Entanglement Islands in Braneworlds

Based on: arxiv: 2006.02438, JHEP 09 (2020) 121;

arxiv: 2012.04671;

arxiv: 2103.07477;

arxiv: 2103.xxxxx.

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Plan:

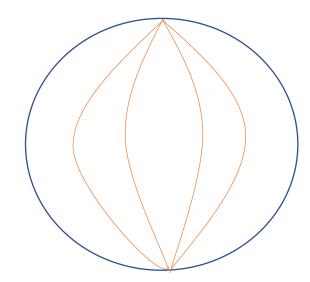
- Brief Review of Karch-Randall (KR) Brane-world (Karch and Randall 01)
- An Information Paradox (Hartman, Maldacena 13; Mathur 14; Almheri, Mahajan and Maldacena 19)
- Quantum Extremal Surface and Double Holography (Almheri, Mahajan and Santos 19)
- Analytic Resolution of the Paradox in KR (HG and Karch 20)
- Massive Graviton and Islands
- An Attempt to Massless: Information Transfer with a Gravitating Bath (HG, Karch, Perez, Raju, Randall, Riojas and Shashi 20)
- A Parallel Story in de Sitter (HG, Nomura and Sun to appear)

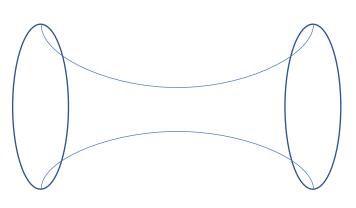
KR Brane-world and Its Three Descriptions

(Karch and Randall 01, a recent generalization: Karch and Randall 20)

Basics of AdS:

1) AdS₅ space:
$$-X_0^2 - X_1^2 + X_2^2 + X_3^2 + X_4^2 + X_5^2 = -L^2$$





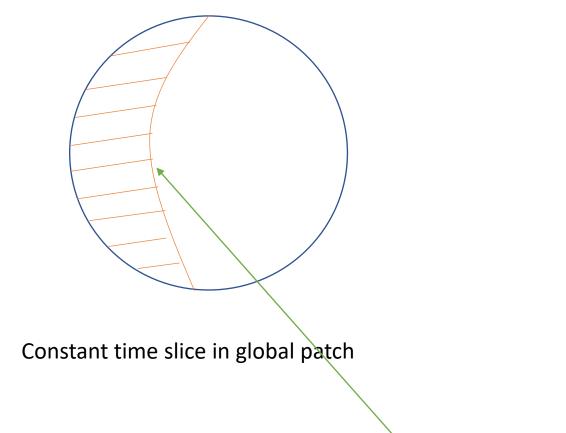
AdS₅ in embedding space

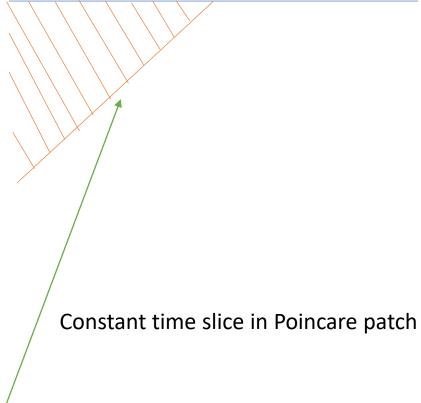
AdS₅ constant time slice in global coordinate

2)
$$AdS_5/AdS_4: -X_0^2-X_1^2+X_2^2+X_3^2+X_4^2=-L^2 \cosh^2(r)$$

 $X_5^2=L^2 \sinh^2(r)$

• Karch-Randall Brane:





Karch-Randall Brane AdS₄

$$S = \frac{1}{16\pi G} \int d^5x \, \sqrt{-g} (R - 2\Lambda) - \frac{1}{8\pi G} \int_{brane} d^4x \, \sqrt{-h} \, (K + T)$$

• Three Descriptions of KR Brane-world:

