

A simple electronic ladder model harboring Z_4 parafermions

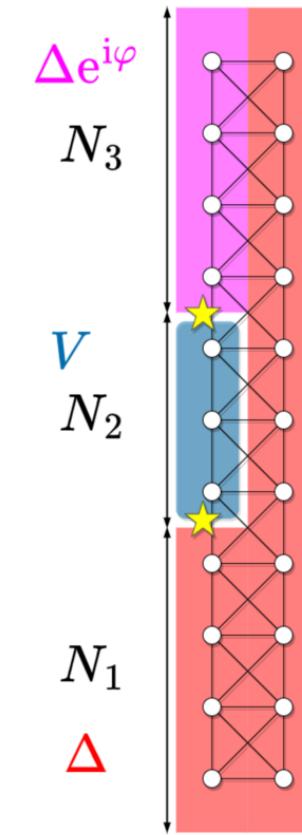
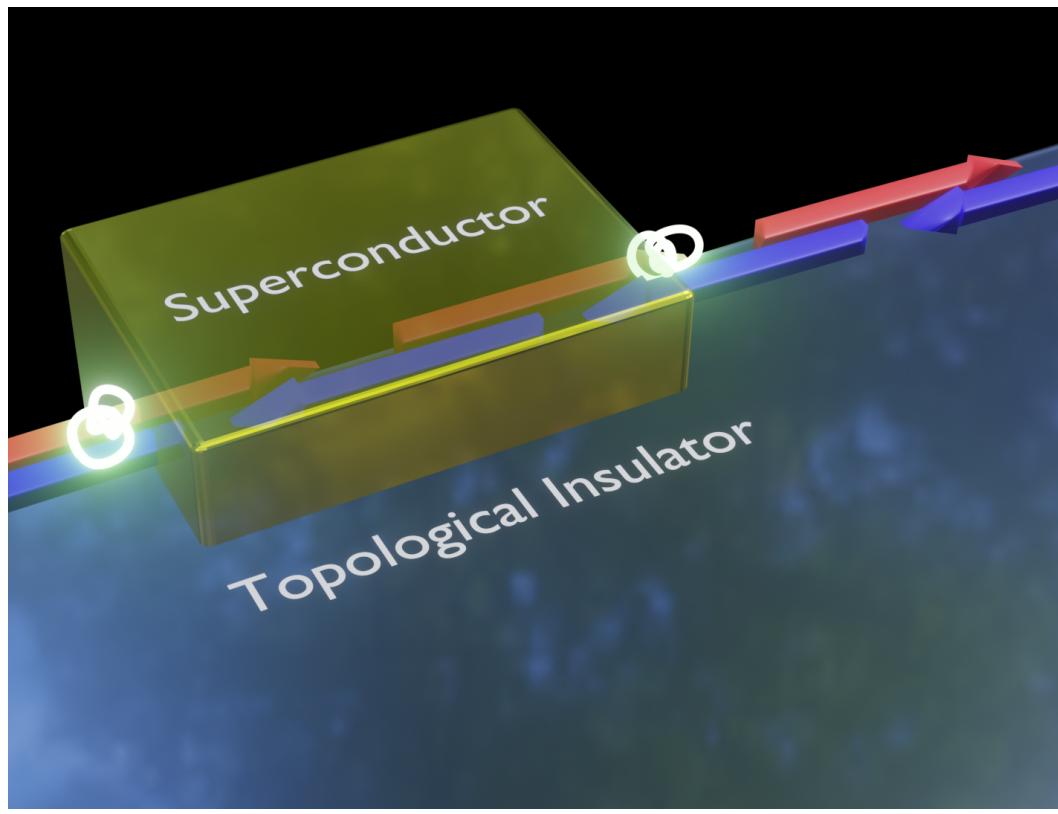


ELTE
EÖTVÖS LORÁND
TUDOMÁNYEGYETEM



Lendület
program

QNL Quantum Information
National Laboratory
HUNGARY



UNKP
Új Nemzeti
Kiválóság Program

Wigner
MŰEGYETEM 1782

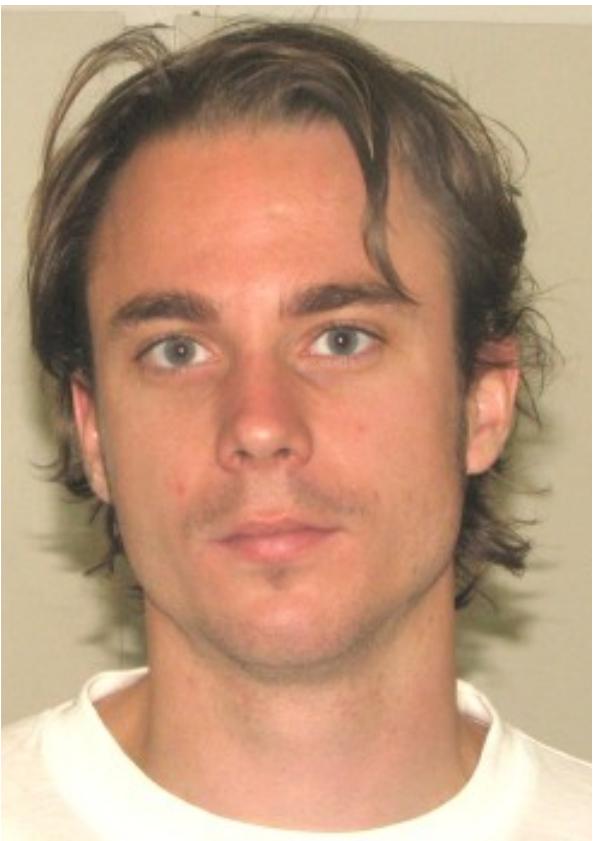
László Oroszlány

Department of Physics of Complex Systems, Eötvös Loránd University
Wigner Research Centre for Physics

The team



Botond Osváth
ELTE



Gergely Barcza
Wigner FK



Balázs Dóra
BME



Örs Legeza
Wigner FK

Outline

- What are Parafermions ? Why should you care?
- Where did people look for parafermions?
- Where we look for parafermions ?
- Where one could look for parafermions ?

Clock models and parafermions

$$H = -J \sum_{p=1}^{L-1} \hat{\sigma}_p^\dagger \hat{\sigma}_{p+1} - f \sum_{p=1}^L \hat{\tau}_p + \text{h.c.}$$

N=3 Clock model

$$\sigma = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \Omega & 0 \\ 0 & 0 & \Omega^2 \end{pmatrix}$$

$$\tau = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\Omega = e^{i2\pi/N} = \omega^2$$

Clock models and parafermions

$$H = -J \sum_{p=1}^{L-1} \hat{\sigma}_p^\dagger \hat{\sigma}_{p+1} - f \sum_{p=1}^L \hat{\tau}_p + \text{h.c.}$$

\Updownarrow

$$H = -J \sum_{p=1}^{L-1} \omega \hat{\alpha}_{2p}^\dagger \hat{\alpha}_{2p+1} - f \sum_{p=1}^L \omega \hat{\alpha}_{2p-1}^\dagger \hat{\alpha}_{2p}$$

Jordan-Wigner

$$\hat{\alpha}_{2p-1} = \hat{\sigma}_p \prod_{q < p} \hat{\tau}_q$$

$$\hat{\alpha}_{2p} = -\omega \hat{\tau}_p \hat{\sigma}_p \prod_{q < p} \hat{\tau}_q$$

N=3 Clock model

$$\sigma = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \Omega & 0 \\ 0 & 0 & \Omega^2 \end{pmatrix}$$

$$\tau = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\Omega = e^{i2\pi/N} = \omega^2$$

Parafermion

$$\hat{\alpha}_p^N = \hat{1}, \hat{\alpha}_p^\dagger = \hat{\alpha}_p^{N-1}$$

$$\hat{\alpha}_p \hat{\alpha}_q = \Omega^{\text{sign}(q-p)} \hat{\alpha}_q \hat{\alpha}_p$$

Clock models and parafermions

$$H = -J \sum_{p=1}^{L-1} \hat{\sigma}_p^\dagger \hat{\sigma}_{p+1} - f \sum_{p=1}^L \hat{\tau}_p + \text{h.c.}$$

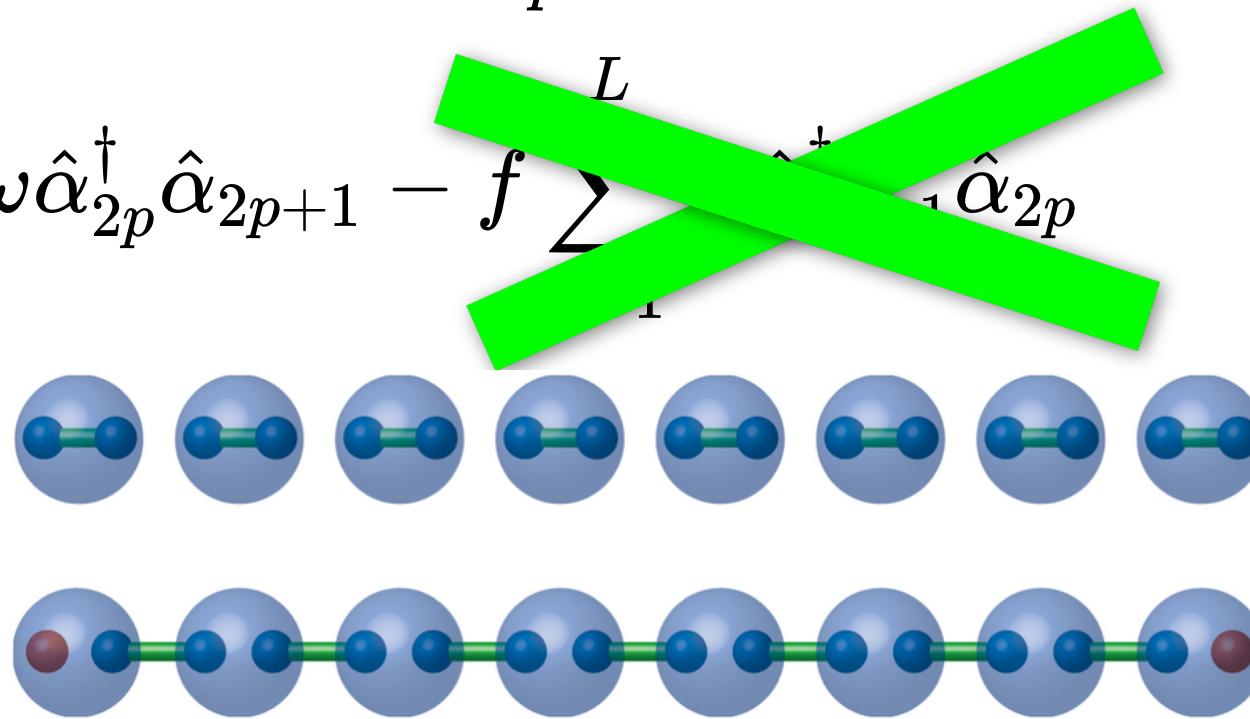
\Updownarrow

$$H = -J \sum_{p=1}^{L-1} \omega \hat{\alpha}_{2p}^\dagger \hat{\alpha}_{2p+1} - f \sum_{p=1}^L \hat{\alpha}_{2p} + \text{h.c.}$$

Jordan-Wigner

$$\hat{\alpha}_{2p-1} = \hat{\sigma}_p \prod_{q < p} \hat{\tau}_q$$

$$\hat{\alpha}_{2p} = -\omega \hat{\tau}_p \hat{\sigma}_p \prod_{q < p} \hat{\tau}_q$$



$f = 0 \rightarrow$ parafermions at the edges, $\hat{\alpha}_1$ & $\hat{\alpha}_{2L}$, absent from the Hamiltonian!

The missing two parafermions encode an N-fold degenerate subspace!

N=3 Clock model

$$\sigma = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \Omega & 0 \\ 0 & 0 & \Omega^2 \end{pmatrix}$$

$$\tau = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\Omega = e^{i2\pi/N} = \omega^2$$

Parafermion

$$\hat{\alpha}_p^N = \hat{1}, \hat{\alpha}_p^\dagger = \hat{\alpha}_p^{N-1}$$

$$\hat{\alpha}_p \hat{\alpha}_q = \Omega^{\text{sign}(q-p)} \hat{\alpha}_q \hat{\alpha}_p$$

Majoranas vs. parafermions

- Majorana modes can be potentially realized in non-interacting systems. (*i.e.* mean-field description is sufficient)

Majoranas vs. parafermions

- Majorana modes can be potentially realized in non-interacting systems. (*i.e.* mean-field description is sufficient)
- With braiding alone, Majorana modes can realize nontrivial unitary operations, but no entangling qbit gates.

Majoranas vs. parafermions

- Majorana modes can be potentially realized in non-interacting systems. (*i.e.* mean-field description is sufficient)
- With braiding alone, Majorana modes can realize nontrivial unitary operations, but no entangling qbit gates.
- Parafermions need interaction. (*i.e.* mean-field description is not sufficient)

Majoranas vs. parafermions

- Majorana modes can be potentially realized in non-interacting systems. (*i.e.* mean-field description is sufficient)
- With braiding alone, Majorana modes can realize nontrivial unitary operations, but no entangling qbit gates.
- Parafermions need interaction. (*i.e.* mean-field description is not sufficient)
- \mathbb{Z}_{even} parafermions can realize entangling gates just with braiding!

Majoranas vs. parafermions

- Majorana modes can be potentially realized in non-interacting systems. (*i.e.* mean-field description is sufficient)
- With braiding alone, Majorana modes can realize nontrivial unitary operations, but no entangling qbit gates.
- Parafermions need interaction. (*i.e.* mean-field description is not sufficient)
- \mathbb{Z}_{even} parafermions can realize entangling gates just with braiding!
- \mathbb{Z}_{odd} parafermions: route to universality

Parafermion signatures

- Robustness against disorder
- Highly (>2) degenerate groundstate
- Localized zero-energy excitations
- Nontrivial (fractional) Josephson effect

$$\mathbb{Z}_n \rightarrow 2n\pi \text{ periodic}$$

Z_4 parafermions from ordinary fermions

Hamiltonian in fermion language ...

$N=4$ clock model/
parafermion chain



each site
has 4 states

spinful electron
in 1D wire

\mathbb{Z}_4 parafermions from ordinary fermions

Hamiltonian in fermion language ...

