

Probabilistic Runge-Kutta methods for uncertainty quantification of numerical errors in geometric integration

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Outline

- 1 Motivation
- 2 Geometric numerical integration
- 3 Probabilistic methods for ODEs
- 4 Bayesian inverse problems

Motivation – Chaotic equations

Consider Lorenz equation (atmospheric convection)

$$\begin{aligned}x' &= \sigma(y - x), & x(0) &= -10, \\y' &= x(\rho - z) - y, & y(0) &= -1, \\z' &= xy - \beta z, & z(0) &= 40.\end{aligned}$$

For $\rho = 28$, $\sigma = 10$, $\beta = 8/3$ **chaotic behaviour**.

\implies Numerical integration gives **unreliable solutions**.

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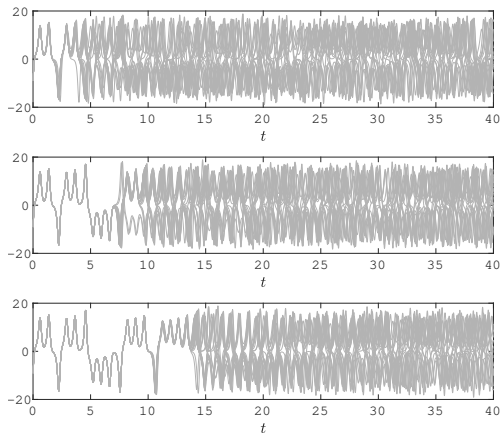
For $\rho = 28$, $\sigma = 10$, $\beta = 8/3$ **chaotic behaviour**.

\implies Numerical integration gives **unreliable solutions**.

Goal: Understand reliability of numerical solutions.

Idea: Perturb the initial data e.g. with Gaussian noise on $x(0)$

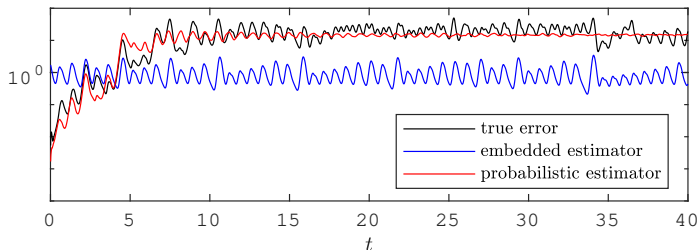
Motivation – Chaotic equations



Solutions of the Lorenz system (x component) – different perturbations

Which one has the correct magnitude wrt numerical error?

Motivation – Error estimators



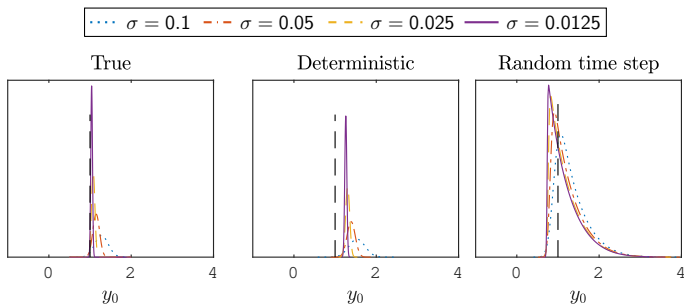
Error estimators for Lorenz given by

Probabilistic solution: use $\|\text{Var } Y_n\|^{1/2}$ where Y_n is a probabilistic family

Classical embedded couple: Local errors don't show the true behaviour!

Goal: A posteriori error estimator $\text{err}_n \approx \|\text{Var } Y_n\|^{1/2} \rightsquigarrow$ Work in progress!

Motivation – Bayesian inverse problems



Posterior distributions (analytic) on y_0 for $y' = -y$, $y(0) = y_0$. One observation corrupted by noise $\mathcal{N}(0, \sigma^2)$ with truth y_0^* .

For $\sigma \rightarrow 0$:

- **True solution:** Posterior converging to $\delta_{y_0^*} \rightarrow$ good
- **Runge-Kutta:** Posterior converging to Dirac delta on wrong value \rightarrow bad
- **Probabilistic method:** Posterior variance \approx numerical error

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Notation

Autonomous dynamical system, function $f: \mathbb{R}^d \rightarrow \mathbb{R}^d$ and the ODE

$$y' = f(y), \quad y(0) = y_0.$$

Flow of the equation φ_t

$$y(t) = \varphi_t(y_0).$$

One-step method (e.g. Runge Kutta): numerical flow Ψ_h

$$y_{n+1} = \Psi_h(y_n).$$

Geometric numerical integration

Goal

Develop numerical methods that preserve geometric properties of certain dynamical systems.