(1) 5d Einstein Gravity in AdS₅ with a Karch-Randall Brane:



holography 2

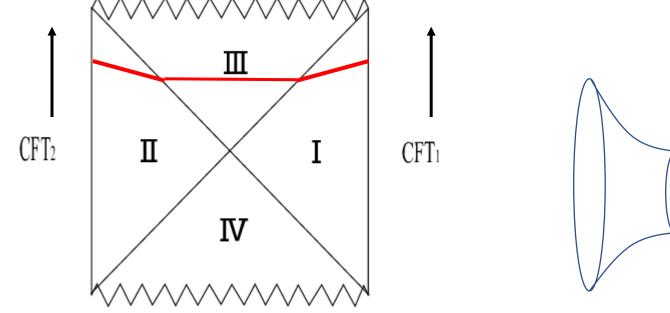
(2) 4d (holographic) CFT Coupled to a 4d Gravitating (holographic) CFT via Transparent Boundary Condition:

(3) 4d Boundary (holographic) Conformal Field Theory (BCFT):

An Information Paradox

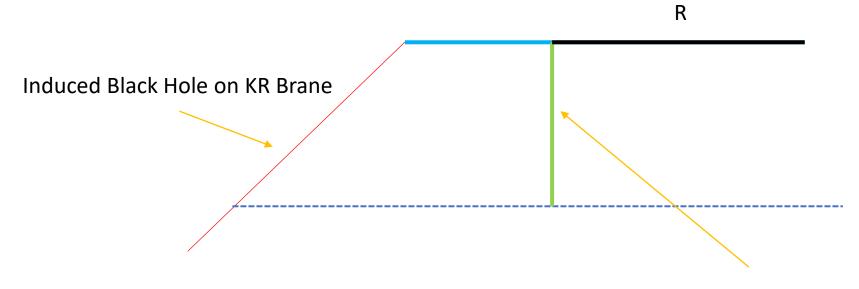
(Hartman, Maldacena 13; Mathur 14; Almheri, Mahajan and Maldacena 19)

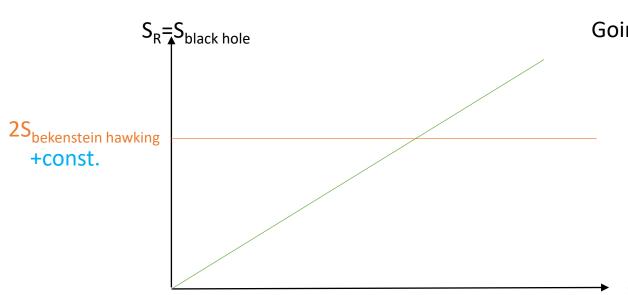
• Eternal Black Hole in AdS:





• The Paradox:





Hartman-Maldacena Surface, Going Through ER Bridge, Linear Growing with Time

Quantum Extremal Surface and Double Holography

(3) 4d Boundary (holographic) Conformal Field Theory (BCFT):



$$S_R = -Tr \rho_R log \rho_R$$

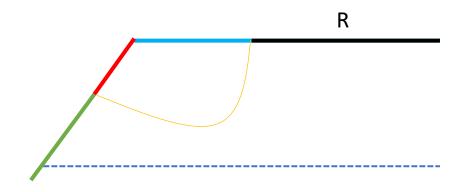
(2) 4d (holographic) CFT Coupled to a 4d Gravitating (holographic) CFT via Transparent





$$S_R = S_{gen}(EW(R)) = S_{gen}(R U I)$$

(1) 5d Einstein Gravity in AdS₅ with a Karch-Randall Brane:



$$S_R = S_{gen}(R U I) = S_{gen}(EW(R U I)) = \frac{A(\partial (EW(R U I)))}{4G_E}$$

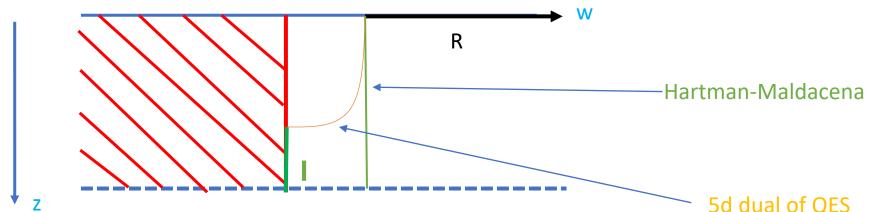
Analytic Resolution of the Paradox in KR

(HG and Karch 20)

