

A Search for Sterile Neutrinos at the NO ν A Far Detector

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ABSTRACT

NO ν A is the current United States flagship long-baseline neutrino experiment designed to study the properties of neutrino oscillations. It consists of two functionally identical detectors each located 14.6 mrad off the central axis from the Fermilab NuMI neutrino beam. The Near Detector is located 1 km downstream from the beam source, and the Far Detector is located 810 km away in Ash River, Minnesota. This long baseline, combined with the ability of the NuMI facility to switch between nearly pure neutrino and anti-neutrino beams, allows NO ν A to make precision measurements of neutrino mixing angles, potentially determine the neutrino mass hierarchy, and begin searching for CP violating effects in the lepton sector. However, NO ν A can also probe more exotic scenarios, such as oscillations between the known active neutrinos and new sterile species.

This thesis showcases the first search for sterile neutrinos in a 3 + 1 model at NO ν A. The analysis presented searches for a deficit in the rate of neutral current events at the Far Detector using the Near Detector to constrain the predicted spectrum. The comparison between the observed and predicted spectra is translated into a measurement of the expanded PMNS mixing matrix elements, $|U_{\mu 4}|^2$ and $|U_{\tau 4}|^2$, assuming a value of $\Delta m_{41}^2 \sim O(1 \text{ eV}^2)$. This analysis was performed using data taken between February 2014 and May 2015 corresponding to 3.52×10^{20} protons on target. The best fit values for the matrix elements were $|U_{\mu 4}|^2 = 0.xy \pm a.bc$ and $|U_{\tau 4}|^2 = 0.vw \pm d.ef$, consistent with the no sterile neutrino hypothesis. At the end of this thesis there is a short discussion of future sensitivity improvements using a larger dataset.

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THIS IS THE DEDICATION.

Acknowledgments

These people were cool.

1

A Brief History of Neutrinos

1.1 INTRODUCTION

The neutrino was first postulated by Wolfgang Pauli as a possible explanation for the continuous spectrum of electrons emitted from nuclear β decay [1]. This decay was originally thought to be the emission of an electron from an atom, resulting in a different nucleus, via the process,

$$N \rightarrow N' + e \tag{1.1}$$

where N and N' are the parent and daughter nuclei, respectively. In a two body decay such as this, the momenta and energies of the outgoing particles are exactly constrained. Pauli's new particle explained the continuous spectrum of electron energy via a modified decay process:

$$N \rightarrow N' + e + \nu \tag{1.2}$$

where ν is the outgoing neutral particle. Pauli's original proposal called the new particle the neutron, but this name was later used to name the massive neutral nucleon discovered by Chadwick in 1932 [2]. Three years after Pauli's idea, Fermi proposed a model for nuclear β decay that included the new particle, which he coined the neutrino, or little neutral one [3].

1.2 FIRST DETECTION OF NEUTRINOS

Twenty years passed from Fermi's model proposal before neutrinos were discovered experimentally. Fred Reines and Clyde Cowan made the discovery by placing a detector near a nuclear reactor as a source of neutrinos and observing inverse β decay [4, 5]. The neutrinos observed were anti-electron neutrinos, thus the following was the observed process.

$$p + \bar{\nu}_e \rightarrow n + e^+ \quad (1.3)$$

Reines earned the Nobel Prize in Physics in 1995 for the detection of the neutrino.

In 1962, the muon neutrino was discovered at Brookhaven National Laboratory using the first neutrino beam [6] in a scheme still used in neutrino experiments today. The beam was generated by colliding protons with a target, producing pions that decayed into muons and muon neutrinos. The resultant beam then passed through thick steel, absorbing everything but the neutrinos. Leon Lederman, Melvin Schwartz, and Jack Steinberger won the Nobel Prize in Physics in 1988 for the discovery of the muon neutrino.

The last generation of neutrino, the tau neutrino, was discovered at Fermilab by the DONUT collaboration in 2000 [7].

1.3 EVIDENCE OF NEUTRINO OSCILLATIONS

Pontecorvo first postulated neutrino oscillations between neutrinos and anti-neutrinos, analogous to K^0/\bar{K}^0 oscillations, in 1957 [8]. Nothing came of the proposal immediately, but the idea was later revived and modified to solve the solar neutrino problem. The physics community initially viewed neu-

trino oscillations with skepticism and believed the experiments to be flawed, but over time oscillations have become an unmistakable and accepted phenomenon.

The solar neutrino problem was born from a large discrepancy between the theoretical and observed number of neutrinos produced by the sun. Neutrinos were used as a study for solar models because photons take a thousand years to escape the dense nuclear plasma to the surface of the sun, but neutrinos are unimpeded. The models, which have been confirmed today, describe a somewhat complicated chain of nuclear reactions, many of which produce neutrinos. Each individual process contributes a neutrinos in a different energy spectrum, but all of the neutrinos are created as electron neutrinos.

The experimental observations and theoretical predictions were both published in 1968. Ray Davis designed an experiment underground in the South Dakota Homestake mine consisting of a tank of an ultra pure chlorine cleaning solution capable of neutrino capture via the process



The argon atoms could be collected and counted for a direct measurement of the neutrino flux [9]. Meanwhile, John Bahcall precisely calculated the expected neutrino flux [10], and the observed rate was found to be about one third of the predicted rate. Pontecorvo revived his theory with the modification of allowing ν_e to ν_μ oscillations [11], but the idea was still not taken seriously and it was another 20 years before the solar neutrino problem was confirmed.

Beginning in 1989, multiple experiments with different methodologies confirmed the solar neutrino problem. Kamiokande, a water Cherenkov detector, measured a rate deficit in 1989 [12]. Two experiments measured solar neutrinos via the reaction



and measured a similar deficit, SAGE in 1991 [13] and GALLEX in 1992 [14]. With results from three different experimental methods all showing similar rate deficits, the solar neutrino problem could no longer be relegated to an experimental error.

Evidence soon emerged for oscillations with atmospheric neutrinos as well. These neutrinos are produced when cosmic rays collide with particles in the atmosphere and decay, predominantly via the following channels.

$$\begin{aligned}\pi^{+/-} &\rightarrow \mu^{+/-} + \nu_\mu/\bar{\nu}_\mu \\ \mu^{+/-} &\rightarrow e^{+/-} + \nu_e/\bar{\nu}_e + \bar{\nu}_\mu/\nu_\mu\end{aligned}\quad (1.6)$$

Thus, the expected ratio of muon family neutrinos to muon family neutrinos was expected to be 2. Kamiokande measured this ratio in 1992 and found the ratio to be much closer to 1 [15]. Furthermore, the ratio seemed to be dependent on zenith angle, with the measurement being nearly 2 for neutrinos coming from directly overhead, and dropping as the angle increased. Super-Kamiokande (or Super-K), the successor to Kamiokande, improved upon this measurement in 1998 [16], providing the most definitive evidence of neutrino oscillations to that point.

A resolution to the solar neutrino problem did not have to wait much longer with detector technologies capable of discerning different neutrino interaction types. SNO was designed as a heavy water (D_2O) Cherenkov detector experiment to be sensitive to both the flux of electron neutrinos and the flux of all neutrinos. In 2002, it released results for these measurements, finding what was then the expected deficit in electron neutrino flux, but a total flux consistent with the standard solar model, see figure 1.1 [17]. With this result, neutrino oscillations were confirmed, and subsequent experiments now measure oscillation parameters with precision.

1.4 POSSIBLE EVIDENCE OF STERILE NEUTRINOS

Several experiments have reported results that can be explained by oscillations between more than three neutrinos. The frequency of oscillations is controlled by the mass splitting between neutrino mass states, Δm^2 , and will be considered in more detail in chapter 2. Historically, solar and atmospheric neutrino oscillations were discovered first and were well explained by oscillations between the standard model's three active neutrinos, and with three neutrinos, there are only two independent mass splittings. Fur-

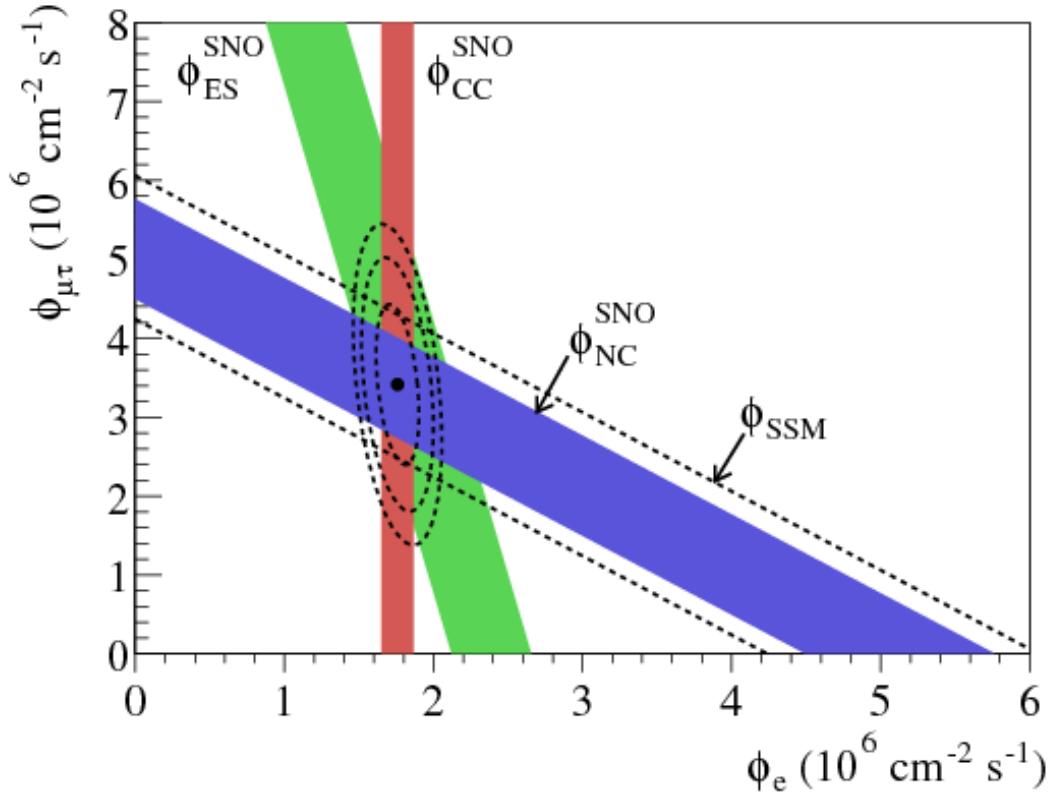


Figure 1.1: The measurement of different event rates at SNO [17]. The red band represents ν_e CC interactions with the deuterium neutron, an interaction only sensitive to electron neutrinos. The blue band represents neutral current scattering off of the deuterium nucleus, an interaction sensitive to the total neutrino flux. The green band represents elastic scattering of the neutrino off the deuterium electron, an interaction sensitive to all neutrino flavors, but not completely independent of neutrino flavor. The dashed straight lines represent the flux prediction by the standard solar model. The point represents the best fit for the flux of electron neutrinos and the flux for the combined muon and tau neutrinos.

thermore, it is known that there can only be three active neutrinos. However, there have been some experimental results that cannot be explained by oscillations based on the solar and atmospheric mass splitting scales, but they can be explained by oscillations based on a third independent mass splitting.

The number of active neutrinos is constrained by measurements of the width of the Z boson. LEP has measured the number of active neutrinos to be 2.984 ± 0.008 [18], so the discoveries of the ν_e , ν_μ , and ν_τ leave no room for new active neutrinos. Strictly speaking, there could be other active neutrinos if they had mass greater than half the mass of the Z boson so the Z could not decay to them, but while current evidence hinting at a sterile neutrino would suggest a mass splitting much larger than the solar and atmospheric scales, it is still much smaller than the mass of the Z.

The first evidence for an additional neutrino came from the Liquid Scintillator Neutrino Detector, or LSND, in 1995. This experiment searched for $\bar{\nu}_e$ appearance in a $\bar{\nu}_\mu$ beam. When it found an excess of events, it reported a measurement of a mass splitting between neutrino states Δm^2 of $O(1 \text{ eV}^2)$ [19], a measurement incompatible with the mass splittings measured in both the atmospheric and solar oscillation experiments, suggesting the addition of at least one more neutrino. However, based on the result from LEP, this new neutrino cannot couple to the Z boson, hence the suggestion for a sterile neutrino.

Searches for sterile neutrinos can be categorized in one of three channels: $(\bar{\nu}_\mu) \rightarrow (\bar{\nu}_e)$ appearance, $(\bar{\nu}_e)$ disappearance, and $(\bar{\nu}_\mu)$ disappearance. The LSND result was from the appearance channel. The Mini-BooNE experiment at Fermilab searched for both ν_e appearance in a ν_μ beam and $\bar{\nu}_e$ appearance in an $\bar{\nu}_\mu$ beam. While they first reported no event excess in 2007 [20], their more results show excesses in both modes [21] that could be consistent with some sterile neutrino models. Not all experiments have found evidence of sterile neutrinos, however. The most recent results from KARMEN in 2002 [22] and NOMAD in 2003 [23], both measuring the appearance channel, showed no evidence of oscillations at the same mass scale as LSND.

Experiments measuring the $(\bar{\nu}_e)$ disappearance channel have also reported some indication for a neutrino at the same mass scale as the appearance channel. The so-called Gallium anomaly comes from a measured deficit of events in GALLEX [24] and Sage [25] when placing a radioactive source inside the detectors for detector response measurements. On the other hand, reactor experiments such as Bugey [26] reported no oscillations and RENO [27] and Daya Bay [28] reported only oscillations consistent with standard three flavor oscillations. Furthermore, T2K, an accelerator experiment not unlike NO ν A, performed a short baseline analysis at its Near Detector and also found no evidence of sterile oscillations at 95% confidence.

To date, experiments probing the $(\bar{\nu}_\mu)$ disappearance channel have universally shown no evidence of sterile neutrinos. This includes an atmospheric neutrino measurement by Super-K [29] and a long baseline analysis by MINOS [30]. This list of experimental evidence for sterile neutrinos (or lack thereof) is not intended to be exhaustive, but it should be clear that there is not yet a scientific consensus on the existence of sterile oscillations despite the variety of past and current searches.

NO ν A is a long baseline neutrino experiment with a near detector, thus capable of performing both short baseline and long baseline sterile neutrino analyses; this dissertation focuses on the long baseline. It measures events generated by the NuMI beam, a mostly pure ν_μ beam, and thus measures the ν_μ disappearance channel. (In the future, the beam will switch to $\bar{\nu}_\mu$ allowing for measurement in the $\bar{\nu}_\mu$ disappearance channel.) The analysis presented here searches for a deficit in the number of neutral current events at the far detector, using the near detector data to constrain the predicted spectrum. Neutral currents are insensitive to the flavors of the standard 3 active neutrinos, so a rate deficit would point to the existence of oscillations between active and sterile neutrinos.

2

Theory of Neutrino Oscillations

The idea of neutrino oscillations was first proposed by Pontecorvo in 1957 [8], but his proposal described oscillations between neutrinos and anti-neutrinos. In 1962, after the discovery of the muon neutrino, Maki, Nakagawa, and Sakata proposed the theory that described oscillations between neutrino flavors due to differing neutrino flavor and mass eigenstates [31]. This chapter describes the modern formalism in detail and uses natural units where $\hbar = c = 1$, except where otherwise noted.

2.1 THE PMNS MATRIX

In the Standard Model, neutrinos only interact via the W and Z bosons as shown by the Feynman diagrams in figure 2.1. From these diagrams, it is clear that neutrinos always interact in a definite flavor eigenstate, $|\nu_\alpha\rangle$. Furthermore, when a neutrino is produced from a W boson, the flavor is always deter-

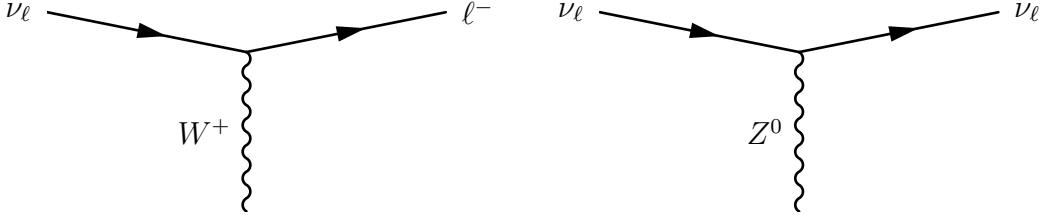


Figure 2.1: Standard Model Weak interactions involving a neutrino. Left: Charged Current interaction. Right: Neutral Current interaction.

mined by the associated charged lepton shown in equation 2.1.

$$\begin{pmatrix} \nu_e \\ e \end{pmatrix}, \quad \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}, \quad \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix} \quad (2.1)$$

On the other hand, neutrinos propagate through spacetime with a definite mass, $|\nu_i\rangle$ an eigenstate of the free Hamiltonian. The flavor states can be written as a superposition of the mass states via

$$|\nu_\alpha\rangle = \sum_{i=1}^n U_{\alpha i}^* |\nu_i\rangle, \quad (2.2)$$

where n is the number of neutrinos, and U is the unitary PMNS matrix, named after Pontecorvo, Maki, Nakagawa, and Sakata. The PMNS matrix is unitary, and would reduce to the identity matrix if neutrinos did not oscillate between flavor states. Yet since it does provide the mechanism for flavor transitions, it can be thought of as analogous to the quark sector CKM matrix.

2.2 VACUUM OSCILLATIONS

In this section, the basics of neutrino oscillations are developed by considering oscillations in a vacuum. The neutrinos are treated as plane waves, as in [32], with the assumption that the neutrino is actually

localized in space put in by hand. A careful, rigorous analysis treating the neutrinos as plane waves in [33] reproduces the same results.