First integral of motion $I: \mathbb{R}^d \rightarrow \mathbb{R}$

$$I(\varphi_t(y_0)) = I(y_0), \quad \forall t > 0.$$

Example: quadratic first integral, given $S \in \mathbb{R}^{d \times d}$, $v \in \mathbb{R}^d$

$$I(y) = y^\top S y + v^\top y,$$

conserved by all Gauss collocation methods (e.g., **implicit midpoint**, ...).

Theorem (Polynomial first integrals)

No Runge-Kutta method can conserve all polynomial first integrals of degree $\text{Deg}(I) \geq 3$.

Geometric numerical integration

Goal

Develop numerical methods that preserve geometric properties of certain dynamical systems.

Hamiltonian systems: Given $Q: \mathbb{R}^{2d} \rightarrow \mathbb{R}$, define

$$y'(t) = J^{-1} \nabla Q(y), \quad y(0) = y_0$$
$$J = \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix}, \quad I \text{ identity in } \mathbb{R}^{d \times d}$$

The flow φ_t is **symplectic**

$$\varphi'_t(y)^\top J \varphi'_t(y) = J \implies \text{Conservation of volumes}$$

Symplectic numerical methods

$$\Psi'_h(y)^\top J \Psi'_h(y) = J$$

Geometric numerical integration

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Theorem

For a symplectic integrator of order q , there exist $C, \kappa > 0$, independent of h such that

$$\mathbb{E}|Q(y_n) - Q(y_0)| \leq C_1 h^q,$$

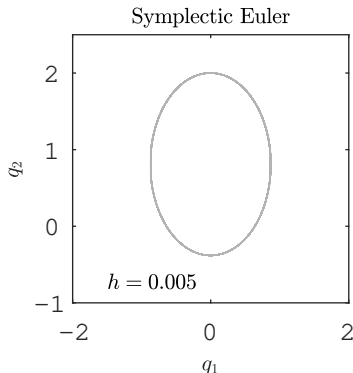
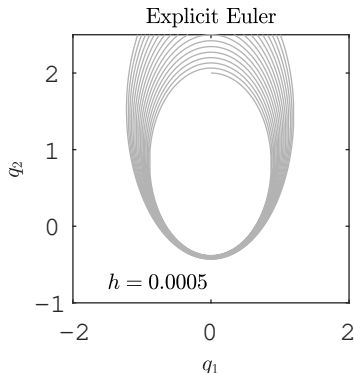
for time intervals of length $nh = \mathcal{O}(e^{\kappa/h})$.

Geometric numerical integration

Example

Two-body problem (planetary orbits), $y = (p, q)^\top \in \mathbb{R}^4$

$$Q(p, q) = \frac{1}{2}(p_1^2 + p_2^2) - \frac{1}{\sqrt{q_1^2 + q_2^2}}$$

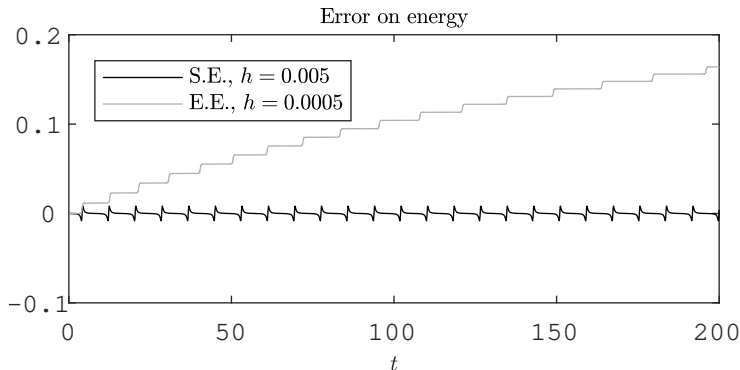


Geometric numerical integration

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Geometric numerical integration

Given Hamiltonian $Q(p, q)$.