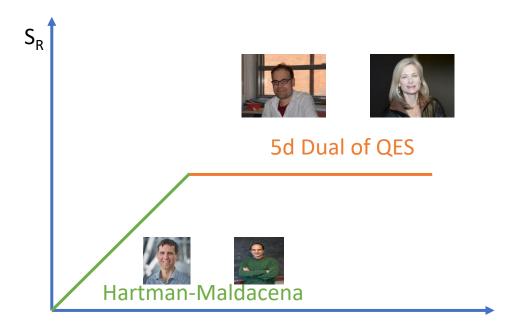
- Problem (Almheri, Mahajan and Santos 19):
 - The bulk geometry in general is very complicated and requires numerical GR
- Analytic Model (HG and Karch 20):

Simple Observation: If the brane is tensionless then there is no back reaction at all!

$$AdS_{5}Plannar\ Black\ Hole:\ ds^{2} = \frac{1}{z^{2}}(-h(z)dt^{2} + \frac{dz^{2}}{h(z)} + dw^{2} + dx^{2} + dy^{2}),\ h(z) = \frac{1}{1 - \frac{z^{4}}{z_{H}^{4}}}$$



The Page Curve:



Massive Graviton and Islands

(HG and Karch 20)

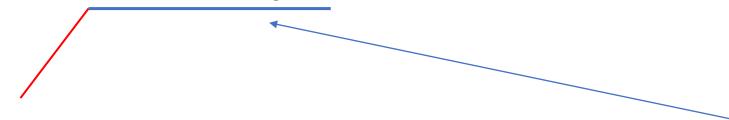
- Massive Graviton in KR:
 - (2) 4d (holographic) CFT Coupled to a 4d Gravitating (holographic) CFT via a Transparent Boundary Condition.

$$\partial_{\alpha}T^{\alpha\beta} \neq 0$$



one-loop correction of graviton is not protected from being zero and this gives nonzero mass correction (Porrati 02)

(1) 5d Einstein Gravity in AdS₅ with a Karch-Randall Brane:



The massless KK mode of 5d graviton is not normalizable because part of the conformal boundary is not cutoff.

General Lessons:

- a) As long as a bulk gravitational theory is coupled to a boundary bath the gravitational theory is modified at one-loop level.
- b) Existing higher dimensional calculations of the Page curve of black hole radiation using AdS/CFT are all in the context of massive gravity.

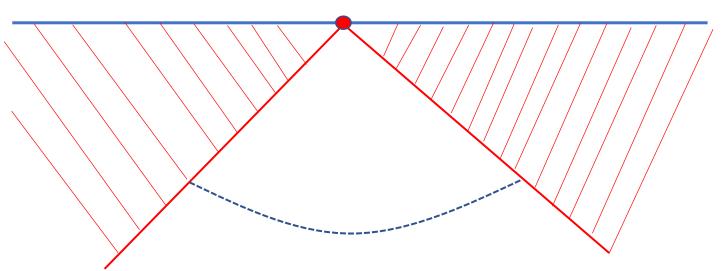
(massive gravity —— Page curve)

C) Comment on massless gravity:

The massless mode can be recovered by cutting off the leftover part of the conformal boundary by another KR brane. But this makes the bath gravitating. And we see a constant Page curve for radiation in this context. (HG, Karch, Perez, Raju, Randall, Riojas and Shashi 20)

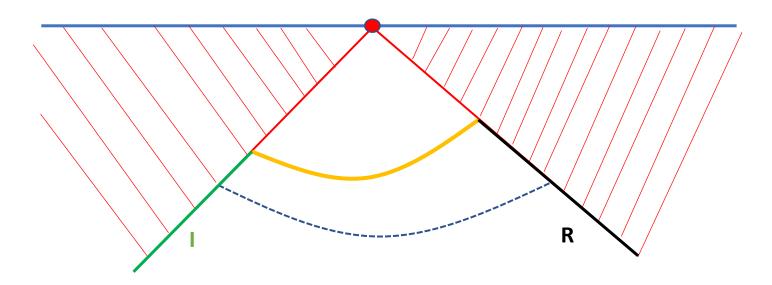
An Attempt to Massless: Information Transfer with a Gravitating Bath

(HG, Karch, Perez, Raju, Randall, Riojas and Shashi 20)