Consider a neutrino in a state of definite flavor α at time $t = 0$, $|\nu(0)\rangle = |\nu_\alpha\rangle$. This state is in a superposition of mass eigenstates. The time evolution of this neutrino is simply the time evolution of the individual mass states. In a vacuum, this adds a phase factor to each mass state.

$$|\nu_\alpha(t)\rangle = \sum_i U_{\alpha i}^* e^{-i(E_i t - \mathbf{p}_i \cdot \mathbf{x})} |\nu_i\rangle \quad (2.3)$$

With the neutrino at position $\mathbf{x} = L$ at time t , the dot product evaluates to $\mathbf{p}_i \cdot \mathbf{x} = p_i L$. Eq 2.3 can then simplified by making use of the fact that neutrinos are ultra-relativistic, allowing for several assumptions. First, the time, t , is replaced by the distance, L . Next, the energy of each mass state is approximated to be the same energy, $E_i = E$. Last, the momentum is expanded as $p_i = \sqrt{E^2 - m_i^2} \approx E - m_i^2/2E$. With these assumptions, equation 2.3 simplifies as:

$$|\nu_\alpha(L)\rangle = \sum_i U_{\alpha i}^* e^{-im_i^2 L/2E} |\nu_i\rangle. \quad (2.4)$$

The mass eigenstate inside the sum is then re-expressed in terms of flavor eigenstates using the inverse of equation 2.2 and unitarity of U .

$$|\nu_\alpha(L)\rangle = \sum_{\alpha'} \sum_i U_{\alpha i}^* U_{\alpha' i} e^{-im_i^2 L/2E} |\nu'_\alpha\rangle. \quad (2.5)$$

Equation 2.5 can then be used to find the probability that the original neutrino in flavor state α has transitioned (or survived) as flavor state β . First, the matrix element $\langle \nu_\alpha | \nu_\beta(L) \rangle$ is computed.

$$\langle \nu_\alpha | \nu_\beta(L) \rangle = \sum_{\beta'} \sum_i U_{\beta' i}^* U_{\beta i} e^{-im_i^2 L/2E} \langle \nu_\alpha | \nu'_{\beta'} \rangle = \sum_i U_{\alpha i}^* U_{\beta i} e^{-im_i^2 L/2E} \quad (2.6)$$

The last equality in equation 2.6 follows from the orthogonality of individual flavor eigenstates. The

probability of the flavor transition is then the square of this matrix element.

$$P(\nu_\alpha \rightarrow \nu_\beta) = |\langle \nu_\alpha | \nu_\beta(L) \rangle|^2 = \sum_{i,j} U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^* e^{-i(m_i^2 - m_j^2)L/2E} \quad (2.7)$$

It is standard to rewrite the mass squared difference as $\Delta m_{ij}^2 \equiv m_i^2 - m_j^2$. Equation 2.7 is then manipulated using the properties of unitary matrices.

$$\begin{aligned} P(\nu_\alpha \rightarrow \nu_\beta) &= \sum_{i,j} U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^* + \sum_{i,j} U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^* (e^{-i\Delta m_{ij}^2 L/2E} - 1) \\ &= \delta_{\alpha\beta} + \sum_{i,j} U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^* (e^{-i\Delta m_{ij}^2 L/2E} - 1) \end{aligned} \quad (2.8)$$

The remaining summed term is further simplified making use of two facts. When $i = j$, the complex phase is 0 as $\Delta m_{ii}^2 = 0$, and thus these terms vanish. Second, the terms with $i < j$ are complex conjugates of those with $i > j$, and $z + z^* = 2\Re(z)$ for any complex number z .

$$P(\nu_\alpha \rightarrow \nu_\beta) = \delta_{\alpha\beta} + 2 \sum_{i>j} \Re \left[U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^* (e^{-i\Delta m_{ij}^2 L/2E} - 1) \right] \quad (2.9)$$

Both pieces of this term are split into their real and imaginary parts, and simplified using the trigonometric identity $\cos 2\theta - 1 = -2\sin^2 \theta$. Defining $\mathcal{U} \equiv U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^* (e^{-i\Delta m_{ij}^2 L/2E} - 1)$ and $\phi \equiv \Delta m_{ij}^2 L/2E$:

$$\Re(\mathcal{U}) = \Re \left[U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^* (e^{-i\Delta m_{ij}^2 L/2E} - 1) \right] \quad (2.10)$$

$$= \Re \{ [\Re(U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^*) + i\Im(U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^*)] [-2\sin^2(\phi/2) - i\sin \phi] \} \quad (2.11)$$

$$= -2\Re(U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^*) \sin^2(\phi/2) + \Im(U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^*) \sin \phi \quad (2.12)$$

Inserting the expression from equation 2.12 into equation 2.9, we find:

$$\begin{aligned} P(\nu_\alpha \rightarrow \nu_\beta) = \delta_{\alpha\beta} & - 4 \sum_{i>j} \Re(U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^*) \sin^2 \Delta_{ij} \\ & + 2 \sum_{i>j} \Im(U_{\alpha i} U_{\beta i}^* U_{\beta j} U_{\alpha j}^*) \sin 2\Delta_{ij}, \end{aligned} \quad (2.13)$$

where $\Delta_{ij} \equiv \Delta m_{ij}^2 L / 4E = 1.267 \Delta m_{ij}^2 L (\text{km}) / E (\text{GeV})$. It can now be seen that the distance the neutrino travels, its energy, and the different mass splittings all affect the frequency of oscillation. Ideally, neutrino oscillations would be studied by having neutrinos with a fixed energy profile (preferably monoenergetic) and varying the baseline. However, neutrino detectors are incredibly large, so in practice the baseline is fixed and the oscillation probability is studied as a function of neutrino energy.

For the case of survival probability, $\alpha = \beta$ and equation 2.13 simplifies further. The imaginary piece from equation 2.13 drops out, as

$$\Im(U_{\alpha i} U_{\alpha i}^* U_{\alpha j} U_{\alpha j}^*) = \Im(|U_{\alpha i}|^2 |U_{\alpha j}|^2) = 0. \quad (2.14)$$

The survival probability is then given by:

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - 4 \sum_{i>j} |U_{\alpha i}|^2 |U_{\alpha j}|^2 \sin^2 \Delta_{ij}. \quad (2.15)$$

Due to the combined influence of mass splitting, oscillation baseline, and neutrino energy on the oscillation probability, it is often the case that only one term contributes to the sums in equations 2.13 and 2.15. The two neutrino approximation can be instructive in this instance. For this model, the mixing matrix simplifies to the two dimensional rotation matrix:

$$U = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}. \quad (2.16)$$

As this matrix is entirely real, the imaginary piece of equation 2.13 drops out. Plugging the matrix elements into the remaining term directly and simplifying slightly, we find the following forms for the

survival and appearance probabilities.

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right) \quad (2.17)$$

$$P(\nu_\alpha \not\rightarrow \nu_\alpha) = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right) \quad (2.18)$$

From these equations it is clear that the mixing matrix parameters control the amplitude of neutrino oscillations. For small angles, most neutrinos will not change flavor, while larger angles can cause most of the neutrinos to change flavor. The case where $\theta = 45^\circ$ is called maximal mixing as at specific baseline lengths the probability of oscillation becomes 1.

2.3 STANDARD 3-FLAVOR OSCILLATIONS

The Standard Model includes three neutrinos, so the PMNS matrix is 3×3 in this picture. Explicitly expanding equation 2.2, U takes the following form:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}. \quad (2.19)$$

The PMNS matrix can be parametrized in terms of 3 real mixing angles, θ_{ij} and a complex phase, δ , called the CP phase. Following the convention from the Particle Data Group [34], the expanded matrix takes the form

$$\begin{aligned} U &= \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \\ &= \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (2.20) \end{aligned}$$

where $c_{ij} \equiv \cos \theta_{ij}$ and $s_{ij} \equiv \sin \theta_{ij}$.

With three neutrinos, the expanded forms of equations 2.13 and 2.15 can still balloon into unwieldy messes. Fortunately, based on current knowledge of the mass splittings, it is usually the case that only one mass splitting scale matters and other terms can be dropped. Figure 2.2 shows a schematic of the mass splittings. For historic reasons, Δm_{21}^2 is known as the solar mass splitting and the larger mass splitting is called the atmospheric mass splitting. The atmospheric mass splitting is about 30 times the solar mass splitting. The sign of the solar mass splitting is known, while that of the atmospheric mass splitting is not. A positive value of Δm_{32}^2 is called the normal hierarchy; a negative value is called the inverted hierarchy.

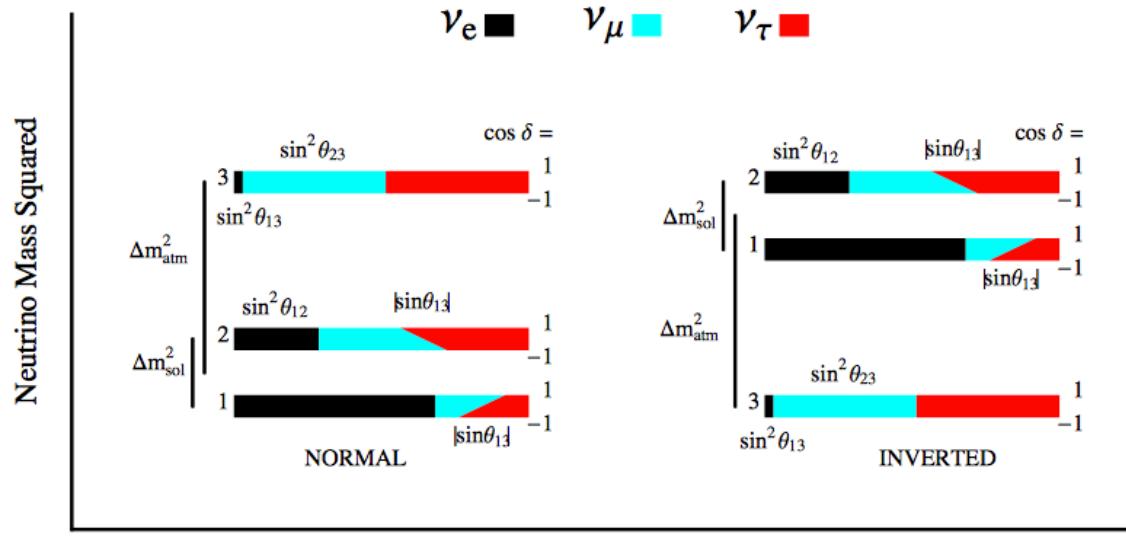


Figure 2.2: A schematic of the mass splittings between the three known neutrino mass states and how much they couple to each of the flavor states [35].

Two oscillation probabilities that are of interest to NO ν A are the muon neutrino survival probability and electron neutrino appearance from a muon neutrino beam. Since $|\Delta m_{21}^2|$ is so much smaller than $|\Delta m_{32}^2|$, the solar oscillation baseline is much longer, thus the oscillation probability is first dominated by terms containing Δm_{32}^2 . This is the case for NO ν A. Furthermore, the probability can be simplified by making the assumption that $|\Delta m_{32}^2| \approx |\Delta m_{31}^2|$. Under these conditions, the survival proba-

bility of muon neutrinos is calculated as follows:

$$P(\nu_\mu \rightarrow \nu_\mu) \approx 1 - 4|U_{\mu 3}|^2(|U_{\mu 1}|^2 + |U_{\mu 2}|^2)\sin^2 \Delta_{32} \quad (2.21)$$

$$\approx 1 - 4s_{23}^2(1 - s_{13}^2)(c_{23}^2 + s_{23}^2s_{13}^2)\sin^2 \Delta_{32} \quad (2.22)$$

$$\approx 1 - 4s_{23}^2c_{23}^2\sin^2 \Delta_{32} + 4s_{23}^2s_{13}^2(c_{23}^2 - s_{23}^2)\sin^2 \Delta_{32} \quad (2.23)$$

$$= 1 - \sin^2 2\theta_{23}\sin^2 \Delta_{32} + 4\sin^2 \theta_{23}\sin^2 \theta_{13}\cos^2 2\theta_{23}\sin^2 \Delta_{32} \quad (2.24)$$

Between equations 2.22 and 2.23, the term proportional to s_{13}^4 was dropped using the current knowledge that s_{13}^2 is small [34]. Note that if θ_{13} were 0, then equation 2.24 would reduce to equation 2.17, the two neutrino survival probability.

The full 3 flavor electron neutrino appearance from muon neutrino oscillation probability is often written in the form [36]:

$$P(^{(\bar{\nu}_\mu)} \rightarrow ^{(\bar{\nu}_e)}) = P_{atm} + 2\sqrt{P_{atm}}\sqrt{P_{sol}} (\cos \delta \cos \Delta_{32} (+) \sin \delta \sin \Delta_{32}) + P_{sol} \quad (2.25)$$

where

$$\sqrt{P_{atm}} \equiv \sin \theta_{23} \sin 2\theta_{13} \sin \Delta_{32} \quad (2.26)$$

$$\sqrt{P_{sol}} \equiv \cos \theta_{23} \sin 2\theta_{12} \sin \Delta_{21} \quad (2.27)$$

where the approximation $|\Delta m_{32}^2| \approx |\Delta m_{31}^2|$ has been made and higher order terms of s_{13}^2 been dropped. For an experiment at a short enough baseline such as NO ν A, the P_{sol} term is negligible as it depends on a higher order term of the solar mass splitting. The cross term is also not the dominant effect as it also depends upon the solar mass splitting, but it demonstrates interesting behavior. The $\cos \delta$ term is CP conserving, but the $\sin \delta$ term exhibits CP violation. This is why δ is called the CP violating phase angle.

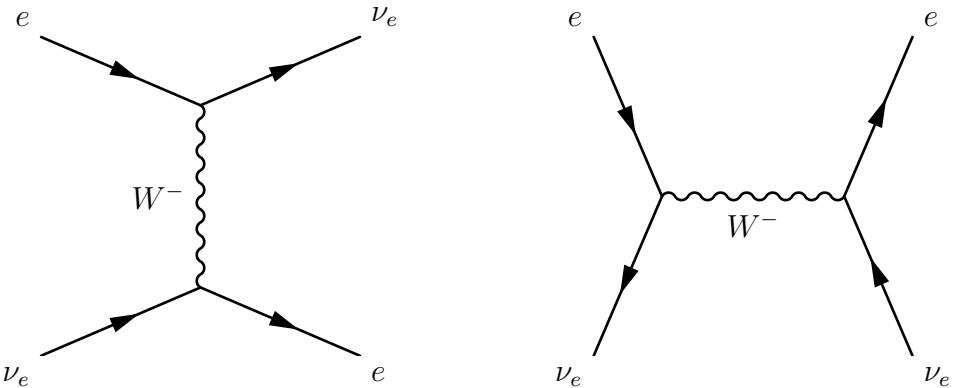


Figure 2.3: Coherent forward scattering interactions involved in the MSW effect. Left: Scattering of electron neutrinos on electrons. Right: Scattering of anti-electron neutrinos on electrons.

2.4 MATTER EFFECTS

So far, the oscillation formalism has been developed only considering neutrinos in a vacuum. However, most neutrino oscillation experiments involve neutrinos traveling through matter, be it the Sun or the Earth. This affects the oscillation probabilities in a process called the Mikheyev-Smirnov-Wolfenstein effect, or MSW effect. The phenomenon was first proposed by Wolfenstein in 1978 [37]; Mikheyev and Smirnov built upon that work in 1985 [38] as a possible solution for the solar neutrino problem.

The MSW effect is the coherent forward scattering of neutrinos off of the electrons in ordinary matter, a channel only available to electron flavor neutrinos and anti-neutrinos. Figure 2.3 illustrates the interactions. The electrons contribute an additional potential term, $V_e = \pm\sqrt{2}G_F N_e$, where G_F is Fermi's constant, N_e is the electron number density, the positive sign is for neutrinos, and the negative for anti-neutrinos. Neutrinos also forward scatter off the neutrons and protons in matter via neutral current interactions, but this only provides an overall phase as all neutrino flavors participate in these interactions equally. The matter induced potential adds an additional term to the Schrödinger equation, affecting the time evolution of the flavor states and thus changing the oscillation probabilities.

The following derivation will consider the MSW effect in the case of two neutrino flavors. The time

evolution of the flavor states is written as follows:

$$i \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \left[U \begin{pmatrix} \frac{m_1^2}{2E} & 0 \\ 0 & \frac{m_2^2}{2E} \end{pmatrix} U^\dagger + \begin{pmatrix} \pm V_e & 0 \\ 0 & 0 \end{pmatrix} \right] \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} \quad (2.28)$$

Inserting the 2 flavor PMNS matrix from equation 2.16, applying some trigonometric identities, and dropping common diagonal terms, equation 2.28 simplifies to

$$i \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \frac{1}{4E} \begin{pmatrix} -\Delta m_{21}^2 \cos 2\theta \pm 4EV_e & \Delta m_{21}^2 \sin 2\theta \\ \Delta m_{21}^2 \sin 2\theta & \Delta m_{21}^2 \cos 2\theta \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}. \quad (2.29)$$

The diagonal terms are dropped because they can be absorbed by the phase convention of the neutrino states. This Hamiltonian can be re-diagonalized with another unitary transformation, $H_M = U_M^\dagger H U_M$, with the following results:

$$H_M = \frac{1}{2} \begin{pmatrix} -\frac{\Delta m_M^2}{2E} & 0 \\ 0 & \frac{\Delta m_M^2}{2E} \end{pmatrix} \quad (2.30)$$

$$U_M = \begin{pmatrix} \cos \theta_M & \sin \theta_M \\ -\sin \theta_M & \cos \theta_M \end{pmatrix}, \quad (2.31)$$

where

$$\sin 2\theta_M \equiv \frac{\sin 2\theta}{A_M} \quad (2.32)$$

$$\Delta m_M^2 \equiv \Delta m_{21}^2 A_M \quad (2.33)$$

$$A_M \equiv \sqrt{\left(\cos 2\theta \mp \frac{2EV_e}{\Delta m_{21}^2} \right)^2 + \sin^2 2\theta}, \quad (2.34)$$

and now the negative sign in A_M is for neutrinos and the positive sign for anti-neutrinos. As the electron number density goes to 0, so too does V_e and the vacuum solution is recovered.

From the form of this solution, it can be seen that the Hamiltonian takes the same form as that in vac-

uum oscillations, but with modified effective masses. Likewise, U_M has the same form as the 2 neutrino PMNS matrix, so θ_M can be considered the effective mixing angle. In the absence of neutrino oscillations (when $\theta = 0$), matter effects cannot “create” them. However, even for small angles θ , the matter effect can create a resonant effect pushing the effective mixing angle, θ_M , maximally to 45° . This occurs when the term in parenthesis in the definition of A_M is 0 (equation 2.34).

$$N_e^{res} = \frac{\Delta m_{21}^2 \cos 2\theta}{2\sqrt{2}G_F E} \quad (2.35)$$

In the case of 3 neutrinos, the same procedure is followed to diagonalize the Hamilton and obtain effective values for the various oscillation parameters. The effects are considerably more complicated, but the general effect is the same—matter changes the effective neutrino mass and alters the oscillation probability curves differently for neutrinos and anti-neutrinos. Under the same conditions that were used to calculate $P(\langle \bar{\nu}_\mu \rangle \rightarrow \langle \bar{\nu}_e \rangle)$ in section 2.3, the results can be simplified to a few basic replacements [39].

$$P(\langle \bar{\nu}_\mu \rangle \rightarrow \langle \bar{\nu}_e \rangle) = P_{atm}^M + 2\sqrt{P_{atm}^M} \sqrt{P_{sol}^M} (\cos \delta \cos \Delta_{32} \text{ (+)} \sin \delta \sin \Delta_{32}) + P_{sol}^M \quad (2.36)$$

This is exactly the same form as equation 2.25. The interesting effects are seen with how P_{atm}^M and P_{sol}^M differ from their respective vacuum counterparts.

$$\sqrt{P_{atm}^M} \equiv \sin \theta_{23} \sin 2\theta_{13} \frac{\sin(\Delta_{31} - aL)}{\Delta_{31} - aL} \Delta_{31} \quad (2.37)$$

$$\sqrt{P_{sol}^M} \equiv \cos \theta_{23} \sin 2\theta_{12} \frac{\sin(aL)}{aL} \Delta_{21} \quad (2.38)$$

Here, $a \equiv \pm G_F N_e / \sqrt{2}$ where the positive sign is for neutrinos and the negative sign for anti-neutrinos. For the Earth, $|a| \approx 1/3500 \text{ km}$.

