

\$ Supplement Note 10: Fourier Series Expansion

As mentioned in **Example 5.3** of **Lecture Notes**, the set of functions

$$B \equiv \{1, \cos x, \cos 2x, \cos 3x, \dots, \sin x, \sin 2x, \sin 3x, \dots\}, \quad (1)$$

may be considered as *vectors* in the (infinitely dimensional) vector space $L^2[-\pi, \pi]$ consisting of *periodic sectionally smooth functions* with a period of 2π .

$f \in L^2[-\pi, \pi]$ implies that

1. $f(x)$ is *continuous and differentiable* everywhere in the interval $[-\pi, \pi]$ except at a finite number of discontinuities,
2. for any integer n , the following periodic relation is satisfied:

$$f(x + 2n\pi) = f(x). \quad (2)$$

The ordinary summation of functions and multiplication by scalars on the functions are taken as the *addition* and *scalar multiplication* operators for vectors in this space. In particular, any linear superposition of vectors in B also belongs to the vector space $L^2[-\pi, \pi]$. (*This is the basic property of linear vector space!*).

$$f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx) \Rightarrow f \in L^2[-\pi, \pi], \quad (3)$$

where a_n and b_n are arbitrary constants.

Further, (**Example 5.3** of **Lecture Notes**), a *scalar product* for two functions f, g in the vector space $L^2[-\pi, \pi]$ can be defined by

$$\langle f, g \rangle \equiv \int_{-\pi}^{\pi} f(x)g(x)dx. \quad (4)$$

With this scalar product, the vectors in B are orthogonal with each other:

$$\langle \cos nx, 1 \rangle \equiv \int_{-\pi}^{\pi} \cos nx \, dx = 0, \quad n = 1, 2, \dots \quad (5a)$$

$$\langle \sin nx, 1 \rangle \equiv \int_{-\pi}^{\pi} \sin nx \, dx = 0, \quad n = 1, 2, \dots \quad (5b)$$

$$\begin{aligned} \langle \cos nx, \sin mx \rangle &\equiv \int_{-\pi}^{\pi} \cos nx \sin mx \, dx = \int_{-\pi}^{\pi} \frac{\sin(m+n)x + \sin(m-n)x}{2} dx \\ &= 0, \quad n, m = 1, 2, \dots \end{aligned} \quad (5c)$$

$$\begin{aligned} \langle \sin nx, \sin mx \rangle &\equiv \int_{-\pi}^{\pi} \sin nx \sin mx \, dx = \int_{-\pi}^{\pi} \frac{\cos(m-n)x - \cos(m+n)x}{2} dx \\ &= 0, \quad n \neq m \text{ are integers} \end{aligned} \quad (5d)$$

$$\begin{aligned} \langle \cos nx, \cos mx \rangle &\equiv \int_{-\pi}^{\pi} \cos nx \cos mx \, dx = \int_{-\pi}^{\pi} \frac{\cos(m-n)x + \cos(m+n)x}{2} dx \\ &= 0, \quad n \neq m \text{ are integers} \end{aligned} \quad (5e)$$

The normalization of these vectors are given by:

$$\langle \sin nx, \sin nx \rangle \equiv \int_{-\pi}^{\pi} \sin^2 nx \, dx = \int_{-\pi}^{\pi} \frac{1 - \cos 2x}{2} dx = \pi \quad (6a)$$

$$\langle \cos nx, \cos nx \rangle \equiv \int_{-\pi}^{\pi} \cos^2 nx \, dx = \int_{-\pi}^{\pi} \frac{1 + \cos 2x}{2} dx = \pi \quad (6b)$$

$$\langle 1, 1 \rangle \equiv \int_{-\pi}^{\pi} dx = 2\pi. \quad (6c)$$

Thus the vectors in B can be taken as basis vectors for a particular subspace of the vector space $L^2[-\pi, \pi]$. The inverse of (3) would imply that *this subspace is the whole vector space $L^2[-\pi, \pi]$* . This fact is provided by the *Fourier theorem* which states that: *A periodic sectionally smooth function $f(x)$ with a period of 2π can be expressed as a Fourier series expansion as in (3):*

$$f \in L^2[-\pi, \pi] \quad \Rightarrow \quad f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx), \quad (7)$$

where the coefficients a_n and b_n are constants. Using the orthogonality conditions (5) and the normalization conditions (6) of the *basis vectors B* , the *Fourier coefficients a_n and b_n* may be obtained by taking a scalar product of f with each *basis*

vector in B .

$$\begin{aligned} \langle f, 1 \rangle &= \langle a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx), 1 \rangle \\ &= a_0 \langle 1, 1 \rangle = 2\pi a_0 \end{aligned} \quad (8a)$$

$$\begin{aligned} \langle f, \cos nx \rangle &= \langle a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx), \cos nx \rangle \\ &= a_n \langle \cos nx, \cos nx \rangle = \pi a_n, \quad n \neq 0 \end{aligned} \quad (8b)$$

$$\begin{aligned} \langle f, \sin nx \rangle &= \langle a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx), \sin nx \rangle \\ &= b_n \langle \sin nx, \sin nx \rangle = \pi b_n, \quad n \neq 0 \end{aligned} \quad (8c)$$