N=4 clock model/
parafermion chain



each site
has 4 states

spinful electron
in 1D wire

$$H = H^{(2)} + H^{(4)} + H^{(6)}$$

$$H^{(2)} = -J \sum_{\sigma,j} \left[c_{\sigma,j}^\dagger c_{\sigma,j+1} - i c_{-\sigma,j}^\dagger c_{\sigma,j+1}^\dagger \right] + h.c.,$$

$$\begin{aligned} H^{(4)} = & -J \sum_{\sigma,j} \left[c_{\sigma,j}^\dagger c_{\sigma,j+1} (-n_{-\sigma,j} - n_{-\sigma,j+1}) \right. \\ & + c_{\sigma,j}^\dagger c_{-\sigma,j+1} i (n_{-\sigma,j} + n_{\sigma,j+1}) \\ & + c_{-\sigma,j}^\dagger c_{\sigma,j+1}^\dagger i (n_{\sigma,j} + n_{-\sigma,j+1}) \\ & \left. + c_{\sigma,j}^\dagger c_{\sigma,j+1}^\dagger (n_{-\sigma,j} - n_{-\sigma,j+1}) \right] + h.c., \end{aligned}$$

$$\begin{aligned} H^{(6)} = & -J \sum_j \left[-2i c_{\sigma,j}^\dagger c_{-\sigma,j+1} (n_{-\sigma,j} n_{\sigma,j+1}) \right. \\ & \left. - 2i c_{-\sigma,j}^\dagger c_{\sigma,j+1}^\dagger (n_{\sigma,j} n_{-\sigma,j+1}) \right] + h.c. \end{aligned}$$

\mathbb{Z}_4 parafermions from ordinary fermions

Hamiltonian in fermion language ...

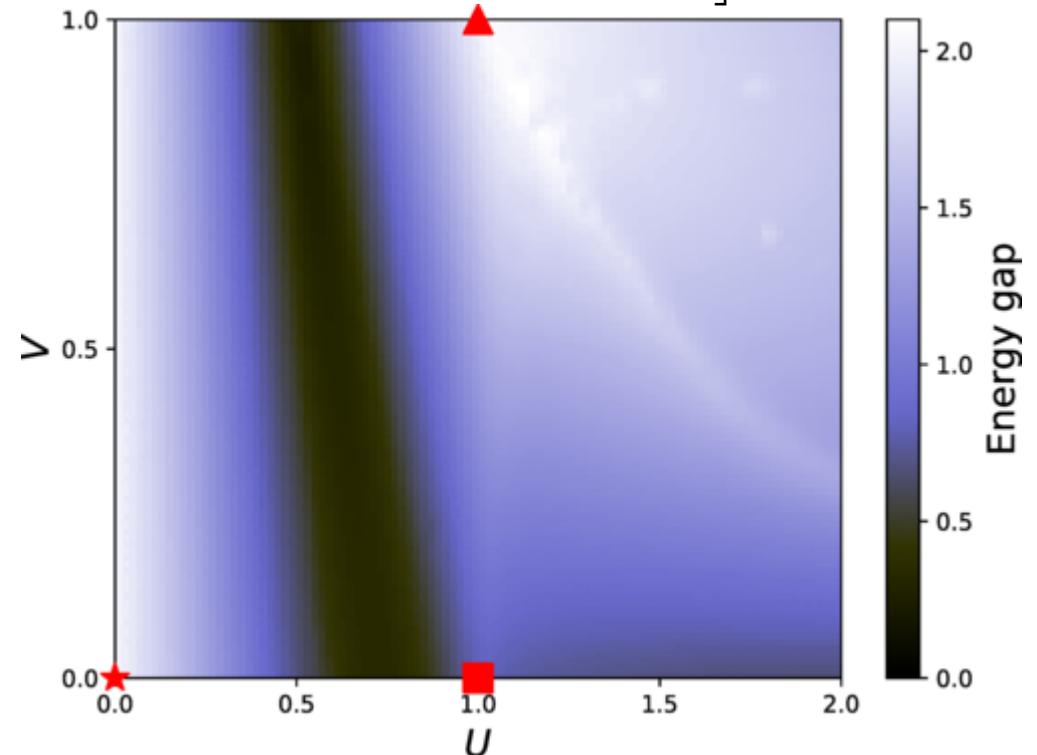
$N=4$ clock model/
parafermion chain



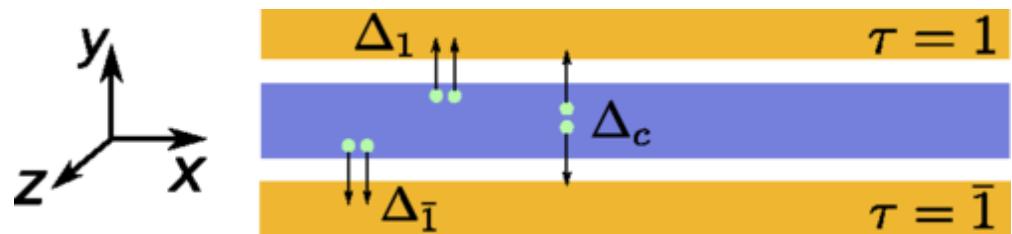
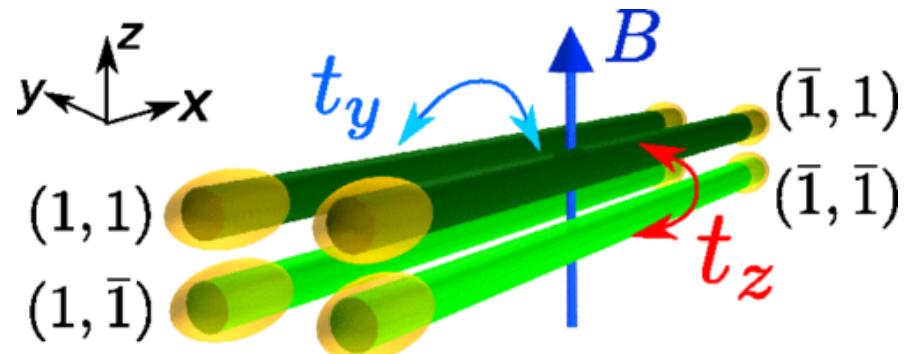
each site
has 4 states

spinful electron
in 1D wire

$$\bar{H}(U, V) = H^{(2)} + U [V (H^{(4)} + H^{(6)}) + (1 - V) \bar{H}^{(4)}]$$
$$\bar{H}^{(4)} = -J \sum_{\sigma, j} \left[c_{\sigma, j}^\dagger c_{\sigma, j+1} (-n_{-\sigma, j} - n_{-\sigma, j+1}) \right. \\ \left. + c_{\sigma, j}^\dagger c_{\sigma, j+1}^\dagger (n_{-\sigma, j} - n_{-\sigma, j+1}) \right] + h.c.$$



Possible experimental blueprints

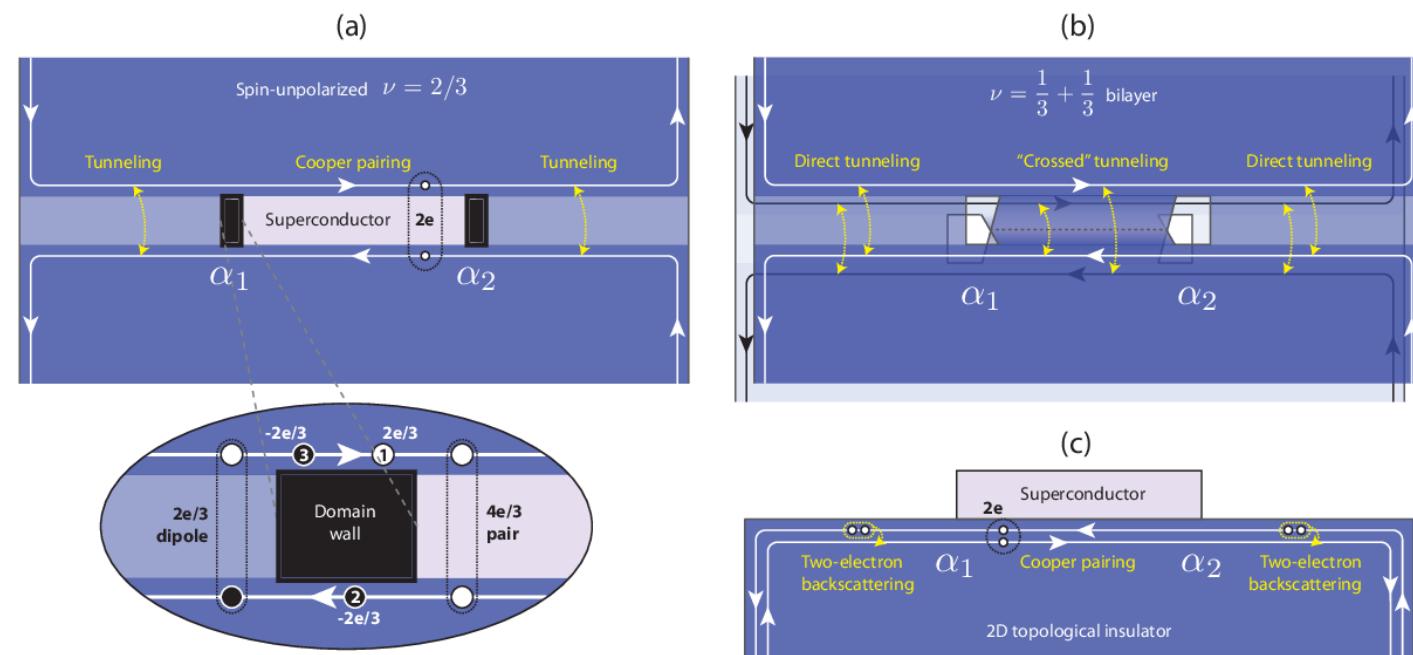
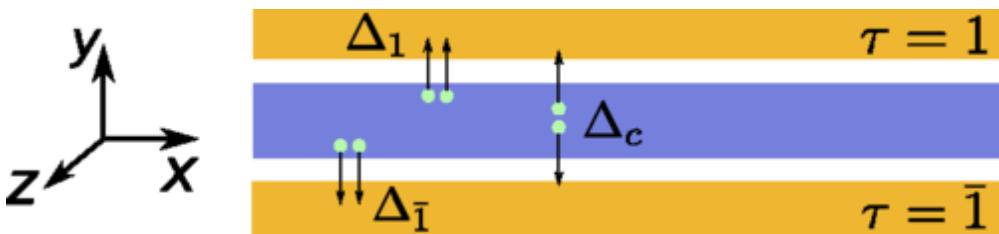
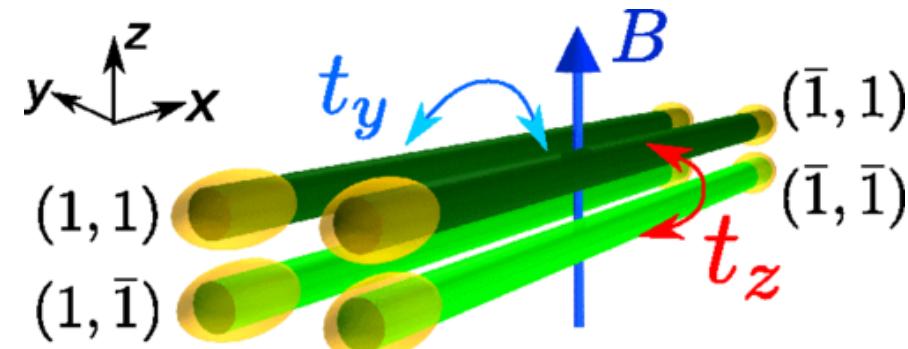


J. Klinovaja and D. Loss

Phys. Rev. Lett. **112**, 246403 (2014)

Phys. Rev. B **90**, 045118 (2014)

Possible experimental blueprints



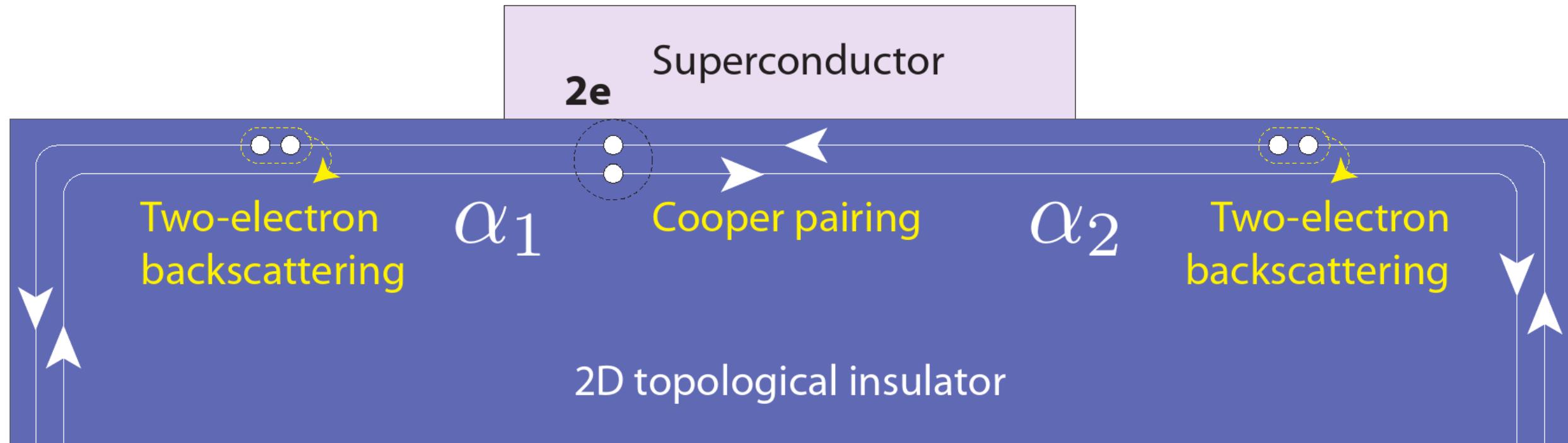
J. Klinovaja and D. Loss

Phys. Rev. Lett. **112**, 246403 (2014)

Phys. Rev. B **90**, 045118 (2014)

J. Alicea, P. Fendley Annu. Rev. Condens. Matter Phys. **7**, 119 (2016.)

Parafermions at TI edge



F. Zhang, C. L. Kane, Phys. Rev. Lett., **113**, 036401 (2014).

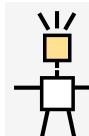
C. P. Orth *et al.* Phys. Rev. B, **91**, 081406 (2015).