The combined effect that appears in equations 2.36, 2.37, and 2.38 due to the presence of matter plays an interesting role in the search for CP violation. The MSW effect by itself mimics CP violation as it alters oscillation probabilities for neutrinos and anti-neutrinos differently. Depending on the value of δ

that nature has chosen, the differences in oscillation probabilities due to the CP violation angle and the MSW effect can either compound or cancel out.

2.5 CURRENT MEASUREMENTS

Most of the free parameters in the PMNS matrix have been measured by various solar, atmospheric, accelerator, and reactor neutrino experiments. However, any given neutrino experiment does not have sensitivity to all of the oscillation parameters. Instead, experiments are sensitive to specific angles based on their baseline and the energies of the neutrinos they observe. Solar neutrino experiments, such as GALLEX, SAGE, Super-K, and SNO, measure neutrinos with energies on the order of several MeV after a very long baseline, and are most sensitive to θ_{12} and Δm_{21}^2 . Due to the strong MSW effect within the Sun, solar neutrino experiments also determined the ordering of mass states ν_1 and ν_2 ; ν_2 is defined as the heavier state. Atmospheric neutrino experiments, such as Super-K, SNO, and MINOS, measure neutrinos generated by cosmic ray collisions with the Earth's atmosphere, and are sensitive to θ_{23} and Δm_{32}^2 .

Reactor neutrino experiments, such as Chooz, Double Chooz, RENO, and Daya Bay, measure $\bar{\nu}_e$ generated by nearby nuclear reactors. Like solar neutrinos, reactor neutrinos have energies on the order of a few MeV. By measuring these neutrinos with a short baseline ($O(1\text{ km})$), the 2 neutrino approximation is valid, so the oscillation probability can be approximated as:

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 1 - \sin^2 2\theta_{13} \sin^2 \Delta_{31}. \quad (2.39)$$

Daya Bay made the first nonzero measurement of θ_{13} in 2012, reporting a value of $\sin^2 2\theta_{13} = 0.092 \pm 0.016 \text{ (stat)} \pm 0.005 \text{ (syst)}$ after taking just 55 days of data [40]. This result excluded a zero value for θ_{13} at 5.2σ and is shown in figure 2.4. Since that result, the limits have only continued to improve, and the leading measurement still comes from Daya Bay.

Accelerator neutrino experiments, such as MINOS, T₂K, and NO ν A, begin with a beam of nearly pure ${}^{\langle}\bar{\nu}_{\mu}{}^{\rangle}$ and search for both a disappearance of ${}^{\langle}\bar{\nu}_{\mu}{}^{\rangle}$ and appearance of other neutrino flavors. These experiments are sensitive to θ_{13} , θ_{23} , Δm_{32}^2 , and δ . The experiments that have the largest matter effect

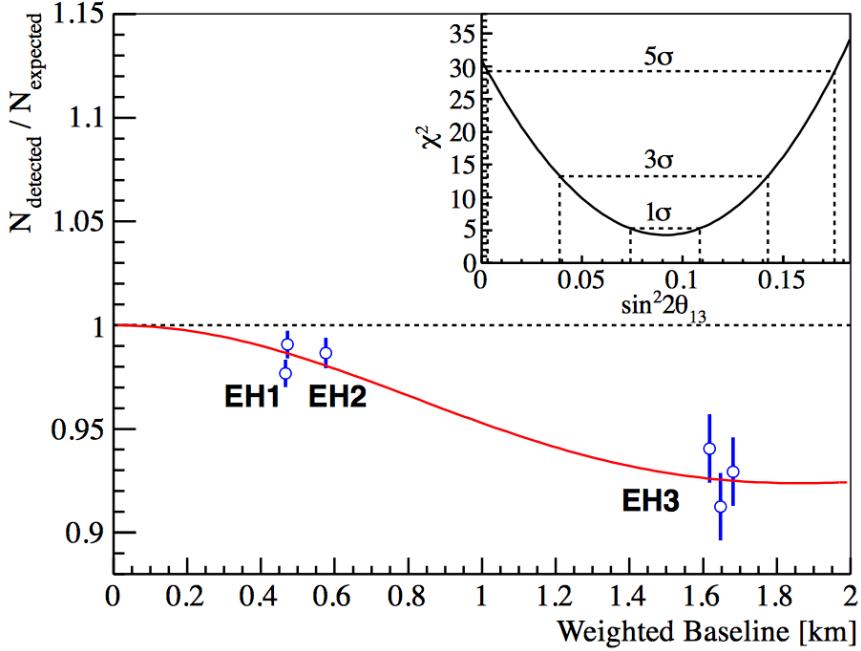


Figure 2.4: First measurement of θ_{13} from Daya Bay [40]. The points show the ratio of observed to expected events assuming $\theta_{13} = 0$. Each point is a measurement at a different detector. The inset in the upper corner shows the χ^2 value vs value of $\sin^2 2\theta_{13}$, and excludes $\theta_{13} = 0$ at greater than 5σ .

are the most sensitive to δ and the mass hierarchy.

Global fits to the combined data of these (and other) neutrino experiments have been performed and summarized in [34, 41]; the best fit values are shown in table 2.1. While most of the parameters have been measured with good precision, there are still a few lingering questions. From the table it is clear that a much better measurement on the CP violation angle is needed. The mass hierarchy still needs to be definitively measured as well. The other main question is whether θ_{23} is maximal, and if not, whether it is in the lower or upper octant.

Current and next generation reactor experiments have a great outlook to answer these outstanding questions. Making use of the MSW effect, the ν_e and $\bar{\nu}_e$ appearance results from experiments like T2K and NO ν A could simultaneously measure δ , the mass hierarchy, and θ_{23} octant. The prospects for NO ν A to make these measurements are shown in figure 2.5. The first analysis results from NO ν A [42, 43] were published after taking about 10% of the experiments design statistics and already show promise. The NO ν A measurement for δ , shown in figure 2.6, provide a hint toward the normal hierarchy and

Table 2.1: Current status of best fit oscillation parameters, from [34, 41]. The last column shows the allowed values within a 3σ range, with the exception of δ , which is shown at a 2σ range. This is because the current global best fit for δ still allows the full range from 0 to 2π at 3σ . NO ν A should vastly improve the limits on δ .

Parameter	Best-Fit ($\pm 1\sigma$)		3σ Range
Δm_{21}^2 [10^{-5} eV 2]	$7.54^{+0.26}_{-0.22}$		6.99 – 8.18
$ \Delta m^2 $ [10^{-3} eV 2]	NH	2.43 ± 0.06	2.23 – 2.61
	IH	2.38 ± 0.06	2.19 – 2.56
$\sin^2 \theta_{12}$		0.308 ± 0.017	0.259 – 0.359
$\sin^2 \theta_{23}$	NH	$0.437^{+0.033}_{-0.023}$	0.374 – 0.628
	IH	$0.455^{+0.039}_{-0.031}$	0.380 – 0.641
$\sin^2 \theta_{13}$	NH	$0.0234^{+0.0020}_{-0.0019}$	0.0176 – 0.0295
	IH	$0.0240^{+0.0019}_{-0.0022}$	0.0178 – 0.0298
δ/π (2σ range)	NH	$1.39^{+0.38}_{-0.27}$	(0.00 – 0.16) \oplus (0.86 – 2.00)
	IH	$1.31^{+0.29}_{-0.33}$	(0.00 – 0.02) \oplus (0.70 – 2.00)

eliminate portions of δ space at 90% confidence.

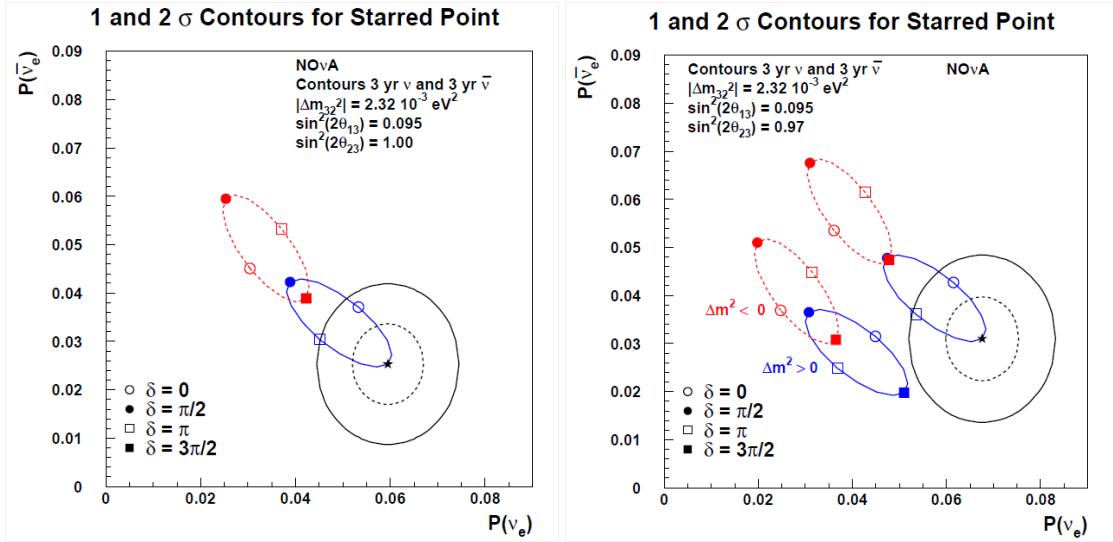


Figure 2.5: Probability of ν_e appearance versus $\bar{\nu}_e$ appearance at NO ν A. The blue ellipses are for the normal hierarchy; the red ellipses are for the inverted hierarchy. The starred points show a possible measurement NO ν A could make. The matter effect can either constructively or destructively combine with the CP violation effect. A larger matter effect, further separates the two mass hierarchy ellipses. This corresponds to neutrinos passing through more matter. On the left, θ_{23} is assumed to be 45° for maximal mixing, purely showcasing the interference between the matter and CP violation effects. On the right, θ_{23} is non-maximal, showing how the dependence on θ_{23} affects both ellipses in the same way.

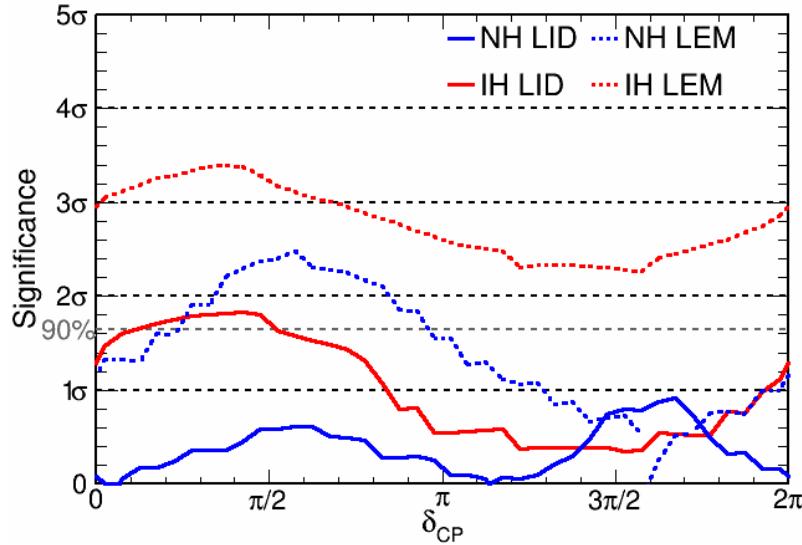


Figure 2.6: First measurement of δ from NO ν A [42]. The plot shows the significance of the difference between the observed and predicted number of events as a function of delta. NO ν A used a primary and secondary selection technique, the primary (secondary) technique is shown as the solid (dotted) line. The secondary selection disfavors the inverted hierarchy for all values of δ .

2.6 STERILE NEUTRINOS

The analysis presented in this dissertation considers sterile neutrinos in a $3+1$ model. This adds a fourth neutrino flavor state, ν_s , and mass state, ν_4 , and the PMNS matrix becomes a 4×4 matrix. Given that standard three flavor oscillations explain solar, atmospheric, and long-baseline oscillations so well, this fourth state must be mostly sterile, i.e., $|U_{\alpha 4}| \ll 1$ for $\alpha \in \{e, \mu, \tau\}$.

The mass splitting suggested by the experimental evidence discussed in section 1.4 is $\Delta m^2 \sim O(1 \text{ eV}^2)$. If the fourth neutrino mass state were the lightest, then all of the other mass states would have absolute masses $m \gtrsim 1 \text{ eV}$. However, the sum of the active neutrino masses is well constrained by cosmology; the most recent Planck results reported $\sum m_\nu < 0.23 \text{ eV}$ [44]. Consequently, it is assumed that the additional neutrino mass state is heavier than the other three. Furthermore, since $\Delta m_{4i}^2 \gg |\Delta m_{ji}^2|$ for $i, j < 4$, it is assumed that $\Delta_{41} \approx \Delta_{42} \approx \Delta_{43}$.

As mentioned above, adding a fourth neutrino state expands the PMNS matrix into a 4×4 matrix. It can be written as the product of individual (complex) rotation matrices, as in the second line of equation

2.20. This adds an additional 3 real rotation angles and 2 complex phases. It is a matter of convention on how exactly to add these mixing angles into the PMNS matrix, but as physical results cannot be affected by a particular parametrization, it is best to choose the most convenient one. Following the parametrization in [45], U takes the form:

$$U = R_{34}C_{24}R_{23}R_{14}C_{13}C_{12}, \quad (2.40)$$

where R_{ij} is a real matrix parametrized by a single mixing angle, θ_{ij} , and C_{ij} is a complex matrix parametrized by a real mixing angle θ_{ij} and complex phase δ_{ij} . This parametrization is convenient because as was the case in section 2.3, the solar mass scale Δ_{21} is negligible, thus mixing angles parametrized by C_{12} drop out of the oscillation probabilities. Based on this, equation 2.40 can be expanded in terms of the mixing angles:

$$U = \begin{bmatrix} U_{e1} & U_{e2} & c_{14}s_{13}e^{-i\delta_{13}} & s_{14} \\ U_{\mu 1} & U_{\mu 2} & -s_{14}s_{13}e^{-i\delta_{13}}s_{24}e^{-i\delta_{24}} + c_{13}s_{23}c_{24} & c_{14}s_{24}e^{-i\delta_{24}} \\ U_{\tau 1} & U_{\tau 2} & -s_{14}c_{24}s_{34}s_{13}e^{-i\delta_{13}} - c_{13}s_{23}s_{34}s_{24}e^{i\delta_{24}} + c_{13}c_{23}c_{34} & c_{14}c_{24}s_{34} \\ U_{s1} & U_{s2} & -s_{14}c_{24}c_{34}s_{13}e^{-i\delta_{13}} - c_{13}s_{23}c_{34}s_{24}e^{i\delta_{24}} - c_{13}c_{23}s_{34} & c_{14}c_{24}c_{34} \end{bmatrix} \quad (2.41)$$

Above, only the columns that appear in the oscillation probabilities have been expanded. Before using these expanded forms, however, it is useful to write the oscillation probability formulae in terms of the matrix elements first.

For NO ν A, which starts with a nearly pure ν_μ beam, the two most important oscillation probabilities are the ν_μ survival probability, $P(\nu_\mu \rightarrow \nu_\mu)$ and the probability that a ν_μ will not oscillate to a sterile neutrino, $1 - P(\nu_\mu \rightarrow \nu_s)$. Starting with equations 2.15 and 2.13, these probabilities are calculated as

follows, using the unitarity of U to remove elements found in its the first two columns.

$$\begin{aligned}
P(\nu_\mu \rightarrow \nu_\mu) &\approx 1 - 4|U_{\mu 4}|^2(|U_{\mu 1}|^2 + |U_{\mu 2}|^2 + |U_{\mu 3}|^2) \sin^2 \Delta_{41} \\
&\quad - 4|U_{\mu 3}|^2(|U_{\mu 1}|^2 + |U_{\mu 2}|^2) \sin^2 \Delta_{31} \\
&\approx 1 - 4|U_{\mu 4}|^2(1 - |U_{\mu 4}|^2) \sin^2 \Delta_{41} \\
&\quad - 4|U_{\mu 3}|^2(1 - |U_{\mu 3}|^2 - |U_{\mu 4}|^2) \sin^2 \Delta_{31}
\end{aligned} \tag{2.42}$$

$$\begin{aligned}
1 - P(\nu_\mu \rightarrow \nu_s) &= 1 + 4 \sum_{i>j} \Re(U_{\mu i} U_{si}^* U_{sj} U_{\mu j}^*) \sin^2 \Delta_{ij} \\
&\quad - 2 \sum_{i>j} \Im(U_{\mu i} U_{si}^* U_{sj} U_{\mu j}^*) \sin 2\Delta_{ij} \\
&\approx 1 + 4\Re(C_{41,42,43}) \sin^2 \Delta_{41} + 4\Re(C_{31,32}) \sin^2 \Delta_{31} \\
&\quad - 2\Im(C_{41,42,43}) \sin 2\Delta_{41} - 2\Im(C_{31,32}) \sin 2\Delta_{31}
\end{aligned} \tag{2.43}$$

In equation 2.43, $C_{41,42,43}$ and $C_{31,32}$ are sums of products of PMNS matrix elements.

$$\begin{aligned}
C_{41,42,43} &= U_{\mu 4} U_{s4}^* (U_{s1} U_{\mu 1}^* + U_{s2} U_{\mu 2}^* + U_{s3} U_{\mu 3}^*) \\
&= U_{\mu 4} U_{s4}^* (-U_{s4} U_{\mu 4}^*) \\
&= -|U_{\mu 4}|^2 |U_{s4}|^2
\end{aligned} \tag{2.44}$$

Note that this term is real, thus the term with $\Im(C_{41,42,43})$ in equation 2.43 drops out.

$$\begin{aligned}
C_{31,32} &= U_{\mu 3} U_{s3}^* (U_{s1} U_{\mu 1}^* + U_{s2} U_{\mu 2}^*) \\
&= -U_{\mu 3} U_{s3}^* (U_{s4} U_{\mu 4}^* + U_{s3} U_{\mu 3}^*) \\
&= -|U_{\mu 3}|^2 |U_{s3}|^2 - U_{\mu 3} U_{\mu 4}^* U_{s3}^* U_{s4} \\
&= -|U_{\mu 3}|^2 |U_{s3}|^2 - \Re(U_{\mu 3} U_{\mu 4}^* U_{s3}^* U_{s4}) - i\Im(U_{\mu 3} U_{\mu 4}^* U_{s3}^* U_{s4}) \\
&= -|U_{\mu 3}|^2 |U_{s3}|^2 - \Re(U_{\mu 3}^* U_{\mu 4} U_{s3} U_{s4}^*) + i\Im(U_{\mu 3}^* U_{\mu 4} U_{s3} U_{s4}^*)
\end{aligned} \tag{2.45}$$

The last equality in 2.45 was done so that all terms in the oscillation probability are subtracted from 1.

