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particle moving on the astroid at constant frequency

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In parametric Cartesian equations, the astroid can be represented by

$$x = a \cos^3 \omega t, \quad y = a \sin^3 \omega t,$$

where  $a > 0$  is a known constant,  $\omega > 0$  is the constant angular frequency, and  $t \in [0, \infty)$  is the time parameter. Thus the position vector of a particle, moving over the astroid, is

$$\mathbf{r} = a \cos^3 \omega t \mathbf{i} + a \sin^3 \omega t \mathbf{j},$$

and its velocity

$$\mathbf{v} = -3a\omega \sin \omega t \cos^2 \omega t \mathbf{i} + 3a\omega \sin^2 \omega t \cos \omega t \mathbf{j},$$

where  $\{\mathbf{i}, \mathbf{j}\}$  is a reference basis. Hence for the particle speed we have

$$v = 3a\omega \sin \omega t \cos \omega t.$$

From the last two equations we get the tangent vector

$$\mathbf{T} = -\sin \omega t \mathbf{i} + \cos \omega t \mathbf{j},$$

and by using the well known formula <sup>1</sup>

$$\left\| \frac{d\mathbf{T}}{dt} \right\| = \frac{v}{\rho},$$

$\rho > 0$  being the radius of curvature at any instant  $t$ , we arrive to the useful equation

$$v = \omega \rho.$$

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<sup>1</sup>By applying the chain rule,

$$\left\| \frac{d\mathbf{T}}{dt} \right\| = \left\| \frac{d\mathbf{T}}{ds} \right\| \left| \frac{ds}{dt} \right| = \left\| \frac{\mathbf{N}}{\rho} \right\| v = \frac{v}{\rho},$$

by Frenet-Serret.  $\mathbf{N}$  is the normal vector.