J. Alicea, P. Fendley Annu. Rev. Condens. Matter Phys. **7**, 119 (2016.)

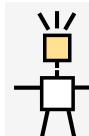
bosonised models

goal: microscopic model + DMRG

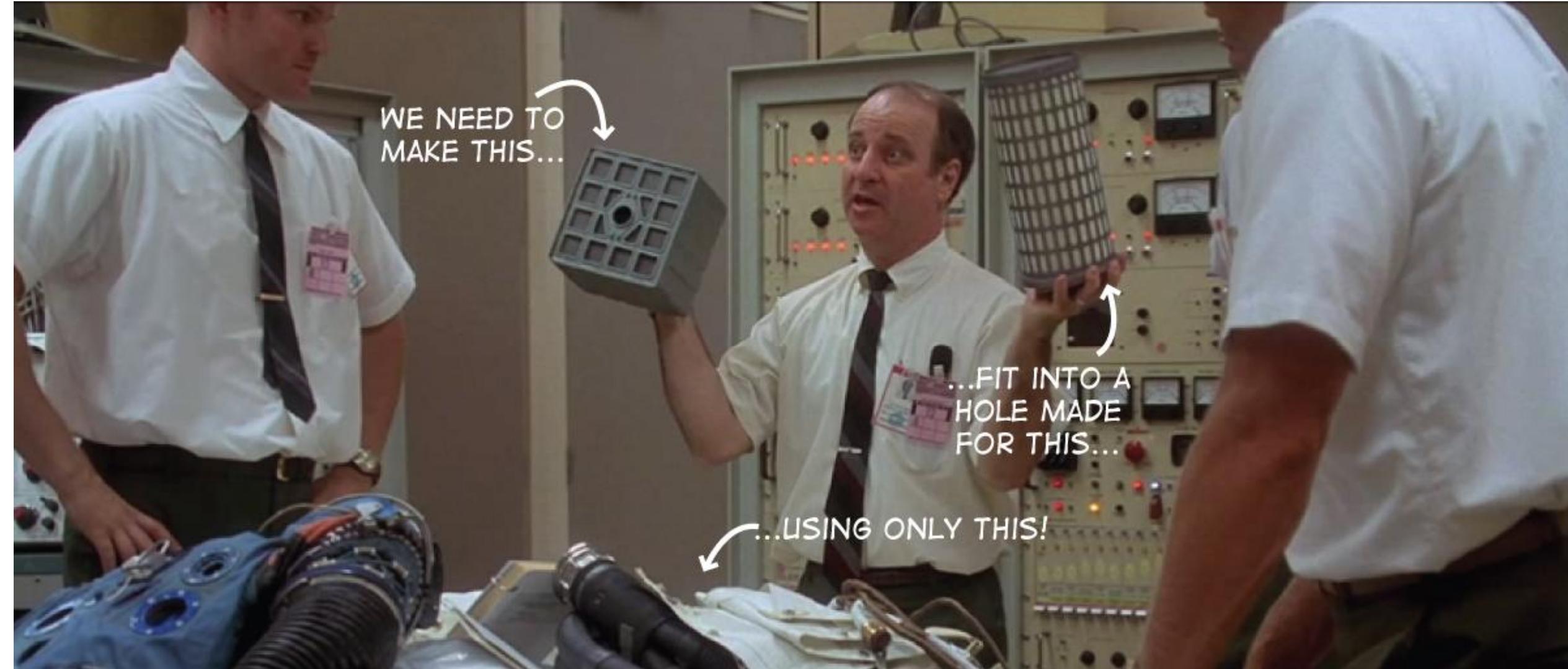
ITensor



Budapest DMRG



A model for 2DTI that can be digested by DMRG?



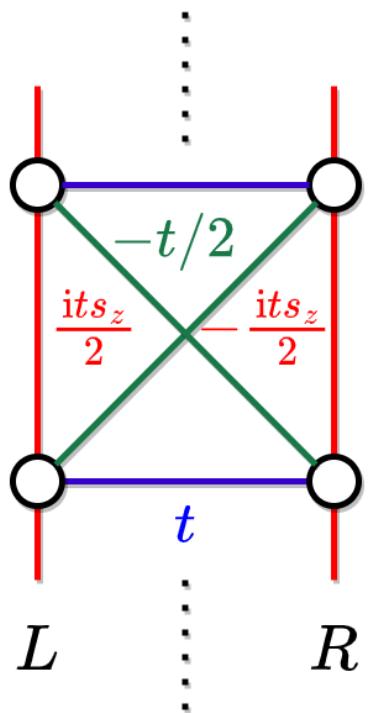
The model

$$H = H_k + H_{sc} + H_{int}$$

The model

$$H = H_k + H_{sc} + H_{int}$$

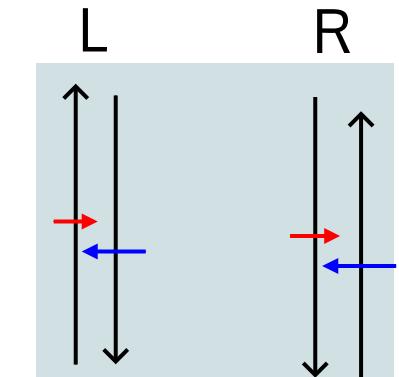
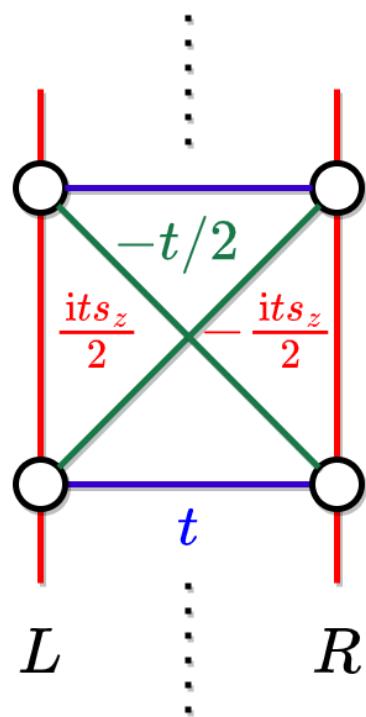
$$H_k = \sum_{m\sigma} \begin{pmatrix} c_{mL\sigma}^\dagger & c_{mR\sigma}^\dagger \end{pmatrix} \begin{pmatrix} -\mu & t \\ t & -\mu \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix}$$
$$-\frac{t}{2} \sum_{m\sigma} \left[\begin{pmatrix} c_{m+1L\sigma}^\dagger & c_{m+1R\sigma}^\dagger \end{pmatrix} \begin{pmatrix} i\sigma & 1 \\ 1 & -i\sigma \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix} + \text{h.c.} \right]$$



The model

$$H = H_k + H_{sc} + H_{int}$$

$$H_k = \sum_{m\sigma} \begin{pmatrix} c_{mL\sigma}^\dagger & c_{mR\sigma}^\dagger \end{pmatrix} \begin{pmatrix} -\mu & t \\ t & -\mu \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix} \\ -\frac{t}{2} \sum_{m\sigma} \left[\begin{pmatrix} c_{m+1L\sigma}^\dagger & c_{m+1R\sigma}^\dagger \end{pmatrix} \begin{pmatrix} i\sigma & 1 \\ 1 & -i\sigma \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix} + \text{h.c.} \right]$$



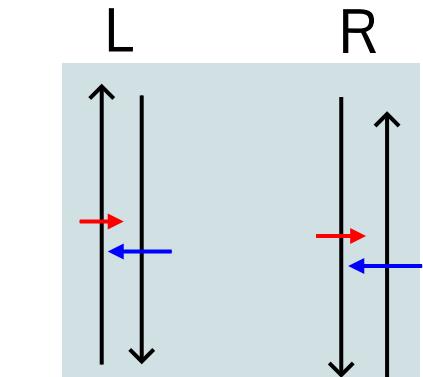
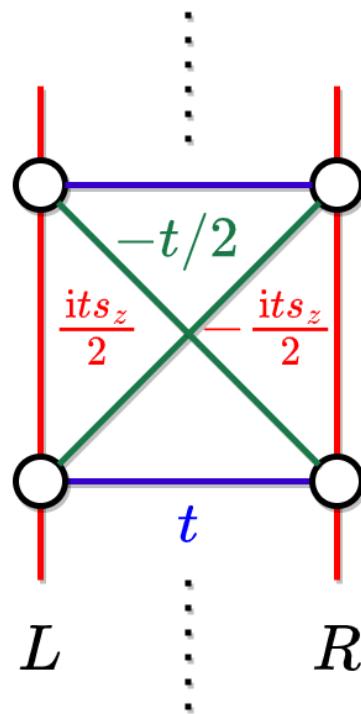
- two "disconnected edges" $\zeta = L, R$

$$H_k \approx ps_z\zeta_z$$

The model

$$H = H_k + H_{sc} + H_{int}$$

$$H_k = \sum_{m\sigma} \begin{pmatrix} c_{mL\sigma}^\dagger & c_{mR\sigma}^\dagger \end{pmatrix} \begin{pmatrix} -\mu & t \\ t & -\mu \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix} - \frac{t}{2} \sum_{m\sigma} \left[\begin{pmatrix} c_{m+1L\sigma}^\dagger & c_{m+1R\sigma}^\dagger \end{pmatrix} \begin{pmatrix} i\sigma & 1 \\ 1 & -i\sigma \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix} + \text{h.c.} \right]$$



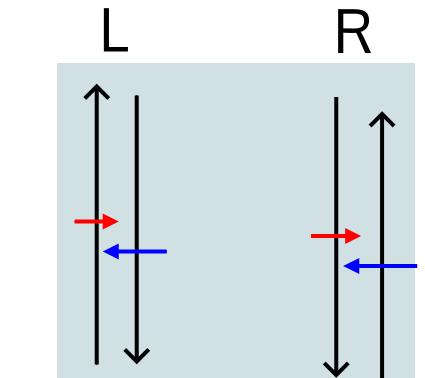
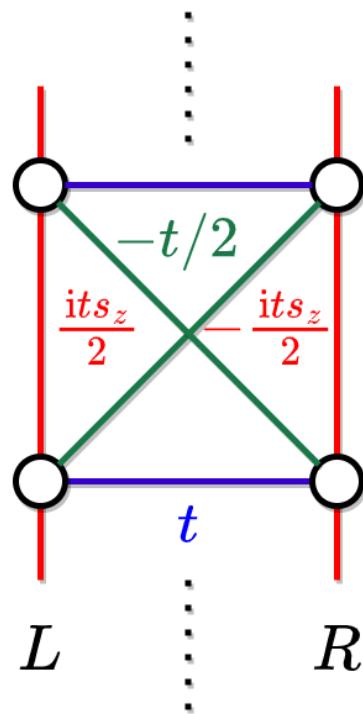
$$H_k \approx ps_z\zeta_z$$

- $$H_{sc} = \sum_{m\zeta} \Delta_{m\zeta} [c_{m\zeta\uparrow}^\dagger c_{m\zeta\downarrow}^\dagger + \text{h.c.}]$$
- $$H_{int} = \sum_{m\zeta} V_{m\zeta} [c_{m\zeta\uparrow}^\dagger c_{m\zeta\downarrow}^\dagger c_{(m+1)\zeta\uparrow}^\dagger c_{(m+1)\zeta\downarrow} + \text{h.c.}]$$
- **two "disconnected edges" $\zeta = L, R$**
 - **explicit superconductivity and interactions**

The model

$$H = H_k + H_{sc} + H_{int}$$

$$H_k = \sum_{m\sigma} \begin{pmatrix} c_{mL\sigma}^\dagger & c_{mR\sigma}^\dagger \end{pmatrix} \begin{pmatrix} -\mu & t \\ t & -\mu \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix} - \frac{t}{2} \sum_{m\sigma} \left[\begin{pmatrix} c_{m+1L\sigma}^\dagger & c_{m+1R\sigma}^\dagger \end{pmatrix} \begin{pmatrix} i\sigma & 1 \\ 1 & -i\sigma \end{pmatrix} \begin{pmatrix} c_{mL\sigma} \\ c_{mR\sigma} \end{pmatrix} + \text{h.c.} \right]$$



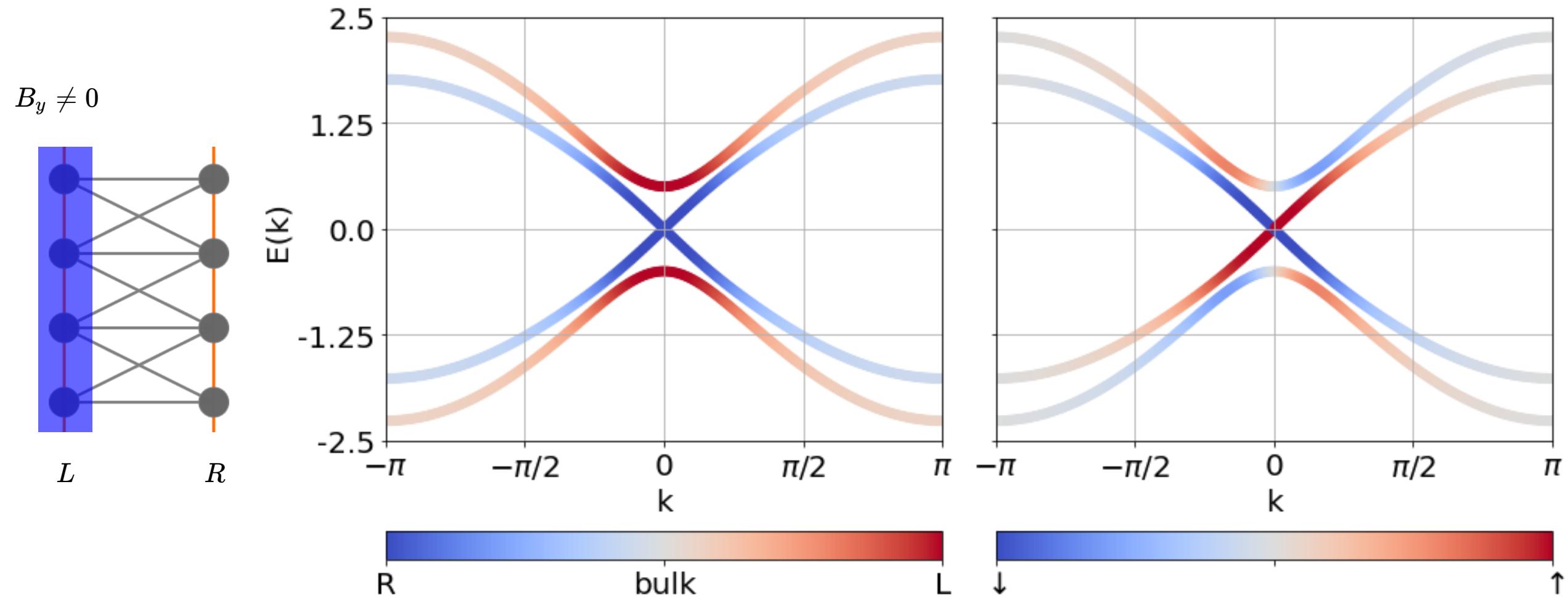
$$H_k \approx ps_z\zeta_z$$

$$H_{sc} = \sum_{m\zeta} \Delta_{m\zeta} [c_{m\zeta\uparrow}^\dagger c_{m\zeta\downarrow}^\dagger + \text{h.c.}]$$

$$H_{int} = \sum_{m\zeta} V_{m\zeta} [c_{m\zeta\uparrow}^\dagger c_{m\zeta\downarrow}^\dagger c_{(m+1)\zeta\uparrow}^\dagger c_{(m+1)\zeta\downarrow} + \text{h.c.}]$$

