

A recursive parameterisation of unitary matrices

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Abstract

A simple recursive scheme for parameterisation of n -by- n unitary matrices is presented. The n -by- n matrix is expressed as a product containing the $(n-1)$ -by- $(n-1)$ matrix and a unitary matrix that contains the additional parameters needed to go from $n-1$ to n . The procedure is repeated to obtain recursion formulas for n -by- n unitary matrices.

1 The parameterisation

It has been known since a long time that unitary transformations play a central role in physics. An excellent example is Wigner's paper of 1939 [1] which has had a great impact in the development of physics and is still important.

It is also a known fact that a general n -by- n unitary matrix $X^{(n)}$ may be expressed as a product of three unitary matrices,

$$X^{(n)} = \Phi^{(n)}(\vec{\alpha}) V^{(n)} \Phi^{(n)}(\vec{\beta}) \quad (1)$$

where the matrices Φ are diagonal unitary matrices,

$$\Phi^{(n)}(\vec{\alpha}) = \begin{pmatrix} e^{i\alpha_1} & & & \\ & e^{i\alpha_2} & & \\ & & \ddots & \\ & & & e^{i\alpha_n} \end{pmatrix} \quad (2)$$

$\Phi(\vec{\beta})$ is defined analogously; the α 's and β 's being real. The matrix $X^{(n)}$ has n^2 real parameters. In the following, for simplicity, the word parameter stands for real parameter. The quantities $\vec{\alpha}$ and $\vec{\beta}$ take care of $2n - 1$ parameters of $X^{(n)}$ because only the sums $\alpha_i + \beta_j$ enter, where i and j run from 1 to n . The remaining $(n - 1)^2$ parameters reside in the non-trivial matrix $V^{(n)}$ which is the subject of this study.

We start by putting $V^{(1)} = 1$, whereby $X^{(1)} = e^{i(\alpha_1 + \beta_1)}$ is the most general one-by-one unitary "matrix". For $n \geq 2$, we write the matrix $V^{(n)}$ in the form

$$V^{(n)} = \begin{pmatrix} V^{(n-1)} + (1 - c_n) |A^{(n-1)}\rangle\langle B^{(n-1)}| & s_n |A^{(n-1)}\rangle \\ s_n \langle B^{(n-1)}| & c_n \end{pmatrix} \quad (3)$$

Here we have introduced an angle denoted by θ_n and have used the common notation $s_n = \sin\theta_n$, $c_n = \cos\theta_n$. The complex vectors " $A^{(n-1)}$ " and " $B^{(n-1)}$ " have each $n-1$ components, i.e.,

$$|A^{(n-1)}\rangle = \begin{pmatrix} a_1^{(n-1)} \\ a_2^{(n-1)} \\ \vdots \\ a_{n-1}^{(n-1)} \end{pmatrix}, \quad |B^{(n-1)}\rangle = \begin{pmatrix} b_1^{(n-1)} \\ b_2^{(n-1)} \\ \vdots \\ b_{n-1}^{(n-1)} \end{pmatrix} \quad (4)$$

Furthermore,

$$\langle B^{(n-1)}| = (b_1^{(n-1)*}, b_2^{(n-1)*}, \dots, b_{n-1}^{(n-1)*}) \quad (5)$$

and

$$(|A^{(n-1)}\rangle \langle B^{(n-1)}|)_{ij} \equiv a_i^{(n-1)} b_j^{(n-1)*} \quad (6)$$

$A^{(n-1)}$ and $B^{(n-1)}$ are not arbitrary but are required to satisfy the conditions

$$\langle A^{(n-1)}|A^{(n-1)}\rangle = 1, \quad |B^{(n-1)}\rangle = -V^{(n-1)\dagger}|A^{(n-1)}\rangle \quad (7)$$

whereby

$$\langle B^{(n-1)}|B^{(n-1)}\rangle = 1, \quad |A^{(n-1)}\rangle = -V^{(n-1)}|B^{(n-1)}\rangle \quad (8)$$

