taking curve fits to slices of data.

Then the distances between these minima were plotted.

$$\Delta E = \left[ 1 + \frac{\lambda}{L} (2n - 1) \right] E_a$$
$$\sigma = \frac{k_B T}{p \lambda}$$
$$E_a = \Delta E_{(n=\frac{1}{2})}$$



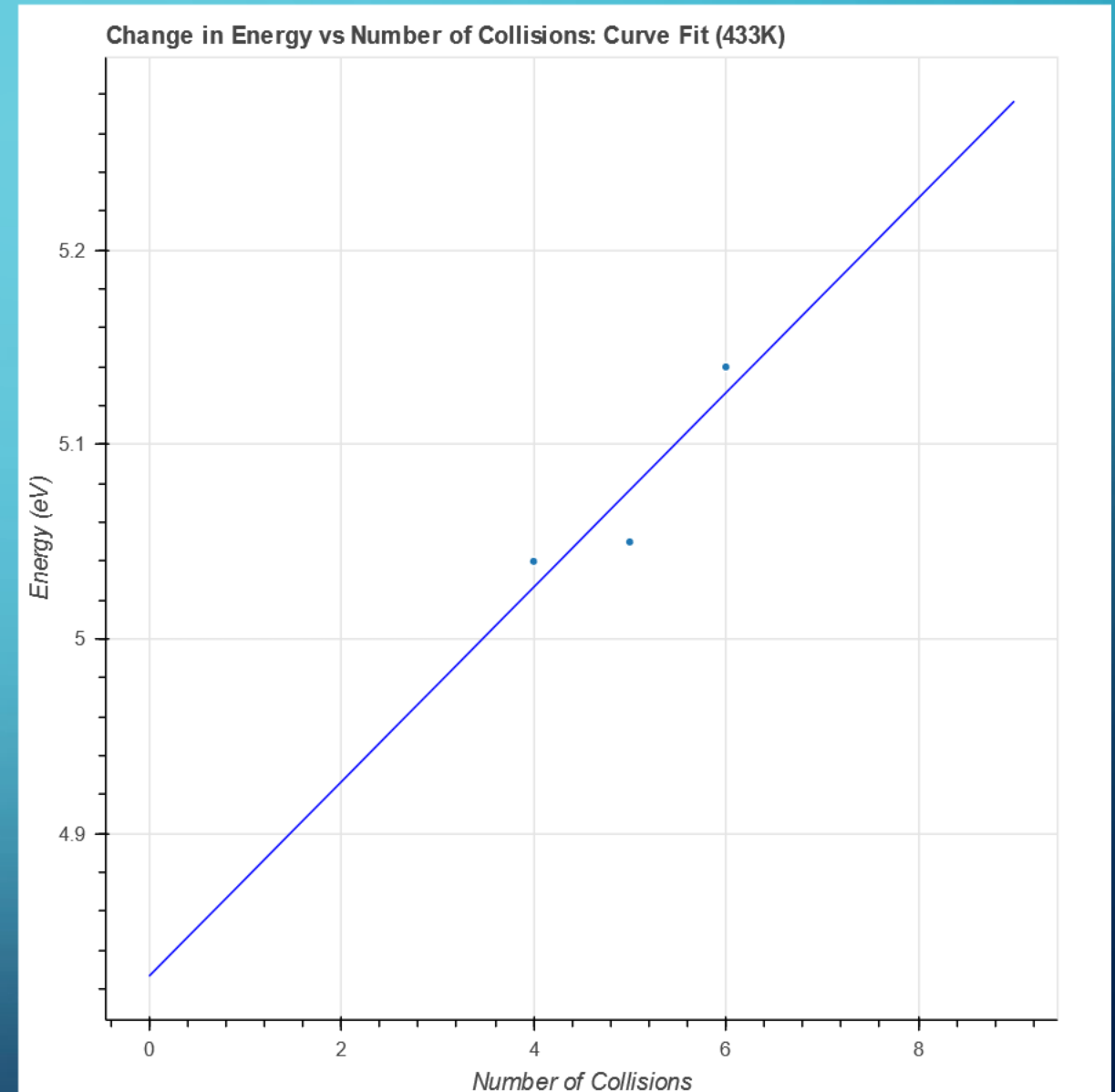
# FINDING THE CROSS-SECTION FOR COLLISIONS

- Cross-section is sensitive to movement of particles, therefore varying with temperature:

$$\sigma = \frac{k_B T}{8.7 * 10^9 \frac{3110[K]}{T} [Pa] \lambda}$$

- The mean free path is taken from the slope of the function:  $\frac{d\Delta E}{dn} = \frac{2\lambda}{L} E_a$

Temperature [K]:	433	447	462	476
$\lambda$ [ $10^{-6}\text{m}$ ]:	41.2	1.59	6.41	12.2
Uncertainty:	9.24	8.90	8.96	9.12
$\sigma$ [ $10^{-18}\text{m}^2$ ]:	2.54	40.5	6.16	2.11
Uncertainty [ $\text{m}^2$ ]:	0.569	22.7	8.62	1.58





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