Using these results, equation 2.43 becomes

$$\begin{aligned} 1 - P(\nu_\mu \rightarrow \nu_s) &\approx 1 - 4|U_{\mu 4}|^2|U_{s 4}|^2 \sin^2 \Delta_{41} \\ &- 4(|U_{\mu 3}|^2|U_{s 3}|^2 + \Re(U_{\mu 3}^* U_{\mu 4} U_{s 3} U_{s 4}^*)) \sin^2 \Delta_{31} \\ &- 2\Im(U_{\mu 3}^* U_{\mu 4} U_{s 3} U_{s 4}^*) \sin 2\Delta_{31}. \end{aligned} \quad (2.46)$$

Equations 2.42 and 2.46 are the probabilities that a ν_μ will survive as a ν_μ and survive as one of the active neutrino flavors, respectively. These can now be written in terms of the mixing angles. However, several approximations make the expressions much more bearable. First, based on the best value reported in section 2.5, the value of s_{13} is taken as 0. This has the fortunate side effect of removing any dependence on δ_{13} as well. Second, the value of s_{14} is also taken as 0, based on reactor and other constraints, e.g. [28, 46]. Note that in the relevant matrix elements, s_{13} and s_{14} always appear together as a product, further justifying the approximation. The matrix element products that appear in the oscillation probabilities are calculated below.

$$|U_{\mu 4}|^2 = s_{24}^2 \quad (2.47)$$

$$|U_{s 4}|^2 = c_{24}^2 c_{34}^2 \quad (2.48)$$

$$|U_{\mu 3}|^2 = s_{23}^2 c_{24}^2 \quad (2.49)$$

$$|U_{s 3}|^2 = c_{23}^2 s_{34}^2 + s_{23}^2 s_{24}^2 c_{34}^2 + 2c_{23}s_{23}s_{24}c_{34}s_{34} \cos \delta_{24} \quad (2.50)$$

$$U_{\mu 3}^* U_{\mu 4} U_{s 3} U_{s 4}^* = -s_{23} c_{24}^2 s_{24} c_{34} (s_{23} s_{24} c_{34} + c_{23} s_{34} e^{-i\delta_{24}}) \quad (2.51)$$

$$\Re(U_{\mu 3}^* U_{\mu 4} U_{s 3} U_{s 4}^*) = -s_{23} c_{24}^2 s_{24} c_{34} (s_{23} s_{24} c_{34} + c_{23} s_{34} \cos \delta_{24}) \quad (2.52)$$

$$\Im(U_{\mu 3}^* U_{\mu 4} U_{s 3} U_{s 4}^*) = c_{23} s_{23} c_{24}^2 s_{24} c_{34} s_{34} \sin \delta_{24} \quad (2.53)$$

These results can be substituted into equations 2.42 and 2.46.

$$\begin{aligned}
P(\nu_\mu \rightarrow \nu_\mu) &\approx 1 - 4|U_{\mu 4}|^2(1 - |U_{\mu 4}|^2)\sin^2 \Delta_{41} \\
&\quad - 4|U_{\mu 3}|^2(1 - |U_{\mu 3}|^2 - |U_{\mu 4}|^2)\sin^2 \Delta_{31} \\
&\approx 1 - 4s_{24}^2(1 - s_{24}^2)\sin^2 \Delta_{41} \\
&\quad - 4s_{23}^2 c_{24}^2(1 - s_{23}^2 c_{24}^2 - s_{24}^2)\sin^2 \Delta_{31} \\
&\approx 1 - 4c_{24}^2 s_{24}^2 \sin^2 \Delta_{41} - 4c_{24}^4 s_{23}^2(1 - s_{23}^2)\sin^2 \Delta_{31} \\
&\approx 1 - \sin^2 2\theta_{24} \sin^2 \Delta_{41} - c_{24}^4 \sin^2 2\theta_{23} \sin^2 \Delta_{31} \tag{2.54}
\end{aligned}$$

$$\begin{aligned}
1 - P(\nu_\mu \rightarrow \nu_s) &\approx 1 - 4|U_{\mu 4}|^2|U_{s 4}|^2 \sin^2 \Delta_{41} \\
&\quad - 4(|U_{\mu 3}|^2|U_{s 3}|^2 + \Re(U_{\mu 3}^* U_{\mu 4} U_{s 3} U_{s 4}^*)) \sin^2 \Delta_{31} \\
&\quad - 2\Im(U_{\mu 3}^* U_{\mu 4} U_{s 3} U_{s 4}^*) \sin 2\Delta_{31} \\
&\approx 1 - 4c_{24}^2 s_{24}^2 c_{34}^2 \sin^2 \Delta_{41} \\
&\quad - 4[s_{23}^2 c_{24}^2(c_{23}^2 s_{34}^2 + s_{23}^2 s_{24}^2 c_{34}^2 + 2c_{23} s_{23} s_{24} c_{34} s_{34} \cos \delta_{24}) \\
&\quad \quad - s_{23} c_{24}^2 s_{24} c_{34}(s_{23} s_{24} c_{34} + c_{23} s_{34} \cos \delta_{24})] \sin^2 \Delta_{31} \\
&\quad - 2c_{23} s_{23} c_{24}^2 s_{24} c_{34} s_{34} \sin \delta_{24} \sin 2\Delta_{31} \\
&\approx 1 - c_{34}^2 \sin^2 2\theta_{24} \sin^2 \Delta_{41} \\
&\quad - 4[c_{23}^2 s_{23}^2 c_{24}^2 s_{34}^2 - s_{23}^2 c_{24}^2 s_{24}^2 c_{34}^2(1 - s_{23}^2) \\
&\quad \quad + c_{23} s_{23} c_{24}^2 s_{24} c_{34} s_{34}(2s_{23}^2 - 1) \cos \delta_{24}] \sin^2 \Delta_{31} \\
&\quad - \frac{1}{4} c_{24} \sin 2\theta_{23} \sin 2\theta_{24} \sin 2\theta_{34} \sin \delta_{24} \sin 2\Delta_{31} \\
&\approx 1 - c_{34}^2 \sin^2 2\theta_{24} \sin^2 \Delta_{41} - c_{24}^2(s_{34}^2 - s_{24}^2 c_{34}^2) \sin^2 2\theta_{23} \sin^2 \Delta_{31} \\
&\quad - \frac{1}{4} c_{24} \sin 2\theta_{23} \sin 2\theta_{24} \sin 2\theta_{34} \sin \delta_{24} \sin 2\Delta_{31} \tag{2.55}
\end{aligned}$$

In the last equality, the term proportional to $\cos \delta_{24}$ was dropped noting that $2s_{23}^2 - 1 \approx 0$. Even if

it is not exact, there is already an additional factor of $s_{24}s_{34}$. Both s_{24} and s_{34} must be small in order to satisfy $|U_{\mu 4}|, |U_{\tau 4}| \ll 1$, thus the term proportional to $s_{24}s_{34}(2s_{23}^2 - 1)$ can be considered higher order than the remaining terms.

The effect of each of the oscillation parameters can be seen in the figures below. Each page shows the ν_μ survival probability followed by the active neutrino survival probability and uses the exact oscillation probabilities. Figures 2.7-2.11 show how the mixing parameters that appear in equations 2.54 and 2.55 affect the oscillation probabilities. Figures 2.12-2.14 are similar but show the much smaller effects of the parameters that dropped out of these equations due to approximations. (Solar mixing parameters are not shown.) The individual plots show different probabilities curves while varying a single oscillation parameter over a set of representative values; the mixing parameters that are not varied in a particular figure are held fixed to the values in table 2.2. In the figures, the bottom axis shows the neutrino L/E , thus the plots show oscillation probabilities at both the near and far detectors. The top axis shows the neutrino energy at each NO ν A detector, the dashed line separates the Near Detector (ND) on the left from the Far Detector (FD) on the right. Finally, the grey bands show the neutrino energies each detector is sensitive to; a darker color means a higher intensity of neutrinos at that energy.

Table 2.2: The four flavor oscillation parameters assumed for studying the effects of different parameters on the oscillation probabilities.

Oscillation Parameter	Value
ρ	2.84 g/cm^3
Δm_{21}^2	$7.53 \times 10^{-5} \text{ eV}^2$
Δm_{32}^2	$2.44 \times 10^{-3} \text{ eV}^2$
Δm_{41}^2	0.5 eV^2
$\sin^2 2\theta_{12}$	0.846
$\sin^2 2\theta_{13}$	0.085
θ_{23}	$\pi/4$
θ_{14}	0
θ_{24}	10°
θ_{34}	35°
δ_{13}	0
δ_{24}	0

For the particular analysis presented in the rest of this dissertation, it is assumed that there are no Near

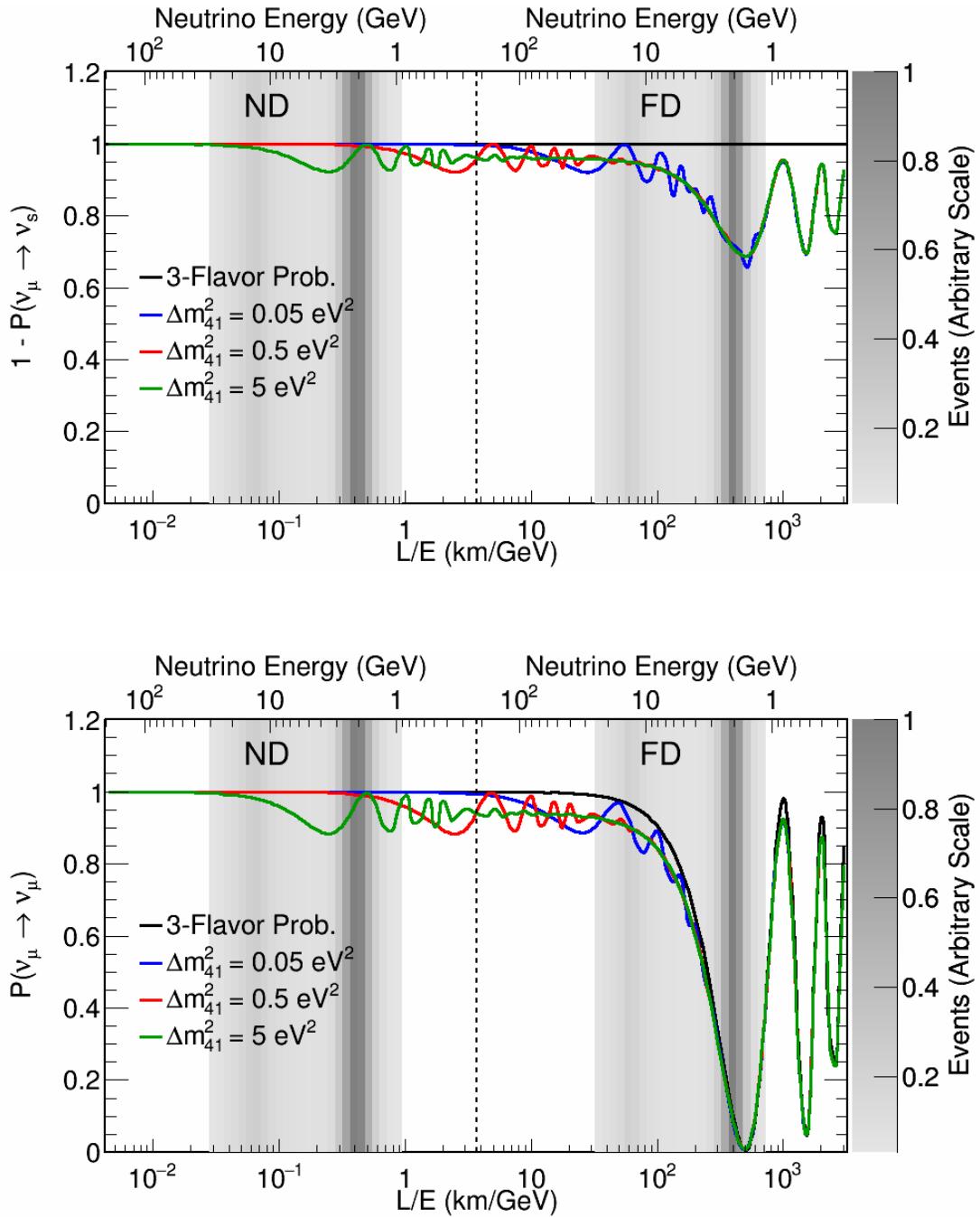


Figure 2.7: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of Δm_{41}^2 . Above $\Delta m_{41}^2 = 0.5 \text{ eV}^2$, oscillations start to affect the ND. The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

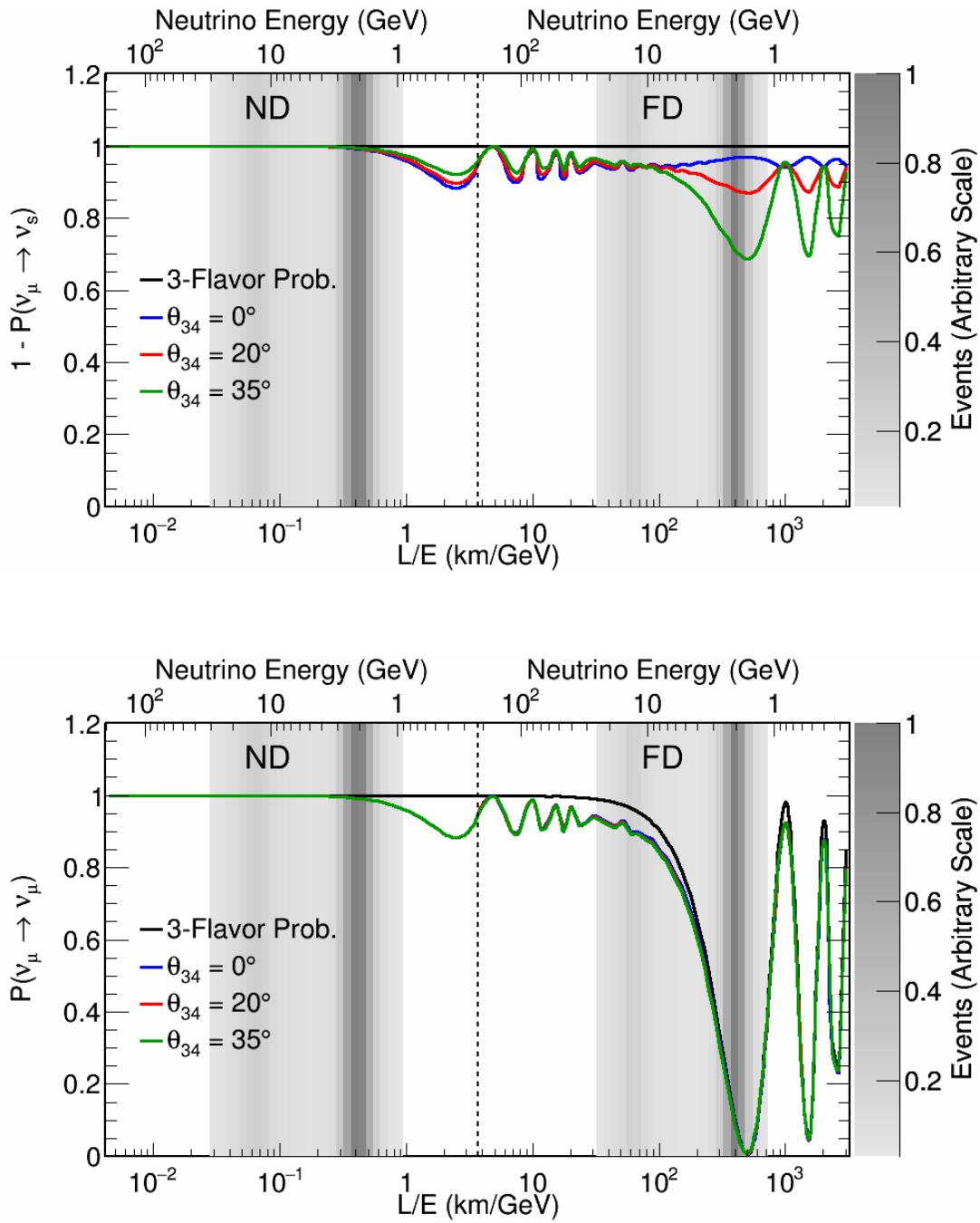


Figure 2.8: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of θ_{34} . The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

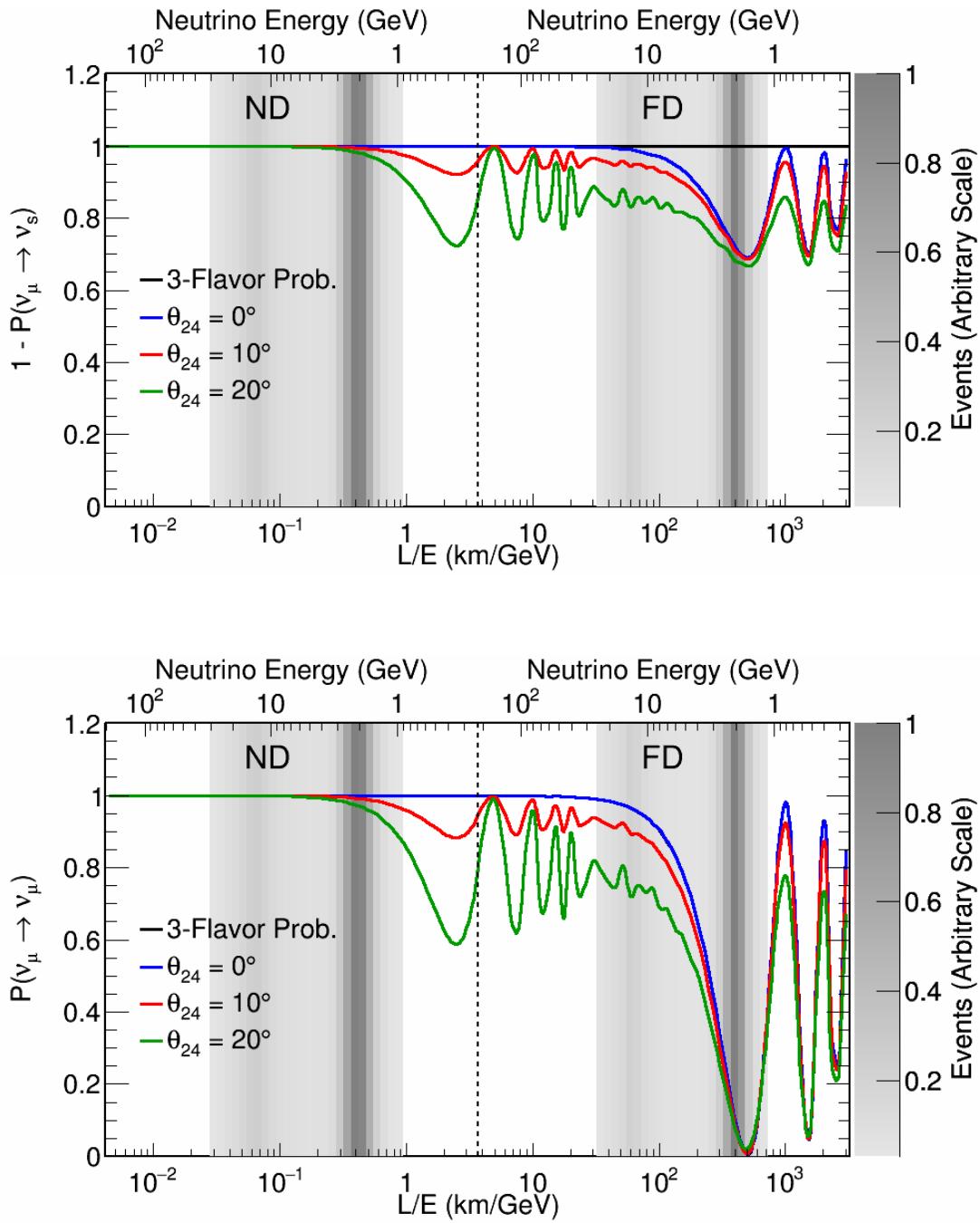


Figure 2.9: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of θ_{24} . The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