- **two "disconnected edges" $\zeta = L, R$**
- **explicit superconductivity and interactions**
- **time reversal symmetry**

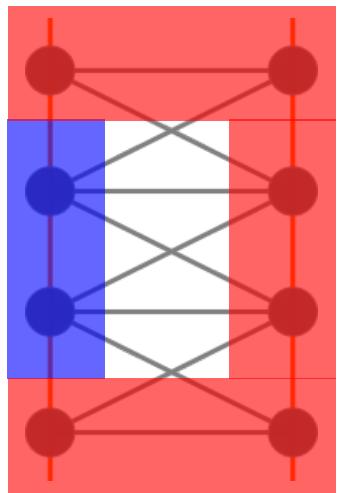
Single particle spectrum



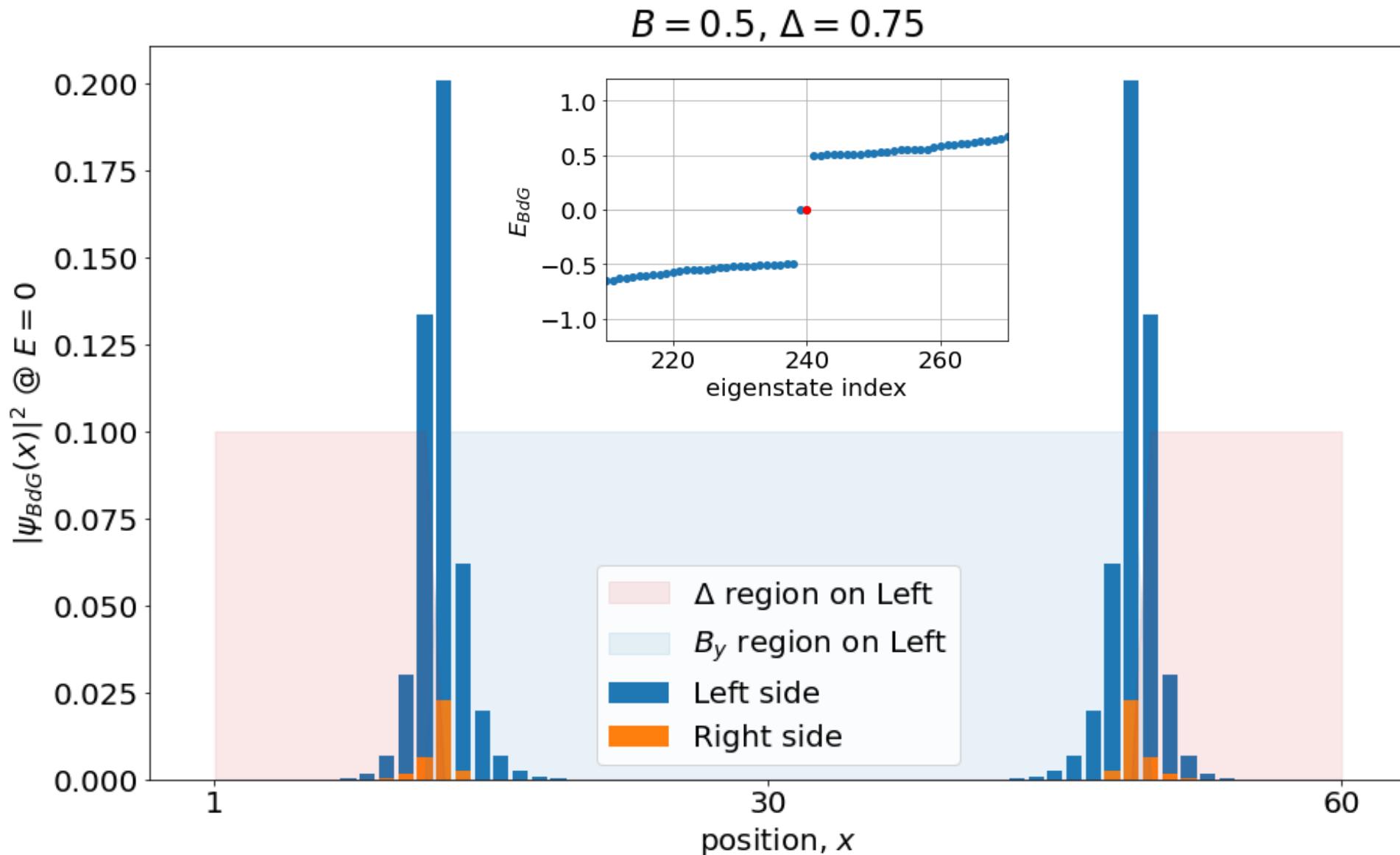
small B_y on the left for better visibility

We still have Majoranas !

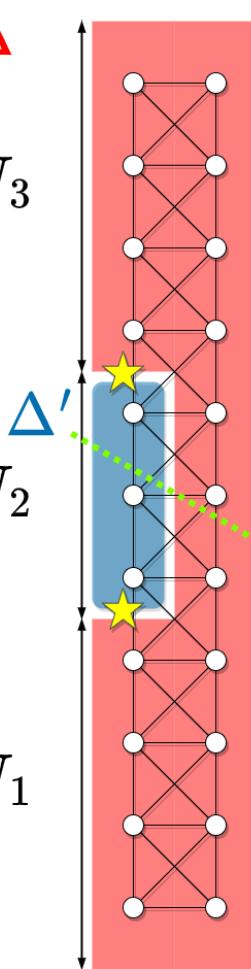
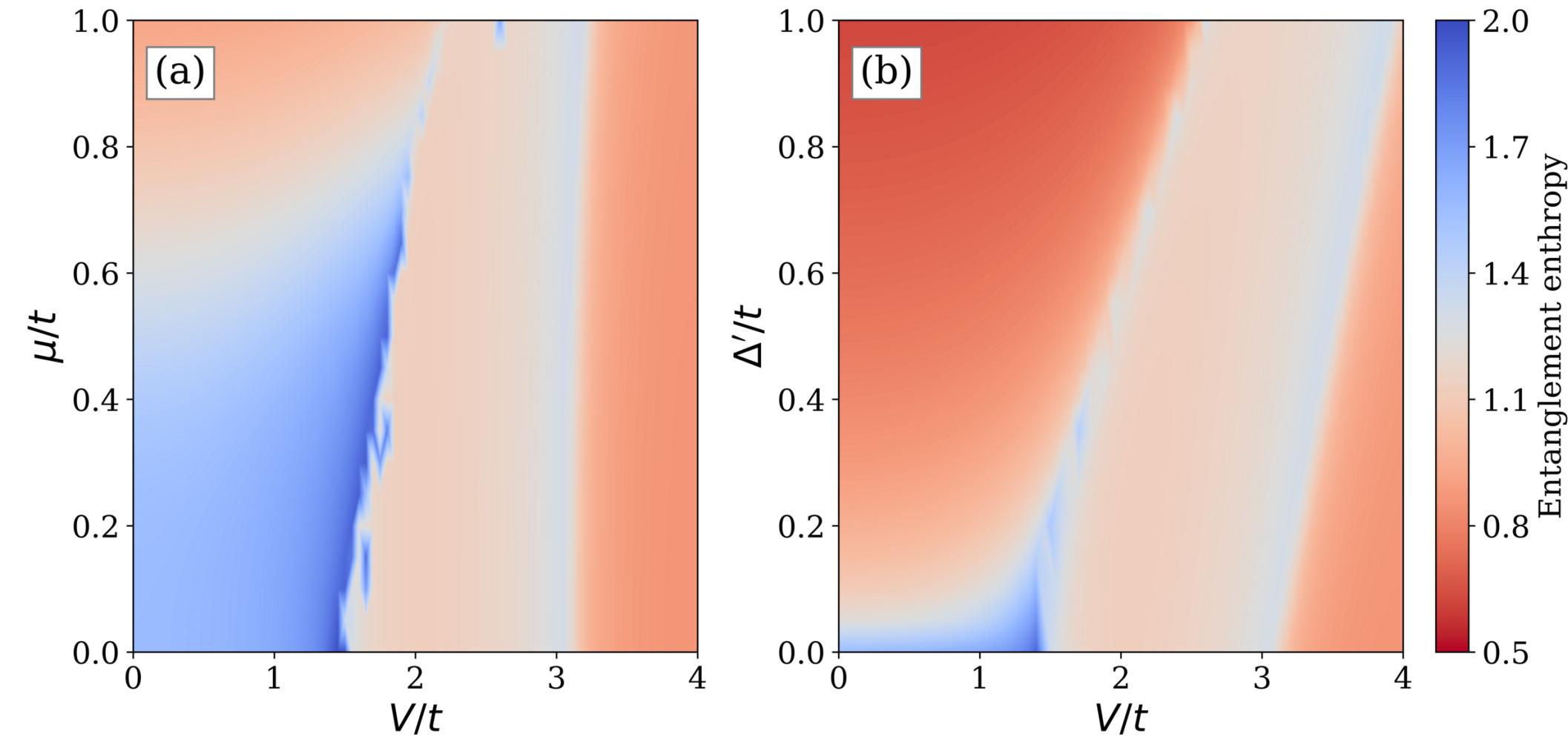
$$B_y \neq 0 \quad \Delta \neq 0$$



