

## planetmath.org

Math for the people, by the people.

## derivation of integral representations of Jacobi $\vartheta$ functions

 ${\bf Canonical\ name} \quad {\bf Derivation Of Integral Representations Of Jacobivar theta Functions}$ 

Date of creation 2013-03-22 14:39:54 Last modified on 2013-03-22 14:39:54

Owner rspuzio (6075) Last modified by rspuzio (6075)

Numerical id 26

Author rspuzio (6075) Entry type Derivation Classification msc 33E05 By rearranging the Fourier series of  $\cos(ux)$ , one obtains the series

$$\frac{\pi \cos(ux)}{2u \sin(\pi u)} = \frac{1}{2u^2} + \sum_{n=1}^{\infty} (-1)^n \frac{\cos(nx)}{u^2 - n^2}$$

This equation which is valid for all real values of x such that  $-\pi \leq x \leq \pi$  and all non-integral complex values of u. By comparison with the convergent series  $\sum_{n=0}^{\infty} 1/n^2$ , it follows that this series is absolutely convergent. Note that this series may be viewed as a Mittag-Leffler partial fraction expansion.

Let y be a positive real number. Multiply both by  $2ue^{-yu^2}$  and integrate.

$$\int_{i-\infty}^{i+\infty} \frac{\pi \cos(ux)e^{-yu^2}}{\sin(\pi u)} dv = 2 \int_{i-\infty}^{i+\infty} e^{-yu^2} \left[ \frac{1}{2u^2} + \sum_{n=0}^{\infty} (-1)^n \frac{\cos(nx)}{u^2 - n^2} \right] u du$$

Because of the exponential, the integrand decays rapidly as  $u \to i \pm \infty$  provided that  $\Re u > 0$ , and hence the integral converges absolutely. Make a change of variables  $v = u^2$ 

$$= \int_{P} e^{-yv} \left[ \frac{1}{2v} + \sum_{n=1}^{\infty} (-1)^{n} \frac{\cos(nx)}{v - n^{2}} \right] dv$$

The contour of integration P is a parabola in the complex v-plane, symmetric about the real axis with vertex at v = -1, which encloses the real axis. Its equation is  $\Re v + 1 = 2(\Im v)^2$ 

Let  $S_m$  (m is an integer) be the straight line segment joining the points  $v = (i + m + 1/2)^2$  and  $v = (i - m - 1/2)^2$ . Along this line segment, we may bound the integrand in absolute value as follows:

$$\left| \sum_{n=1}^{\infty} (-1)^n \frac{\cos(nx)}{v - n^2} \right| \le \sum_{n=1}^{\infty} \frac{(-1)^n}{|v - n^2|} \le \sum_{n=1}^{\infty} \frac{(-1)^n}{|v_m - n^2|}$$

where  $v_m = m^2 + m - 3/4$  is the point of intersection of  $S_m$  with the real axis. To proceed further, we break up the last summation into two parts.

Since the squares closest in absolute value to  $v_m$  are  $m^2$  and  $(m+1)^2 = m^2 + 2m + 1$ , it follows that  $|v_m - n^2| \ge |m - 3/4|$  for all m, n. Hence, we have

$$\sum_{i=1}^{2m} \frac{1}{|v_m - n^2|} \le \frac{2m}{m - 3/4} \le 8$$

When n > 2m, we have  $n^2 \ge (2m+1)^2 = 4m^2 + 4m + 1 > 4m^2 + 4m - 3 = 4v_m$ . Hence,  $|n^2 - v_m| > 3n^2/4$  and

$$\sum_{n=2m+1}^{\infty} \frac{1}{|v_m - n^2|} < \frac{4}{3} \sum_{n=2m+1}^{\infty} \frac{1}{n^2} < \frac{4}{3} \sum_{n=1}^{\infty} \frac{1}{n^2} = \frac{2\pi}{9}$$

Finally  $1/(2v_m) < 1/2$  since  $v_m > 1$  when  $m \ge 1$ . Also,  $|e^{-yv}| = e^{-y\Re v} - e^{-yv_m} < e^{-ym^2}$ . From these observations, we conclude that

$$\left| \int_{S_m} e^{-yv} \left[ \frac{1}{2v} + \sum_{n=1}^{\infty} (-1)^n \frac{\cos(nx)}{v - n^2} \right] dv \right| < e^{-ym^2} \left( 1 + 8 + \frac{2\pi}{9} \right) \int_{S_m} dv = (4m + 2) \left( 9 + \frac{2\pi}{9} \right) e^{-ym^2}$$

Note that this quantity approaches 0 in the limit  $m \to \infty$ .

Let  $P_m$  be the arc of the parabola P bounded by the endpoints of  $S_m$ . Together,  $S_m$  and  $P_m$  form a closed contour which encloses poles of the integrand. Hence, by the residue theorem, we have

$$\int_{P_m} e^{-yv} \left[ \frac{1}{2v} + \sum_{n=1}^{\infty} (-1)^n \frac{\cos(nx)}{v - n^2} \right] dv + \int_{S_m} e^{-yv} \left[ \frac{1}{2v} + \sum_{n=1}^{\infty} (-1)^n \frac{\cos(nx)}{v - n^2} \right] dv =$$

$$2\pi i \sum_{n=1}^{\infty} (-1)^n \cos(nx) e^{-n^2 y}$$

Taking the limit  $m \to \infty$  we obtain

$$\int_{P} e^{-yv} \left[ \frac{1}{2v} + \sum_{n=1}^{\infty} (-1)^n \frac{\cos(nx)}{v - n^2} \right] dv = 2\pi i \left( \frac{1}{2} + \sum_{n=1}^{\infty} (-1)^n \cos(nx) e^{-n^2 y} \right)$$

