Derived equations for the polarization/sensitive FP microscopy

The electric fields passing through an anisotropic sample encounter polarization and the magnitude and phase will be changed. This phenomenon can be interpreted as:

$$\begin{pmatrix} E'_{x} \\ E'_{y} \end{pmatrix} = \begin{pmatrix} J_{xx} & J_{xy} \\ J_{yx} & J_{yy} \end{pmatrix} \begin{pmatrix} E_{x} \\ E_{y} \end{pmatrix} (1)$$

The Jones matrix includes the corresponding matrices for polarizer, retarder and rotator

$$J' = J(-\theta) J_R J(\theta) (2)$$

$$J' = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} p_x & 0 \\ 0 & p_y \end{pmatrix} \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix}$$
(3)

J has been defined for linear polarizer here.

Accordingly, the Jones matrix for a rotated polarizer is:

$$J_{p} = \begin{pmatrix} p_{x}cos^{2}\theta + p_{y}sin^{2}\theta & (p_{x} - p_{y})sin\theta cos\theta \\ (p_{x} - p_{y})sin\theta cos\theta & p_{x}sin^{2}\theta + p_{y}cos^{2}\theta \end{pmatrix} (4)$$

Therefore, for an ideal linear polarizer (each channel in this paper), it simplifies as:

$$J_p(\theta) = \begin{pmatrix} \cos^2 \theta & \sin \theta \cos \theta \\ \sin \theta \cos \theta & \sin^2 \theta \end{pmatrix} (5)$$

By considering a retardation phase as δ , the Jones matrix of a retarder in the left/handed circular polarization satate, will be:

$$\begin{pmatrix} E'_{x} \\ E'_{y} \end{pmatrix} = \begin{pmatrix} e^{\frac{-j\delta}{2}} & 0 \\ 0 & e^{\frac{j\delta}{2}} \end{pmatrix} \begin{pmatrix} E_{x} \\ E_{y} \end{pmatrix}, (5)$$

$$J_{R}(\delta) = \begin{pmatrix} e^{\frac{-j\delta}{2}} & 0 \\ 0 & e^{\frac{j\delta}{2}} \end{pmatrix}$$

It means that the total phase shift δ is halve into two orthogonally polarized components by assigning a phase advance of $\delta/2$ to the fast y axis and phase retardation of $-\delta/2$ to the slow x axis. It should be considered that these values have inverse sign in the case of right-handed circular polarization.

According to (2),

$$J(\delta,\theta) = \begin{pmatrix} \cos\left(\frac{\delta}{2}\right) - j\sin\left(\frac{\delta}{2}\right)\cos\left(2\theta\right) & -j\sin\left(\frac{\delta}{2}\right)\sin\left(2\theta\right) \\ -j\sin\left(\frac{\delta}{2}\right)\sin\left(2\theta\right) & \cos\left(\frac{\delta}{2}\right) + j\sin\left(\frac{\delta}{2}\right)\cos\left(2\theta\right) \end{pmatrix} = \begin{pmatrix} J_{xx} & J_{xy} \\ J_{yx} & J_{yy} \end{pmatrix}$$

Intensity is generally defined as $I' = E^*$. E

And hence, we have:

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Based upon the paper, $E_{in} = \sqrt{\frac{I}{2}} [1 \ i]^T$ that for the LH-CP should be $E_{in} = \sqrt{\frac{I}{2}} [1 \ -i]^T$. Otherwise equations of 5-8 are not verified.

For 0 deg:

$$E_x = J_{xx} + j J_{xy}, E_y = 0$$

$$I_0 = \left(\frac{I}{2}\right) \cdot |E_x|^2 = (I/2)(1 - \sin\delta\sin(2\theta))$$

Similarly for the other polarization states, we have:

For 45 deg:

$$I_{45} = \left(\frac{I}{2\sqrt{2}}\right) \left| \binom{J_{xx} + J_{yx} - j(J_{xy} + J_{yy})}{J_{xx} + J_{yx} - j(J_{xy} + J_{yy})} \right|^2 = (I/2)(1 + \sin\delta\cos(2\theta))$$

For 90 deg:

$$E_x = 0$$
, $E_y = J_{yx} - j J_{yy}$
 $I_{90} = \left(\frac{I}{2}\right) \cdot \left|E_y\right|^2 = (I/2)(1 + \sin\delta\sin(2\theta))$

For 135 deg:

$$I_{135} = \left(\frac{I}{2\sqrt{2}}\right) \left| \begin{pmatrix} (J_{xx} - J_{yx}) + j(J_{xy} - J_{yy}) \\ -(J_{xx} - J_{yx}) - j(J_{xy} - J_{yy}) \end{pmatrix} \right|^2 = (I/2)(1 - \sin\delta\cos(2\theta))$$

Until now, we do not have information about retardation (δ) and rotation (θ) . Furthermore, the intensity of incoming wave (I) has not been defined. To remove the dependency to I, two auxiliary variables are defined as v_1 and v_2 .

$$\begin{split} I_{90} + I_0 &= I \\ I_{135} + I_{45} &= I \\ I_{90} - I_0 &= I(sin\delta\sin(2\theta)) \\ I_{45} - I_{135} &= I(sin\delta\cos(2\theta)) \\ v_1 &= (I_{90} - I_0)/(I_{90} + I_0) \,, \qquad v_2 &= (I_{45} - I_{135})/(I_{45} + I_{135}) \end{split}$$

$$v_1^2 + v_2^2 = \sin \delta^2$$

$$\frac{v_1}{v_2} = \tan(2\theta)$$

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