

# Quantum Wasserstein distance based on an optimization over separable states

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G. Tóth<sup>1,2,3,4,5</sup>



and J. Pitrik<sup>5,6,7</sup>



<sup>1</sup>University of the Basque Country UPV/EHU, Bilbao, Spain

<sup>2</sup>EHU Quantum Center, University of the Basque Country UPV/EHU, Spain

<sup>3</sup>Donostia International Physics Center (DIPC), San Sebastián, Spain

<sup>4</sup>IKERBASQUE, Basque Foundation for Science, Bilbao, Spain

<sup>5</sup>Wigner Research Centre for Physics, Budapest, Hungary

<sup>6</sup>Alfréd Rényi Institute of Mathematics, Budapest, Hungary

<sup>7</sup>Department of Analysis, Budapest University of  
Technology and Economics, Budapest, Hungary

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# Outline

## 1 Motivation

- Connecting Wasserstein distance to entanglement theory?

## 2 Background

- Quantum Wasserstein distance
- Quantum Fisher information

## 3 Wasserstein distance and separable states

- Quantum Wasserstein distance based on an optimization over separable states
- Relation to entanglement conditions

# Motivation

- Many distance measures are maximal for orthogonal states.
- Recently, the Wasserstein distance appeared, which is different and this makes it very useful.
- For the quantum case, surprisingly, the self-distance can be nonzero.
- Can we connect these to entanglement theory?

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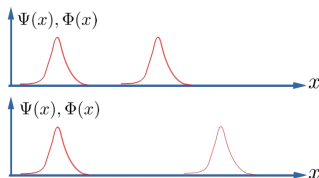
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## Wasserstein distance and separable states

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# An important property of the Wasserstein distance

- Many distance measures are maximal for orthogonal states, e.g., for the following state-pairs.



- In the second example, the two states are further apart from each other, based on common sense.
- Wasserstein distance can recognize this since it is related to the "cost of moving sand from a distribution to the other one."
- It can be used for machine learning.

# Quantum Wasserstein distance

- **Definition.**—The square of the distance between two quantum states described by the density matrices  $\varrho$  and  $\sigma$  is

$$D_{\text{DPT}}(\varrho, \sigma)^2 = \frac{1}{2} \min_{\varrho_{12}} \sum_{n=1}^N \text{Tr}[(H_n^T \otimes \mathbb{1} - \mathbb{1} \otimes H_n)^2 \varrho_{12}],$$

s. t.

$$\begin{aligned}\varrho_{12} &\in \mathcal{D}, \\ \text{Tr}_2(\varrho_{12}) &= \varrho^T, \\ \text{Tr}_1(\varrho_{12}) &= \sigma,\end{aligned}$$

where  $\mathcal{D}$  is the set of density matrices, and  $H_n$  are Hermitian matrices.

- Note the relation to the representability problem.

G. De Palma and D. Trevisan, Quantum optimal transport with quantum channels, *Ann. Henri Poincaré* 22, 3199 (2021).

# Self-distance can be nonzero (unlike in the classical case)

- The self-distance of a state is

$$D_{\text{DPT}}(\varrho, \varrho)^2 = \sum_{n=1}^N I_{\varrho}(H_n),$$

where the Wigner-Yanase skew information is defined as

$$I_{\varrho}(H) = \text{Tr}(H^2 \varrho) - \text{Tr}(H \sqrt{\varrho} H \sqrt{\varrho}).$$

- This connects connects Wasserstein distance and quantum metrology.

G. De Palma and D. Trevisan, Quantum optimal transport with quantum channels, *Ann. Henri Poincaré* 22, 3199 (2021).



# Wasserstein distance between a pure state

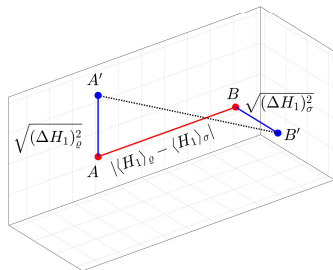
$\varrho = |\Psi\rangle\langle\Psi|$  and a mixed state  $\sigma$

- The distance is given as

$$\begin{aligned} D_{\text{DPT}}(\varrho, \sigma)^2 \\ = \frac{1}{2} \sum_{n=1}^N [(\Delta H_n)^2_{\varrho} + (\Delta H_n)^2_{\sigma} + (\langle H_n \rangle_{\varrho} - \langle H_n \rangle_{\sigma})^2], \end{aligned}$$

see the following figure, where  $(\Delta H_n)^2$  is the variance.