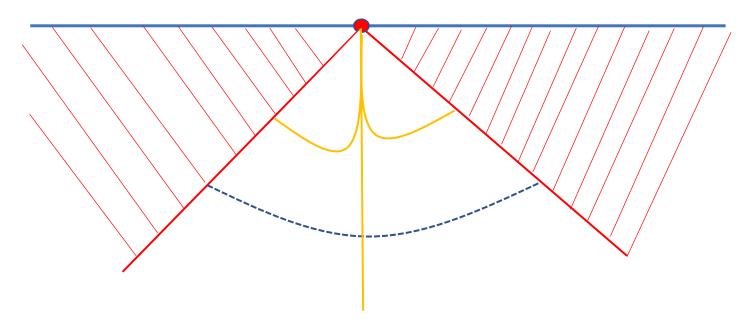
- The bulk geometry is AdS₅ black string which is foliated by AdS₄ black holes.
- We embed two Karch-Randall branes into the bulk which rescues the massless graviton mode.
- Bath is gravitating so there is no freedom to specify the radiation at the fine-grained level.

A Dynamical Principle and a Page Curve



- $S_R = S_{gen}(R \cup I) = S_{gen}(EW(R \cup I)) = \frac{A(\partial (EW(R \cup I)))}{4G_5}$ and now we show minimize over both ∂R and ∂I .
- The RT surface is black string horizon. Therefore, we have a constant Page curve. Consistent with Maldacena 01 and Laddha, Prabhu, Raju and Shrivastava 20.

Another Question



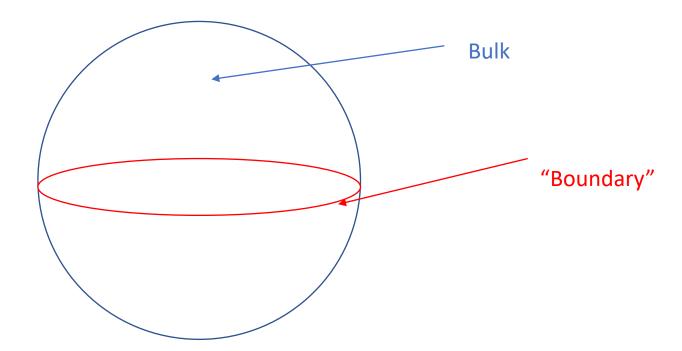
- An internal bipartition of the Hilbert space into L and R sectors.
- Three candidates entangling surfaces for $S_{L/R}$.
- Indeed, we can have a time dependent Page curve if the gravity on the brane is weak enough.
- We have an explicit CFT computation. (Geng, Lüst, Mishra and Wakeham to appear)

An Information Paradox in de Sitter and Its Resolution

(HG, Nomura and Sun to appear)

The DS/dS Correspondence:

Quantum gravity in a (d+1)-dimensional de Sitter spacetime is dual to a field theory system living on a d-dimensional de Sitter space. The field theory system consists of two UV-cutoff conformal field theories coupled to dynamical gravity.



An Information Paradox

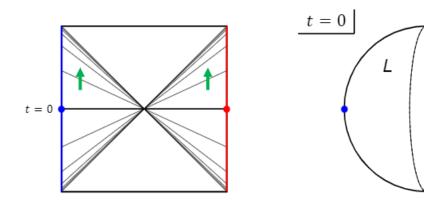


FIG. 1. The Penrose diagram of dS spacetime in the extended static patch (left) and its spatial section at t=0 (right). The blue and red dots represent the north and south poles at t=0. The green arrows represent the direction of time evolution.

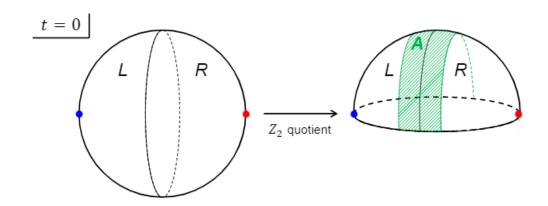


FIG. 2. The \mathbb{Z}_2 orbifolding removes a half of dS_d^2 . The region A, of which we calculate the entanglement entropy, at t=0 is depicted as the green shaded region on the right panel.