Going back to the beginning of the proof, where the integral on the left hand side was expressed as an integral with respect to u, we obtain

$$\int_{i-\infty}^{i+\infty} \frac{\pi \cos(ux)e^{-yu^2}}{\sin(\pi u)} dv = 2\pi i \left(\frac{1}{2} + \sum_{n=1}^{\infty} (-1)^n \cos(nx)e^{-n^2y}\right)$$

Making a change of variables  $x=2z,y=-i\pi\tau$  and tidying up some, we obtain

$$\int_{i-\infty}^{i+\infty} \frac{\cos(2uz)e^{i\pi\tau u^2}}{\sin(\pi u)} \, dv = i\left(1 + 2\sum_{n=1}^{\infty} (-1)^n e^{i\pi n^2\tau} \cos(2nz)\right) = i\vartheta_4(z|\tau)$$

Because of the initial assumption about the Fourier series, we only know that this formula is valid when  $\tau$  is purely imaginary with strictly positive imaginary part and z is real and  $\pi/2 < z < \pi/2$ . However, we can use analytic continuation to extend the domain of its validity. On the one hand, the theta function on the right-hand side is analytic for all z and all  $\tau$  such that  $\Im \tau > 0$ .

On the other hand, I claim that the integral on the left hand side is also an analytic function of z and  $\tau$  whenever  $\Im \tau > 0$ . To validate this claim, we need to examine the behaviour of the integrand as  $u \to i \pm \infty$ . The contribution of the denominator is bounded;

$$\left| \frac{1}{\sin \pi u} \right| < c$$

for some constant c whenever  $\Im u = 1$ . The absolute value of the cosine in the numerator is easy to bound:

$$|\cos(2uz)| < e^{2|u||z|}$$

To bound the remaining term, let us examine the argument of the exponential carefully:

$$\Im(\tau u^2) = 2\Re\tau \Re u + \Im\tau(\Re u)^2 - \Im\tau = \Im\tau \left( \left(\Re u + \frac{\Re\tau}{\Im\tau}\right)^2 - 1 - \left(\frac{\Re\tau}{\Im\tau}\right)^2 \right)$$

Therefore, if  $|\Re u| > 1 + 3|\Re \tau|/(\Im \tau)$ , it will be the case that  $\Im(\tau u^2) \ge \Im \tau (\Re u)^2/9$ , and so

$$\left| e^{i\pi\tau u^2} \right| = e^{-\pi\Im(\tau u^2)} \le e^{-\pi\Im\tau\,(\Re u)^2/9}$$

Taken together, the estimates of the last paragraph imply that

$$\left| \int_{i+R}^{i+\infty} \frac{\cos(2uz)e^{i\pi\tau u^2}}{\sin(\pi u)} \right| < c \int_{i+R}^{i+\infty} e^{2|u||z| - \pi\Im\tau \,(\Re u)^2/9}$$

when  $R > 1 + 3|\Re \tau|/(\Im \tau)$ . If we impose the further conditions

$$R > \frac{180|z|}{\pi \, \Im \tau} \qquad R^2 > \frac{180|z|}{\pi \, \Im \tau} \qquad ,$$

it will be the case that

$$2|u||z| - \pi\Im\tau (\Re u)^2/9 < 2\Re u |z| + 2|z| - \pi\Im\tau (\Re u)^2/9 <$$

$$(2\Re u |z| - \pi\Im\tau (\Re u)^2/180) + (2|z| - \pi\Im\tau (\Re u)^2/180) - \pi\Im\tau (\Re u)^2/10 < -\pi\Im\tau (\Re u)^2/10 ,$$

and hence

$$\left| \int_{i+R}^{i+\infty} \frac{\cos(2uz)e^{i\pi\tau u^2}}{\sin(\pi u)} \, du \right| < c \int_{i+R}^{i+\infty} e^{-\pi\Im\tau \, (\Re u)^2/10} \, du < \frac{5c}{\pi\,\Im\tau} Re^{-\pi\Im\tau \, R^2/10}$$

Likewise, under the same restriction on R,

$$\left| \int_{i-\infty}^{i-R} \frac{\cos(2uz)e^{i\pi\tau u^2}}{\sin(\pi u)} \, du \right| < c \int_{i+R}^{i+\infty} e^{-\pi\Im\tau \, (\Re u)^2/10} \, du < \frac{5c}{\pi\,\Im\tau} Re^{-\pi\Im\tau \, R^2/10}$$

Since the contour of integration is compact and the integrand is analytic in a neighborhood of the contour,

$$\int_{i-R}^{i+R} \frac{\cos(2uz)e^{i\pi\tau u^2}}{\sin(\pi u)} du$$

will be an analytic function of z and  $\tau$ . Suppose that z and  $\tau$  are restricted to bounded regions of the complex plane and that, furthermore,  $Im\tau$  is positive and bounded away from zero. Then the inequalities of the last paragraph imply that the integral converges uniformly as  $R \to \infty$ , and hence

$$\int_{i-\infty}^{i+\infty} \frac{\cos(2uz)e^{i\pi\tau u^2}}{\sin(\pi u)} \, du$$

is an analytic function of u and z in the domain  $\Im \tau > 0$ .

Thus, by the fundamental theorem of analytic continuation, we may conclude that

$$\int_{i-\infty}^{i+\infty} \frac{\cos(2uz)e^{i\pi\tau u^2}}{\sin(\pi u)} \, dv = i\left(1 + 2\sum_{n=1}^{\infty} (-1)^n e^{i\pi n^2 \tau} \cos(2nz)\right) = i\vartheta_4(z|\tau)$$

throughout this domain.

Finally, integral representations of the remaining three theta functions may be easily obtained from this one by adding the appropriate half-quasiperiods to z.