# Wasserstein distance between a pure state $\varrho = |\psi\rangle\langle\psi|$ and a mixed state $\sigma$ II



- $N = 1$  with operator  $H_1$ .
- The quantum Wasserstein distance equals  $1/2$  times the usual Euclidean distance between  $A'$  and  $B'$ .

GT and J. Pitrik, Quantum Wasserstein distance based on an optimization over separable states, arXiv:2209.09925.

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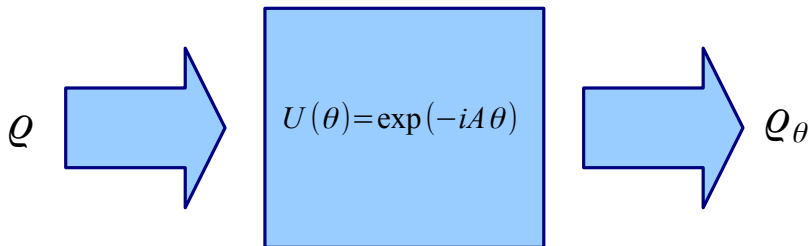
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# Quantum metrology

- Fundamental task in metrology



- We have to estimate  $\theta$  in the dynamics

$$U = \exp(-iA\theta).$$

# The quantum Fisher information

- Cramér-Rao bound on the precision of parameter estimation

$$(\Delta\theta)^2 \geq \frac{1}{mF_Q[\varrho, A]},$$

where  $F_Q[\varrho, A]$  is the **quantum Fisher information**, and  $m$  is the number of independent repetitions.

- The quantum Fisher information is

$$F_Q[\varrho, A] = 2 \sum_{k,l} \frac{(\lambda_k - \lambda_l)^2}{\lambda_k + \lambda_l} |\langle k|A|l \rangle|^2,$$

where  $\varrho = \sum_k \lambda_k |k\rangle\langle k|$ .

# Formula based on convex roofs

The quantum Fisher information is the convex roof of the variance times four

$$F_Q[\varrho, A] = 4 \min_{\{p_k, |\psi_k\rangle\}} \sum_k p_k (\Delta A)_{\psi_k}^2,$$

where

$$\varrho = \sum_k p_k |\psi_k\rangle\langle\psi_k|.$$

GT, D. Petz, Phys. Rev. A 87, 032324 (2013); S. Yu, arXiv1302.5311 (2013);

GT, I. Apellaniz, J. Phys. A: Math. Theor. 47, 424006 (2014).

# Formula based on concave roofs

The variance is the concave roof of itself

$$(\Delta A)^2_{\varrho} = \max_{\{p_k, |\psi_k\rangle\}} \sum_k p_k (\Delta A)^2_{\psi_k},$$

where

$$\varrho = \sum_k p_k |\psi_k\rangle\langle\psi_k|.$$

GT, D. Petz, Phys. Rev. A 87, 032324 (2013);

GT, I. Apellaniz, J. Phys. A: Math. Theor. 47, 424006 (2014).

# A single relation for the QFI and the variance

For any decomposition  $\{p_k, |\psi_k\rangle\}$  of the density matrix  $\varrho$  we have

$$\frac{1}{4}F_Q[\varrho, A] \leq \sum_k p_k (\Delta A)_{\psi_k}^2 \leq (\Delta A)_{\varrho}^2,$$

where the upper and the lower bounds are both **tight**.

- Note that

$$\frac{1}{4}F_Q[\varrho, A] \leq (\Delta A)_{\varrho}^2,$$

where for pure states we have an equality.

- The QFI is strongly related to the variance.



# Formula based on an optimization in the two-copy space

- Two-copy formulation for the variance

$$(\Delta H)^2_\Psi = \text{Tr}(\Omega|\Psi\rangle\langle\Psi| \otimes |\Psi\rangle\langle\Psi|),$$

where we define the operator

$$\Omega = H^2 \otimes \mathbb{1} - H \otimes H.$$

We can reformulate the convex roof as

$$\begin{aligned} F_Q[\varrho, H] &= \min_{\varrho_{12}} && 4\text{Tr}(\Omega\varrho_{12}), \\ &\text{s. t.} && \varrho_{12} \in \mathcal{S}', \\ &&& \text{Tr}_2(\varrho_{12}) = \varrho. \end{aligned}$$

Here  $\mathcal{S}'$  is the set of symmetric separable states.

