# **GENERAL PHYSICS PH1110**

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# 1. MECHANICS

### 1.5 DYNAMICS OF RIGID BODIES

- 1 RIGID BODY
- 2 Center of mass
- 3 Translation and rotation
- 4 ROTATIONAL ENERGY AND MOMENT OF INERTIA
- 5 Torque of a force
- 6 Angular momentum
- 7 Torque and angular acceleration
- 8 Torque and angular momentum

**Mechanics** ▷ Dynamics of rigid bodies

1. Rigid body

**Mechanics** > Dynamics of rigid bodies

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**Mechanics** ▷ Dynamics of rigid bodies

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  - ▶ Some results studied in this chapter can be applied for solid bodies which are not rigid.

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2. Center of mass (CM)

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**Mechanics** ▷ Dynamics of rigid bodies

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### **Mechanics** > Dynamics of rigid bodies

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The <u>CM</u> of a system is the point where the weighted relative position of the distributed mass sums to zero.

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• For a rigid body, the general motion is the combination of the translational and rotational motions.

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- Therefore, in studying a pure translational motion, we only need to study the motion of the CM of the body.

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The vectors  $\overrightarrow{\omega}$  and  $\overrightarrow{\beta}$  lie on the z-axis; direction of  $\overrightarrow{\omega}$  depends on the sign of  $\dot{\varphi}$ ; direction of  $\vec{\beta}$  depends on the sign of  $\ddot{\varphi}$ .

$$\hat{\mathbf{p}} \times \hat{\mathbf{p}} = \hat{\mathbf{z}}$$
,

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,

$$\hat{\mathbf{\phi}} \times \hat{\mathbf{z}} = \hat{\mathbf{\rho}},$$

$$\hat{\mathbf{z}} \times \hat{\boldsymbol{\rho}} = \hat{\boldsymbol{\phi}}$$

$$\hat{\mathbf{o}} \times \hat{\mathbf{o}} = \hat{\mathbf{z}}$$

$$\hat{\boldsymbol{\varphi}} \times \hat{\mathbf{z}} = \hat{\boldsymbol{\rho}},$$

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### **Mechanics** ▷ Dynamics of rigid bodies

### 3. Translation and rotation

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,

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### **Mechanics** ▷ Dynamics of rigid bodies

## 3. Translation and rotation

• The unit vectors  $(\hat{\rho}, \hat{\phi}, \hat{z})$  form a right-handed triad:

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,

$$\mathbf{\hat{z}}\times\mathbf{\hat{\rho}}=\mathbf{\hat{\phi}}$$

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,

$$\hat{\mathbf{\phi}} \times \hat{\mathbf{\phi}} = 0,$$

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• Let  $\vec{\rho} \equiv \rho \hat{\rho}$ .

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$$\mathbf{\hat{z}} \times \mathbf{\hat{z}} = 0$$

• Let  $\vec{\rho} \equiv \rho \hat{\rho}$ . Now, you can easily prove the following relations.

$$\vec{\mathbf{v}} = \vec{\boldsymbol{\omega}} \times \vec{\boldsymbol{\rho}}$$

$$\vec{a}_t = \vec{\beta} \times \vec{\rho}$$

$$\vec{\mathbf{a}}_{n} = \vec{\mathbf{v}} \times \vec{\boldsymbol{\omega}}$$

• The unit vectors ( $\hat{\rho}$ ,  $\hat{\phi}$ ,  $\hat{z}$ ) form a right-handed triad:

$$\hat{\mathbf{p}} \times \hat{\mathbf{\phi}} = \hat{\mathbf{z}}$$

$$\hat{\mathbf{\phi}} \times \hat{\mathbf{z}} = \hat{\mathbf{\rho}}$$

$$\hat{\mathbf{z}} \times \hat{\boldsymbol{\rho}} = \hat{\boldsymbol{\phi}}$$

$$\hat{\mathbf{p}} \times \hat{\mathbf{p}} = 0$$
,

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$$\vec{\omega} = \frac{\vec{\rho} \times \vec{\mathbf{v}}}{\rho^2}$$

