

Direct Fidelity Estimation

Ahmed Aman

August 2025

What is DFE and Why?

- DFE is a new technique to estimate if an arbitrary pure quantum state is close to a desired ideal state.
- It's crucial for verifying the quality of state preparation in quantum computation.
- This estimation uses a few "constant" numbers of Pauli measurements, selected randomly based on an importance-weighted rule.
- DFE is faster and requires fewer resources than traditional Quantum State Tomography (QST).
- It allows for the validation of a quantum state's closeness to a desired state without needing to reconstruct the full state itself.
- It can be applied to a wider range of states than the Entanglement Witness method.

Setting of the Protocol

- **Hilbert Space** \mathcal{H} with $\dim(\mathcal{H}) = d = 2^n$.

Assumptions

- 1 The ideal state ρ is pure.
- 2 We can measure n-qubit Pauli observables, which are tensor products of single-qubit Pauli operations.
- 3 There are no assumptions of additional structure or symmetry on ρ .

Fidelity Estimation

The fidelity between the ideal state ρ and the actual state σ is given by:

$$F(\rho, \sigma) = \text{tr}([\rho^{1/2} \sigma \rho^{1/2})^{1/2}) = \text{tr}(\rho \sigma) \quad (1)$$

This can be expressed in terms of Pauli expectation values.

- Let W_k ($k = 1, \dots, d^2$) be all possible Pauli operators.
- We define the characteristic function:

$$\mathcal{X}_\rho(k) = \text{tr}(\rho W_k / d^{1/2}) \quad (2)$$

- The fidelity can then be written as:

$$F(\rho, \sigma) = \sum_k \mathcal{X}_\rho(k) \mathcal{X}_\sigma(k) \quad (3)$$

Estimation is performed by randomly selecting Pauli operators according to an importance-weighted rule, with the probability of selecting W_k given by:

$$p(k) = (\mathcal{X}_\rho(k))^2 \quad (4)$$

DFE Protocol Summarized

- 1) Generate ℓ independent and identically distributed random samples k_1, \dots, k_ℓ from the importance sampling distribution $p(k_i)$, where $\ell = \lceil 1/(\epsilon^2 \delta) \rceil$.
- 2) For each k_i , measure the observable W_{k_i} a total of m_i times, where m_i is defined as:

$$m_i = \left\lceil \frac{2}{(\mathcal{X}_\rho(k_i))^2 d \ell \epsilon^2} \ln \frac{2}{\delta} \right\rceil$$

- 3) For each measurement i , use the outcomes $A_{ij} \in \{-1, 1\}$ to calculate the estimator \tilde{X}_i :

$$\tilde{X}_i = \frac{1}{m_i \sqrt{d} \mathcal{X}_\rho(k_i)} \sum_{j=1}^{m_i} A_{ij}$$

- 4) Calculate the final fidelity estimate \tilde{Y} :

$$\tilde{Y} = \frac{1}{\ell} \sum_{i=1}^{\ell} \tilde{X}_i$$

- 5) Return \tilde{Y} as the fidelity estimator. This provides an unbiased estimate of $F(\rho, \sigma)$ with an accuracy of 2ϵ and a confidence level of $1 - 2\delta$.

Calculation of Required Sample Counts

This section details the derivation for the number of measurement settings, ℓ , and the number of measurements per setting, m_i .

- The calculations rely on two fundamental concentration inequalities: Chebyshev's inequality for ℓ and Hoeffding's inequality for m_i .
- The total number of copies of the state σ required is $m = \sum_{i=1}^{\ell} m_i$.
- While this depends on the random choices of k_i , we can bound its expectation value.

The expected number of total copies can be shown to be:

$$(m) \leq 1 + \frac{2d}{\epsilon^2 \delta} + \frac{2d}{\epsilon^2} \log(2/\delta) \quad (5)$$

Numerics

- The goal was to recreate the graph from the DFE paper.
- Parameters used: $\epsilon = 0.05$ and $\delta = 0.05$, which results in $\ell = 8000$.
- The number of qubits n is 4.