GT, T. Moroder, and O. Gühne, Evaluating convex roof entanglement measures, Phys. Rev. Lett. 114, 160501 (2015); GT, D. Petz, Phys. Rev. A 87, 032324 (2013).

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# Quantum Wasserstein distance based on an optimization over separable states

- **Definition**—We can also define

$$D_{\text{DPT},\text{sep}}(\varrho, \sigma)^2 = \frac{1}{2} \min_{\varrho_{12}} \sum_{n=1}^N \text{Tr}[(H_n^T \otimes \mathbb{1} - \mathbb{1} \otimes H_n)^2 \varrho_{12}],$$

s. t.  $\varrho_{12} \in \mathcal{S},$   
 $\text{Tr}_2(\varrho_{12}) = \varrho^T,$   
 $\text{Tr}_1(\varrho_{12}) = \sigma,$

where  $\mathcal{S}$  is the set of separable states.

GT and J. Pitrik, Quantum Wasserstein distance based on an optimization over separable states, arXiv:2209.09925.

# Quantum Wasserstein distance based on an optimization over separable states II

- For two-qubits, **it is computable numerically** with semidefinite programming.
- For systems of larger dimensions, one can obtain a **very good lower bound** based on an optimization over states with a positive partial transpose (PPT).
- Even better lower bounds can be obtained.

P. Horodecki, Phys. Lett. A 232, 333 (1997);

A. Peres, Phys. Rev. Lett. 77, 1413 (1996);

A. C. Doherty, P. A. Parrilo, and F. M. Spedalieri, Phys. Rev. A 69, 022308 (2004).

# Self-distance

- The self-distance for  $N = 1$  is

$$D_{\text{DPT,sep}}(\varrho, \varrho)^2 = \frac{1}{4} F_Q[\varrho, H_1].$$

- Note that

$$I_{\varrho}(A) \leq \frac{1}{4} F_Q[\varrho, A] \leq (\Delta A)_{\varrho}^2.$$

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# Entanglement of $\varrho_{12}$

- In general,

$$D_{\text{DPT,sep}}(\varrho, \sigma) \geq D_{\text{DPT}}(\varrho, \sigma).$$

- If the relation

$$D_{\text{DPT,sep}}(\varrho, \sigma) > D_{\text{DPT}}(\varrho, \sigma),$$

holds, then the optimal  $\varrho_{12}$  for  $D_{\text{DPT}}(\varrho, \sigma)$  is entangled.

- Allowing an entangled  $\varrho_{12}$  decreases the cost!
- Thus, an entangled  $\varrho_{12}$  can be cheaper than a separable one.

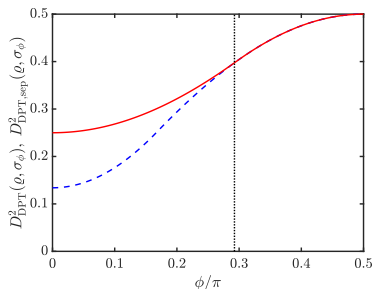
# Comparison of the two types of Wasserstein distance

- Let us consider the distance between two single-qubit mixed states

$$\varrho = \frac{1}{2}|1\rangle\langle 1|_x + \frac{1}{2} \cdot \frac{\mathbb{1}}{2},$$

and

$$\sigma_\phi = e^{-i\frac{\sigma_y}{2}\phi} \varrho e^{+i\frac{\sigma_y}{2}\phi}.$$





# Bounds on the distance

- **Entanglement condition:** Let us choose a set of  $H_n$  such that

$$\frac{1}{2} \sum_n \langle (H_n^T \otimes \mathbb{1} - \mathbb{1} \otimes H_n)^2 \rangle \geq \text{const.}$$

holds for separable states.

- E. g.,  $\{H_n\} = \{j_x, j_y, j_z\}$  and "const." =  $j$ .

- If the inequality

$$D_{\text{DPT}}(\varrho, \sigma) < \text{const.}$$

holds, then the the optimal  $\varrho_{12}$  for  $D_{\text{DPT}}(\varrho, \sigma)$  is entangled.

- Then, we will have a **minimal distance**

$$D_{\text{DPT,sep}}(\varrho, \sigma) \geq \text{const.}$$

# Summary

- For the quantum Wasserstein distance, we restrict the optimization to separable states.
- Then, the self-distance equals the quantum Fisher information over four.
- We found a fundamental connection from quantum optimal transport to quantum entanglement theory and quantum metrology.

G. Tóth and J. Pitrik, [arXiv:2209.09925](https://arxiv.org/abs/2209.09925).

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