We can easily check that if the matrix $V^{(n-1)}$, in Eq.(3), is unitary so is $V^{(n)}$. In order for $V^{(n)}$ to be the most general n -by- n unitary matrix, modulus the phase matrices Φ , it must have the required number of parameters. The vector $A^{(n-1)}$, having $n-1$ complex components, would seem to represent $2(n-1)$ parameters. But it has only $2(n-2)$ because it is normalised and its overall phase can be absorbed into the matrices Φ , i.e., the transformation

$$|A^{(n-1)}\rangle \rightarrow e^{i\eta}|A^{(n-1)}\rangle \quad (9)$$

yields

$$|B^{(n-1)}\rangle \rightarrow e^{i\eta}|B^{(n-1)}\rangle \quad (10)$$

and

$$V^{(n)} \rightarrow \Phi(0, 0, \dots, e^{-i\eta})V^{(n)}\Phi(0, 0, \dots, e^{i\eta}) \quad (11)$$

The parameter counting, therefore, goes as follows. On the LHS of Eq.(3) we need to have $(n-1)^2$ parameters. On the RHS, we have $(n-2)^2$ from $V^{(n-1)}$ and $2(n-2)$ from the vector $A^{(n-1)}$. Thus, together with the angle θ_n , the number of parameters is $(n-2)^2 + 2(n-2) + 1$ which equals $(n-1)^2$ as required.

We may use the relation (7) between $A^{(n-1)}$ and $B^{(n-1)}$ to rewrite the matrix $V^{(n)}$ in terms of either $A^{(n-1)}$ or $B^{(n-1)}$. In terms of $A^{(n-1)}$, we have

$$\begin{aligned} V^{(n)} &= \begin{pmatrix} 1 - (1 - c_n)|A^{(n-1)}\rangle \langle A^{(n-1)}| & s_n|A^{(n-1)}\rangle \\ -s_n \langle A^{(n-1)}| & c_n \end{pmatrix} \begin{pmatrix} V^{(n-1)} & 0 \\ 0 & 1 \end{pmatrix} \\ &\equiv A_{n,n-1} \begin{pmatrix} V^{(n-1)} & 0 \\ 0 & 1 \end{pmatrix} \end{aligned} \quad (12)$$

While writing the matrix in terms of $B^{(n-1)}$ yields

$$\begin{aligned} V^{(n)} &= \begin{pmatrix} V^{(n-1)} & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 - (1 - c_n)|B^{(n-1)}\rangle\langle B^{(n-1)}| & -s_n|B^{(n-1)}\rangle \\ s_n\langle B^{(n-1)}| & c_n \end{pmatrix} \\ &\equiv \begin{pmatrix} V^{(n-1)} & 0 \\ 0 & 1 \end{pmatrix} B_{n,n-1} \end{aligned} \quad (13)$$

These relations allow for a systematic construction of unitary matrices order by order. By repeating the above procedure for the matrix $V^{(n-1)}$ in terms of $A^{(n-2)}$ and $B^{(n-2)}$, and following down the chain we find the recursion formulas that we are looking for,

$$V^{(n)} = A_{n,n-1}A_{n,n-2}\dots A_{n,2}A_{n,1} \quad (14)$$

$$V^{(n)} = B_{n,1}B_{n,2}\dots B_{n,n-2}B_{n,n-1} \quad (15)$$

The matrices $A_{n,n-1}$ and $B_{n,n-1}$ were previously defined in Eqs.(12) and (13). For $j < n - 1$ we have

$$A_{n,j} = \begin{pmatrix} \begin{pmatrix} 1 - (1 - c_{j+1})|A^{(j)}\rangle\langle A^{(j)}| & s_{j+1}|A^{(j)}\rangle \\ -s_{j+1}\langle A^{(j)}| & c_{j+1} \end{pmatrix} & 0 \\ 0 & I_{n-j-1} \end{pmatrix} \quad (16)$$

$$B_{n,j} = \begin{pmatrix} \begin{pmatrix} 1 - (1 - c_{j+1})|B^{(j)}\rangle\langle B^{(j)}| & -s_{j+1}|B^{(j)}\rangle \\ s_{j+1}\langle B^{(j)}| & c_{j+1} \end{pmatrix} & 0 \\ 0 & I_{n-j-1} \end{pmatrix} \quad (17)$$