Thus we obtain

$$a_0 = \frac{1}{2\pi} \langle f, 1 \rangle = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx \quad (9a)$$

$$a_n = \frac{1}{\pi} \langle f, \cos nx \rangle = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos nx dx, \quad n \neq 0 \quad (9b)$$

$$b_n = \frac{1}{\pi} \langle f, \sin nx \rangle = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin nx dx, \quad n \neq 0 \quad (9c)$$

Example 1 : Let us consider the Fourier series expansion (7) of the periodic square-well function

$$f(x) = \begin{cases} 0 & -\pi < x < 0 \\ 1 & 0 < x < \pi \end{cases} \quad f(x + 2n\pi) = f(x), n = 1, 2, 3, \dots \quad (10)$$

The Fourier coefficients may be obtained from (9):

$$a_0 = \frac{1}{2\pi} \int_0^{\pi} dx = \frac{1}{2} \quad (11a)$$

$$a_k = \frac{1}{\pi} \int_0^{\pi} \cos(kx) dx = \frac{1}{\pi} \left[\frac{\sin kx}{k} \right]_0^{\pi} = 0 \quad (11b)$$

$$b_k = \frac{1}{\pi} \int_{-\pi}^{\pi} \sin(kx) dx = -\frac{1}{\pi} \left[\frac{\cos kx}{k} \right]_0^{\pi} = \frac{1 - (-1)^k}{k\pi} \quad (11c)$$

All the coefficients a_k 's for non-zero k and b_k 's with even values of k vanish. By writing the odd k 's as $2m + 1$ the Fourier series expansion of the periodic square-well may be written as:

$$f(x) = \frac{1}{2} + \sum_{m=0}^{\infty} \frac{2}{(2m+1)\pi} \sin(2m+1)x \quad (12)$$

This is coded in *FourierSeries.py* and the result is shown in Figure 1. As can be seen from the figure, the Fourier Series expansion (12) converges to the periodic square-well at all points except at the discontinuities. Two interesting observations worth to mention here:

1. At the discontinuity the Fourier converges to the mean of the two limits of the original function from both sides.

$$a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx) = \frac{1}{2} \left\{ \lim_{x \rightarrow 0^+} f(x) + \lim_{x \rightarrow 0^-} f(x) \right\} \quad (13)$$

2. There are overshoots near the edges of the discontinuities from the partial sums of the Fourier Series expansion (12). *Although the positions of the overshoot all move ever closer to the discontinuities as the number of terms increases, the magnitudes of the overshoot remain essentially constant.* This is known as *Gibbs phenomena*.

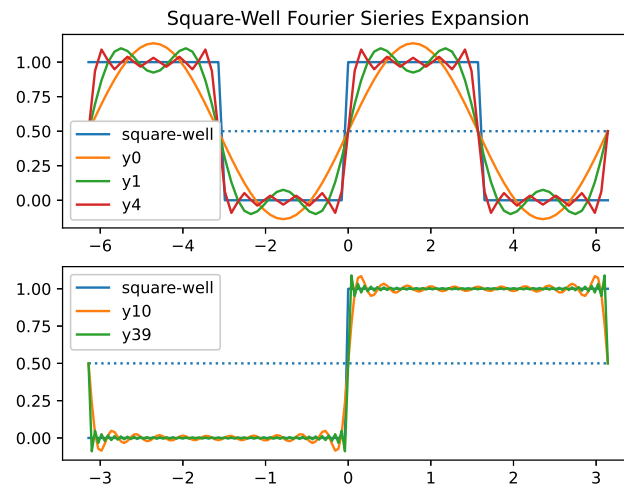


Figure 1: Fourier Series Expansion of Periodic Square-Well

```
# FourierSeries.py
# Square-Well Fourier Series Expansion
```

```

import numpy as np
import matplotlib.pyplot as plt
Pi= np.pi
Pi2, twoOpi = 2*Pi, 2/Pi
x = np.linspace(- Pi2, Pi2, 81)
nx = len(x)
z = np.zeros(nx)
z[:21] = 1
z[40:61] = 1
plt.figure()
plt.subplot(2,1,1)
plt.title('Square-Well Fourier Series Expansion')
plt.plot(x,z,label='square-well')
plt.hlines(0.5, -Pi2,Pi2,linestyles='dotted')
iout= [0,1,4]
y = 0.5*np.ones(nx)
for i in range(5):
    n = 2*i+ 1
    xx = n*x
    fac = twoOpi/n
    y = y + fac*np.sin(xx)
    if i in iout:
        s = 'y'+str(i)
        plt.plot(x,y,label=s )
plt.legend()
###plot including higher terms and finer grid points
x = np.linspace(- Pi, Pi, 161)
nx = len(x)
z = np.zeros(nx)
z[80:] = 1
print(z[80])
iout= [10,39]
plt.subplot(2,1,2)

```

```

plt.plot(x,z,label='square-well')
plt.hlines(0.5, -Pi,Pi,linestyles='dotted')
y = 0.5*np.ones(nx)
for i in range(40):
    n = 2*i+ 1
    xx = n*x
    fac = two0pi/n
    y = y + fac*np.sin(xx)
    if i in iout:
        s = 'y'+str(i)
        plt.plot(x,y,label=s )
plt.legend()
fig = plt.gcf()
fig.savefig('FourierSeries.eps', format='eps')
plt.show()

```