States Investigated

- Haar Random States
- W States
- GHZ States

Errors Applied (10%)

- Depolarizing Noise
- Phase Damping Noise
- Amplitude Damping Noise
- General Amplitude Damping

Fidelity Estimates Distribution Comparison

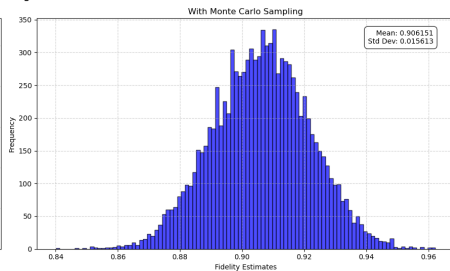
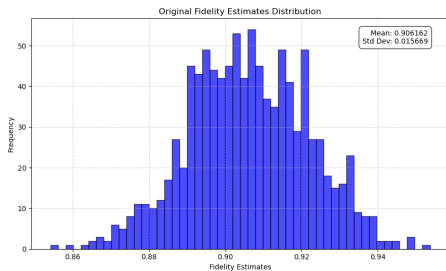
Noise Type: Depolarizing 10%

state=Haar, n=4 qubits

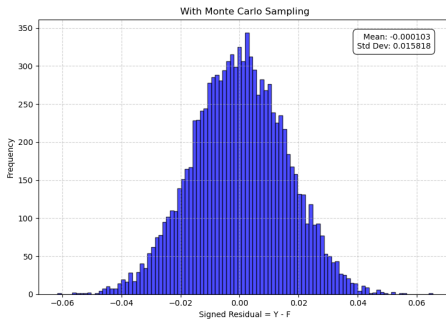
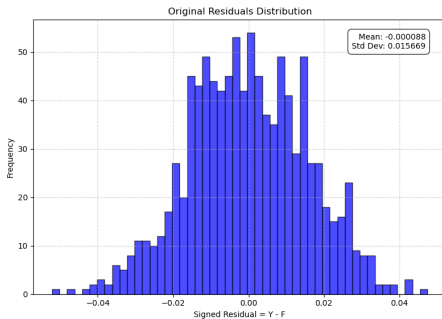
$\epsilon=0.05$, $\delta=0.05$

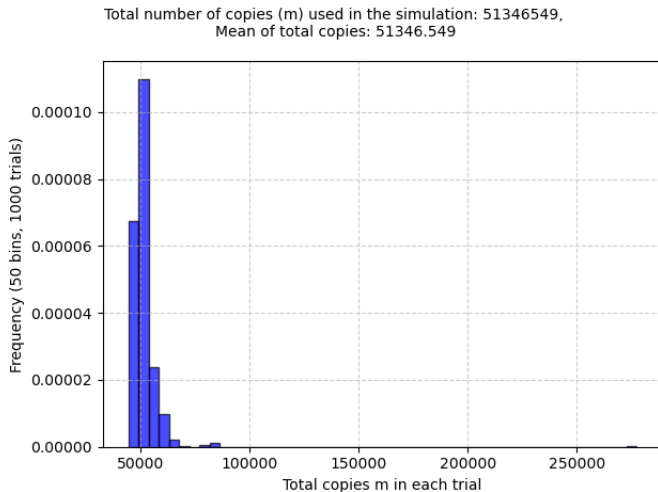
True Fidelity: 0.90626

Absolute Mean Difference between Original and Monte Carlo: 0.000109



Residuals Distribution Comparison
Noise Type: Depolarizing 10%
state=Haar, n=4 qubits
 $\epsilon=0.05$, $\delta=0.05$
Absolute Mean Difference: 0.000015