Symplectic Euler method – order 1

$$\begin{aligned}p_{n+1} &= p_n - hQ_q(p_{n+1}, q_n), \\q_{n+1} &= q_n + hQ_p(p_{n+1}, q_n).\end{aligned}$$

Störmer-Verlet scheme – order 2

$$\begin{aligned}p_{n+1/2} &= p_n - \frac{h}{2}Q_q(p_{n+1/2}, q_n), \\q_{n+1} &= q_n + \frac{h}{2}(Q_p(p_{n+1/2}, q_n) + Q_p(p_{n+1/2}, q_{n+1})), \\p_{n+1} &= p_n - \frac{h}{2}Q_q(p_{n+1/2}, q_{n+1}).\end{aligned}$$

Geometric numerical integration

Given **separable** Hamiltonian $Q(p, q) = Q_1(p) + Q_2(q)$.

Symplectic Euler method – order 1, explicit

$$\begin{aligned}p_{n+1} &= p_n - hQ'_2(q_n), \\ q_{n+1} &= q_n + hQ'_1(p_{n+1}).\end{aligned}$$

Störmer-Verlet scheme – order 2, explicit

$$\begin{aligned}p_{n+1/2} &= p_n - \frac{h}{2}Q'_2(q_n), \\ q_{n+1} &= q_n + hQ'_1(p_{n+1/2}), \\ p_{n+1} &= p_n - \frac{h}{2}Q'_2(q_{n+1}).\end{aligned}$$

Several examples of separable Hamiltonians (**Two-body problem**, ...)

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Probabilistic methods for ODEs

Filtering methods for ODEs: fix a prior on $y(t)$ (Gaussian process), update with evaluations of $f(y)$.

- Kersting and Hennig (2016); Kersting et al. (2018)
- Chkrebtii et al. (2016)
- Schober et al. (2014, 2018)
- ...

Randomised methods for ODEs: random perturbation of deterministic numerical solutions \rightarrow sampling

- Conrad et al. (2016)
- Lie et al. (2017)
- Abdulle and Garegnani (2018)
- ...

Additive noise method

Stochastic process $\{Y_n\}_{n=1,2,\dots}$ with recurrence

$$Y_{n+1} = \underbrace{\Psi_h(Y_n)}_{\text{deterministic}} + \underbrace{\xi_n(h)}_{\text{random}}.$$

Main assumption: For $p > 1$ and $Q \in \mathbb{R}^{d \times d}$

$$\xi_n(h) \stackrel{\text{iid}}{\sim} \mathcal{N}(0, Qh^{2p+1}).$$

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Properties

If Ψ_h is of order q and for $\Phi: \mathbb{R}^d \rightarrow \mathbb{R}$ smooth

- Strong convergence: $\mathbb{E}\|y(hn) - Y_n\| \leq Ch^{\min\{p,q\}},$
- Weak convergence: $|\Phi(y(hn)) - \mathbb{E}\Phi(Y_n)| \leq Ch^{\min\{2p,q\}},$
- Good qualitative behavior in Bayesian inverse problems.

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Issues

- Robustness: $\Psi_h(Y_{n-1}) > 0 \not\Rightarrow \mathbb{P}(Y_n < 0) = 0$,
- Geometric properties are not conserved from Ψ_h .

Random time steps

Intrinsic noise: Random time-stepping Runge-Kutta (RTS-RK)

$$Y_{n+1} = \Psi_{H_n}(Y_n),$$

Main assumption: $\{H_n\}_{n=0,1,\dots}$ iid such that for $h, C > 0$ and $p > 1$

$$H_n > 0 \text{ a.s.}, \quad \mathbb{E} H_n = h, \quad \text{Var } H_n = Ch^{2p+1}.$$

Example: $H_n \stackrel{\text{iid}}{\sim} \mathcal{U}(h - h^{p+1/2}, h + h^{p+1/2})$.

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Properties – Geometric

- Conservation of (polynomial) first integrals is inherited by Ψ_h ,
- Flow map is symplectic if Ψ_h is symplectic,
- Long-time conservation of energy in Hamiltonian systems.

