

Fórmulas relativistas

$$\vec{p} = \frac{m\vec{u}}{\sqrt{1-u^2/c^2}} = m\gamma\vec{u} = m\vec{U}$$

donde  $\vec{U} = \frac{d\vec{x}}{dt}$

Ahora, determinaremos  $\mathcal{E}(u)$

La conexión entre  $K, K'$  para las velocidades  $\textcircled{2}$

$$\textcircled{1} \quad (u_x)_x = \frac{-c\beta \sin\theta'}{\gamma(1-\beta^2 \cos\theta')} \quad (u_z)_z = \frac{c\beta(1-\cos\theta')}{1-\beta^2 \cos\theta'}$$

para  $\theta'$  pequeños  $\cos\theta' \approx 1 - \frac{\theta'^2}{2}$   $\textcircled{3}$

$$(u_z)_z = \frac{c\beta[1 - (1 - \theta'^2/2)]}{[1 - \beta^2(1 - \theta'^2/2)]} \quad \text{Usando } \textcircled{3} \text{ en } \textcircled{2}$$

Usando  $\textcircled{3}$  en  $\textcircled{1}$  y tomando el cuadrado

$$(u_x)_x^2 = \frac{c^2 \beta^2 \theta'^2}{\gamma^2 (1-\beta^2)^2 (1 + \frac{\theta'^2}{2(1-\beta^2)})^2} \approx \frac{c^2 \beta^2 \theta'^2}{\gamma^2 (1-\beta^2)^2} + \mathcal{O}(\theta'^4)$$

$$(u_z)_z^2 = \frac{c^2 \beta^2 (\theta'^2/2)^2}{[1 - \beta^2 + \frac{\beta^2 \theta'^2}{2}]^2} \sim \mathcal{O}(\theta'^4) + \mathcal{O}(\theta'^2)$$

$$(u_a)^2 = (u_x)_x^2 + (u_z)_z^2 = \left[ \frac{c^2 \beta^2 \theta'^2}{(1-\beta^2)^2} \right] = u^2$$

$$(u_a)^2 = u^2 \quad \leftarrow$$

De la misma forma  $(u_x)_x, (u_z)_z$

$$\Rightarrow u_a^2 = (u_x)_x^2 + (u_z)_z^2 = u_a^2 - \frac{u^2}{\gamma_a^3}$$

$$\gamma_a = \frac{1}{\sqrt{1-u_a^2/c^2}}$$

Ejercicio Propuesto

Regresamos a la conservación de la energía visto en  $K$

$$\mathcal{E}(u_a) + \mathcal{E}(0) = \mathcal{E}(u_c) + \mathcal{E}(u_d)$$

$$\text{donde } u_c^2 = \frac{(u_a)^2}{\gamma_a^3}$$

Expandiendo

$$\mathcal{E}(u_c(n)) = \mathcal{E}(u_c(n)) \Big|_{n=0} + \left[ \frac{\partial \mathcal{E}(u_c)}{\partial u_c^2} \frac{\partial u_c^2}{\partial n} \right]_{n=0} \cdot n + \mathcal{O}(n^2)$$

$$\mathcal{E}(u_c(n)) = \mathcal{E}(u_a) + n \left[ \frac{\partial \mathcal{E}(u_a)}{\partial u_c^2} \frac{\partial u_c^2}{\partial n} \right]_{n=0}$$

$$\mathcal{E}(u_d(n)) = \mathcal{E}(u_d(n)) \Big|_{n=0} + n \left[ \frac{\partial \mathcal{E}(u_d)}{\partial u_d^2} \frac{\partial u_d^2}{\partial n} \right]_{n=0} + \dots$$

$$\mathcal{E}(u_d(n)) = \mathcal{E}(0) + n \left[ \frac{\partial \mathcal{E}(u_d)}{\partial u_d^2} \frac{\partial u_d^2}{\partial n} \right]_{n=0} + \dots$$

Usando estas aproximaciones en la conservación de la energía

$$\mathcal{E}(u_a) + \mathcal{E}(0) = \mathcal{E}(u_a) + n \left[ \frac{\partial \mathcal{E}(u_a)}{\partial u_c^2} \frac{\partial u_c^2}{\partial n} \right]_{n=0} + \mathcal{E}(0) + n \left[ \frac{\partial \mathcal{E}(u_d)}{\partial u_d^2} \frac{\partial u_d^2}{\partial n} \right]_{n=0}$$

$$0 = -\frac{u^2}{\gamma_a^3} \left( \frac{\partial \mathcal{E}(u_a)}{\partial u_c^2} \right)_{n=0} + n \left[ \frac{\partial \mathcal{E}(u_d)}{\partial u_d^2} \cdot 1 \right]_{n=0}$$

$$0 = -\frac{1}{\gamma_a^3} \frac{\partial \mathcal{E}(u_a)}{\partial u_a^2} + \left( \frac{\partial \mathcal{E}(u_d)}{\partial u_d^2} \right)_{n=0}$$

$$0 = -\frac{1}{\gamma_a^3} \frac{\partial \mathcal{E}(u_a)}{\partial u_a^2} + \frac{\partial \mathcal{E}(0)}{\partial u_d^2}$$

$$\frac{\partial \mathcal{E}(0)}{\partial u^2} = \frac{m}{2}$$

$$\frac{\partial \mathcal{E}(u_a)}{\partial u_a^2} = \frac{m}{2} \gamma_a^3 = \frac{m}{2(1-u_a^2/c^2)^{3/2}}$$