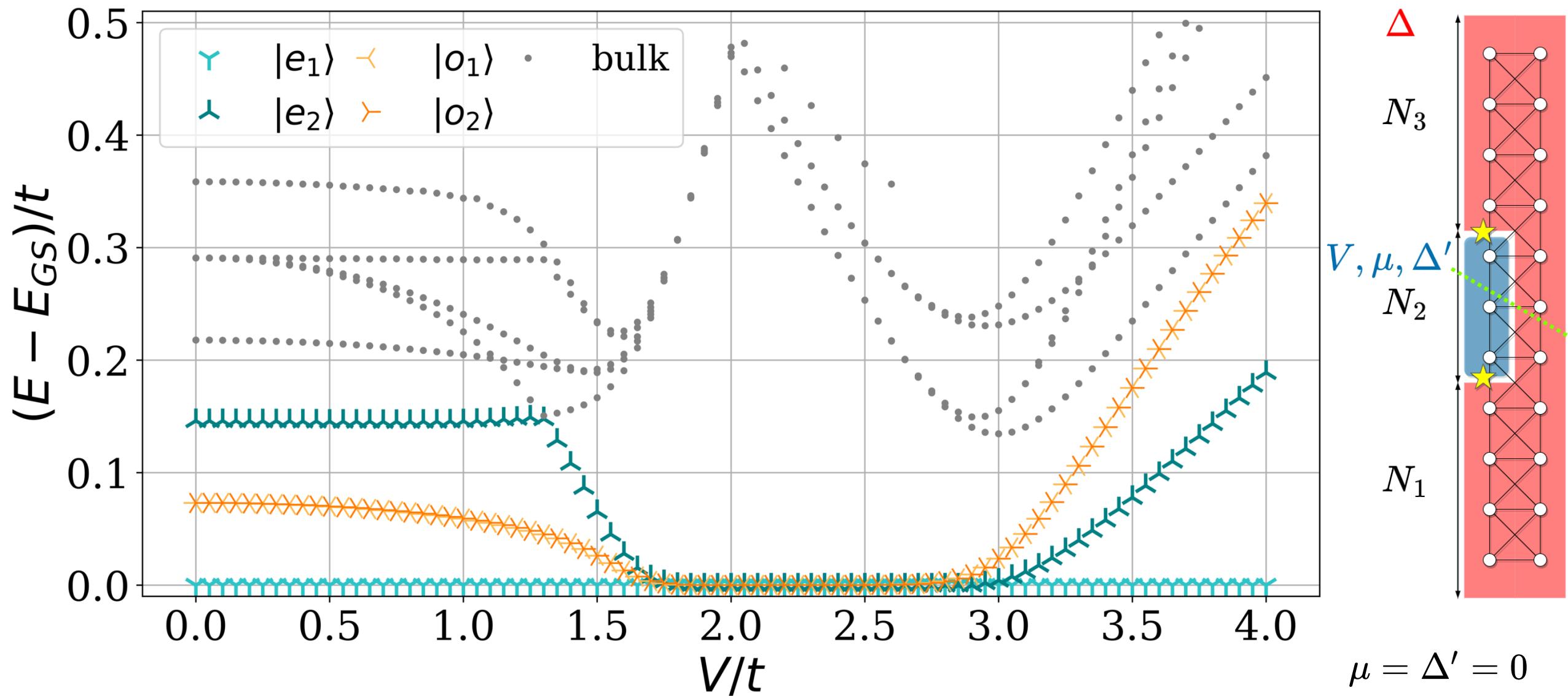
$$V = 0$$



Finite-size DMRG calculations: phase diagram

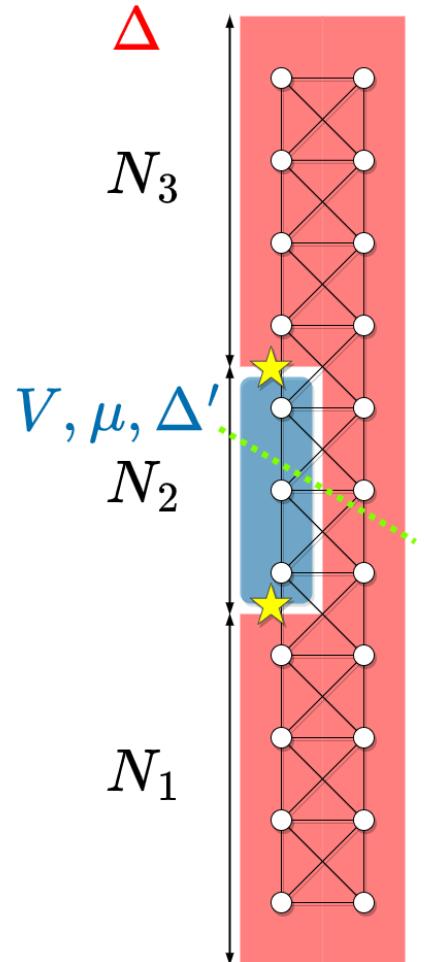


Finite-size DMRG calculations: excitation spectrum



Finite-size DMRG calculations: local quantities

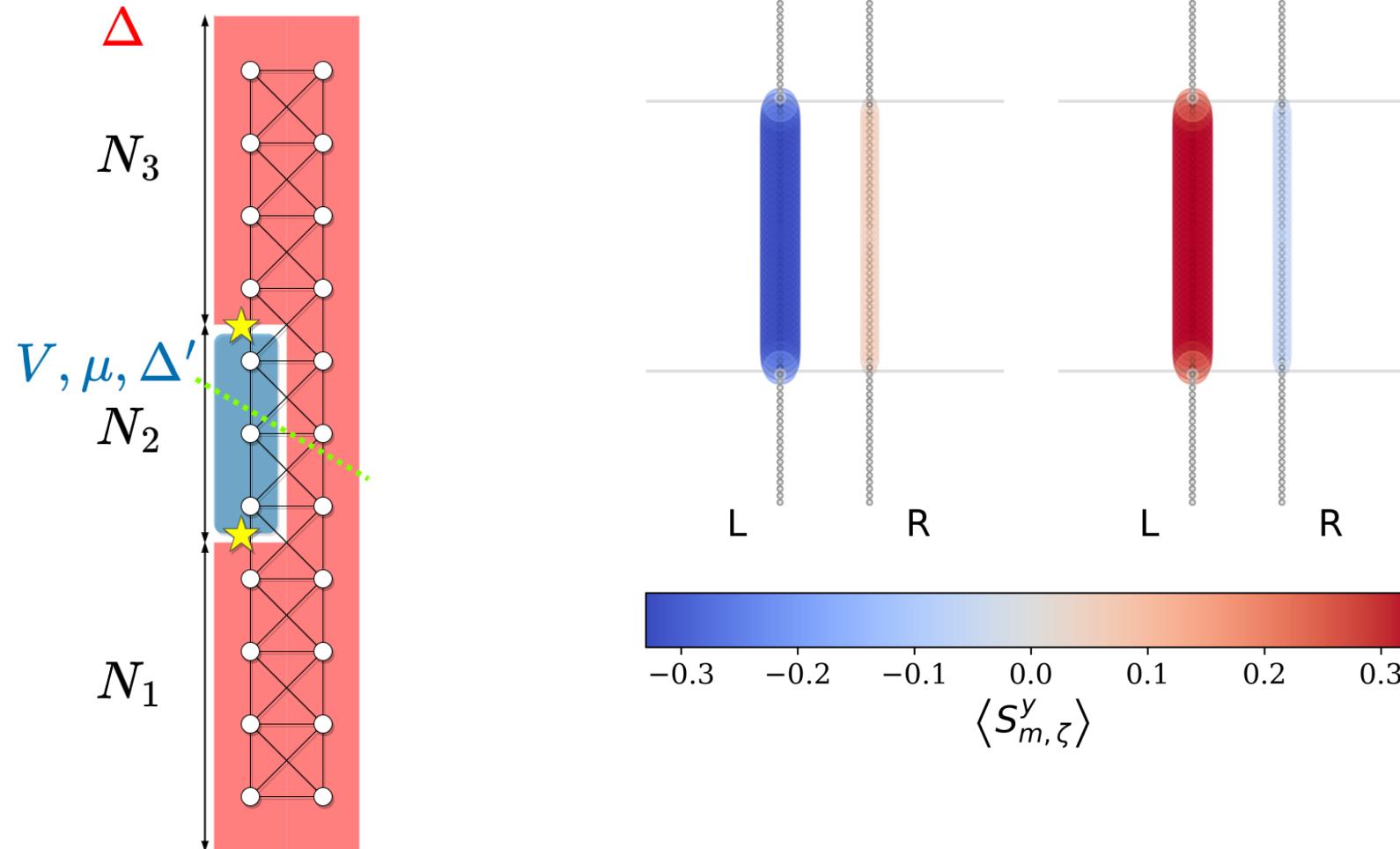
$$\mu = \Delta' = 0, V/t = 2.2$$



$$\langle GS_p | n_i | GS_q \rangle \propto \delta_{pq}$$

Finite-size DMRG calculations: local quantities

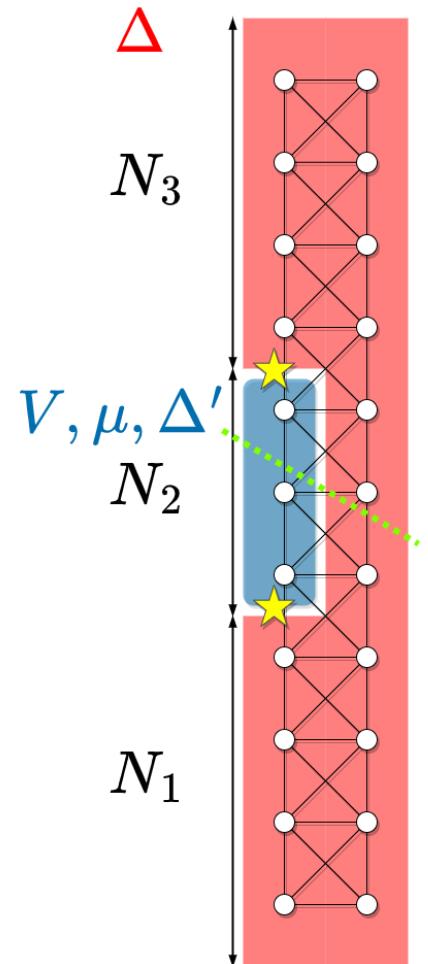
$$\mu = \Delta' = 0, V/t = 2.2$$



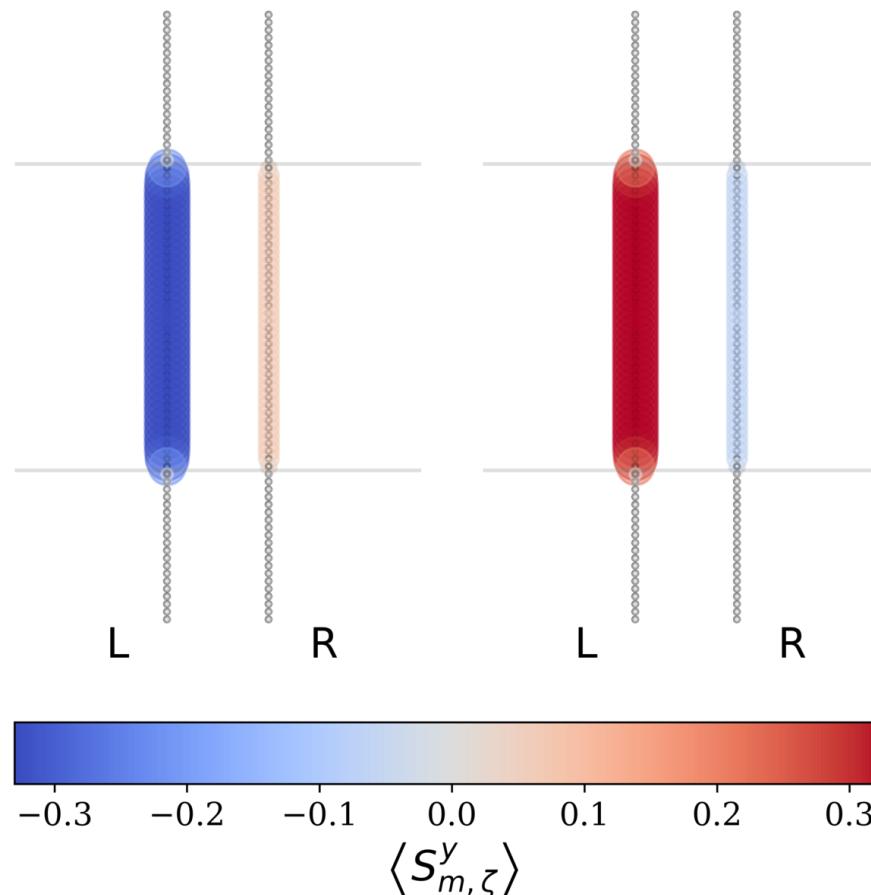
$$\langle GS_p | n_i | GS_q \rangle \propto \delta_{pq}$$

Finite-size DMRG calculations: local quantities

$$\mu = \Delta' = 0, V/t = 2.2$$



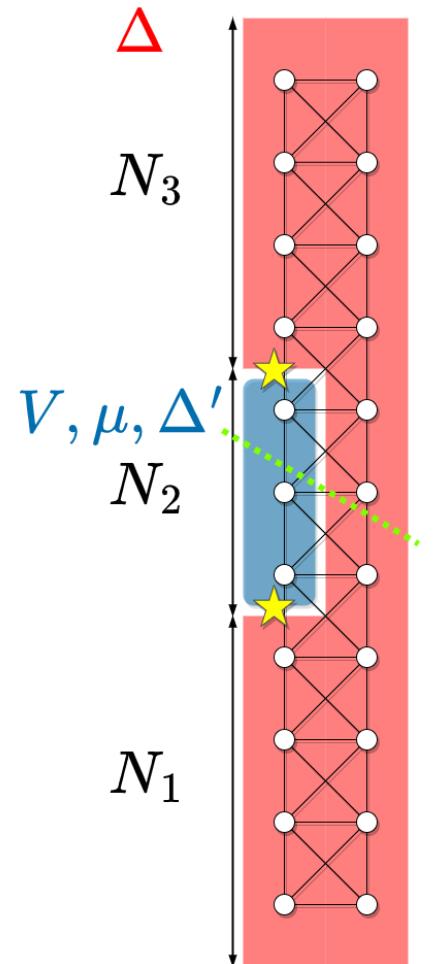
$$\langle GS_p | n_i | GS_q \rangle \propto \delta_{pq}$$



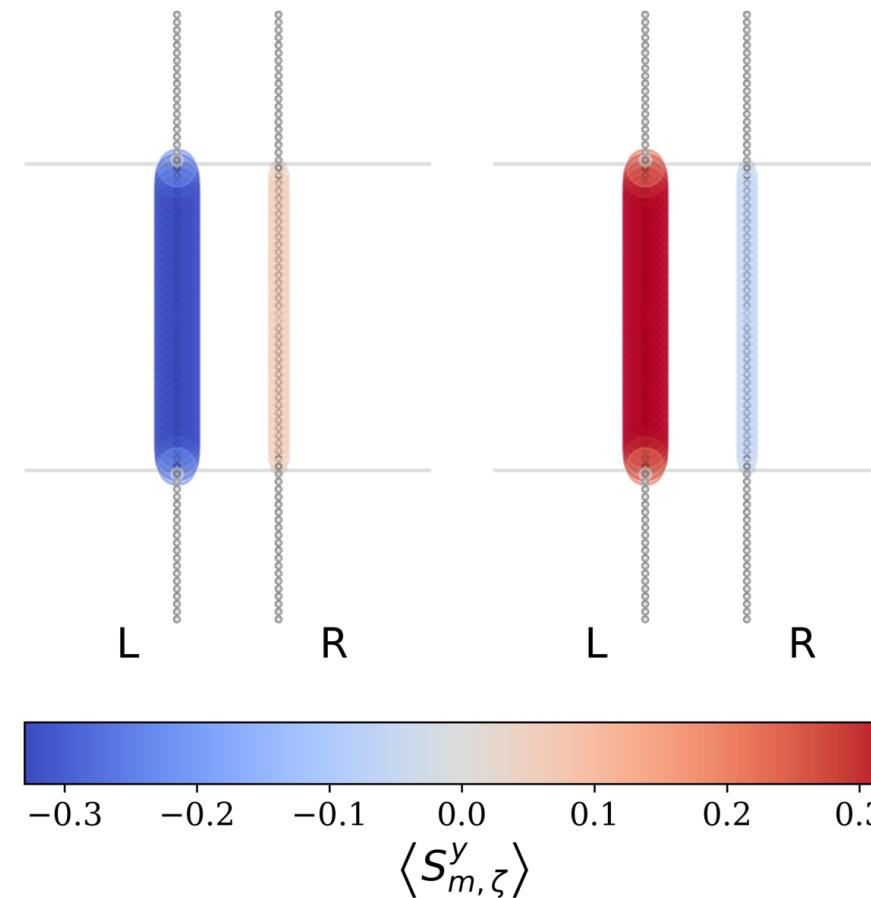
$$\begin{aligned} H_{int} &= \sum_{m\zeta} V_{m\zeta} \left[c_{m\zeta\uparrow}^\dagger c_{m\zeta\downarrow} c_{(m+1)\zeta\uparrow}^\dagger c_{(m+1)\zeta\downarrow} + \text{h.c.} \right] \\ &= \sum_{m,\zeta} \frac{V_{m,\zeta}}{2} \left[S_{m,\zeta}^x S_{m+1,\zeta}^x - S_{m,\zeta}^y S_{m+1,\zeta}^y \right] \end{aligned}$$