Conservation of first integrals – Additive noise

Recall: $Y_{n+1} = \Psi_h(Y_n) + \xi_n(h)$, with $\mathbb{E} \xi_n(h) \xi_n(h)^\top = h^{2p+1} Q$

Linear first integrals: $I(y) = v^\top y$ such that $I(\Psi_h(Y_1)) = I(y_0)$. Then

$$I(Y_1) = v^\top (y_0 + \xi_0(h)) \implies \mathbb{E} I(Y_1) = I(y_0) \text{ iff } \mathbb{E} \xi_0(h) = 0.$$

Quadratic first integrals: $I(y) = y^\top S y$ such that $I(\Psi_h(Y_1)) = I(y_0)$. Then

$$\begin{aligned} I(Y_1) &= I(y_0) + 2\xi_0(h)^\top S \Psi_h(y_0) + \xi_0(h)^\top S \xi_0(h), \\ \implies \mathbb{E} I(Y_1) &= I(y_0) + Q : S h^{2p+1}, \quad (\text{with } \mathbb{E} \xi_0(h) = 0) \end{aligned}$$

Quadratic first integrals are not conserved on average!

Conservation of first integrals

Theorem (Conservation of invariants)

If the Runge-Kutta scheme defined by Ψ_h conserves an invariant $I(y)$ for an ODE, then the RTS-RK method conserves $I(y)$ for the same ODE.

Proof

If $I(\Psi_h(y)) = I(y)$ for any h , then $I(\Psi_{H_0}(y)) = I(y)$ for any value that H_0 can assume.

Symplecticity

Energy $Q: \mathbb{R}^{2d} \rightarrow \mathbb{R}$ and Hamiltonian system

$$y' = J^{-1} \nabla Q(y), \quad J = \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix}.$$

Symplectic integrator Ψ_h of order q .

Theorem (Strong approximation of the Hamiltonian)

There exist positive constants κ, C_1, C_2, C_3 , independent of h such that

$$\mathbb{E}|Q(Y_n) - Q(y_0)| \leq C_1 h^q,$$

for time intervals of length $\mathcal{O}(h^{1-2p})$.

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Idea of the proof

Use classical backward error analysis. Considering that $\mathbb{E} H_k = h$, we have at “first order” the same conservation as deterministic methods. The time interval reduction comes from the time steps’ variance.

Numerical experiments – Geometric properties

Consider the perturbed Kepler equation (model for two-body problem)

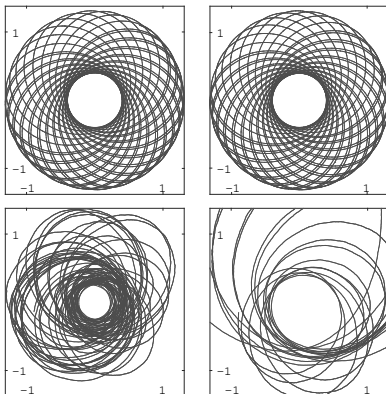
$$\begin{aligned}w_1' &= v_1, & v_1' &= -\frac{w_1}{\|w\|^3} - \frac{\delta w_1}{\|w\|^5}, \\w_2' &= v_2, & v_2' &= -\frac{w_2}{\|w\|^3} - \frac{\delta w_2}{\|w\|^5}.\end{aligned}$$

The **angular momentum** is conserved (quadratic first integral)

$$I(v, w) = w_1 v_2 - w_2 v_1$$

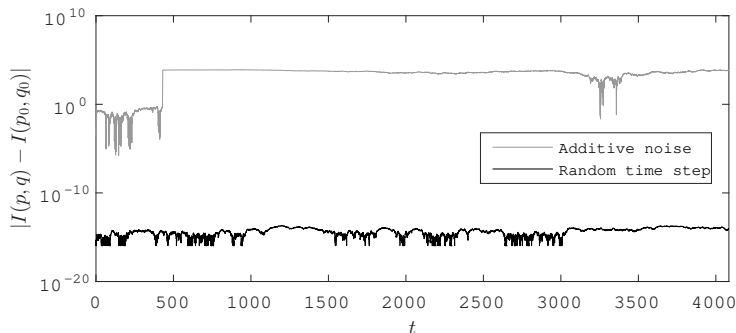
→ employ a Gauss method (implicit midpoint rule).