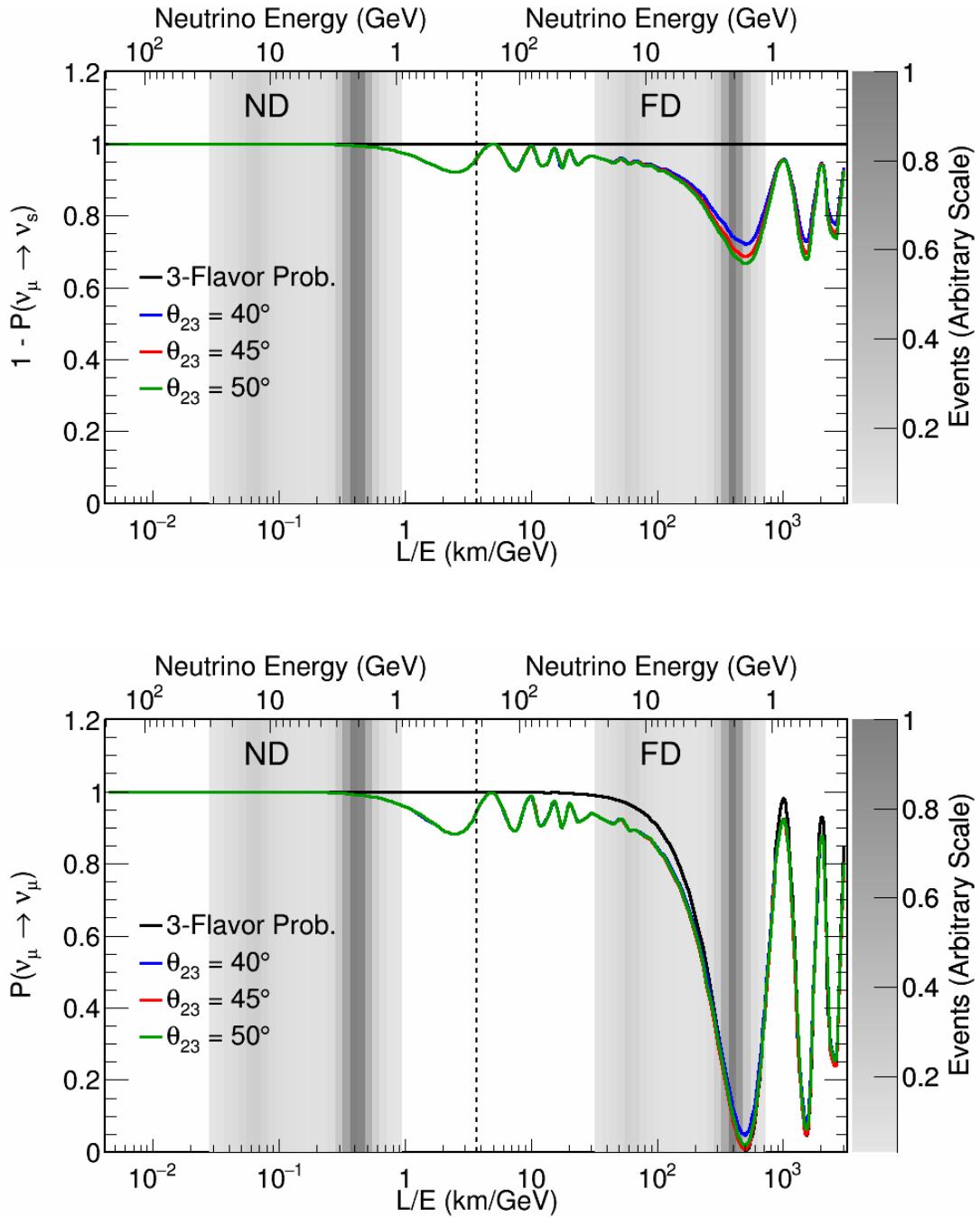


Figure 2.10: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of θ_{23} . Shifts in θ_{23} have a modest effect at the NO ν A peak energy. The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

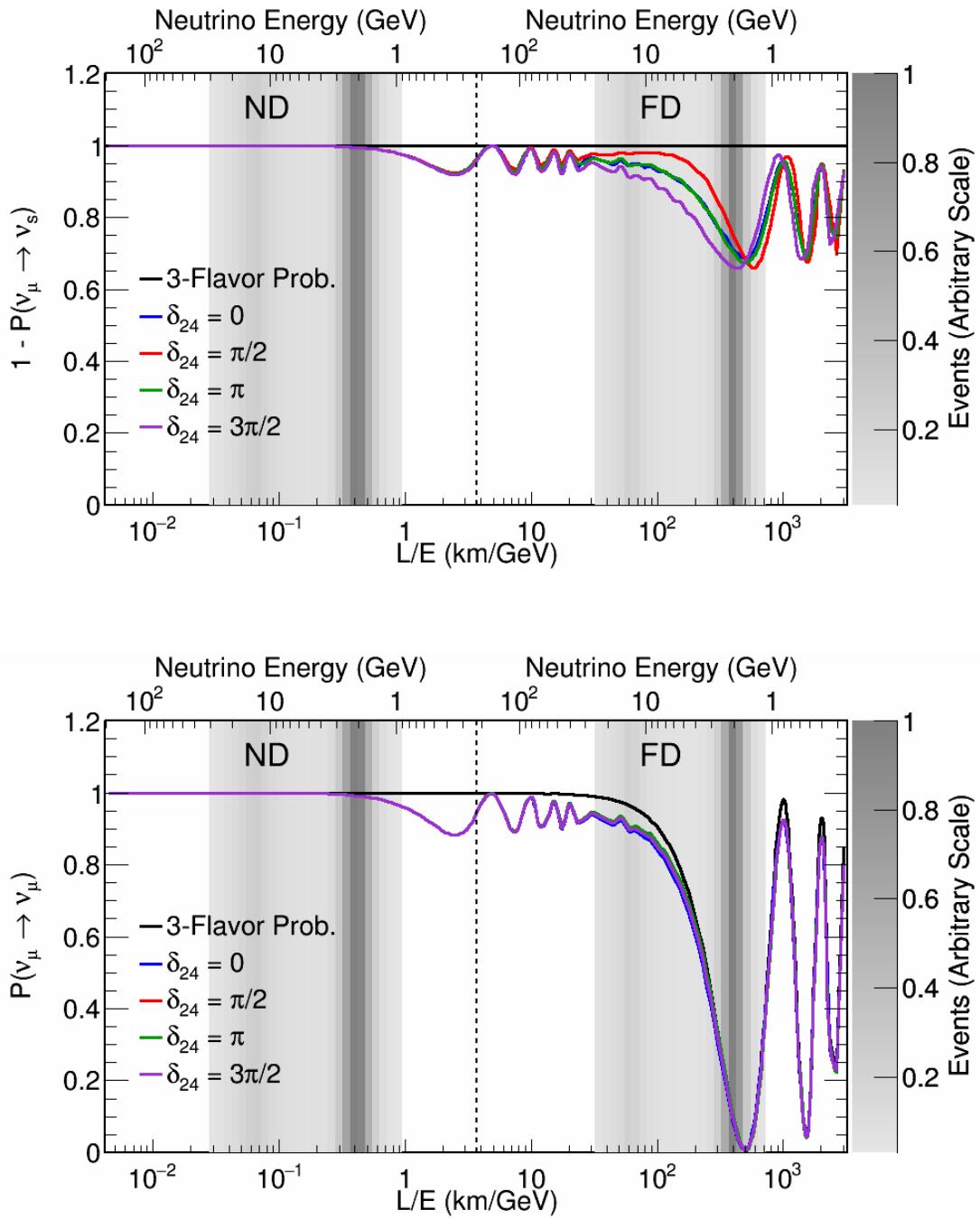


Figure 2.11: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of δ_{24} . The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

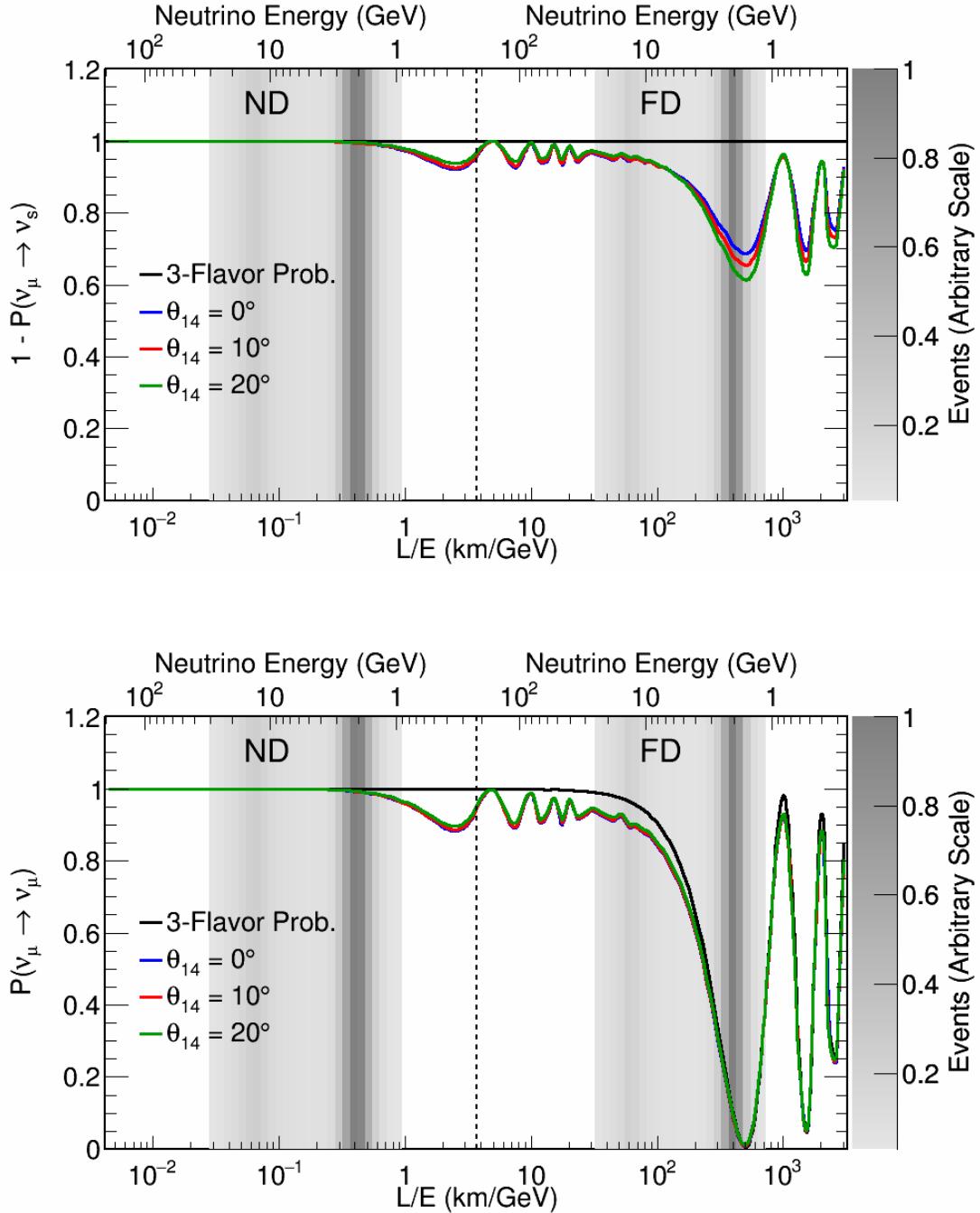


Figure 2.12: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of θ_{14} . The values shown are well outside of experimental bounds, but even these large shifts have only a small effect on the oscillation probabilities that is largely degenerate with θ_{23} . The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

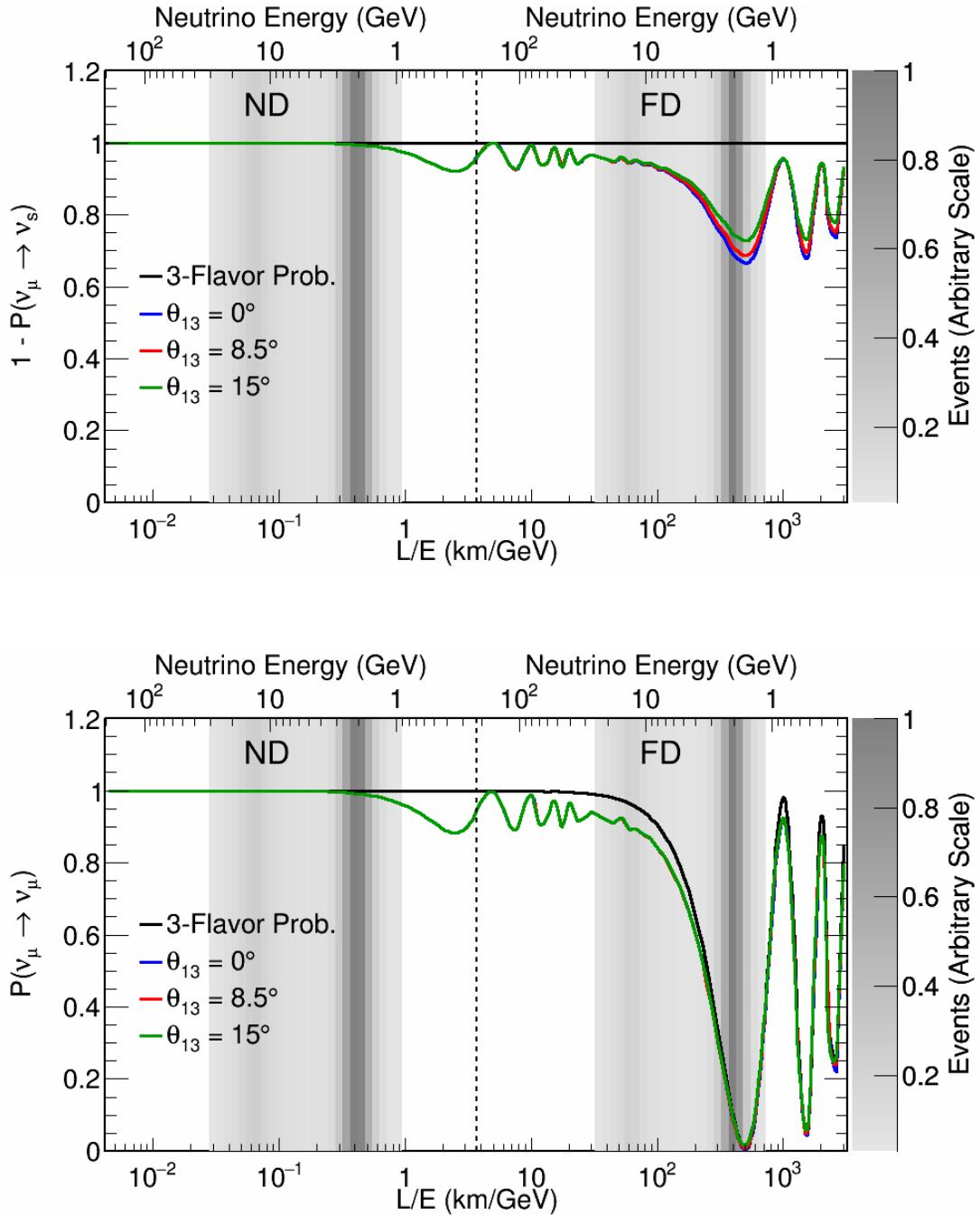


Figure 2.13: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of θ_{13} . The values shown are well outside of experimental bounds, but even these large shifts have only a small effect on the oscillation probabilities that is largely degenerate with θ_{23} . The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

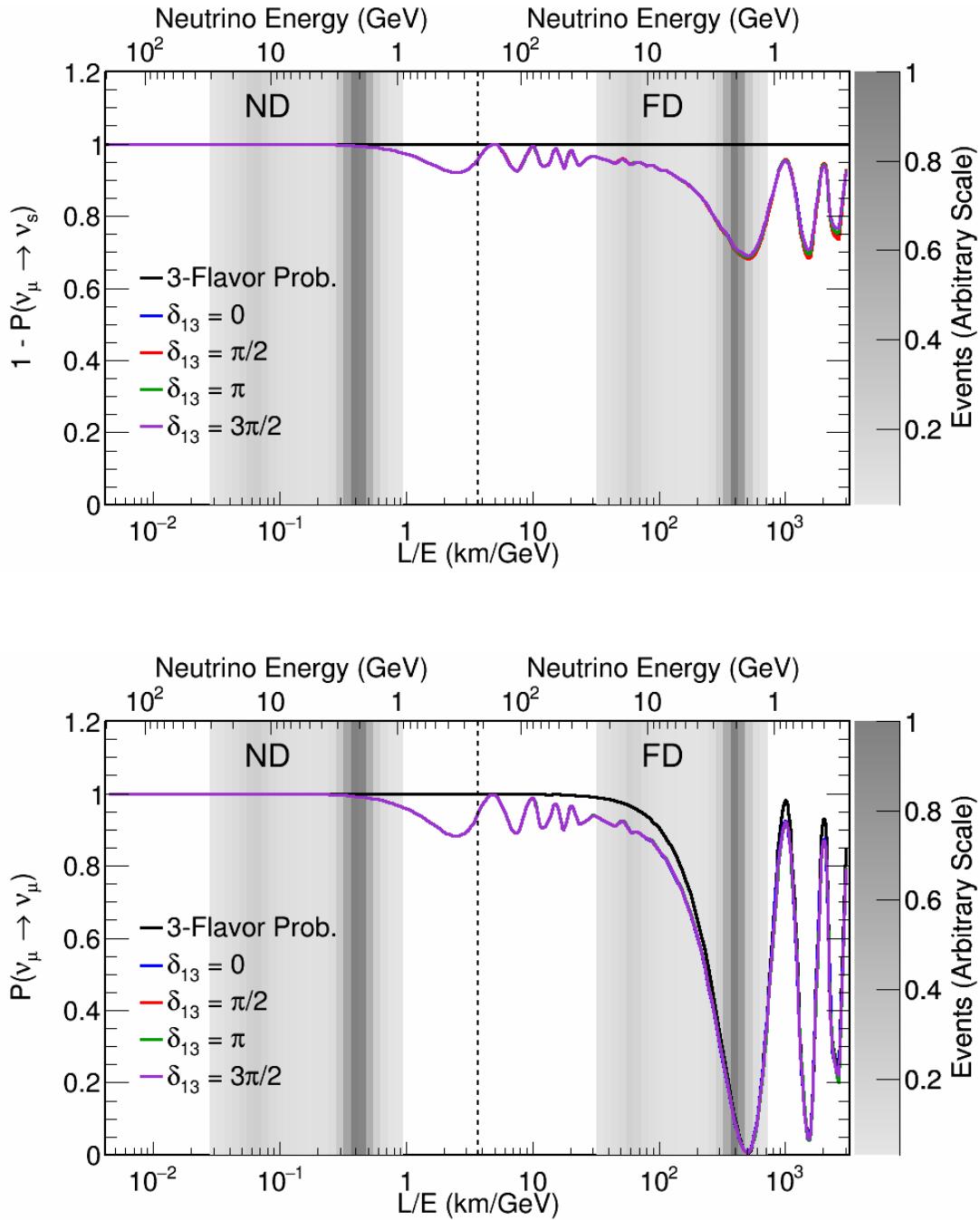


Figure 2.14: Total active neutrino survival probability (top) and ν_μ survival probability (bottom) for different values of δ_{13} . The effects of δ_{13} are negligible on the oscillation probabilities. The grey bands show the neutrino energy spectra observed by NO ν A, with darker colors representing a higher intensity of neutrinos. Parameters not varied in the figures are held fixed to values shown in table 2.2.