$$\vec{\beta} = \frac{\vec{\rho} \times \vec{a}}{\rho^2}$$

**Mechanics** ▷ Dynamics of rigid bodies

#### 3. Translation and rotation

## **Mechanics** > Dynamics of rigid bodies

#### 3. Translation and rotation

$$\beta$$
 = constant,

$$\beta$$
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$$\omega =$$

$$\varphi =$$

$$\beta = \text{constant}, \qquad \omega = \omega_0 + \beta t,$$

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## **Mechanics** ▷ Dynamics of rigid bodies

#### 3. Translation and rotation

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## **Mechanics** ▷ Dynamics of rigid bodies

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,

### **Mechanics** > Dynamics of rigid bodies

#### 3. Translation and rotation

$$\beta = \text{constant},$$
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$$\overline{\omega} = \frac{\omega_0 + \omega}{2}, \quad \varphi - \varphi_0 =$$

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### **Mechanics** > Dynamics of rigid bodies

#### 3. Translation and rotation

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Rotation with constant acceleration

$$\beta = \text{constant}, \qquad \omega = \omega_0 + \beta t, \qquad \varphi = \varphi_0 + \omega_0 t + \frac{1}{2} \beta t^2$$

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- A disc of radius of 10 cm rotates from rest with a constant angular acceleration. It requires 2s for it to rotate through an angular displacement of 60°. Find:
  - (a) the angular acceleration of the disc,
  - (b) its angular velocity at 2s and at 6s,
  - (c) the linear speed at 2s of a point that is at a distance of 7 cm from the center of the disc,

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(d) the distance moved by this point in the first 2s.

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- A disc of radius of 10 cm rotates from rest with a constant angular acceleration. It requires 2s for it to rotate through an angular displacement of 60°. Find:
  - (a) the angular acceleration of the disc,  $[0.525 \text{ rad s}^{-2}]$
  - (b) its angular velocity at 2s and at 6s,
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  - (a) the angular acceleration of the disc,  $[0.525 \text{ rad s}^{-2}]$
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  - (d) the distance moved by this point in the first 2s. [0.074m]

# 1. MECHANICS

## 1.5 DYNAMICS OF RIGID BODIES

- 1 Rigid body
- 2 Center of mass
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- 5 Torque of a force
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**Mechanics** > Dynamics of rigid bodies

- 4. Rotational energy and moment of inertia
- Rotation of a system about a fixed axis (the z-axis).

**Mechanics** > Dynamics of rigid bodies

- **Rotation of a system about a fixed axis** (the z-axis).
  - Kinetic energy:

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$$E_{\mathbf{k}} = \sum_{i=1}^{n} \frac{1}{2} m_i v_i^2$$

- **A Rotation of a system about a fixed axis** (the *z*-axis).
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$$E_{k} = \sum \frac{1}{2} m_{i} v_{i}^{2} = \sum \frac{1}{2} m_{i} (\rho_{i} \omega)^{2}$$

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$$I = \sum m_i \rho_i^2$$

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Here,  $\rho_i$  is the distance of  $m_i$  from the axis.

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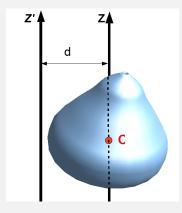
$$I = \int \rho^2 \, \mathrm{d}m$$

**Mechanics** ⊳ Dynamics of rigid bodies

- 4. Rotational energy and moment of inertia
- Parallel-axis theorem

- **Mechanics** > Dynamics of rigid bodies
- 4. Rotational energy and moment of inertia
- Parallel-axis theorem (Huygens-Steiner theorem)

Parallel-axis theorem (Huygens-Steiner theorem)

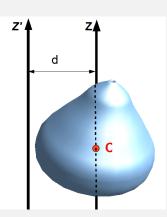


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Parallel-axis theorem (Huygens-Steiner theorem)