Finite-size DMRG calculations: local quantities

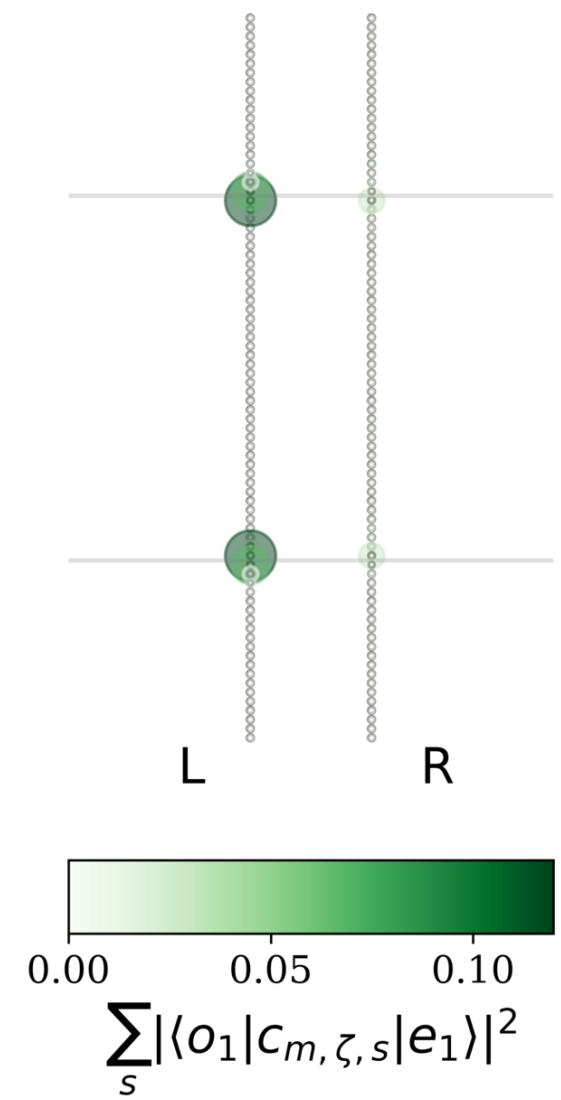
$$\mu = \Delta' = 0, V/t = 2.2$$



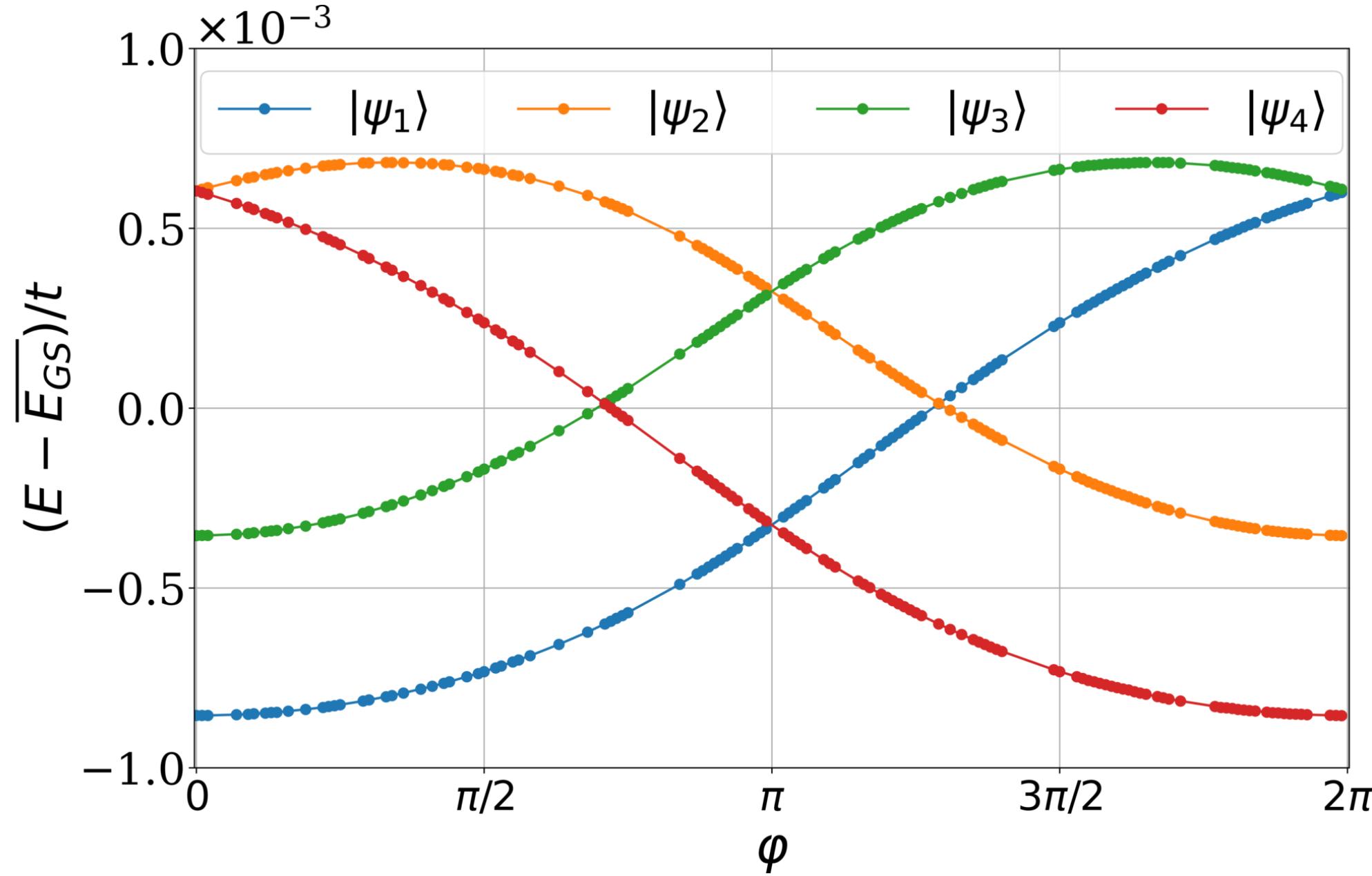
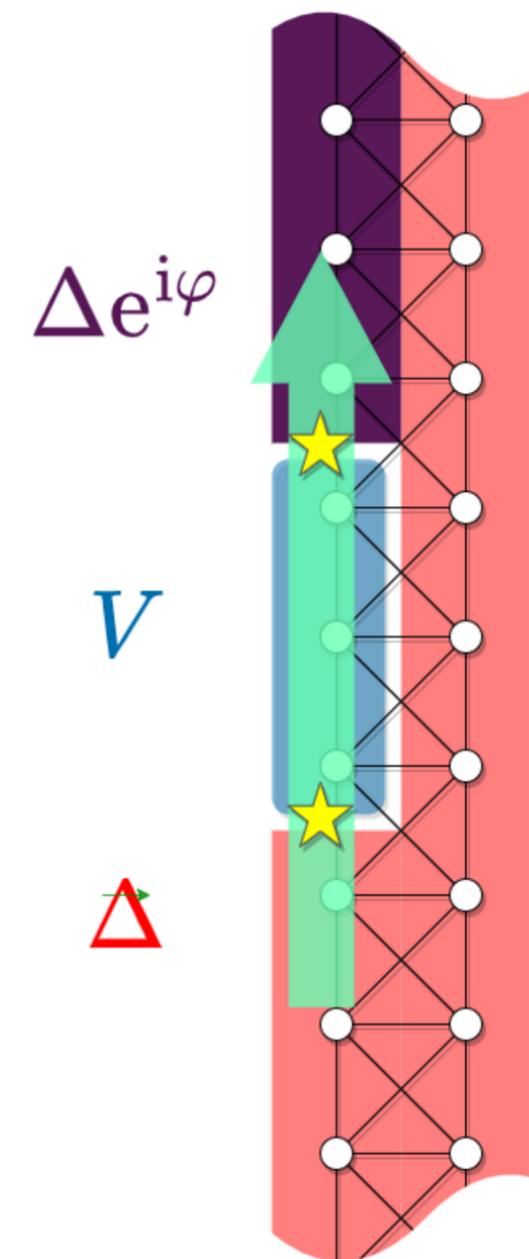
$$\langle GS_p | n_i | GS_q \rangle \propto \delta_{pq}$$



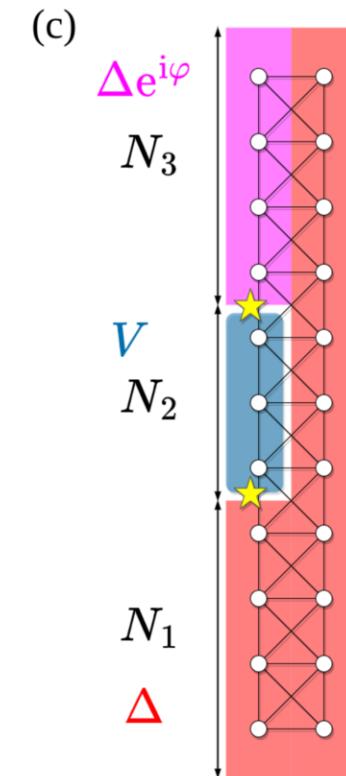
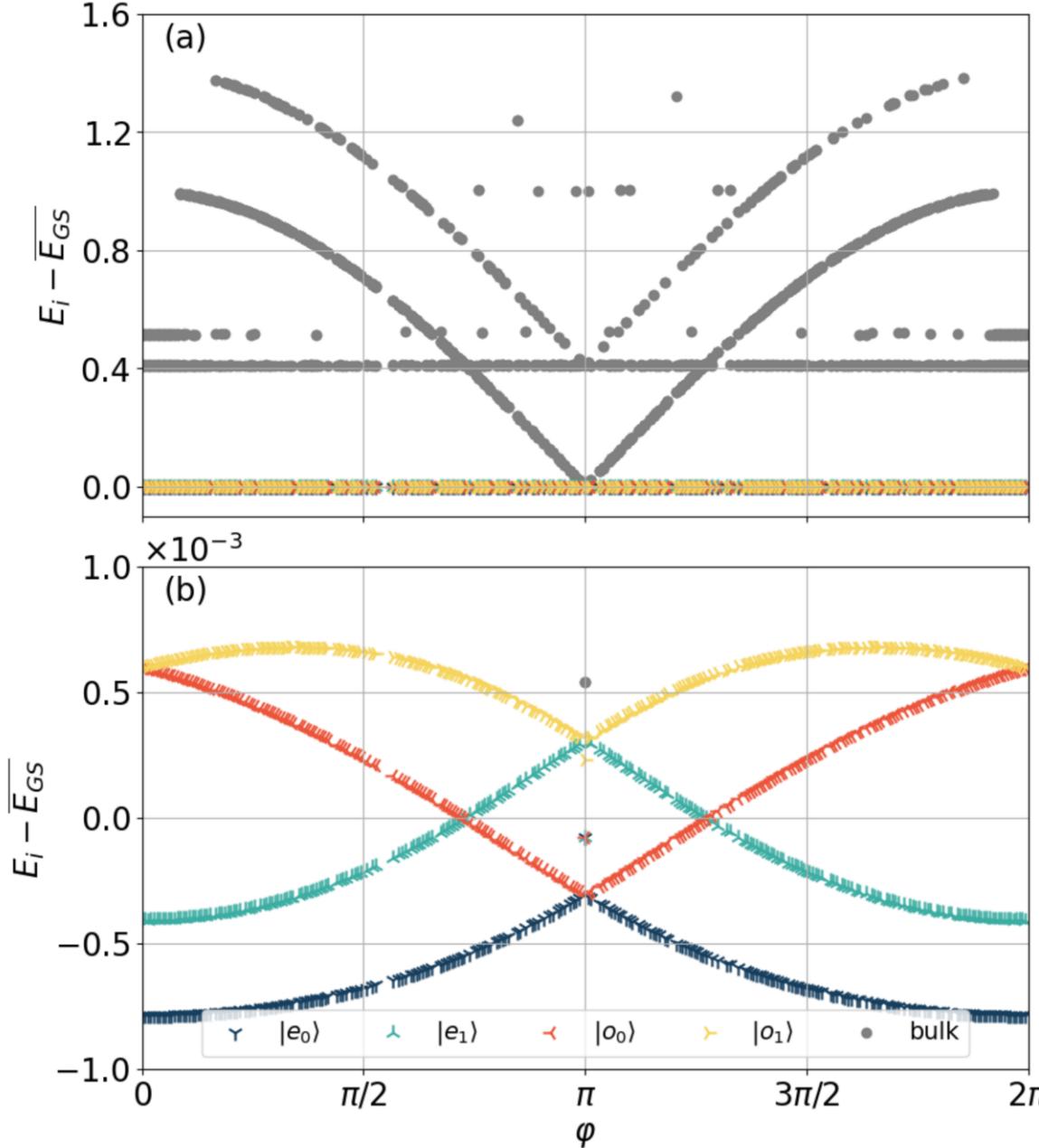
$$\begin{aligned} H_{int} &= \sum_{m\zeta} V_{m\zeta} \left[c_{m\zeta\uparrow}^\dagger c_{m\zeta\downarrow} c_{(m+1)\zeta\uparrow}^\dagger c_{(m+1)\zeta\downarrow} + \text{h.c.} \right] \\ &= \sum_{m,\zeta} \frac{V_{m,\zeta}}{2} \left[S_{m,\zeta}^x S_{m+1,\zeta}^x - S_{m,\zeta}^y S_{m+1,\zeta}^y \right] \end{aligned}$$



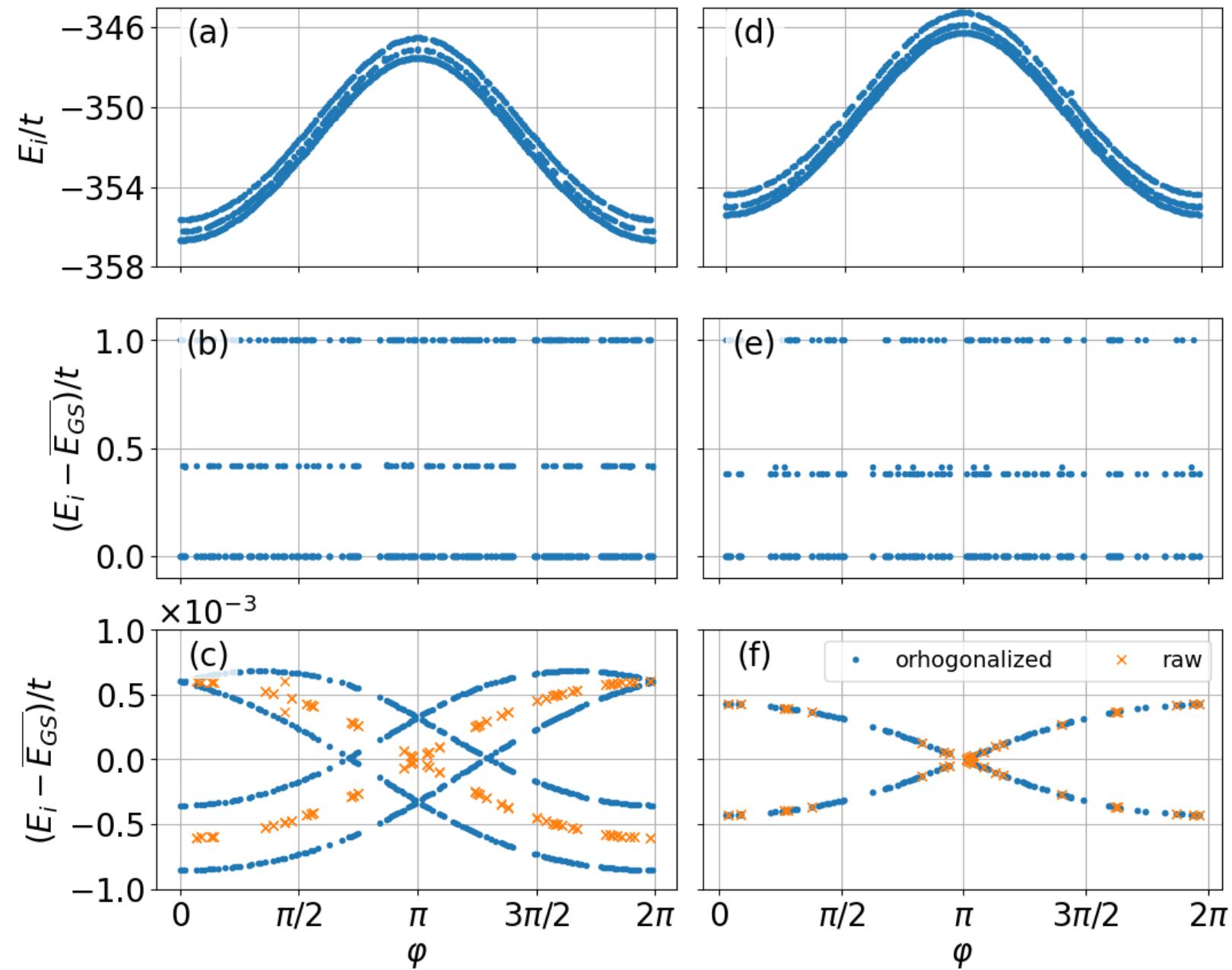
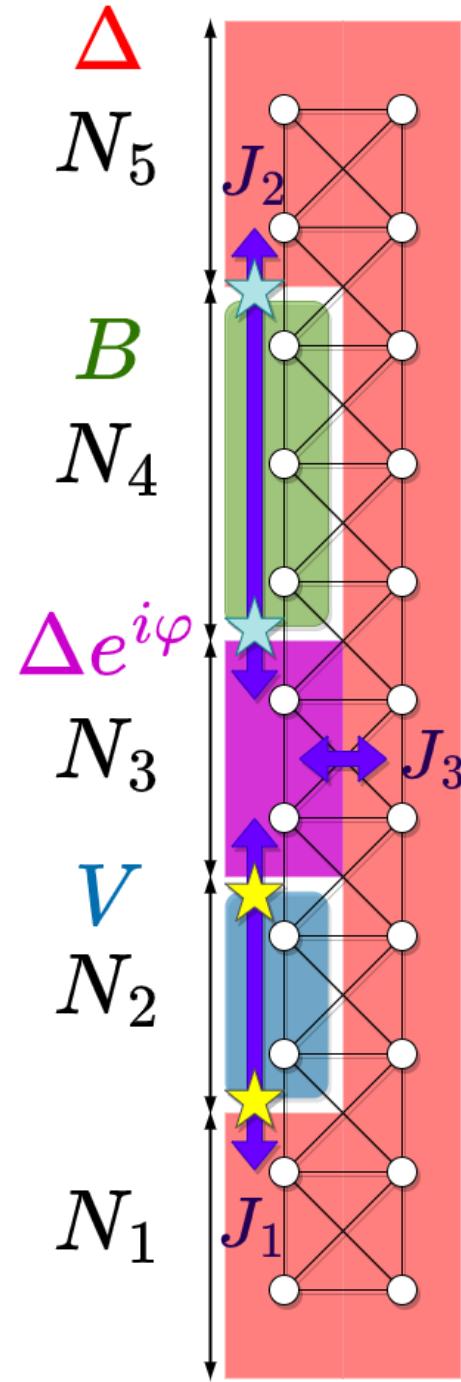
Finite-size DMRG calculations: Josephson spectrum



