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Bessel's equation

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The linear differential equation

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} + (x^2 - p^2)y = 0, \quad (1)$$

in which p is a constant (non-negative if it is real), is called the *Bessel's equation*. We derive its general solution by trying the series form

$$y = x^r \sum_{k=0}^{\infty} a_k x^k = \sum_{k=0}^{\infty} a_k x^{r+k}, \quad (2)$$

due to Frobenius. Since the parameter r is indefinite, we may regard a_0 as distinct from 0.

We substitute (2) and the derivatives of the series in (1):

$$x^2 \sum_{k=0}^{\infty} (r+k)(r+k-1) a_k x^{r+k-2} + x \sum_{k=0}^{\infty} (r+k) a_k x^{r+k-1} + (x^2 - p^2) \sum_{k=0}^{\infty} a_k x^{r+k} = 0.$$

Thus the coefficients of the powers x^r , x^{r+1} , x^{r+2} and so on must vanish, and we get the system of equations

$$\begin{cases} [r^2 - p^2]a_0 = 0, \\ [(r+1)^2 - p^2]a_1 = 0, \\ [(r+2)^2 - p^2]a_2 + a_0 = 0, \\ \dots \\ [(r+k)^2 - p^2]a_k + a_{k-2} = 0. \end{cases} \quad (3)$$

The last of those can be written

$$(r+k-p)(r+k+p)a_k + a_{k-2} = 0.$$

Because $a_0 \neq 0$, the first of those (the indicial equation) gives $r^2 - p^2 = 0$, i.e. we have the roots

$$r_1 = p, \quad r_2 = -p.$$

Let's first look the the solution of (1) with $r = p$; then $k(2p+k)a_k + a_{k-2} = 0$, and thus

$$a_k = -\frac{a_{k-2}}{k(2p+k)}.$$

From the system (3) we can solve one by one each of the coefficients a_1, a_2, \dots and express them with a_0 which remains arbitrary. Setting for k the integer values we get

$$\begin{cases} a_1 = 0, & a_3 = 0, \dots, a_{2m-1} = 0; \\ a_2 = -\frac{a_0}{2(2p+2)}, & a_4 = \frac{a_0}{2 \cdot 4(2p+2)(2p+4)}, \dots, a_{2m} = \frac{(-1)^m a_0}{2 \cdot 4 \cdot 6 \dots (2m)(2p+2)(2p+4) \dots (2p+2m)} \end{cases} \quad (4)$$

(where $m = 1, 2, \dots$). Putting the obtained coefficients to (2) we get the particular solution

$$y_1 := a_0 x^p \left[\frac{x^2}{2(2p+2)} + \frac{x^4}{2 \cdot 4(2p+2)(2p+4)} - \frac{x^6}{2 \cdot 4 \cdot 6(2p+2)(2p+4)(2p+6)} + \dots \right] \quad (5)$$

In order to get the coefficients a_k for the second root $r_2 = -p$ we have to look after that

$$(r_2 + k)^2 - p^2 \neq 0,$$

or $r_2 + k \neq p = r_1$. Therefore

$$r_1 - r_2 = 2p \neq k$$

where k is a positive integer. Thus, when p is not an integer and not an integer added by $\frac{1}{2}$, we get the second particular solution, gotten of (5) by replacing p by $-p$:

$$y_2 := a_0 x^{-p} \left[1 - \frac{x^2}{2(-2p+2)} + \frac{x^4}{2 \cdot 4(-2p+2)(-2p+4)} - \frac{x^6}{2 \cdot 4 \cdot 6(-2p+2)(-2p+4)(-2p+6)} + \dots \right] \quad (6)$$

The power series of (5) and (6) converge for all values of x and are linearly independent (the ratio y_1/y_2 tends to 0 as $x \rightarrow \infty$). With the appointed value

$$a_0 = \frac{1}{2^p \Gamma(p+1)},$$

the solution y_1 is called the *Bessel function of the first kind and of order p* and denoted by J_p . The similar definition is set for the first kind Bessel function of an arbitrary order $p \in \mathbb{R}$ (and \mathbb{C}). For $p \notin \mathbb{Z}$ the general solution of the Bessel's differential equation is thus

$$y := C_1 J_p(x) + C_2 J_{-p}(x),$$

where $J_{-p}(x) = y_2$ with $a_0 = \frac{1}{2^{-p}\Gamma(-p+1)}$.