The moment of inertia I of a system about a given axis is calculated as

$$I = I_{\rm cm} + md^2$$

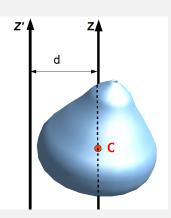


### Parallel-axis theorem (Huygens-Steiner theorem)

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where  $I_{cm}$  is the moment of inertia of the system about the axis which passes through the center of mass of the system,

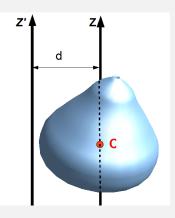


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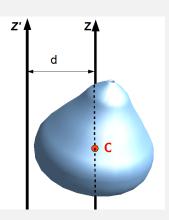
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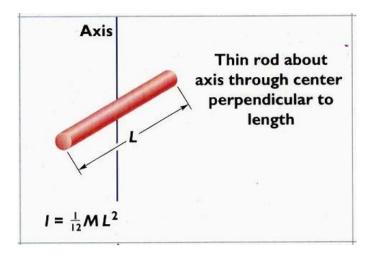
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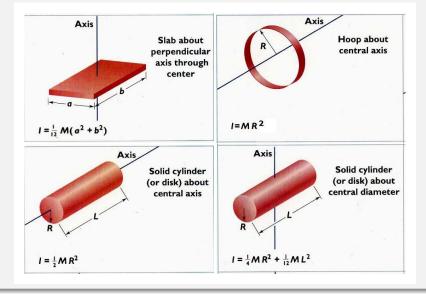
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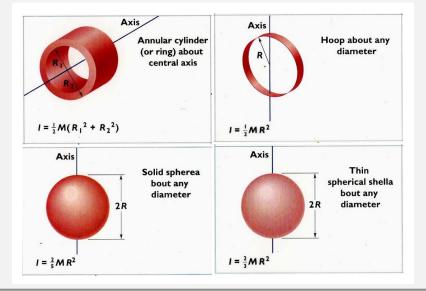


• The proof will be given as a homework.

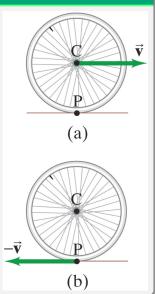


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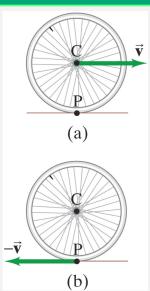




- **Mechanics** ▷ Dynamics of rigid bodies
- 4. Rotational energy and moment of inertia
- Rolling without slipping

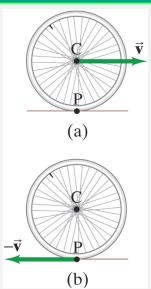


- Rolling without slipping
- (a) A wheel is rolling without slipping.

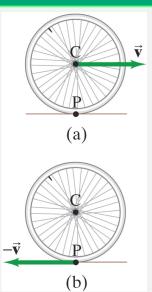


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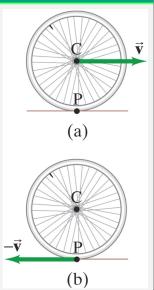
- Rolling without slipping
  - (a) A wheel is rolling without slipping. Point P is instantaneously at rest and point C moves with velocity  $\vec{\mathbf{v}}$ .



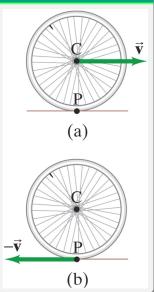
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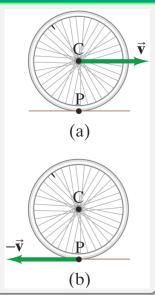
- (a) A wheel is rolling without slipping. Point P is instantaneously at rest and point C moves with velocity  $\vec{\mathbf{v}}$ .
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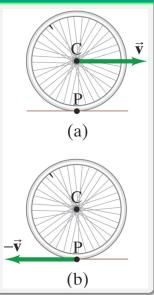
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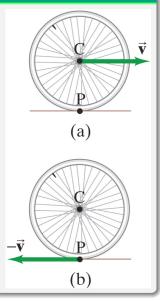