Fidelity Estimates Distribution Comparison

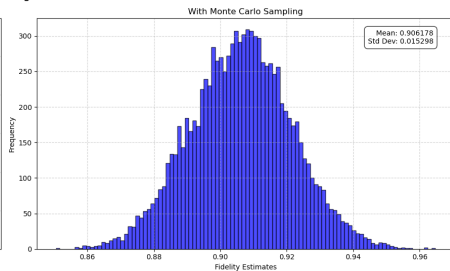
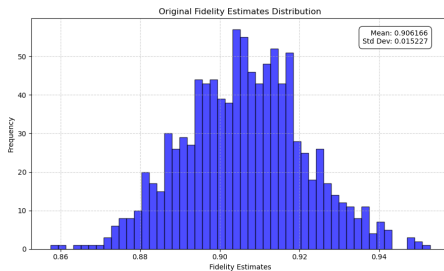
Noise Type: Amplitude 10%

state=Haar, n=4 qubits

$\epsilon=0.05$, $\delta=0.05$

True Fidelity: 0.90626

Absolute Mean Difference between Original and Monte Carlo: 0.000082



Fidelity Estimates Distribution Comparison

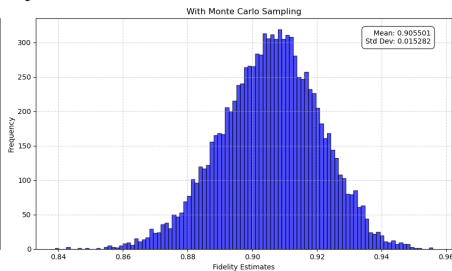
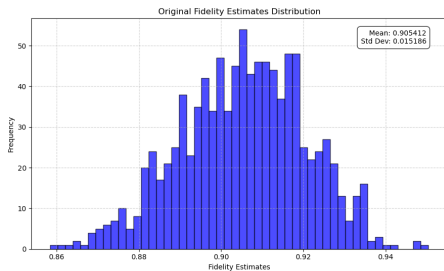
Noise Type: Phase 10%

state=Haar, n=4 qubits

$\epsilon=0.05$, $\delta=0.05$

True Fidelity: 0.90501

Absolute Mean Difference between Original and Monte Carlo: 0.000491



Fidelity Estimates Distribution Comparison

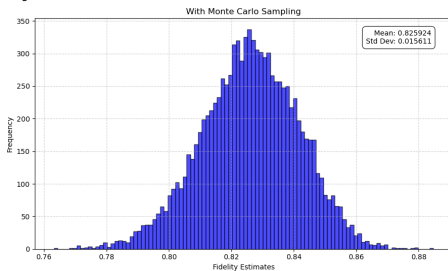
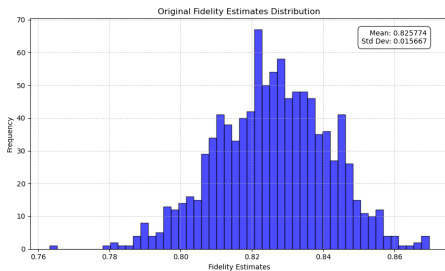
Noise Type: Generalized Amplitude 10%