$\mathcal{E}(u_a) \quad u_a = u$

$$\mathcal{E}(u) = \frac{mc^2}{(1-u^2/c^2)^{1/2}} + [\mathcal{E}(0) - mc^2]$$

$$\int d\mathcal{E}(u) = \int \frac{m}{2(1-u^2/c^2)^{3/2}} du^2$$

$$\mathcal{E}(u) = \frac{mc^2}{(1-u^2/c^2)^{1/2}} + \mathcal{E}(0) - mc^2$$

Se define la energía cinética

$$T(u) = \mathcal{E}(u) - \mathcal{E}(0) = mc^2 \left[ \frac{1}{(1-u^2/c^2)^{1/2}} - 1 \right]$$

$$\lim_{u \ll c} T(u) = mc^2 \left[ 1 + \frac{1}{2} \frac{u^2}{c^2} - 1 \right] = \frac{1}{2} m u^2$$

$$\vec{p} = \frac{m\vec{u}}{\sqrt{1-u^2/c^2}} \quad y \quad \mathcal{E}(u) = \frac{mc^2}{(1-u^2/c^2)^{1/2}} + [\mathcal{E}(0) - mc^2]$$

$$K^0 \rightarrow \gamma\gamma \quad \checkmark \quad E_\gamma + E_\gamma = E_{K^0}$$

$$K^0 \rightarrow \pi^+ \pi^- \quad E_K(u) = E_\pi + E_\pi = 2E_\pi$$

$$E_\pi = \frac{E_{K^0}}{2}$$

$$\text{Pero } E_{\text{cin}} = T_\pi + \mathcal{E}_\pi(0) =$$

$$T_\pi = \frac{E_{K^0}}{2} - \mathcal{E}_\pi(0) \quad \text{Se mide en } \textcircled{1}$$

$$\mathcal{E}_\pi(0) = \frac{E_{K^0}}{2} - T_\pi$$

$$\mathcal{E}(0) = mc^2$$

$$\Rightarrow \mathcal{E}(u) = \frac{mc^2}{\sqrt{1-u^2/c^2}}$$

$$\text{Entonces } \vec{p} = \frac{m\vec{u}}{\sqrt{1-u^2/c^2}} \quad \mathcal{E}(u) = \frac{mc^2}{\sqrt{1-u^2/c^2}}$$

$$\Rightarrow \vec{u} = \frac{c^2 \vec{p}}{E}$$

Se puede construir el 4-vector

$$E/c, \vec{p} \rightarrow \left[ \frac{mc}{\sqrt{1-u^2/c^2}}, \frac{m\vec{u}}{\sqrt{1-u^2/c^2}} \right]$$

$$= [m\gamma c, m\vec{U}]$$

$$\uparrow \quad \uparrow$$

$$p_0, \vec{p}$$

La longitud del 4-vector momento

$$p_0 p_0 - \vec{p} \cdot \vec{p} = p'_0 p'_0 - \vec{p}' \cdot \vec{p}'$$

$$x_0 x_0 - \vec{x} \cdot \vec{x} = x'_0 x'_0 - \vec{x}' \cdot \vec{x}'$$

$$(x_0, \vec{x}) =$$

$$p_0^2 - \vec{p} \cdot \vec{p} = (E/c)^2 - \vec{p} \cdot \vec{p}$$

$$= \frac{m^2 c^4}{(1-u^2/c^2)^2} \cdot \frac{1}{c^2} - \frac{m^2 u^2}{(1-u^2/c^2)^2}$$

$$= m^2 c^2 = p_0'^2 - \vec{p}' \cdot \vec{p}'$$

$m$ : masa en reposo

$$p_0^2 - \vec{p} \cdot \vec{p} = m^2 c^2 \quad \checkmark \quad p_0 = E/c$$

$$\text{También } E^2 = m^2 c^4 + p^2 c^2 \quad \checkmark$$

$$c = 1$$

$$p_0^2 - \vec{p} \cdot \vec{p} = m^2$$

$$E^2 = m^2 + p^2$$

$$[E] = \text{GeV, MeV, eV} \quad [m] = \frac{\text{GeV}}{c^2}, \frac{\text{MeV}}{c^2}, \frac{\text{eV}}{c^2}$$