Numerical experiments – Geometric properties



RTS-RK (first row), Additive noise (second row). Time $0 \leq t \leq 200$ and $200 \leq t \leq 400$ (left and right)

Numerical experiments – Geometric properties



Conservation of the **angular momentum** (quadratic first integral)

Numerical experiments – Geometric properties

Consider the pendulum system, Hamiltonian with energy

$$Q(v, w) = \frac{1}{2}v^2 - \cos(w).$$

Energy is separable: $Q(v, w) = Q_1(v) + Q_2(w)$.

Störmer-Verlet scheme – order 2

$$v_{n+1/2} = v_n - \frac{h}{2}Q_w(v_{n+1/2}, w_n),$$

$$w_{n+1} = w_n + \frac{h}{2}(Q_v(v_{n+1/2}, w_n) + Q_v(v_{n+1/2}, w_{n+1})),$$

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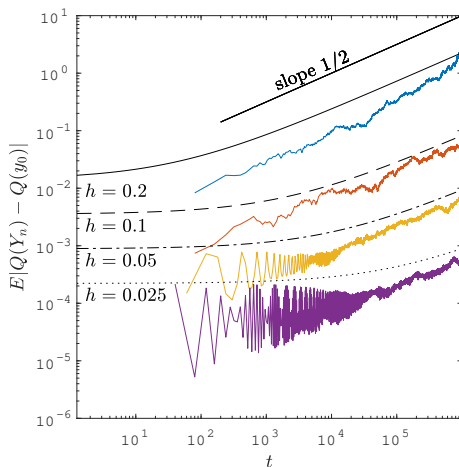
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Störmer-Verlet scheme – order 2, explicit

$$\begin{aligned}v_{n+1/2} &= v_n - \frac{h}{2}Q'_2(w_n), \\w_{n+1} &= w_n + hQ'_1(v_{n+1/2}), \\v_{n+1} &= v_n - \frac{h}{2}Q'_2(w_{n+1}).\end{aligned}$$

Numerical experiments – Geometric properties



Mean error on the Hamiltonian for different values of the time step h .

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Bayesian inverse problems

Goal

Given $\vartheta \in \mathbb{R}^n$, $f_\vartheta: \mathbb{R}^d \rightarrow \mathbb{R}^d$ and the ODE

$$y' = f_\vartheta(y), \quad y(0) = y_{0,\vartheta} \in \mathbb{R}^d,$$

retrieve the true value ϑ^* from observations of $y(t)$, $t > 0$.

Bayesian inverse problems

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retrieve the true value ϑ^* from observations of $y(t)$, $t > 0$.

Bayesian setting: fix prior $\pi_{\text{prior}}(\vartheta)$, consider the forward operator \mathcal{G} and model observations as

$$\underbrace{\mathcal{Y}}_{\text{observations}} = \underbrace{\mathcal{G}(\vartheta^*)}_{\text{forward}} + \underbrace{\varepsilon}_{\text{noise}}, \quad \varepsilon \sim \pi_{\text{noise}},$$

then the **posterior distribution (density)** is

$$\pi(\vartheta \mid \mathcal{Y}) \propto \pi_{\text{prior}}(\vartheta) \pi_{\text{noise}}(\mathcal{Y} - \mathcal{G}(\vartheta)).$$

Bayesian inverse problems

The posterior $\pi(\vartheta \mid \mathcal{Y})$ is not computable, approximate with

$$\pi^h(\vartheta \mid \mathcal{Y}) \propto \pi_{\text{prior}}(\vartheta) \pi_{\text{noise}}(\mathcal{Y} - \mathcal{G}^h(\vartheta)).$$

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Properties

$\pi^h \rightarrow \pi$ for $h \rightarrow 0$ (in the Hellinger distance).

Issue

- π^h concentrated around values “far” from $\vartheta^* \rightarrow$ non-predictive posterior

Bayesian inverse problems

The posterior $\pi(\vartheta \mid \mathcal{Y})$ is not computable, approximate with

$$\pi^{h,\text{RTS}}(\vartheta \mid \mathcal{Y}) \propto \pi_{\text{prior}}(\vartheta) \mathbb{E}^{\mathbf{H}} \pi_{\text{noise}}(\mathcal{Y} - \mathcal{G}^{\mathbf{H}}(\vartheta)),$$

where $\mathbf{H} = (H_0, H_1, \dots)$ is the vector of all time steps chosen in one run.

