

Quantum Information and Computing 2021-2022

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1-D quantum harmonic oscillator-Eigenvalues problem

$$\hat{H} |\psi\rangle = \left[\frac{\hat{p}^2}{2m} + \frac{1}{2m} \omega^2 q^2 \right] |\psi\rangle = E |\psi\rangle$$

- good approximation of any 1-D system around stable equilibrium point
- analytical solution
- **time independent Shrodinger equation** for harmonic potential

Representation: **coordinate basis**

$$\hat{q} \rightarrow \hat{x}, \quad \hat{p} \rightarrow -i\hbar \frac{\partial}{\partial \hat{x}}$$

Natural units for both energy and distance

$$\hbar\omega x \rightarrow x, \quad \sqrt{\frac{\hbar}{m\omega}} E \rightarrow E$$

$$\hat{H} |\psi\rangle = \left[\frac{-i}{2} \frac{\partial}{\partial \hat{x}^2} + \frac{1}{2} \omega^2 x^2 \right] |\psi\rangle = E |\psi\rangle$$

- Eigenvalues $E_n = n + \frac{1}{2}, \quad n = 0, 1, 2, \dots$
- Eigenvectors: $\psi_n(x) = \frac{1}{\sqrt{2^n n!}} \pi^{-\frac{1}{4}} \exp\left(-\frac{x^2}{2}\right) H_n(x)$,
where $H_n(x)$ are the Hermite polynomials

Computational representation: **discretization**

$$x \rightarrow x_k, \quad x_k = L1 + k * \Delta x$$

in which we are building a grid of $N + 1$ points in the interval $[L1, L2]$, with a spacing

$$\Delta x = \frac{L2 - L1}{N}$$

The discretized Shrodinger equation is then, using finite differences method:

$$[\hat{H}|\psi\rangle]_{x=x_k} = \frac{1}{2}\omega^2 x_k^2 + \frac{\psi(x = x_{k+1}) - 2\psi(x = x_k) + \psi(x = x_{k-1}))}{\Delta x^2} + O(\Delta x^2)$$

The representation of the system Hamiltonian in the computational basis ends up to be **tridiagonal** and **symmetric**:

$$H_{i,i+1} = H_{i,i-1} = -\frac{1}{2\Delta x^2}, \quad H_{i,i} = \frac{1}{2\Delta x^2} + \frac{1}{2}\omega^2 x_i^2$$

$$H = \begin{bmatrix} \frac{1}{2\Delta x^2} + \frac{1}{2}\omega^2 x_1^2 & -\frac{1}{2\Delta x^2} & & & 0 \\ -\frac{1}{2\Delta x^2} & \frac{1}{2\Delta x^2} + \frac{1}{2}\omega^2 x_2^2 & -\frac{1}{2\Delta x^2} & & \\ \vdots & -\frac{1}{2\Delta x^2} & \frac{1}{2\Delta x^2} + \frac{1}{2}\omega^2 x_3^2 & -\frac{1}{2\Delta x^2} & \\ 0 & & \ddots & & \\ & & & -\frac{1}{2\Delta x^2} & \frac{1}{2\Delta x^2} + \frac{1}{2}\omega^2 x_{N+1}^2 \end{bmatrix}$$

The compilation and execution of the Fortran script and module exploited are controlled by running:

python execution.py L1 L2 dx

where L1, L2 are the range extrema and dx is the discretization spacing.
In the *Hermitian.f90* we exploit the DSTEVD fortran subroutine for real, symmetric, tridiagonal matrices diagonalization to solve our eigenvalues problem.

```
N = size(h, 1)
LDZ = MAX(1, N)

allocate(eigvalues(N)); allocate(UD(N-1)); allocate(WORK(MAX(1, 2*N-2))); allocate(eigvectors(LDZ, N))

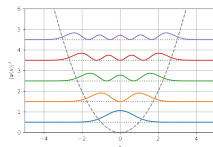
do ii=1, N
    eigvalues(ii) = h(ii, ii)
enddo

do ii=1, N-1
    UD(ii) = h(ii, ii+1)
enddo

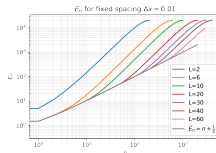
call DSTEVD( 'V', N, eigvalues, UD, eigvectors, LDZ, WORK, INFO )

call check( condition= (INFO.NE.0), &
            msg_type = 'ERROR', &
            msg = 'something went wrong in dstev')
```

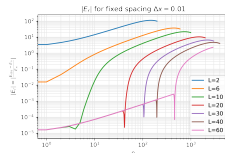
Figure: DSTEVD subroutine call



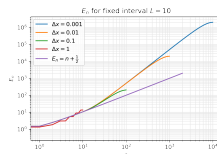
(a)



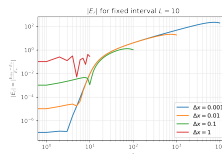
(b)



(c)



(d)



(e)

Figure: Eigenvalues and eigenvectors (2a). Eigenvalues (2b) and their relative error (2c) for fixed spacing Δx and varying L . Eigenvalues (2d) and their relative error (2e) for fixed spacing Δx and varying L .

■ **Correctness**

The curves reporting the absolute value of the relative error keeping respectively Δx and L fixed are shown in the slide before. We can see a strong dependence both on the spacing (finite differences method related error) and both on the interval range (higher energies eigenvalues see the potential as infinite-well shaped)

■ **Stability**

The program returns the same results (up to finite precision) for different runs

■ **Discretization**

There is a strong dependence on the chosen discretization spacing.

■ **Flexibility**

The program responds well to different input values of both Δx and L

■ **Efficiency**

The program exploits the ad-hoc STDEV subroutine for real, tridiagonal, symmetric matrices. Still, there is a strong dependence of the execution time on the number of points in the lattice.