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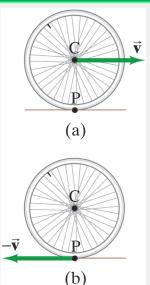
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$$E_{\mathbf{k}} =$$



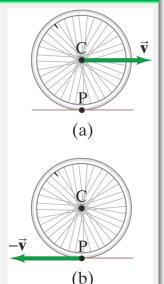
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$$E_{\bf k} = \frac{1}{2} m v^2 + \frac{1}{2} I \omega^2 = \frac{1}{2} \left( m + \frac{I}{R^2} \right) v^2$$



# 1. MECHANICS

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5. Torque of a force		

#### **Mechanics** ⊳ Dynamics of rigid bodies

## 5. Torque of a force

For a body rotating about point O,

#### **Mechanics** ▷ Dynamics of rigid bodies

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#### **Mechanics** > Dynamics of rigid bodies

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#### **Mechanics** > Dynamics of rigid bodies

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- $ightharpoonup \tau = rF\sin\alpha$ , where  $\alpha$  is the angle between  $\overrightarrow{\mathbf{r}}$  and  $\overrightarrow{\mathbf{F}}$ .

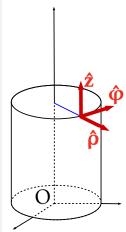
**Mechanics** ▷ Dynamics of rigid bodies

## 5. Torque of a force

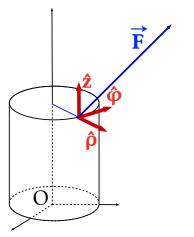
• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):

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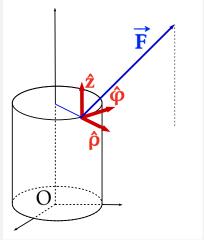
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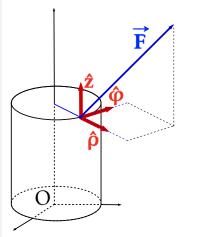
• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):



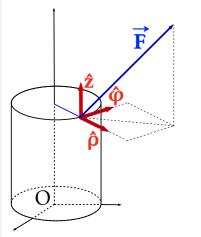
• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):



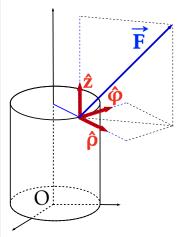
• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):



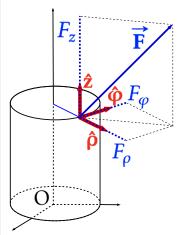
• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):



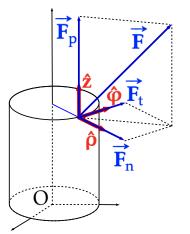
• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):



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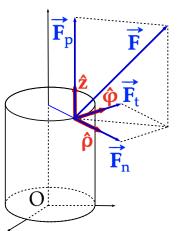


• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):



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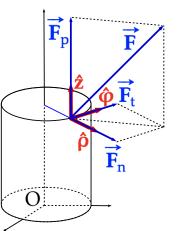
$$\vec{F} = \vec{F}_n + \vec{F}_t + \vec{F}_p$$



• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):

$$\vec{\mathbf{F}} = \vec{\mathbf{F}}_n + \vec{\mathbf{F}}_t + \vec{\mathbf{F}}_p$$

$$\equiv F_{\rho} \hat{\boldsymbol{\rho}} + F_{\varphi} \hat{\boldsymbol{\phi}} + F_z \hat{\boldsymbol{z}}$$

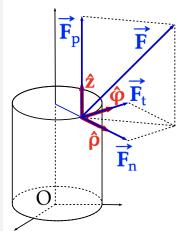


• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):

$$\vec{\mathbf{F}} = \vec{\mathbf{F}}_n + \vec{\mathbf{F}}_t + \vec{\mathbf{F}}_p$$

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•  $\vec{\mathbf{F}}_n$  is perpendicular to the axis

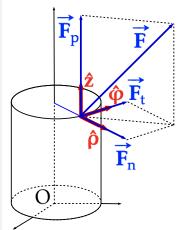


• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):

$$\vec{\mathbf{F}} = \vec{\mathbf{F}}_{n} + \vec{\mathbf{F}}_{t} + \vec{\mathbf{F}}_{p}$$

$$\equiv F_{\rho} \hat{\boldsymbol{\rho}} + F_{\varphi} \hat{\boldsymbol{\phi}} + F_{z} \hat{\boldsymbol{z}}$$

•  $\vec{F}_n$  is perpendicular to the axis and  $\vec{F}_p$  is parallel to the axis.