Bayesian inverse problems

The posterior $\pi(\vartheta \mid \mathcal{Y})$ is not computable, approximate with

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Properties

- $\pi^{h,\text{RTS}} \rightarrow \pi$ for $h \rightarrow 0$ (in the Hellinger distance). Lie et al. (2017)
- “correct” the non-predictive behaviour of deterministic approximations

Warning

- Approximation of $\mathbb{E}^{\mathbf{H}} \pi_{\text{noise}}(\mathcal{Y} - \mathcal{G}^{\mathbf{H}}(\vartheta))$ is required

Numerical experiment – Bayesian inverse problems

Consider the Hénon-Heiles system (motion of a star around a galactic center), Hamiltonian with **energy**

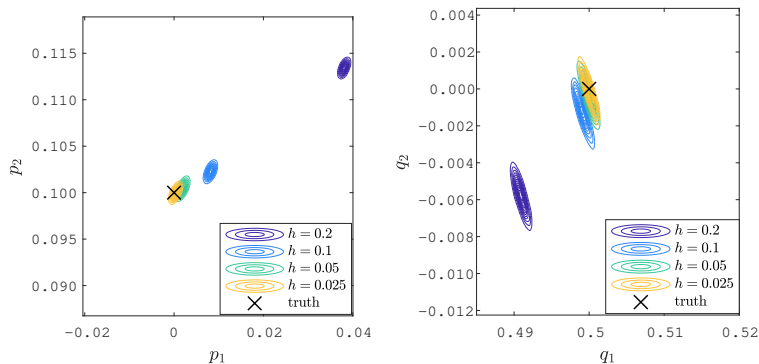
$$E(v, w) = \frac{1}{2} \|v\|^2 + \frac{1}{2} \|w\|^2 + w_1^2 w_2 - \frac{1}{3} w_2^3.$$

Chaotic problem for certain levels of energy.

Goal

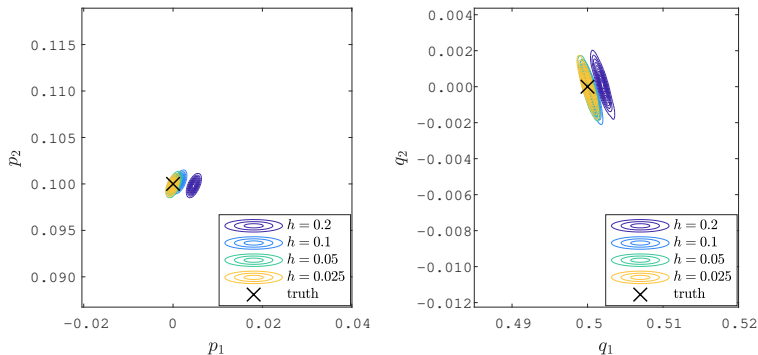
Find posterior $\pi((v_0, w_0) \mid \mathcal{Y})$ over the initial condition from a single observation of $(v(10), w(10))$

Numerical experiment – Bayesian inverse problems



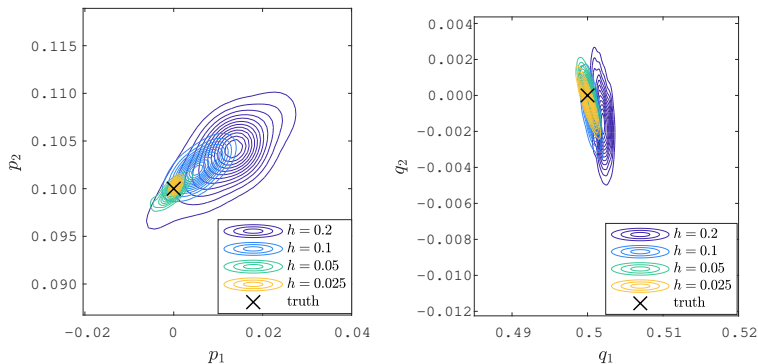
Posterior distributions given by **deterministic Heun method**.

Numerical experiment – Bayesian inverse problems



Posterior distributions given by [deterministic Störmer-Verlet method](#).

Numerical experiment – Bayesian inverse problems



Posterior distributions given by **RTS-RK Störmer-Verlet** method.

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