Detector oscillations. Thus, the analysis will make measurements of θ_{24} , θ_{34} , and these will be translated into measurements of $|U_{\mu 4}|^2$ and $|U_{\tau 4}|^2$, and the results will be valid in the range $0.05 < \Delta m_{41}^2 < 0.5 \text{ eV}^2$.

To make the above measurements, the analysis presented here considers the rate of Neutral Current (NC) events. The rate of Charged Current (CC) events is obviously affected by neutrino oscillations. Given that standard three flavor oscillations already impact the number of CC events, it would be a difficult task to decouple oscillations between the active neutrino flavors and a fourth, sterile flavor. On the other hand, NC interactions are completely insensitive to the active neutrino flavors, and hence to oscillations between them. In the standard three flavor model, then, the total rate of NC events is only determined by the overall flux of neutrinos. Yet sterile neutrinos do not interact via the weak force at all, so if oscillations occur between active and sterile flavor states, this would also manifest in the NC rate as a deficit of events. Therefore, a NC event deficit would be a clear indication of physics outside the standard three neutrino flavor model, and this dissertation will interpret that in the context of oscillations to an extra, sterile, neutrino flavor state.

3

The NO ν A Experiment

The NO ν A experiment, or the NuMI Off-Axis ν_e Appearance experiment, is a long baseline, two detector, neutrino oscillation experiment designed to measure ν_μ to ν_e appearance from the Fermilab NuMI beam. The two functionally identical NO ν A detectors are situated 14.6 mrad off-axis of the beam center, with its Near Detector (ND) located approximately 1 km downstream of the beam target source and its Far Detector (FD) located 810 km from the source in Ash River, MN. The main purpose of the ND is to constrain the neutrino beam energy and composition by measuring the beam very close to its source, i.e., before the neutrinos have had a chance to oscillate. The FD then measures the oscillated neutrinos. This chapter describes the NuMI beam and NO ν A detectors in more detail.

3.1 THE NUMI BEAM

The NuMI beam is a ν_μ beam generated at Fermi National Accelerator Laboratory, or Fermilab, in Batavia, IL, and the source of neutrinos for the NO ν A experiment. The neutrino beam is created by

accelerating protons to 120 GeV, colliding them with a graphite target, and allowing the products to decay to neutrinos.

The NuMI beam originates with the 120 GeV protons, and these are accelerated in the complex shown in figure 3.1. First, negatively charged hydrogen atoms are accelerated to 400 MeV in the Linac. These atoms next enter the Booster synchrotron where the electrons are removed and the protons are accelerated to 8 GeV. The output of from the booster is 13 bunches, 12 which are extracted into the Recycler Ring, each with approximately 4×10^{12} protons. In a procedure called slip stacking, additional bunches are used to double the intensity of each bunch. The first bunches in the Recycler are decelerated slightly while 6 new bunches enter the ring. As the two sets of bunches have slightly different energies, they slip relative to each other. When two bunches overlap they are captured with a special RF pulse, creating 6 larger bunches that are extracted into the Main Injector (MI) and accelerated to 120 GeV. Finally, these protons are directed to a target for neutrino production. At this point, the 6 bunches, or the spill, total about 5×10^{13} protons. With a cycle time of 1.333 s, the NuMI beam reaches 700 kW, which is the most powerful beam in the world [47].

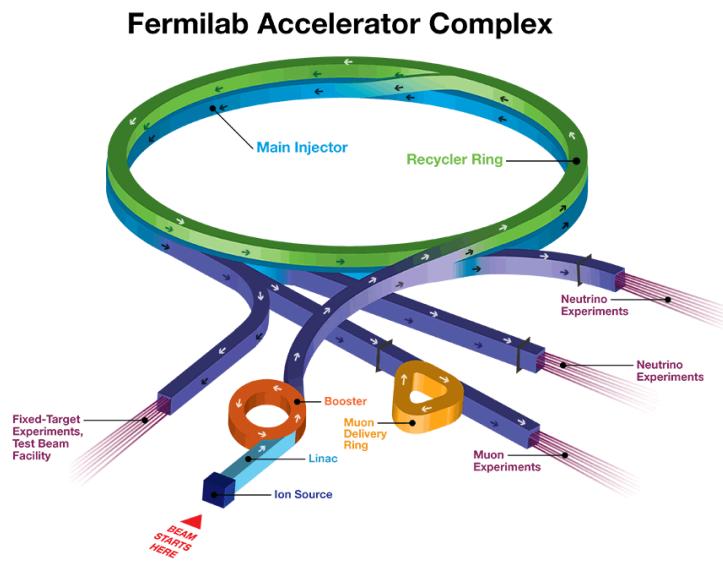


Figure 3.1: Schematic of the Fermilab accelerator complex.

In truth, the above is the design goal of the NuMI beam upgrades. The NuMI beam was originally designed for the MINOS experiment [48], and it was operating at ~ 200 kW with no slip stacking at the time of writing the NO ν A Technical Design Report (TDR) in 2007 [47]. Since then, the Fermilab Accelerator Division has been steadily ramping up the beam power, as shown in figure 3.2. At the time of writing of this dissertation, the NuMI beam was in the midst of its 2016 summer shutdown. Before the beam was powered down, it most recently was running stably at about 560 kW with 6+4 slip stacking, or doubling the proton intensity in just 4 of 6 bunches [49, 50]. However, the NuMI beam briefly reached its design goal of 700 kW with full 6+6 slip stacking during a test on June 13, 2016 [51].

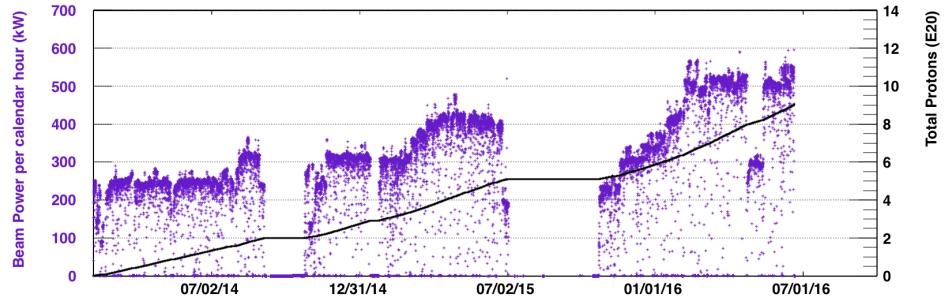


Figure 3.2: Beam power vs time since NO ν A began collecting data. The average beam power is shown as purple points, and the integrated POT is shown in the black curve.

To generation of the actual neutrino beam begins with the extraction of the 120 GeV protons from the MI into the complex shown as a schematic in figure 3.3. The protons are directed 3.3° downward into the earth so the NuMI beam points directly toward the MINOS FD in Soudan, MN. First the protons collide with atoms inside a long, thin target. The target consists of 47 graphite ‘fins’ that are each 20.0 mm long, 6.4 mm wide, and spaced 0.3 mm apart for a total length of 95.4 cm. The beam size at the target is 1 mm \times 1 mm, and the length of the target corresponds to ~ 2 interaction lengths for the protons [48]. NO ν A performs its data accounting by summing the protons on target, or POT. The interactions in the target produce a large number of secondary hadrons consisting mainly of charged pions with a smaller contribution of kaons. These particles next pass through two magnetic, parabolic horns. A current of 200 kA is passed through the horns creating a $1/r$ magnetic field that focuses a one sign of the charged particles and defocuses the opposite sign. The current in the horns can be reversed,

which allows the NuMI beam to run in either neutrino or antineutrino mode. When positively charged particles are focused, they decay into neutrinos, and similar for negatively charged particles and antineutrinos. Next, the focused charged particles travel through a 675 m long, 1 m radius decay pipe where the hadrons decay into neutrinos and tertiary charged leptons [47]. Finally, everything passes through a hadron monitor and absorber followed by three muon monitors spaced between sections of rock, leaving only the neutrinos to reach the ND hall.

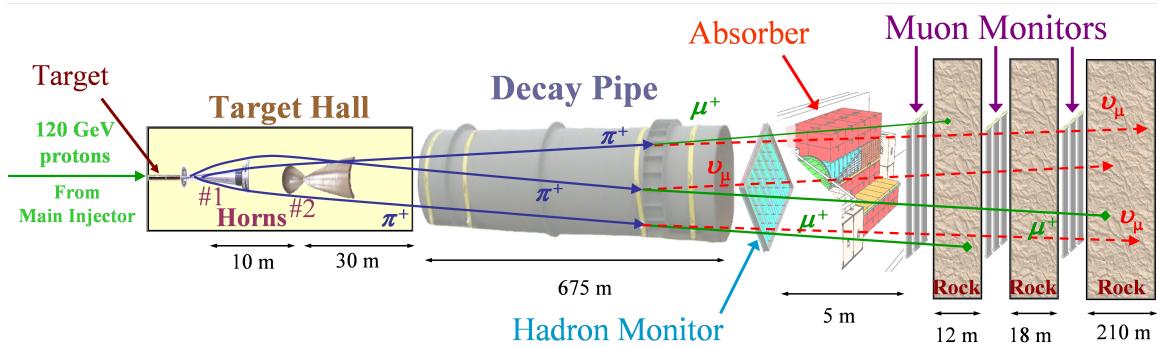


Figure 3.3: Schematic of the NuMI beam.

As a final note for this section, all of the data analyzed in this dissertation was taken during neutrino mode running, so the anti-neutrino mode is not considered further.

3.2 OFF-AXIS DETECTORS

The NO ν A experiment was designed to have its detectors located 14.6 mrad off-axis from the NuMI beam, and this greatly affects the observed neutrino energy spectrum. The main decay mode for charged pions (99.98770%) and kaons (63.56%) is the two body decay into a muon and muon-neutrino [34]. In the center of mass frame of the decaying meson, this is a deterministic, isotropic decay. However, the lab frame is highly boosted resulting in the following neutrino flux and energy spectrum for a detector with cross section area A located a distance L from the decay point.

$$\phi_\nu = \left(\frac{2\gamma_m}{1 + \gamma_m^2 \theta^2} \right)^2 \frac{A}{4\pi L^2} \quad (3.1)$$

$$E_\nu = \left(1 - \frac{m_\mu^2}{m_m^2} \right) \frac{E_m}{1 + \gamma_m^2 \theta^2} \quad (3.2)$$

Above, θ is the angle between the neutrino and meson, and the subscript m refers to the decaying meson, i.e., E_m is the meson energy, m_m is the meson mass, and γ_m is E_m/m_m . The effects of an off-axis location is shown in figure 3.4. There are two competing effects that occur. While the flux of neutrinos is drastically reduced, the neutrino energy is highly constrained, with a net result of a very narrow band beam of neutrinos.

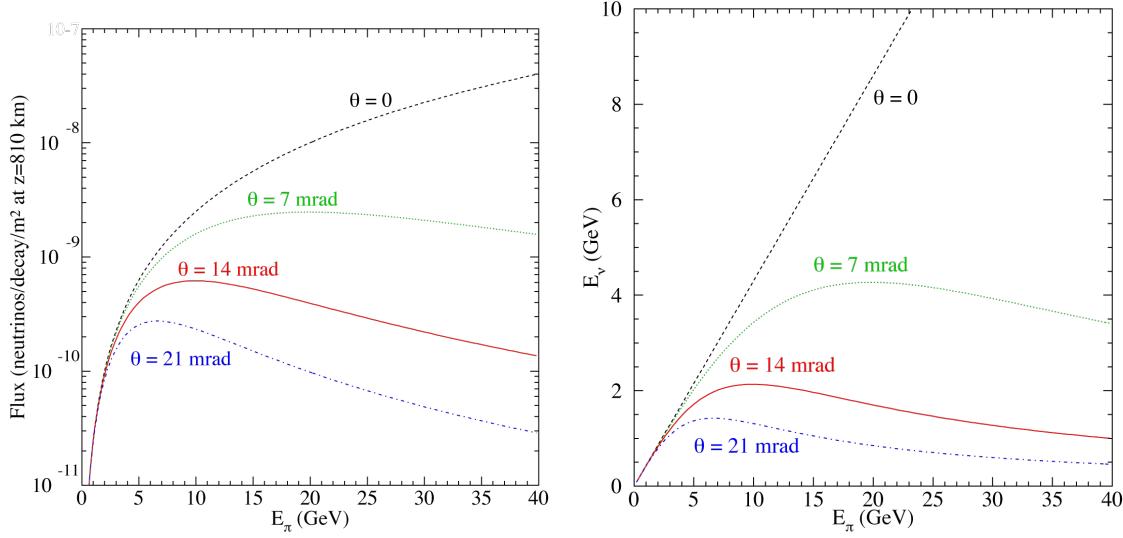


Figure 3.4: The neutrino flux (left) and energy (right) as a function of the decaying pion energy for several different off-axis angles. The on axis values are shown in black, and the red curves show the values for NO ν A. Plots from [47]

The narrow energy band result is shown in figure 3.5, which compares the energy spectrum for several off-axis angles and shows composition of the neutrino beam observed by NO ν A, both for the FD and without considering oscillations. The small contamination of $\bar{\nu}_\mu$, ν_e , and $\bar{\nu}_e$ components come from several effects. The focusing horns discussed in the previous section are unable to defocus the most energetic wrong sign mesons, which mostly contributes to the $\bar{\nu}_\mu$ component through the decay $\pi^- \rightarrow \bar{\nu}_\mu + \mu^-$. In the decay pipe, some of the muons from pion decay also decay, contributing to both the $\bar{\nu}_\mu$ and ν_e components. While most K^+ decay to ν_μ , a small fraction (5.07% [34]) decay via $K^+ \rightarrow \pi^0 + e^+ + \nu_e$. There is a very small $\bar{\nu}_e$ component that is largely due to K_L decay.

NO ν A was designed to be off-axis in an attempt to optimize the amount of observed ν_e appearance, as the experiment name would suggest. At the NO ν A FD, with $L = 810$ km, the neutrino energy range observed due to the particular off-axis angle falls right within the first oscillation maximum for ν_e

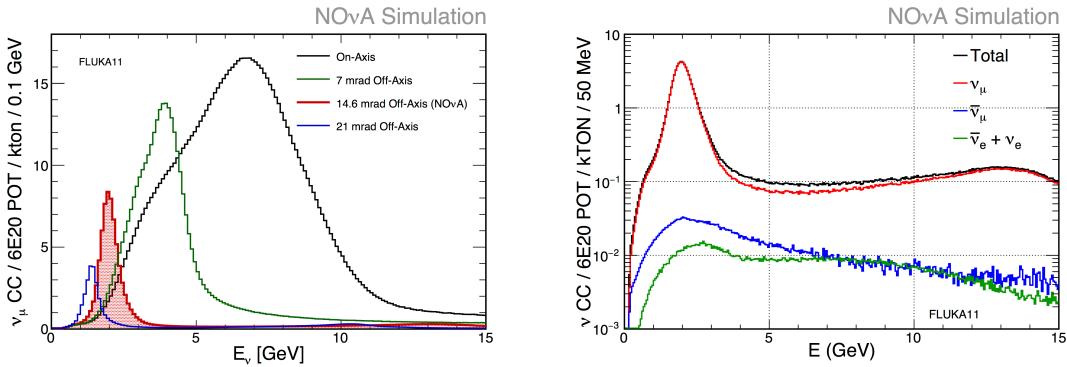


Figure 3.5: The neutrino flux (left) and energy (right) as a function of the decaying pion energy for several different off-axis angles. The on axis values are shown in black, and the red curves show the values for NO ν A.

appearance. Figure 3.6 shows the oscillation probability for ν_e appearance and the event predictions for the dominant components, both as a function of the true neutrino energy. The narrow band beam and first oscillation peaks match quite well, allowing for the best possible measurement of ν_e appearance, one of NO ν A's two main analyses [42]. Furthermore, the probability of ν_μ disappearance (shown only in the event component predictions) highly suppresses the number of expected ν_μ CC events, allowing for precision measurements of ν_μ disappearance, the other main NO ν A analysis [43].

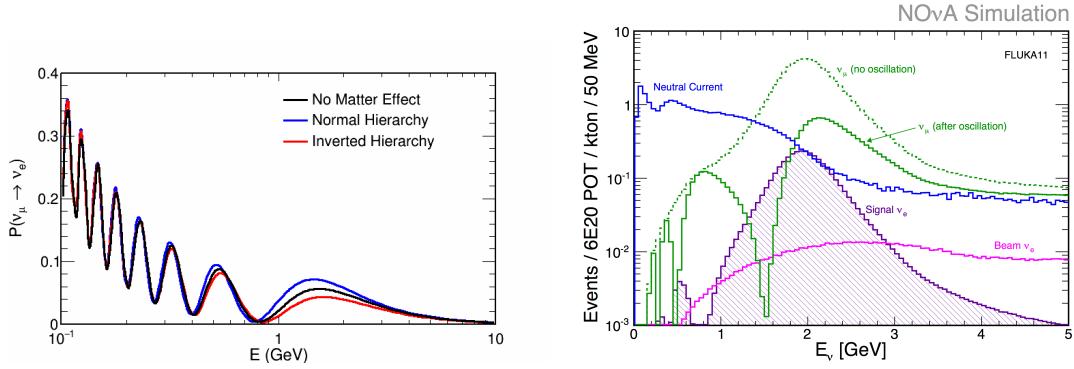


Figure 3.6: Left: the probability of ν_e appearance. The probabilities are calculated assuming standard three flavor oscillations using the relevant values shown in table 2.2, except for δ_{CP} which is here set to $3\pi/2$. Right: The expected number of events as a function of true neutrino energy (with the exception of NCs), for individual components. The NC component is shown as a function of true visible energy $E_{vis} \equiv E_\nu \times y_{Bj}$, where y_{Bj} is the fractional energy loss Bjorken scaling variable.

As with the main NO ν A analyses, the off-axis nature of the experiment has both pros and cons when searching for sterile oscillations via NC disappearance. Obviously a reduced neutrino flux is not great

when searching for neutrino interactions. However, the CC backgrounds are highly suppressed due to oscillations, even with the ν_e appearance. This makes the selection of a relatively pure sample of NC events at the FD a much easier task.