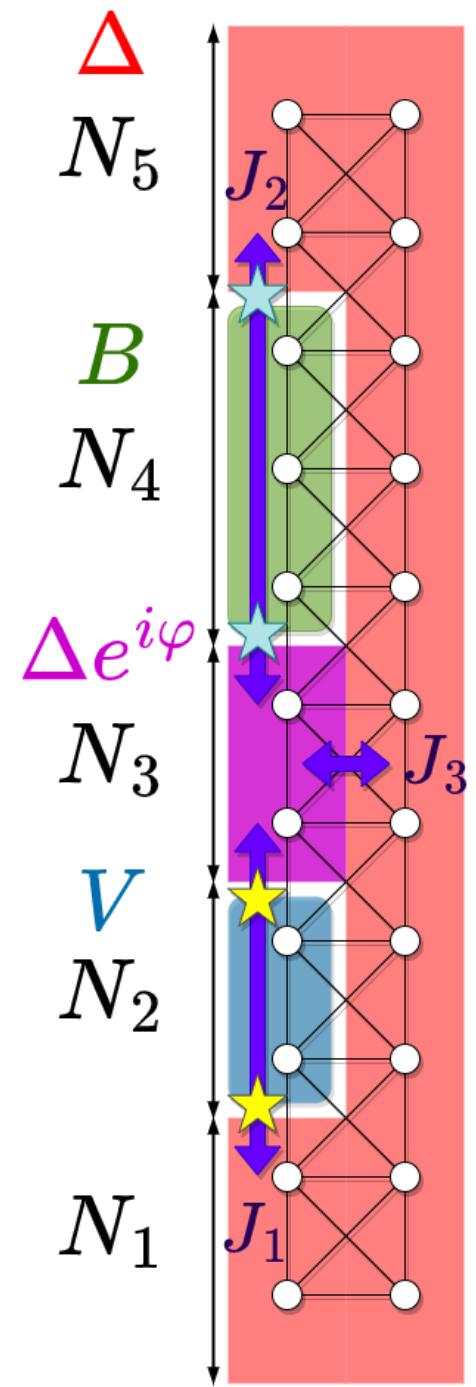
Josephson spectrum.. the details



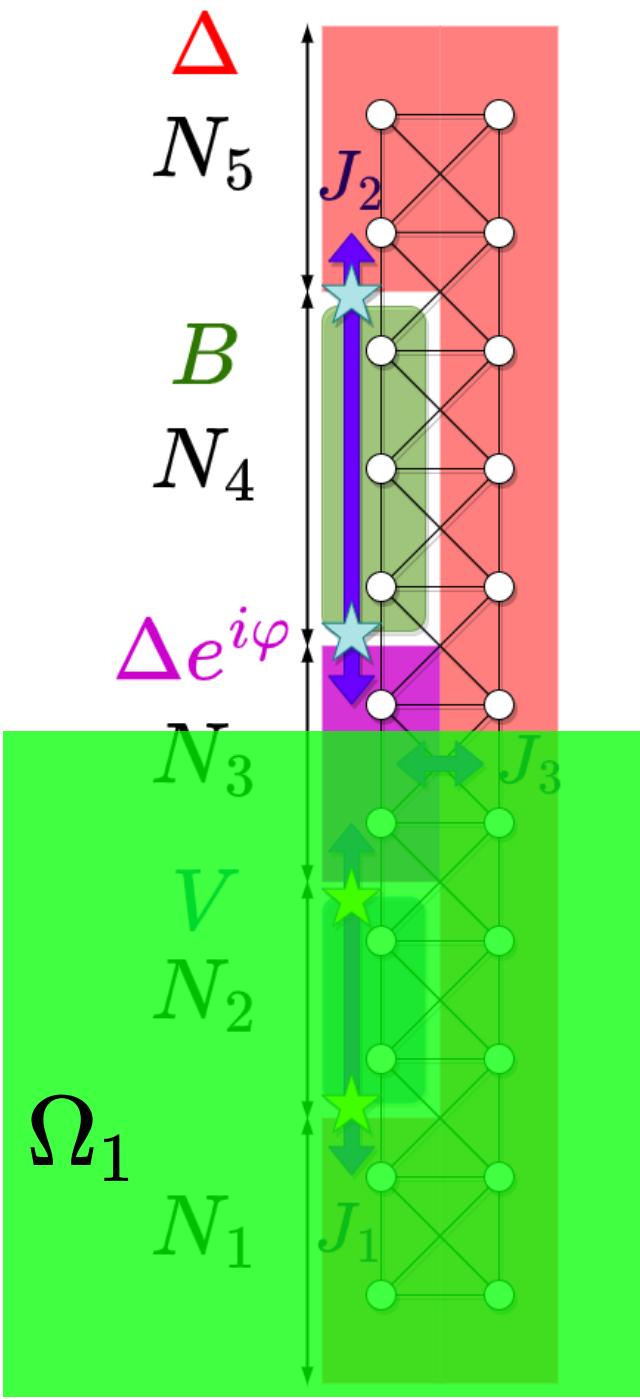
Josephson spectrum.. the details



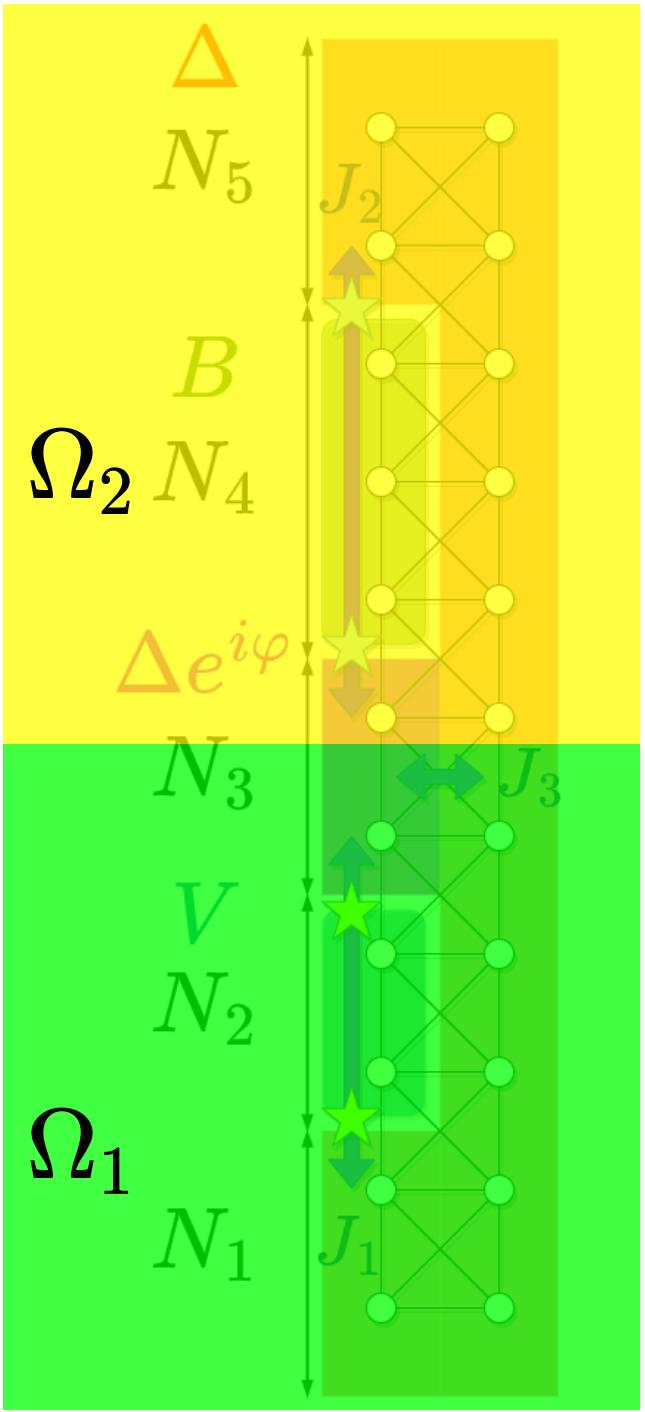
Local parity



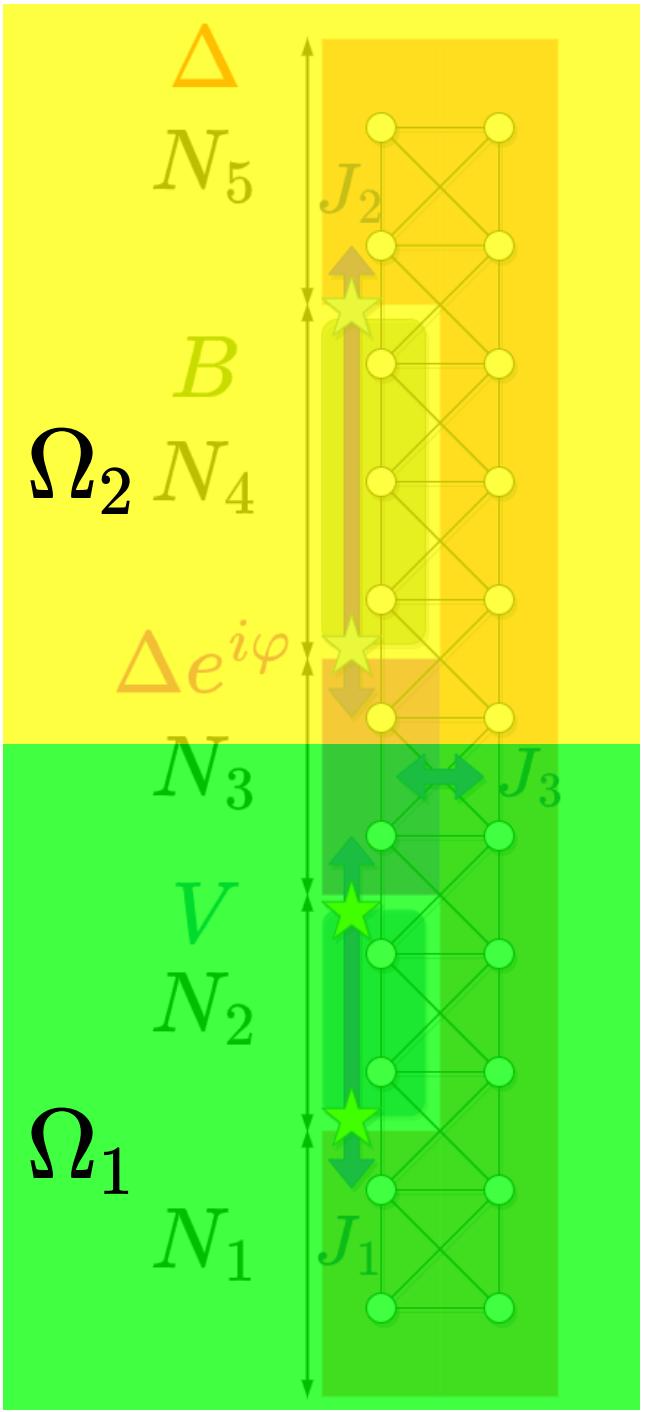
Local parity



Local parity

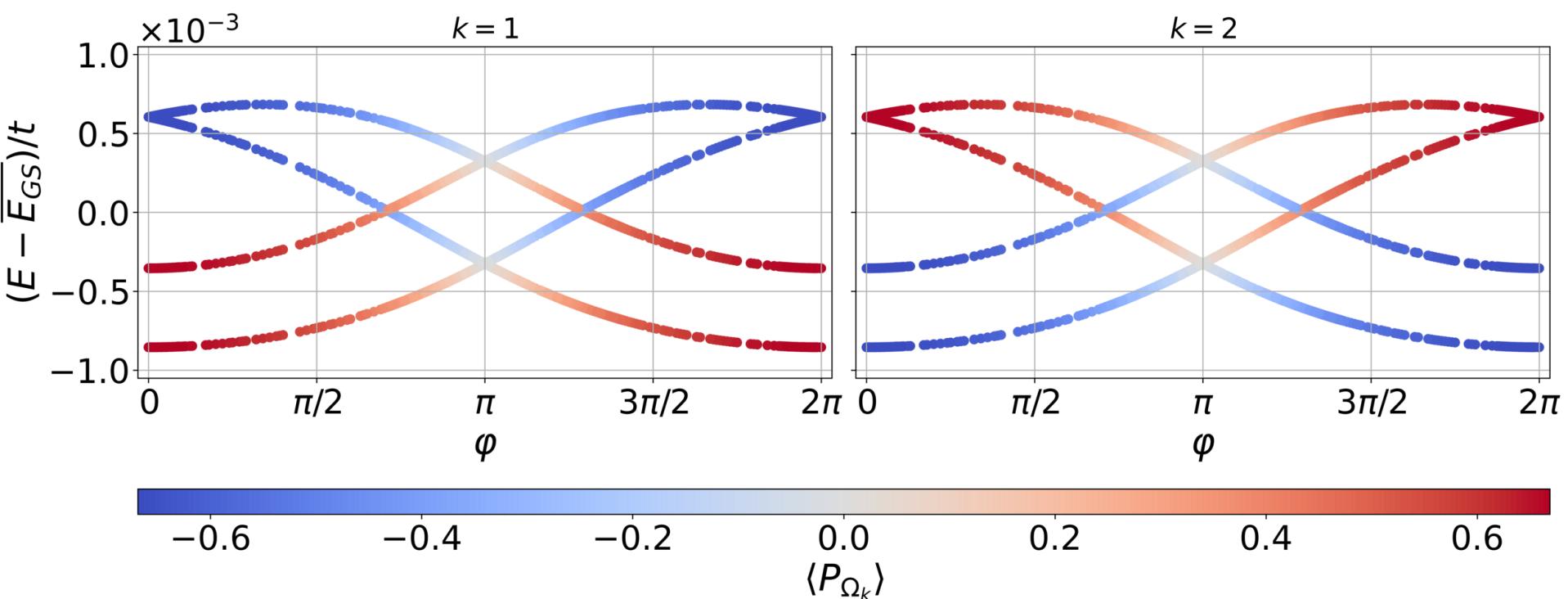
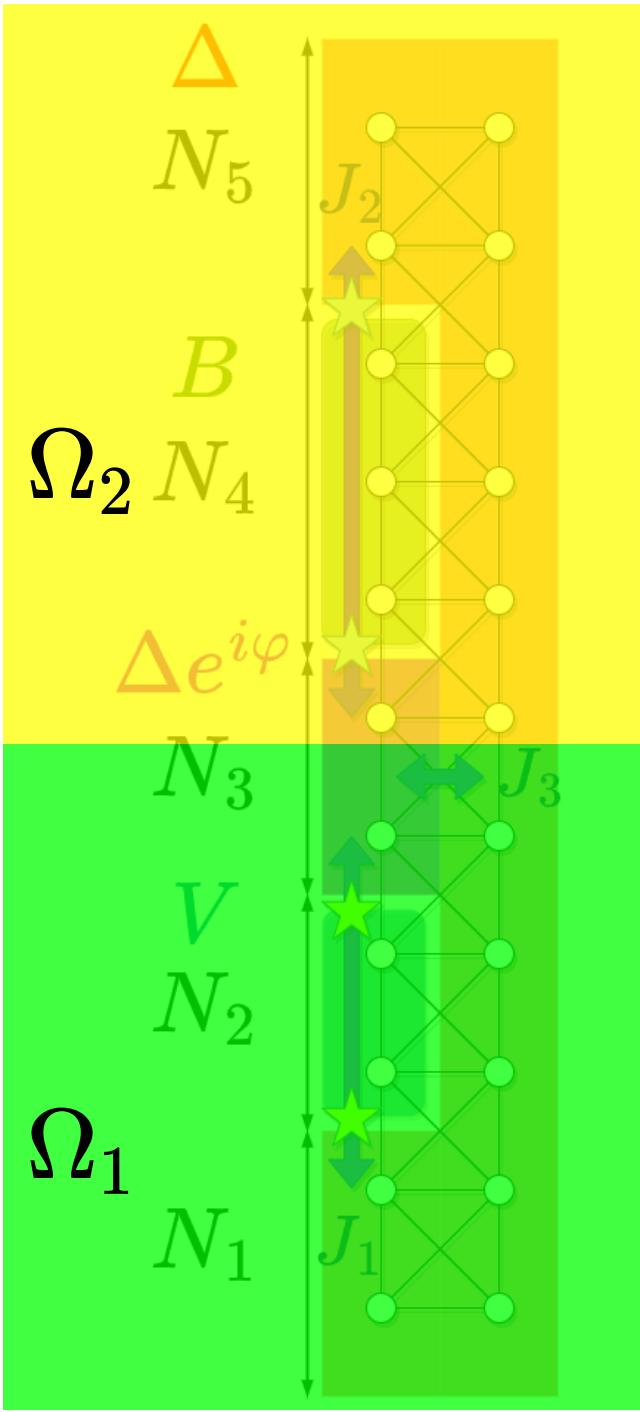


Local parity



$$P_\Omega = \prod_{p \in \Omega, \sigma} (-1)^{n_{p,\sigma}}$$

Local parity



$$P_\Omega = \prod_{p \in \Omega, \sigma} (-1)^{n_{p,\sigma}}$$

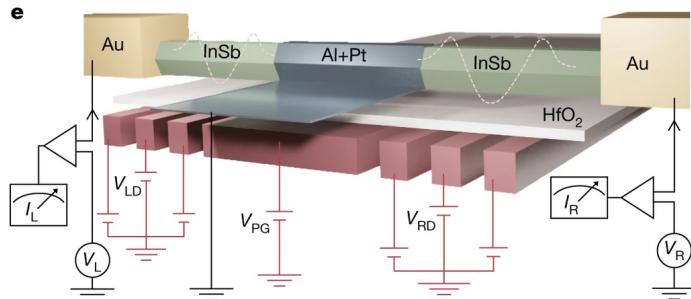
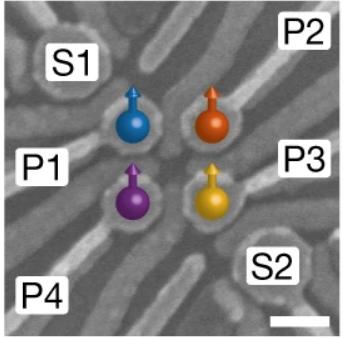
Parafermion signatures

- Robustness against disorder ✓
- Fourfold degenerate groundstate ✓
- Localized zero-energy excitations ✓
- Nontrivial (fractional) Josephson effect ✓

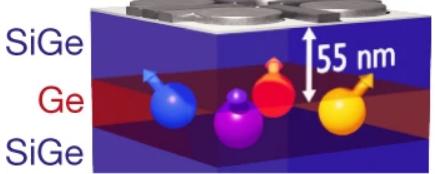
$\mathbb{Z}_4 \rightarrow 8\pi$ periodic