There would be a paradox if S_A is bigger than the minimal Hilbert space dimension of A and \bar{A} or more severely if it is imaginary.

Embedding the Paradox into DS/dS

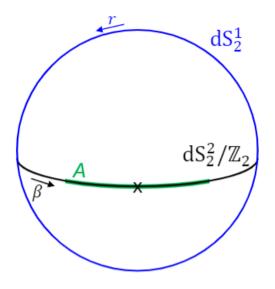
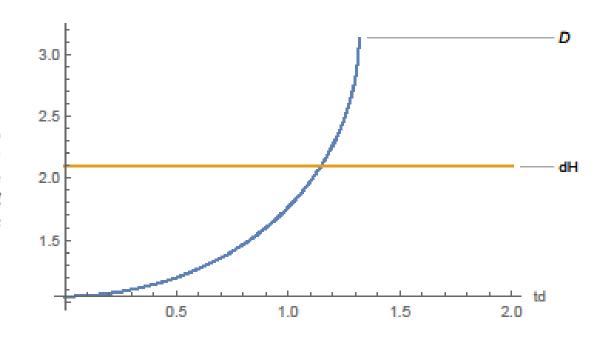


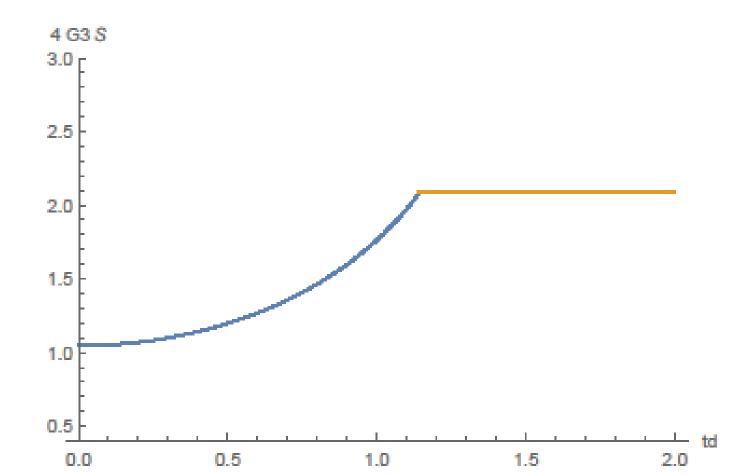
FIG. 3. The spatial section at t = 0 for d = 2. The two intersecting dS systems are depicted as the blue circle (dS_2^1) and black semicircle (dS_2^2/\mathbb{Z}_2) . The green segment represents region A, of which we calculate the entanglement entropy, and the cross represents the horizon at $\beta = \frac{\pi}{2}$. The holographic bulk space (dS_3/\mathbb{Z}_2) is the hemisphere bounded by dS_2^1 .

$$S_{A} = \frac{D}{4G3} = \frac{Arccos(1 - 2cos2(\boldsymbol{\beta}*)cosh2(t))}{4G3}$$



Resolution

The island Ryu-Takayanagi surface is just \bar{A} .



Lessons

- 1) A time-dependent Page curve.
- 2) The graviton is indeed massless.
- 3) The non-factorizability issue in Laddha, Prabhu, Raju and Shrivastava 20 is automatically avoided.

A Question

Is this a fine-grained entropy?

Objection: No, it isn't. Because there is no diffeomorphic invariant way to specify a region in a gravitating spacetime.

Support: Well, the static patch is associated to a specific observer. Hence the diffeomorphic invariance has been broken on the moment we are focusing on a static patch. The observer can always specify a region covariantly using light.

Summary

- A time-dependent Page curve of black hole radiations in AdS requires the gravitational theory to be massive. (HG, Karch 20)
- An attempt to recuse the massless graviton gives wedge holography. No time-dependent Page curve of black hole radiations. But a timedependent Page curve for an internal bipartition of the wedge holographic dual. (HG, Karch, Perez, Raju, Randall, Riojas and Shashi 20)
- A CFT computation of the Page curve for the internal bipartition exists.

(HG, Lüst, Mishra and Wakeham to appear)

Extension of the entanglement island program into de Sitter space.

(HG, Nomura and Sun to appear)

Thank You!