3.3 THE NO ν A DETECTORS

The NO ν A Detectors are highly segmented, liquid scintillating calorimeters. They consist of PVC plastic cells filled with a liquid scintillating solution, and the active scintillating material makes about 65% of the detectors by mass. The two detectors are functionally identical in order to reduce the effect of systematic errors, with the ND used to constrain the neutrino beam flux and energy and the FD used to measure the oscillated neutrino spectrum. This section first describes the general detector design, and later focuses on the specific details of the individual detectors. The information contained here is largely a summary of the NO ν A TDR [47], which contains a much more detailed account of the detectors.

The basic structure of the NO ν A detectors is a set of PVC plastic cells arranged together as planes and alternating in orientation. The cross section of the cells is approximately rectangular with dimensions of 4 cm \times 6 cm. The cell length spans the entire height or width of the particular detector, depending on the orientation of the given plane. The PVC is extruded in groups of 16 cells, as shown in figure 3.7. Two extrusions are glued together to create a ‘module,’ shown in figure 3.8. Many modules are glued together side by side to form a single plane. The planes are oriented such that the NuMI beam direction is approximately parallel to the plane normal vector.



Figure 3.7: A schematic of a PVC cell extrusion. The PVC is extruded along the axis into the page with a length equal to the detector height or width. Figure from [47].

The local detector coordinate systems consider the NuMI beam direction as the Z axis, height (or altitude) as the Y axis, and the X axis is the other remaining. As the NuMI beam points approximately to the north, the direction of positive X happens to point west. (The beam travels in the positive Z direction, and the positive Y direction points to higher elevation.) In this frame, a given plane can either

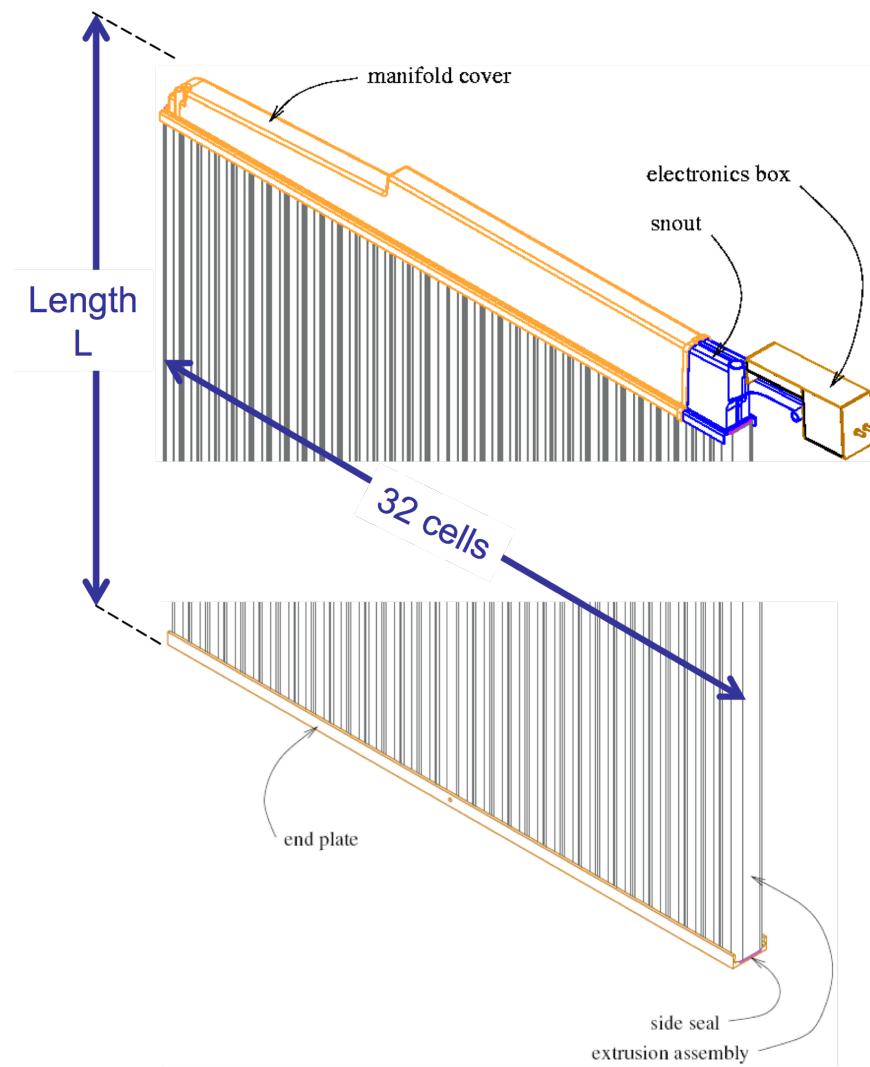


Figure 3.8: A schematic of a module. A full plane is made up of modules connected side by side. Figure from [47].

locate activity in an XZ plane if the cells are in the vertical orientation, or a YZ plane if the cells are in the horizontal orientation. For this reason, the planes alternate between horizontal and vertical orientations to give a full three dimensional view of detector activity.

Each cell is filled with a liquid scintillator solution that consists of several different compounds, each with a specific function. Pseudocumene (1,2,4-Trimethylbenzene) is the scintillant used in the solution. The solution contains two wave shifters, PPO (2,5-Diphenyloxazole) and bis-MIS (1,4-bis(2-methylstyryl)benzene). A small amount of Stadis-425 is used in the solution as an antistatic agent, and a small amount of tocopherol (Vitamin E) is added as an antioxidant. Mineral oil provides the solvent

for each of the above compounds and makes up the majority of the solution. Table 3.1 summarizes the solution contents and lists their mass fractions. The resultant mixture produces scintillation light in the low wavelength UV region, and shifts this to the violet-blue region.

Table 3.1: A summary of each component in the scintillation solution. Adapted from [47].

Compound	Purpose	Mass Fraction
Mineral Oil	Solvent	95.8%
Pseudocumene	Scintillant	4.1%
PPO	Waveshifter	0.091%
bis-MSB	Waveshifter	0.0013%
Stadis-425	Anti-static	0.0003%
Vitamin E	Antioxidant	0.0010%

To collect the scintillation light, the inside of each cell is coated with titanium dioxide and a loop of a plastic wavelength shifting (WLS) fiber rests inside each cell. The TiO_2 is used as a reflective agent to maximize the scintillation light collection. A fluorescent dye in the WLS fiber absorbs this scintillation light and shifts it to the blue-green region. This light is then transmitted through the fiber to electronics for readout. A complete schematic of a cell is shown in figure 3.9.

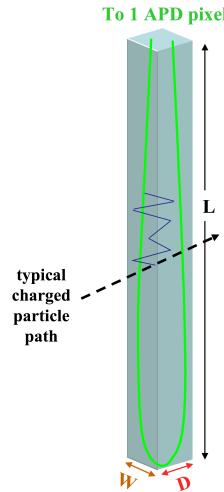


Figure 3.9: A schematic of the inside of a cell. The width and depth are approximately $W \approx 4\text{ cm}$ and $D \approx 6\text{ cm}$, respectively. The green curve is the WLS fiber, and the zigzagging blue line is a path light might take from emission by the scintillator through absorption by the WLS fiber. Figure from [47].

The first component in data readout is collection and amplification by a Hamamatsu avalanche photo-

todiode (APD). A single APD collects the light from all 32 cells in a module, collecting light from both ends of the WLS fiber. Figure 3.10 shows an APD and the ends of the WLS fibers set in a mount to be placed next to the APD. APDs were chosen over PMTs due to their high quantum efficiency, especially for green light. This is shown in figure 3.11 against the light output of the WLS fibers.

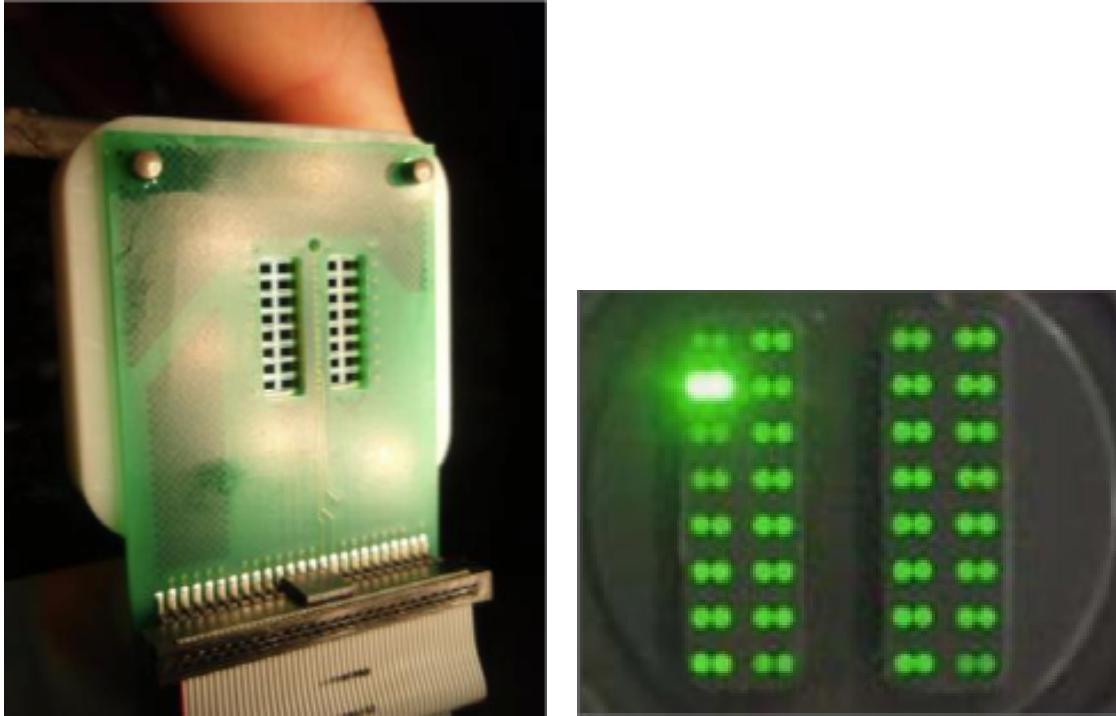


Figure 3.10: Left: An APD. Right: The ends of 32 WLS fibers that will be pressed against the APD face shown on the left.

Each APD is connected directly to a thermoelectric cooler and housed in an electronics box as shown in figure 3.8. The cooler keeps the operating temperature for the APD at -15° to reduce noise. The APDs are held at a voltage of about 425 V to run at a gain of 100, though exact voltages are set individually. In October 2015, the gain for the FD APDs was increased to 150 to be more sensitive to light from the far ends of the fibers. The thermoelectric cooler is part of a larger front-end board (FEB) that is also housed inside the same electronics box as the APD. The FEB thus provides the cooling and voltage for the APD, but it also digitizes the signals from the APD.

The electronics for the NO ν A detectors were installed in 64 plane sections called diblocks. The signals from (up to) 64 FEBs are read out by a single data concentrator module (DCM). Due to the alter-

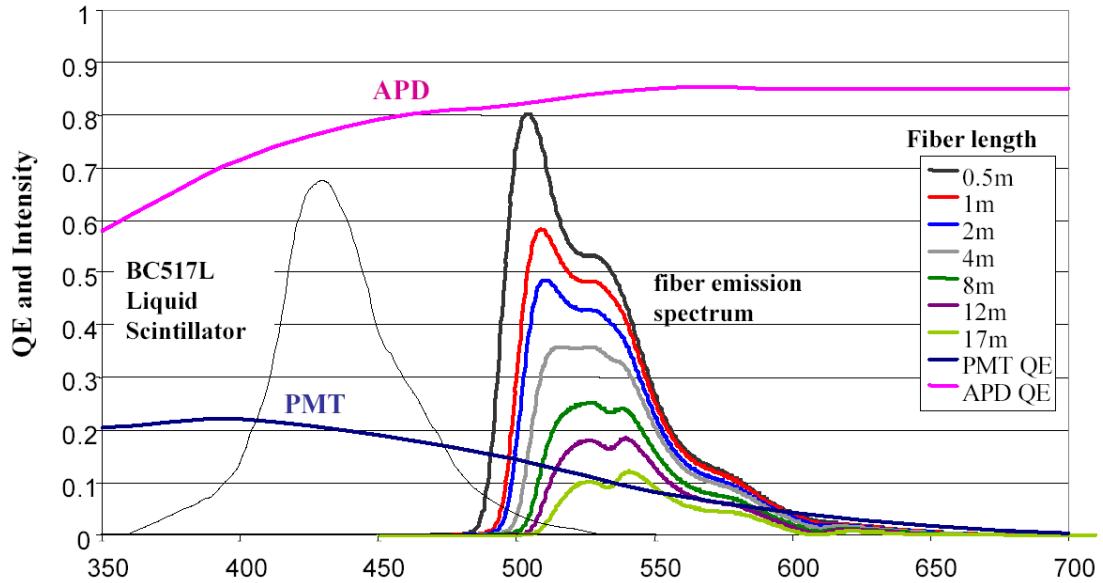


Figure 3.11: Quantum efficiency of APDs as a function of light wavelength. For comparison, the QE of PMTs is shown in navy. The light output of the liquid scintillator is shown in the thin, black curve. The light output for the WLS fibers is also shown for a variety of different lengths through the fiber. The intensities are in arbitrary units. While the blue light is almost fully attenuated through the fiber, the green light is only modestly attenuated. Figure from [47].

nating orientations of the planes, and thus the modules, electronics boxes are found on either the top or one side of the detector. Consequently, a given DCM does not read out one FEB per plane, but rather (up to) two FEBs per plane. The two FEBs are always from adjacent modules, and the modules always share the same X or Y coordinate range based on the module orientation. This diblock structure is important because it allowed for data collection at the FD before the detector was fully commissioned.

Data stored by the DCMs are stored in $50\ \mu\text{s}$ blocks called microslices, which are transferred to a data buffer farm in larger $5\ \text{ms}$ blocks called millislices. Data is stored in the farm for up to 20 minutes before being erased. During this time, different trigger decisions can be applied to keep the data for different purposes such as recording beam activity, searching for good calibration candidates, or searching for exotic physics like monopoles. Data that pass a trigger decision are written to disk and permanently stored, and this mechanism allows the FEBs to collect data continuously.

Detector activity that could be due to a beam event is captured within a $500\ \mu\text{s}$ window. These events are specifically triggered by a signal from Fermilab, corrected for time of flight at the FD. The event window is centered around the $\sim 10\ \mu\text{s}$ beam spill, which allows for considerable drift in the timing systems

without loss of data [36]. Other $500\ \mu\text{s}$ event windows without beam candidates are stored as cosmic trigger data that are useful for predicting the cosmic background event rate.

Data are collected by organizing these events into larger run and subrun file structures. Subruns are populated with events until either a maximum file size is reached, or after an hour of data taking has elapsed, whichever comes first. Runs are comprised of 64 subruns for the FD, and 24 subruns for the ND, and are numbered in succession such that larger numbered runs correspond to later times.

3.3.1 FAR DETECTOR

The NO ν A FD is situated 810 km from the beam target in Ash River, MN. Both the horizontal and vertical planes consist of 12 modules, or 384 cells, for a total cross sectional area of about $15.2\ \text{m} \times 15.2\ \text{m}$. The fully commissioned detector has 14 diblocks (data taking began with as few as 4), or 896 planes, with a total length of about 59.6 m. This means that the detector reads out 344,064 individual cells, or channels. The FD weighs 14 kT, which is about 65% scintillator and 35% PVC plastic.

The detector itself is inside a building that is at the surface, mostly dug underground and partially above ground. The site is at an elevation of about 1220 ft and thus subject to a large background of cosmic rays. A large amount of overburden is used to mitigate this effect. 2.5 ft of precast concrete covers the entire detector and is capped with 1.5 ft of concrete cast in place. Finally, 6" of barite, or barium sulfate, is placed above the concrete and composed of loosely packed rocks. Barite was chosen because it is a high Z material with a high stopping power for cosmic ray photons. Together, the concrete and barite provide around 14 radiation lengths for the cosmic rays. Figure 3.12 shows the FD and the building inside which it is located.

Even with the overburden, it is still expected to see cosmic rays in the FD with a rate over 200 kHz. This is the main reason why NO ν A must use a trigger from the Fermilab accelerator to signal a beam event, but also why taking cosmic trigger data is so important. These data are used to develop and train cuts designed to remove cosmic background. Later, a predicted background is made from an independent sample, namely the data regions in the NuMI trigger events that fall safely outside the smaller beam window.



Figure 3.12: Left: The NO ν A FD. The beam comes from the back of the photograph and out of the page. The red and yellow structure is a pivoter that was used to rotate and put the detector blocks into position and is not part of the detector. A person can be seen halfway up on the right of the detector for scale. Right: The FD building. The detector is beneath the rocky section on the right of the photograph, which only partially shows the barite overburden.

3.3.2 NEAR DETECTOR

The NO ν A ND is situated 1015 m from the beam target, on site at Fermilab. It is composed of two different sections, a fully active region much like the FD and a muon catcher. In the fully active region, each plane consists of 3 modules, or 96 cells, for a total cross sectional area of about 3.8 m \times 3.8 m. The region is made up of 192 planes with a length of about 12.7 m and corresponds to 3 diblocks for the electronics. The NuMI beam first passes through this fully active region, but since it is so much shorter than the FD, the muon catcher is necessary to be able to accurately measure particle energies. For this section, planes of steel, which have an excellent stopping power for muons, are interspersed between the alternating planes. The energy loss of muons passing through steel is well known allowing for accurate event reconstruction. This region is specifically made of a horizontal plane, a 10 cm steel plane, and a vertical plane, repeated 11 times for an additional 22 active planes. However, for economic reasons the steel and planes in this region were recycled from a prototype near detector that was built using the original design of 3 modules \times 2 modules. Thus, the planes in the muon catcher are as wide as the fully active region, but only 2/3 as tall, at about 2.6 m. In total, the ND is 15.8 m in length, has 20,192 cells, and 214 active planes. It weighs 290 T, of which 130 T is scintillator, 81 T is PVC plastic, and 78 T is steel

[52]. Figure 3.13 shows the ND from both beam directions, and table 3.2 summarizes the properties of both detectors.