The explicit expressions for $J_{\pm p}$ are

$$J_{\pm p}(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m \pm p + 1)} \left(\frac{x}{2}\right)^{2m \pm p}, \quad (7)$$

which are obtained from (5) and (6) by using the last for gamma function.

E.g. when $p = \frac{1}{2}$ the series in (5) gets the form

$$y_1 = \frac{x^{\frac{1}{2}}}{\sqrt{2}\Gamma(\frac{3}{2})} \left[1 - \frac{x^2}{2 \cdot 3} + \frac{x^4}{2 \cdot 4 \cdot 3 \cdot 5} - \frac{x^6}{2 \cdot 4 \cdot 6 \cdot 3 \cdot 5 \cdot 7} + \dots \right] = \sqrt{\frac{2}{\pi x}} \left(x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots \right).$$

Thus we get

$$J_{\frac{1}{2}}(x) = \sqrt{\frac{2}{\pi x}} \sin x;$$

analogically (6) yields

$$J_{-\frac{1}{2}}(x) = \sqrt{\frac{2}{\pi x}} \cos x,$$

and the general solution of the equation (1) for $p = \frac{1}{2}$ is

$$y := C_1 J_{\frac{1}{2}}(x) + C_2 J_{-\frac{1}{2}}(x).$$

In the case that p is a non-negative integer n , the “+” case of (7) gives the solution

$$J_n(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! (m+n)!} \left(\frac{x}{2}\right)^{2m+n},$$

but for $p = -n$ the expression of $J_{-n}(x)$ is $(-1)^n J_n(x)$, i.e. linearly dependent on $J_n(x)$. It can be shown that the other solution of (1) ought to be searched in the form $y = K_n(x) = J_n(x) \ln x + x^{-n} \sum_{k=0}^{\infty} b_k x^k$. Then the general solution is $y := C_1 J_n(x) + C_2 K_n(x)$.

Other formulae

The first kind Bessel functions of integer order have the generating function F :

$$F(z, t) = e^{\frac{z}{2}(t - \frac{1}{t})} = \sum_{n=-\infty}^{\infty} J_n(z) t^n \quad (8)$$

This function has an essential singularity at $t = 0$ but is analytic elsewhere in \mathbb{C} ; thus F has the Laurent expansion in that point. Let us prove (8) by using the general expression

$$c_n = \frac{1}{2\pi i} \oint_{\gamma} \frac{f(t)}{(t-a)^{n+1}} dt$$

of the coefficients of Laurent series. Setting to this $a := 0$, $f(t) := e^{\frac{z}{2}(t-\frac{1}{t})}$, $\zeta := \frac{zt}{2}$ gives

$$c_n = \frac{1}{2\pi i} \oint_{\gamma} \frac{e^{\frac{zt}{2}} e^{-\frac{z}{2t}}}{t^{n+1}} dt = \frac{1}{2\pi i} \left(\frac{z}{2}\right)^n \oint_{\delta} \frac{e^{\zeta} e^{-\frac{z^2}{4\zeta}}}{\zeta^{n+1}} d\zeta = \sum_{m=0}^{\infty} \frac{(-1)^m}{m!} \left(\frac{z}{2}\right)^{2m+n} \frac{1}{2\pi i} \oint_{\delta} \zeta^{-m-n-1} e^{\zeta} d\zeta.$$

The paths γ and δ go once round the origin anticlockwise in the t -plane and ζ -plane, respectively. Since the residue of $\zeta^{-m-n-1} e^{\zeta}$ in the origin is $\frac{1}{(m+n)!} = \frac{1}{\Gamma(m+n+1)}$, the <http://planetmath.org/CauchyResidueTheorem> residue theorem gives

$$c_n = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+n+1)} \left(\frac{z}{2}\right)^{2m+n} = J_n(z).$$

This that F has the Laurent expansion (8).

By using the generating function, one can easily derive other formulae, e.g. the of the Bessel functions of integer order:

$$J_n(z) = \frac{1}{\pi} \int_0^{\pi} \cos(n\varphi - z \sin \varphi) d\varphi$$

Also one can obtain the addition formula

$$J_n(x+y) = \sum_{\nu=-\infty}^{\infty} J_{\nu}(x) J_{n-\nu}(y)$$

and the series of cosine and sine:

$$\cos z = J_0(z) - 2J_2(z) + 2J_4(z) - + \dots$$

$$\sin z = 2J_1(z) - 2J_3(z) + 2J_5(z) - + \dots$$

References

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