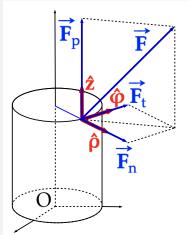
• For a body rotating about a fixed axis, chosen as the *z*-axis (in cylindrical coordinates):

$$\vec{\mathbf{F}} = \vec{\mathbf{F}}_{n} + \vec{\mathbf{F}}_{t} + \vec{\mathbf{F}}_{p}$$

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•  $\vec{F}_n$  is perpendicular to the axis and  $\vec{F}_p$  is parallel to the axis.

These components give no effect to the rotational motion.



• For a body rotating about a fixed axis, chosen as the z-axis (in cylindrical coordinates):

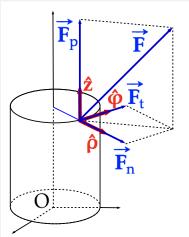
$$\vec{\mathbf{F}} = \vec{\mathbf{F}}_{n} + \vec{\mathbf{F}}_{t} + \vec{\mathbf{F}}_{p}$$

$$\equiv F_{\rho} \hat{\boldsymbol{\rho}} + F_{\varphi} \hat{\boldsymbol{\phi}} + F_{z} \hat{\boldsymbol{z}}$$

•  $\vec{F}_n$  is perpendicular to the axis and  $\vec{F}_p$  is parallel to the axis.

These components give no effect to the rotational motion.

They are cancelled out by the reaction forces from the axis.



**Mechanics** > Dynamics of rigid bodies

## 5. Torque of a force

 $\bullet$  The tangential component  $\overrightarrow{F}_t$ 

### **Mechanics** > Dynamics of rigid bodies

# 5. Torque of a force

• The tangential component  $\vec{F}_t$  is the only one that gives effect to the rotational motion about the *z*-axis.

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$$\vec{\mathbf{F}}_{t} \equiv F_{\varphi} \,\hat{\mathbf{\phi}}, \qquad F_{\varphi} = -F_{x} \sin \varphi + F_{y} \cos \varphi$$

• Therefore, in rotational motion about a fixed axis,

• The tangential component  $\vec{F}_t$  is the only one that gives effect to the rotational motion about the *z*-axis.

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$$\vec{\mathbf{F}}_{t} \equiv F_{\varphi} \,\hat{\mathbf{\phi}}, \qquad F_{\varphi} = -F_{x} \sin \varphi + F_{y} \cos \varphi$$

$$\overrightarrow{\tau} = \overrightarrow{\rho} \times \overrightarrow{F}_t$$

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$$\vec{\boldsymbol{\tau}} = \vec{\boldsymbol{\rho}} \times \vec{\mathbf{F}}_{t} = (\rho \,\hat{\boldsymbol{\rho}}) \times (F_{\varphi} \,\hat{\boldsymbol{\phi}})$$

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$$\vec{\mathbf{F}}_{t} \equiv F_{\varphi} \,\hat{\mathbf{\phi}}, \qquad F_{\varphi} = -F_{x} \sin \varphi + F_{y} \cos \varphi$$

$$\vec{\boldsymbol{\tau}} = \vec{\boldsymbol{\rho}} \times \vec{\mathbf{F}}_{\mathsf{t}} = (\rho \,\hat{\boldsymbol{\rho}}) \times (F_{\varphi} \,\hat{\boldsymbol{\phi}}) \quad \rightarrow \quad \vec{\boldsymbol{\tau}} = \rho F_{\varphi} \,\hat{\mathbf{z}}$$