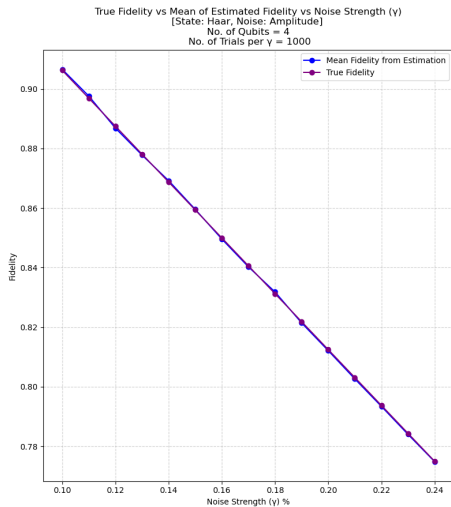
state=Haar, n=4 qubits

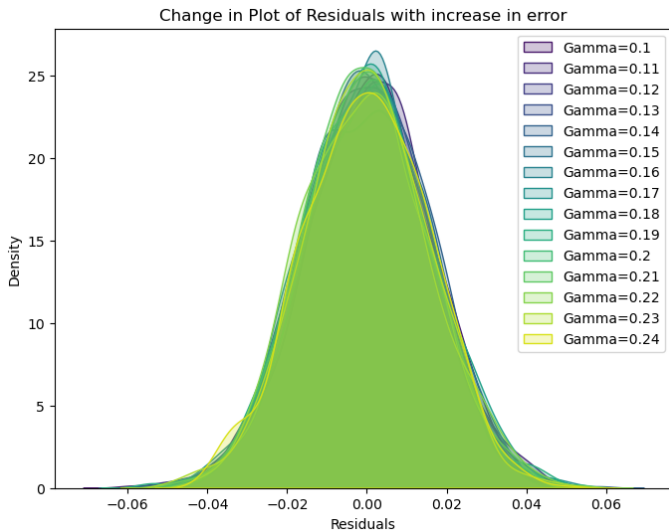
$\epsilon=0.05$, $\delta=0.05$

True Fidelity: 0.82511

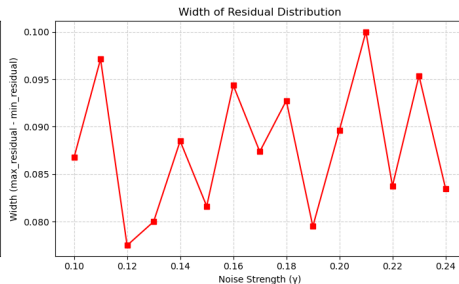
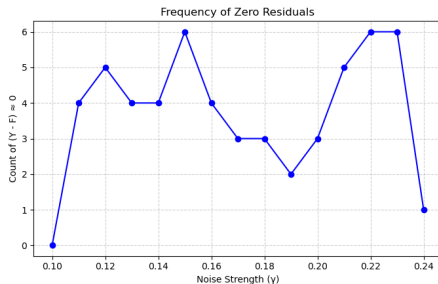
Absolute Mean Difference between Original and Monte Carlo: 0.000814



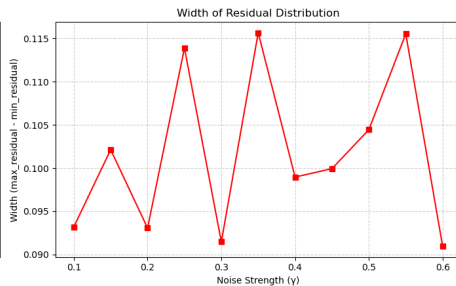
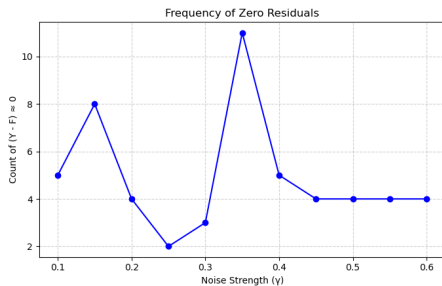




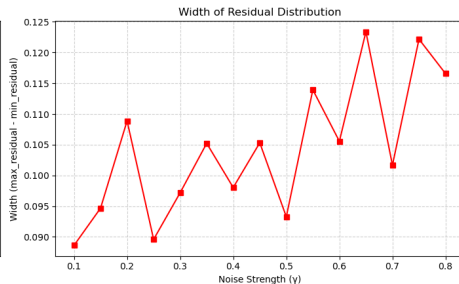
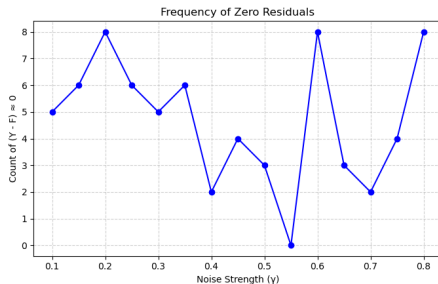
Analysis of Residuals ($Y - F$) vs. Noise Strength (γ)
 [State: W, Noise: Depolarizing]
 No. of Trials per $\gamma = 1000$
 No. of Qubits = 4



Analysis of Residuals ($Y - F$) vs. Noise Strength (γ)
[State: Haar, Noise: Depolarizing]
No. of Trials per $\gamma = 1000$
No. of Qubits = 4



Analysis of Residuals ($Y - F$) vs. Noise Strength (γ)
[State: Haar, Noise: Depolarizing]
No. of Trials per $\gamma = 1000$
No. of Qubits = 4



Other Plots and Code

- **Click here** to check out the other plots, as well as the code for the simulation.