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telegraph equation

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$$f_{rr}'' - af_{tt}'' - bf_t' - cf = 0, (1)$$

where x is distance, t is time and a, b, c are non-negative constants. The equation is a generalised form of the wave equation.

If the initial conditions are $f(x, 0) = f'_t(x, 0) = 0$ and the boundary conditions f(0, t) = g(t), $f(\infty, t) = 0$, then the Laplace transform of the solution function f(x, t) is

$$F(x, s) = G(s)e^{-x\sqrt{as^2 + bs + c}}. (2)$$

In the special case $b^2 - 4ac = 0$, the solution is

$$f(x,t) = e^{-\frac{bx}{2\sqrt{a}}}g(t - x\sqrt{a})H(t - x\sqrt{a}). \tag{3}$$

Justification of (2). Transforming the partial differential equation (1) (x) may be regarded as a parameter) gives

$$F_{xx}''(x, s) - a[s^2F(x, s) - sf(x, 0) - f_t'(x, 0)] - b[sF(x, s) - f(x, 0)] - cF(x, s) = 0,$$

which due to the initial conditions simplifies to

$$F''_{xx}(x, s) = (\underbrace{as^2 + bs + c}_{K^2})F(x, s).$$

The solution of this ordinary differential equation is

$$F(x, s) = C_1 e^{Kx} + C_2 e^{-Kx}.$$

Using the latter boundary condition, we see that

$$F(\infty, s) = \int_0^\infty e^{-st} f(\infty, t) dt \equiv 0,$$

whence $C_1 = 0$. Thus the former boundary condition implies

$$C_2 = F(0, s) = \mathcal{L}\{g(t)\} = G(s).$$

So we obtain the equation (2).

Justification of (3). When the http://planetmath.org/QuadraticFormuladiscriminant of the quadratic equation $as^2+bs+c=0$ vanishes, the http://planetmath.org/Equationroots coincide to $s=-\frac{b}{2a}$, and $as^2+bs+c=a(s+\frac{b}{2a})^2$. Therefore (2) reads

$$F(x, s) = G(s)a^{-x\sqrt{a}(s+\frac{b}{2a})} = e^{-\frac{bx}{2\sqrt{a}}}e^{-x\sqrt{a}s}G(s).$$

According to the delay theorem, we have

$$\mathcal{L}^{-1}\{e^{-ks}G(s)\} = g(t-k)H(t-k).$$

Thus we obtain for $\mathcal{L}^{-1}\{F(x, s)\}$ the expression of (3).