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$$\vec{\mathbf{F}}_{t} \equiv F_{\varphi} \,\hat{\mathbf{\phi}}, \qquad F_{\varphi} = -F_{x} \sin \varphi + F_{y} \cos \varphi$$

 Therefore, in rotational motion about a fixed axis, the torque of a force is defined as the torque of its tangential component.

$$\vec{\boldsymbol{\tau}} = \vec{\boldsymbol{\rho}} \times \vec{\mathbf{F}}_{\mathsf{t}} = (\rho \,\hat{\boldsymbol{\rho}}) \times (F_{\varphi} \,\hat{\boldsymbol{\phi}}) \quad \rightarrow \quad \vec{\boldsymbol{\tau}} = \rho F_{\varphi} \,\hat{\mathbf{z}}$$

• The torque is along the axis of rotation.

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$$\vec{\mathbf{F}}_{t} \equiv F_{\varphi} \,\hat{\mathbf{\phi}}, \qquad F_{\varphi} = -F_{x} \sin \varphi + F_{y} \cos \varphi$$

$$\vec{\boldsymbol{\tau}} = \vec{\boldsymbol{\rho}} \times \vec{\mathbf{F}}_{\mathsf{t}} = (\rho \,\hat{\boldsymbol{\rho}}) \times (F_{\varphi} \,\hat{\boldsymbol{\phi}}) \quad \rightarrow \quad \vec{\boldsymbol{\tau}} = \rho F_{\varphi} \,\hat{\mathbf{z}}$$

- The torque is along the axis of rotation.
- The torque is equal to the product of the perpendicular distance and the tangential component of the force.

**Mechanics** ▷ Dynamics of rigid bodies

## 5. Torque of a force

#### **Mechanics** ▷ Dynamics of rigid bodies

## 5. Torque of a force

$$\vec{\tau}_{O} = \vec{r} \times \vec{F}$$

#### **Mechanics** ▷ Dynamics of rigid bodies

## 5. Torque of a force

$$\vec{\boldsymbol{\tau}}_{\mathrm{O}} = \vec{\mathbf{r}} \times \vec{\mathbf{F}} = (x\,\hat{\mathbf{x}} + y\,\hat{\mathbf{y}} + z\,\hat{\mathbf{z}}) \times (F_x\,\hat{\mathbf{x}} + F_y\,\hat{\mathbf{y}} + F_z\,\hat{\mathbf{z}})$$

$$\vec{\boldsymbol{\tau}}_{\mathrm{O}} = \vec{\mathbf{r}} \times \vec{\mathbf{F}} = (x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}) \times (F_x\hat{\mathbf{x}} + F_y\hat{\mathbf{y}} + F_z\hat{\mathbf{z}})$$

$$\overrightarrow{\boldsymbol{\tau}}_{\mathrm{O}} = (yF_z - zF_y)\hat{\boldsymbol{x}} + (zF_x - xF_z)\hat{\boldsymbol{y}} + (xF_y - yF_x)\hat{\boldsymbol{z}}$$

Generally, for rotational motion about the origin,

$$\overrightarrow{\boldsymbol{\tau}}_{\mathrm{O}} = \overrightarrow{\mathbf{r}} \times \overrightarrow{\mathbf{F}} = (x\,\hat{\mathbf{x}} + y\,\hat{\mathbf{y}} + z\,\hat{\mathbf{z}}) \times (F_x\,\hat{\mathbf{x}} + F_y\,\hat{\mathbf{y}} + F_z\,\hat{\mathbf{z}})$$

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$$\overrightarrow{\boldsymbol{\tau}}_{\mathrm{O}} = (yF_z - zF_y)\,\hat{\mathbf{x}} + (zF_x - xF_z)\,\hat{\mathbf{y}} + (xF_y - yF_x)\,\hat{\mathbf{z}}$$

$$\tau_z = xF_y - yF_x = \rho\cos\varphi F_y - \rho\sin\varphi F_x = \rho(-\sin\varphi F_x + \cos\varphi F_y)$$