Quantum dot arrays for parafermions!!

a



b

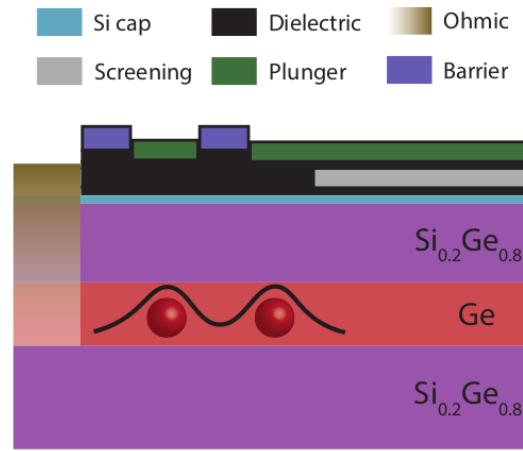
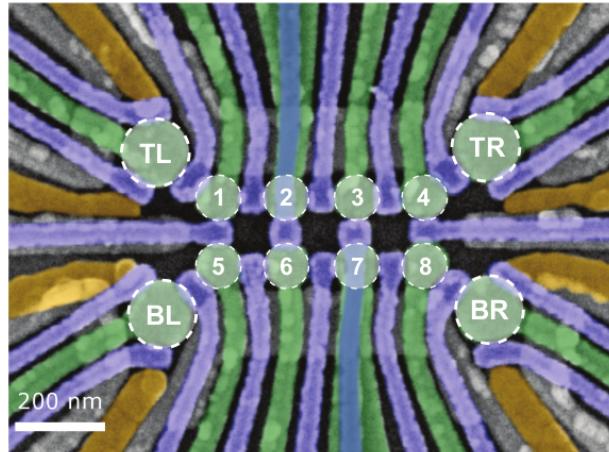


Dvir *et al.*

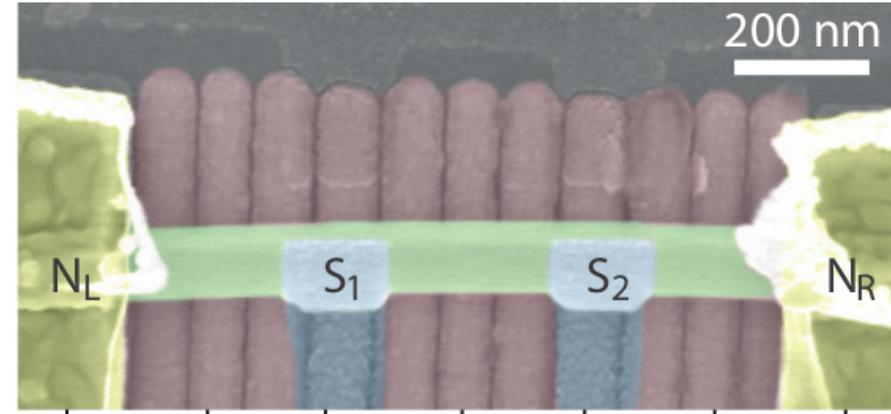
Nature **614**, 445 (2023)

Hendrickx *et al.*

Nature **591**, 580 (2021).

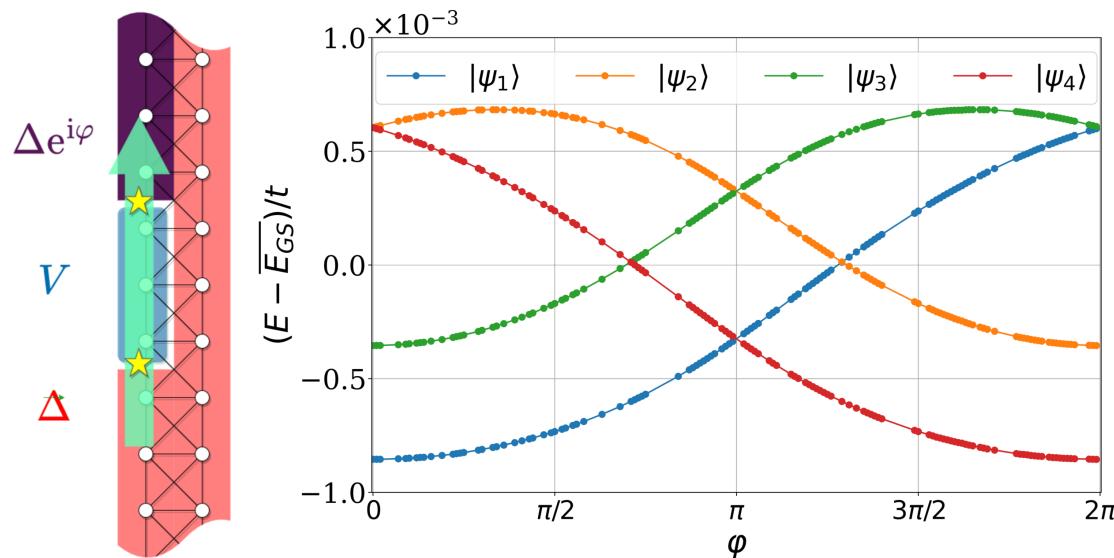
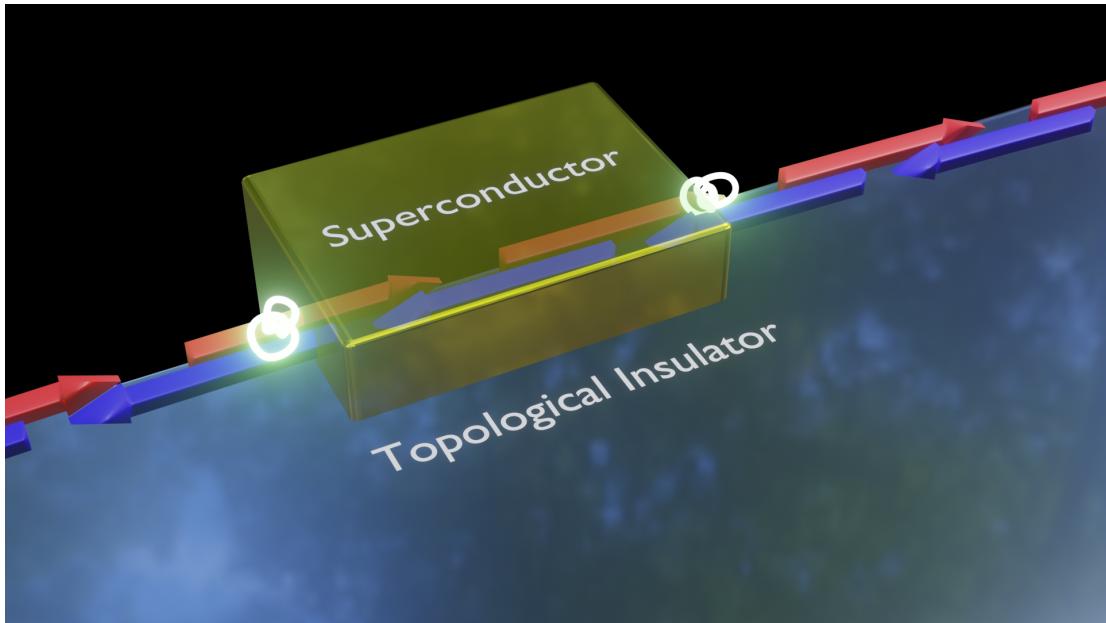


Hsiao *et al.* arXiv:2307.02401 (2023).



Bordin *et al.* arXiv:2306.07696

Summary



- Simple ladder model capable to capture physics at a single edge of a TI.
- Interactions and superconductivity are explicitly taken into account through DMRG calculations.
- Fourfold degeneracy and localized interface states can be realized.
- 8π Josephson spectrum \rightarrow parafermions!
- New realization avenues for parafermions in QD arrays are suggested.

<https://arxiv.org/abs/2311.07359>