Here I_{n-j-1} is the unit matrix of order $n - j - 1$. The two unitary matrices $A_{n,j}$ and $B_{n,j}$ are related by

$$A_{n,j} = \begin{pmatrix} V^{(j)} & 0 \\ 0 & I_{n-j} \end{pmatrix} B_{n,j} \begin{pmatrix} V^{(j)\dagger} & 0 \\ 0 & I_{n-j} \end{pmatrix} \quad (18)$$

2 Simple examples

The simplest case is $n = 2$ for which we take $|A^{(1)}\rangle = 1$ whereby $|B^{(1)}\rangle = -1$. Using $V^{(1)} = 1$ we obtain, from Eqs.(12) and (13)

$$V^{(2)} = A_{2,1} = B_{2,1} = \begin{pmatrix} c_2 & s_2 \\ -s_2 & c_2 \end{pmatrix} \quad (19)$$

This is the familiar rotation matrix $R_2(\theta_2)$, θ_2 being the rotation angle in two dimensions.

The next simplest case is $n = 3$ for which we may either use the "mixed form" or the pure forms. For the mixed form we have

$$V^{(3)} = \begin{pmatrix} R_2(\theta) + (1 - c_3)|A^{(2)}\rangle\langle B^{(2)}| & s_3|A^{(2)}\rangle \\ s_3\langle B^{(2)}| & c_3 \end{pmatrix} \quad (20)$$

Here $R_2(\theta)$ is again the rotation matrix in Eq.(19) and we may put

$$|A^{(2)}\rangle = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix}, \quad |B^{(2)}\rangle = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} \quad (21)$$

From Eq.(7) follows that

$$|A^{(2)}\rangle = -R_2(\theta)|B^{(2)}\rangle, \quad |B^{(2)}\rangle = -R_2(-\theta)|A^{(2)}\rangle \quad (22)$$

Hence these vectors represent two parameters, for example

$$|A^{(2)}\rangle = \begin{pmatrix} \cos\gamma \\ \sin\gamma e^{i\delta} \end{pmatrix} \quad (23)$$

where γ and δ are real. The pure forms are also obtained very simply, for example

$$V^{(3)} = \begin{pmatrix} 1 - (1 - c_3)|A^{(2)}\rangle\langle A^{(2)}| & s_3|A^{(2)}\rangle \\ -s_3\langle A^{(2)}| & c_3 \end{pmatrix} \begin{pmatrix} R_2(\theta_2) & 0 \\ 0 & 1 \end{pmatrix} \quad (24)$$

where $A^{(2)}$ is as defined in Eq.(23). Evidently, depending on the application one has in mind some choices may be more convenient than others. This is demonstrated in Refs.[2] and [3] which deal with the so called quark and lepton mixing matrices. The essential point is that $V^{(3)}$ is described in a rather simple fashion by four parameters as it should be. For the case of $n = 4$ we could, for example, take $A^{(3)}$ to be

$$|A^{(3)}\rangle = \begin{pmatrix} \cos\rho \\ \sin\rho\cos\sigma e^{i\delta_1} \\ \sin\rho\sin\sigma e^{i\delta_2} \end{pmatrix} \quad (25)$$

where ρ , σ , δ_1 and δ_2 are the four parameters needed to define the most general $A^{(3)}$.

In principle, the above recursive procedure may "easily" be extended to much larger n with the help of computers.

References

- [1] E. P. Wigner, Ann. Math. **40** (1939) 149; see also E. P. Wigner, "Symmetries and Reflections" (MIT Press 1970)
- [2] C. Jarlskog, hep-ph/0503199, to appear in Phys. Lett. B
- [3] C. Jarlskog, hep-ph/0504012