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$$\tau_z = \rho F_{\varphi}$$

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$$\overrightarrow{\boldsymbol{\tau}}_{\mathrm{O}} = (yF_z - zF_y)\,\hat{\boldsymbol{x}} + (zF_x - xF_z)\,\hat{\boldsymbol{y}} + (xF_y - yF_x)\,\hat{\boldsymbol{z}}$$

▶ The *z*-component:

$$\tau_z = xF_y - yF_x = \rho\cos\varphi F_y - \rho\sin\varphi F_x = \rho(-\sin\varphi F_x + \cos\varphi F_y)$$

$$\tau_z = \rho F_{\varphi}$$
 (same as the torque of  $\vec{\mathbf{F}}$  about z-axis)

The torque of a force about an axis is equal to the component on the axis of the torque of that force about a point.

## 1. MECHANICS

### 1.5 DYNAMICS OF RIGID BODIES

- 1 Rigid body
- 2 Center of mass
- 3 Translation and rotation
- 4 ROTATIONAL ENERGY AND MOMENT OF INERTIA
- 5 Torque of a force
- 6 Angular momentum
- 7 Torque and angular acceleration
- 8 Torque and angular momentum

## 6. Angular momentum

$$\vec{\mathbf{L}} = \sum \vec{\mathbf{L}}_i$$

$$\overrightarrow{\mathbf{L}} = \sum \overrightarrow{\mathbf{L}}_i = \sum \overrightarrow{\mathbf{r}}_i \times \overrightarrow{\mathbf{p}}_i$$

$$\overrightarrow{\mathbf{L}} = \sum \overrightarrow{\mathbf{r}}_i = \sum \overrightarrow{\mathbf{r}}_i \times \overrightarrow{\mathbf{p}}_i = \sum m_i (\overrightarrow{\mathbf{r}}_i \times \overrightarrow{\mathbf{v}}_i)$$

• For a system rotating about the origin:

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▶ Cylindrical coordinates (unit vectors):  $\vec{\mathbf{r}}_i =$ 

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- ▶ Cylindrical coordinates (unit vectors):  $\vec{\mathbf{r}}_i = \rho_i \hat{\boldsymbol{\rho}} + z_i \hat{\mathbf{z}}$ .
- For a system rotating about the *z*-axis:

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$$\vec{\mathbf{v}}_i = v_i \,\hat{\boldsymbol{\varphi}}$$

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- For a system rotating about the *z*-axis:

$$\vec{\mathbf{v}}_i = v_i \,\hat{\boldsymbol{\varphi}} = \rho_i \omega \,\hat{\boldsymbol{\varphi}}$$

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$$\vec{\mathbf{r}}_{i} \times \vec{\mathbf{v}}_{i} = (\rho_{i} \,\hat{\boldsymbol{\rho}} + z_{i} \,\hat{\mathbf{z}}) \times (\rho_{i} \omega \,\hat{\boldsymbol{\phi}})$$

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$$\vec{\mathbf{L}} = \left(\sum m_{i} \rho_{i}^{2}\right) \omega \,\hat{\mathbf{z}} - \left(\sum m_{i} z_{i} \rho_{i} \,\hat{\boldsymbol{\rho}}\right) \omega$$

$$\overrightarrow{\mathbf{L}} = I\overrightarrow{\boldsymbol{\omega}} - \left(\sum m_i z_i \rho_i \, \hat{\boldsymbol{\rho}}\right) \boldsymbol{\omega}$$

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• If the mass of the system is distributed symmetrically about the *z*-axis

$$\overrightarrow{\mathbf{L}} = I\overrightarrow{\mathbf{w}} - \left(\sum m_i z_i \rho_i \, \hat{\mathbf{\rho}}\right) \omega$$

• If the mass of the system is distributed symmetrically about the *z*-axis then the second term is zero.

$$\overrightarrow{\mathbf{L}} = I\overrightarrow{\boldsymbol{\omega}} - \left(\sum m_i z_i \rho_i \, \hat{\boldsymbol{\rho}}\right) \boldsymbol{\omega}$$

• If the mass of the system is distributed symmetrically about the *z*-axis then the second term is zero.

For a symmetric homogeneous body in pure rotation about its symmetric axis,

$$\overrightarrow{\mathbf{L}} = I\overrightarrow{\boldsymbol{\omega}} - \left(\sum m_i z_i \rho_i \, \hat{\boldsymbol{\rho}}\right) \boldsymbol{\omega}$$

• If the mass of the system is distributed symmetrically about the *z*-axis then the second term is zero.

For a symmetric homogeneous body in pure rotation about its symmetric axis, the <u>total angular momentum</u> is

$$\vec{\mathbf{L}} = I\vec{\boldsymbol{\omega}}$$
.

$$\overrightarrow{\mathbf{L}} = I\overrightarrow{\boldsymbol{\omega}} - \left(\sum m_i z_i \rho_i \, \hat{\boldsymbol{\rho}}\right) \boldsymbol{\omega}$$

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▶ Recall that  $\vec{\mathbf{w}} = \omega \hat{\mathbf{z}}$ 

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▶ Recall that  $\overrightarrow{\mathbf{w}} = \omega \hat{\mathbf{z}} \rightarrow \overrightarrow{\mathbf{L}}$  lies on the axis of rotation.

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- ▶ Recall that  $\overrightarrow{\mathbf{w}} = \omega \hat{\mathbf{z}} \rightarrow \overrightarrow{\mathbf{L}}$  lies on the axis of rotation.
- ▶ In the next parts, we will relate  $\vec{\tau}$  with  $\vec{\beta}$  and  $\vec{L}$ .

# 1. MECHANICS

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6. Torque and angular acceleration

$$\overrightarrow{\tau} =$$

$$\vec{\boldsymbol{\tau}} = \sum \rho_i F_{\varphi i} \, \hat{\mathbf{z}}$$

$$\vec{\boldsymbol{\tau}} = \sum \rho_i F_{\varphi i} \, \hat{\boldsymbol{z}} = \sum \rho_i m_i a_i \, \hat{\boldsymbol{z}}$$

$$\vec{\boldsymbol{\tau}} = \sum \rho_i F_{\varphi i} \, \hat{\boldsymbol{z}} = \sum \rho_i m_i a_i \, \hat{\boldsymbol{z}} = \sum \rho_i^2 m_i \beta \, \hat{\boldsymbol{z}}$$

$$\vec{\boldsymbol{\tau}} = \sum \rho_i F_{\varphi i} \, \hat{\boldsymbol{z}} = \sum \rho_i m_i a_i \, \hat{\boldsymbol{z}} = \sum \rho_i^2 m_i \beta \, \hat{\boldsymbol{z}} = \left(\sum \rho_i^2 m_i\right) \beta \, \hat{\boldsymbol{z}}$$

$$\vec{\boldsymbol{\tau}} = \sum \rho_i F_{\varphi i} \, \hat{\boldsymbol{z}} = \sum \rho_i m_i a_i \, \hat{\boldsymbol{z}} = \sum \rho_i^2 m_i \beta \, \hat{\boldsymbol{z}} = \left(\sum \rho_i^2 m_i\right) \beta \, \hat{\boldsymbol{z}}$$

$$\vec{\tau} = I \vec{\beta}$$

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$$\vec{\tau} = I \vec{\beta}$$
 or  $\vec{\beta} = \frac{\vec{\tau}}{I}$ 

• Net external torque acting on a body rotating about the *z*-axis:

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▶ This is applied to any rigid object, symmetric or not.

# 1. MECHANICS

## 1.5 DYNAMICS OF RIGID BODIES

- 1 Rigid body
- 2 Center of mass
- 3 Translation and rotation
- 4 ROTATIONAL ENERGY AND MOMENT OF INERTIA
- 5 Torque of a force
- 6 Angular momentum
- 7 Torque and angular acceleration
- 8 Torque and angular momentum

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