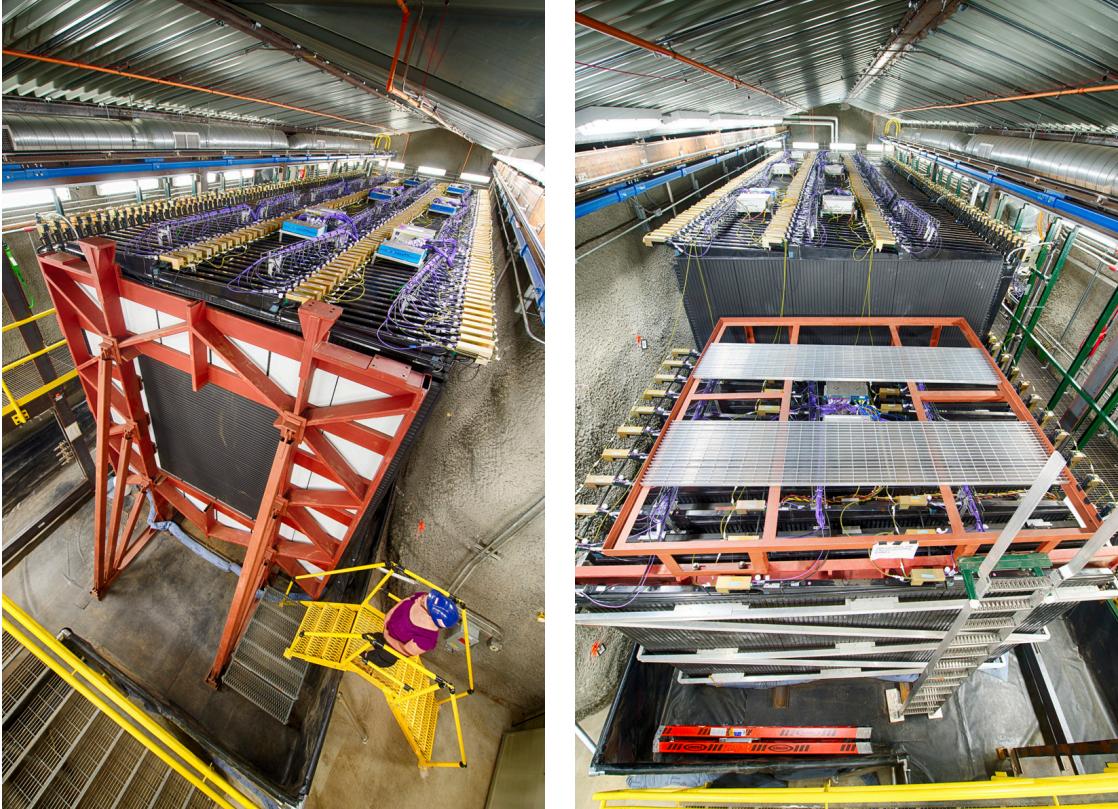


Figure 3.13: The NO ν A ND. Left: View from the side where the beam enters the detector. The red structure is a bookend and not part of the active detector. A person is seen on the right for scale. Right: View from the side where the beam exits the detector. The red and grey structure above the muon catcher is for worker access to this detector region and not part of the active detector.

Unlike the FD, the ND is not a surface detector. It is located 105 m below the surface in the same underground cavern as the MINOS ND, as depicted in figure 3.14. This means that the ND is not subject to a cosmic ray rate like FD. However, given the close proximity to the beam source, the ND can observe multiple neutrino events in every beam spill. This leads to the potential for pileup, so the electronics for the ND were designed to sample the data 4 times faster than the FD to deal with this.

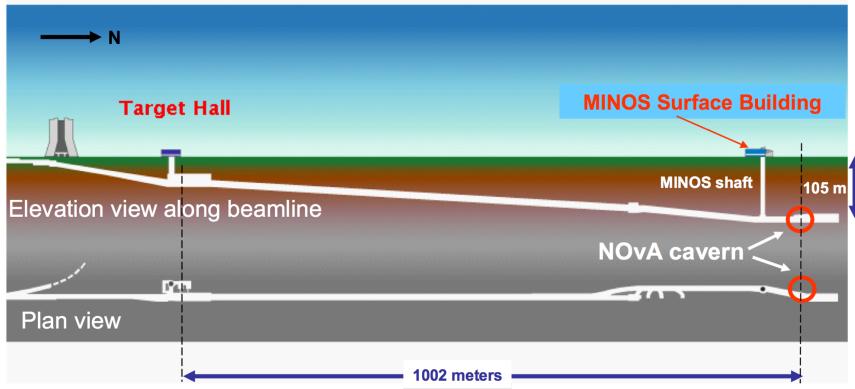


Figure 3.14: A schematic of the NuMI beam line, showing the location of the ND hall. Figure from [47].

Table 3.2: A summary of the physical properties of both detectors. For the ND, a single value means the property is consistent across the entire detector. Otherwise, values in parenthesis are for the muon catcher only. Note that the detector width and height exceed the cell lengths due to inclusion of the end caps on either side of the planes.

Detector Property	FD	ND
Detector Weight (kT)	14	0.3
Number of Diblocks	8	3 (+Muon Catcher)
Number of Planes	896	214
Number of Cells	344064	20192
Detector Length (m)	59.6	15.8
Detector Width (m)	15.8	4.2
Detector Height (m)	15.8	4.2 (2.8)
Cell Depth (cm)	5.6	5.6
Cell Width (cm)	3.6	3.6
Cell Length (m)	15.2	3.8 (3.8H, 2.6V)

4

Experiment Simulation

Simulated data, or Monte Carlo (MC), plays a very important role in the NO ν A analyses. The FD only expects $O(10^2)$ events per year, not nearly enough statistics to fine tune things like selection cuts. Furthermore, it is important that actual parameter measurements use an independent sample of data from that used to train and tune an analysis. MC solves both of these problems by being explicitly different from the real data and having only computer memory limit statistics.

The full simulation chain is very long and comprises multiple steps, which reduces complexity and provides more opportunity for validation. At every step, the information from previous pieces of the simulation chain are kept so it is always possible to reproduce future results and trace any errors that may occur. The simulation chain consists of two main components, simulation of the beam and the detector response. This chapter discusses both of these areas in more detail, and finishes with a discussion of packages designed to study and validate the results.

4.1 FLUX SIMULATION

The first main simulation segment is the simulation of the NuMI beam, or the flux simulation. This section starts with the interaction of protons in the target and ends with experiment independent flux files full of neutrino rays that most importantly contain the flavor, direction, energy, and momentum of the neutrinos.

The flux simulation is performed using the FLUGG package, an interface between FLUKA [53, 54] and Geant4 [55, 56]. The version used for the analysis in this dissertation was FLUGG 2009.3, combining FLUKA2011 and Geant4 v4.9.6.po3(c). The FLUKA package is designed to simulate particle interactions and was used to simulate the proton interactions in the target. Geant4 is a toolkit for simulation of particles propagating through matter, and is used in this context to propagate the target interaction products through a detailed model of the NuMI beam line. This model of the beam includes all of the elements discussed in section 3.1, starting from the target hall all the way through the rock before the detector halls. The second focusing horn, a representative component of the geometry model, is shown in figure 4.1.

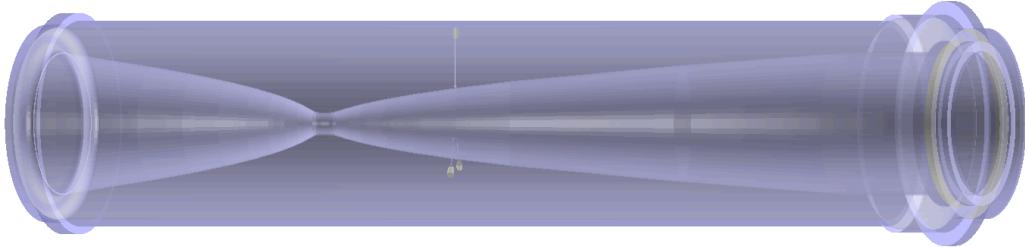


Figure 4.1: Visualization of the geometry used to model the second horn. This image was made during efforts to correctly position the three spider supports seen near the center of the schematic. Figure from [57].

The output of the FLUGG simulation is a set of flugg format flux files. These include the information described at the beginning of this section, but have several other key features as well. Simulating the particle products of every target interaction is extremely expensive on computing resources, so not every particle is tracked. Instead, similar neutrinos are often dropped in lieu of a single entry that is given a greater *importance weight*, one of two weights stored in the flux files. The information about the neutrino parent is also stored in the files, including the particle type, energy, momentum, and decay point.

This affords later simulation steps the ability to “re-decay” a neutrino ray to a specific location and calculate the relative probability that this would occur. The latter value is stored as a *propagation weight*. Retaining the parent information also allows for event reweighting based on hadronic model studies.

The next piece of the flux simulation is largely a repackaging of the flugg format flux files into a unified format, called Dk2Nu [58]. FLUGG is certainly not the only flux generator, and NO ν A is not the only experiment that uses the NuMI beam line. The Dk2Nu format is unified in the sense that it has the same format regardless of flux generator, and valid for any experiment using the NuMI beam.

Dk2Nu format flux files (or simply dk2nu files) are one of the main outputs from the flux simulation. However, one last step is taken to further simplify computation time for later simulation steps. The entries in the dk2nu files are sampled over a ‘window’ that shadows the detector of interest, resulting only in a set of neutrino rays that could create interactions within the detector. The output files from this final step are in the GSsimpleNtpFlux format, or more simply referred to as gsimple files [59].

4.2 DETECTOR SIMULATION

The second main section in the simulation chain is the detector simulation. This section begins with neutrino interactions, models many different detector response effects, and finishes with files of simulated raw hits, or the signals read out by the detector electronics. The latter portion of this section closely follows the NO ν A Simulation Technical Notes [60, 61], which provide more detail and further references.

Neutrino interactions are modeled using the GENIE event generator product [62, 63]. For the analysis in this dissertation, version v2.10.4 was used. GENIE requires as input flux files, a model of the detector geometry, and cross section information, and from this determines if and where a neutrino interaction occurs, the type of interaction, and the kinematics of the interaction products. The generator uses a sophisticated model of the nucleus that allows for quasi-elastic, resonant, or deep inelastic scattering, taking into account any intranuclear scattering that may occur after the initial interaction takes place. The output of a GENIE interaction is a list of primary particles that escape the target nucleus and the relevant kinematic variables that describe each entry.

There is growing evidence [64, 43] that there is a type of nuclear interaction involving a quasi-bound nucleon-nucleon pair. This interaction is coined the meson exchange current, or MEC, and is included as an empirical model in GENIE. These events largely occur in an energy region between quasi-elastic (QE) and resonant (RES) events, vanishing by 5 GeV where QE scattering from free nucleons well model external data. While the inclusion of MEC interactions does not bring the data and MC into complete agreement, it substantially improves the discrepancy; see figure 4.2. Consequently, an additional systematic error is evaluated for the MEC scale, discussed in chapter 8. Furthermore, data used to constrain the empirical model in GENIE was only available for CC interactions, so there are no NC MEC events in the MC [65].

In truth, the actual ‘tuned’ MEC model is applied as a weight on an event level, and there are a couple other effects applied in a similar fashion. The effect of long-range nuclear charge screening, or RPA (random phase approximation) suppression was studied by creating a shifted spectrum using a set of event weights. This shifted spectrum was translated into a systematic error band, assessed in chapter 8. The single non-resonant pion production rate was decreased also using event weights to a level consistent with the most recent data studies [65].

The GENIE output is next fed back to Geant4 for particle propagation through the detector geometry [55, 56]. It models the primary particle interactions within the detector, including energy deposition and secondary interactions. The physical processes considered and modeled are somewhat configurable via the choice of the physics list. At this point in the simulation, it is much more important to accurately track the detailed interactions, so particles are not dropped as in the flux simulation; instead, a more precise model was used here. The specific physics list used was QGSP_BERT_HP, where QGSP specifies that high energy hadrons (> 10 GeV) are modeled with the quark gluon string model, BERT specifies that lower energy hadrons (< 10 GeV) are modeled with the Bertini cascade model, and HP toggles the usage of a high precision neutron model that tracks low energy neutrons (< 20 MeV) [60]. The output from Geant4 at this stage is a list particles involved in the detector interaction and a full suite of information about each one, called the MC truth. The information that is most important to the next simulation step is a list of energy depositions and positions.

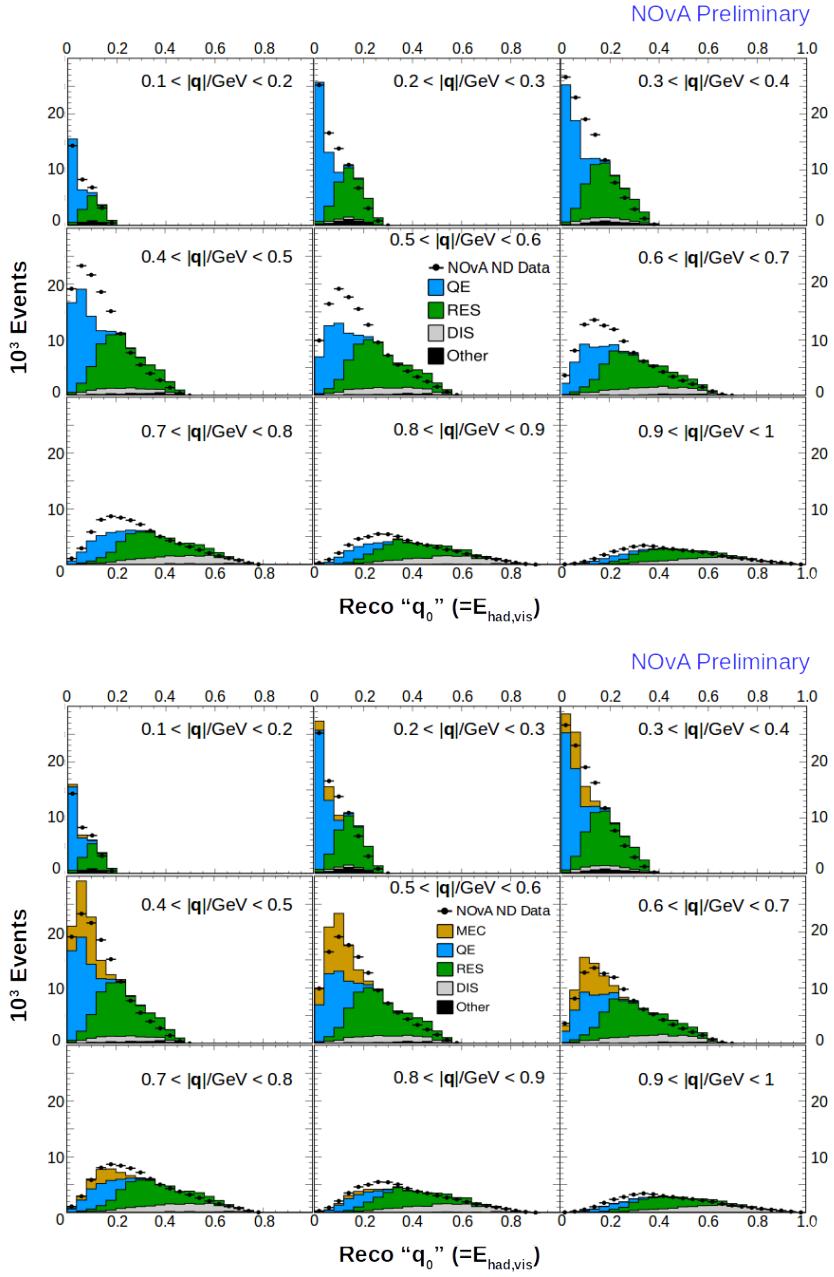


Figure 4.2: Distributions of the reconstructed visible hadronic energy for different ‘slices’ of estimated three momentum transfer, with (right) and without (left) MEC events. The events used for these distributions come entirely from ND selected ν_μ CC events. Plots from [66].

It is worth noting that the simulated events for the Near and far Detectors are generated somewhat differently. The FD expects a very low event rate per beam spill, while the ND expects multiple events per spill. As a result, the FD events are generated one at a time with the simulation calculating the POT needed for the event occur, and ND events are generated assuming a constant POT/spill and a variable number of events. Furthermore, the high rate of interactions at the ND means that it observes many muons that originate from neutrino interactions that begin in the rock outside of the detector. Most of these so called rock events do not result in any detector activity, so tracking them wastes valuable computing time. To deal with this, the ND detector events and single rock events are generated separately, and only the rock events that cause detector activity are kept. The final ND event records are made by overlaying rock events on the detector events. The number of rock singles to overlay is drawn from a Poisson distribution with a mean determined from data [67]. A top down view of the neutrino interaction vertices at the ND is shown in figure 4.3 and shows the level of detail used to model the geometry of the ND hall.

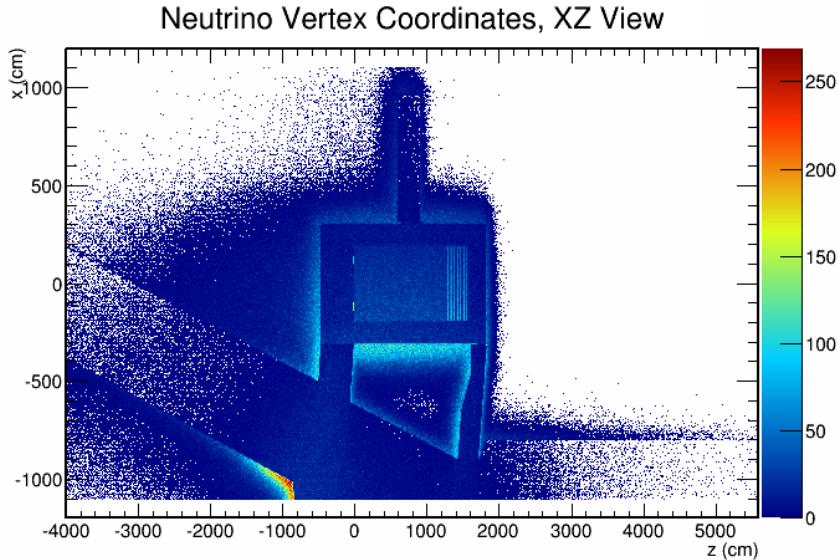


Figure 4.3: Distribution of neutrino interaction vertices in the XZ view at the ND hall. Several features of the geometry model can be seen. The detector sits between ± 200 cm in X and $0 < Z < 1600$ cm. The muon catcher steel planes are easily identified at the back of the detector as the line segments with slightly enhanced interaction rates. The two small blips of greater activity at ± 100 cm in X and 0 cm in Z are the bookend keeping the detector in place. The wide, sloped region extending from 0 on the Z axis to 0 on the X axis is an excavated pathway.

`Geant4` also handles the next piece of the simulation, the conversion between energy deposition and light yield. For low energies, the rate of light produced in liquid scintillator is proportional to the energy deposition, but the light yield begins to quench for particles with high enough energies. The Birks-Chou Law, encapsulated within a `Geant4` module, is used to model the relationship between scintillator light yield rate, $\frac{dL}{dx}$, and particle energy deposition rate, $\frac{dE}{dx}$ [68].

$$\frac{dL}{dx} = L_0 \frac{\frac{dE}{dx}}{1 + k_B \frac{dE}{dx} + k_C \left(\frac{dE}{dx}\right)^2} \quad (4.1)$$

Above, the constants L_0 , k_B , and k_C are dependent on the scintillator material and had to be estimated for NO ν A as no measurement existed for the particular materials used in the experiment. A study was performed comparing the energy deposition at the end of proton tracks in the ND for both data and MC to find parameters that would generate agreement between the two [69]. The results of the study were $k_B = 0.04 \text{ cm/MeV}$ and $k_C = -0.0005 \text{ (cm/MeV)}^2$. The output from this simulation step is a set of energy depositions now encoded as light yield called FLSHits, or fiber in liquid scintillator hits.

The next part of the simulation simulates the capture of the scintillation photons, their transport through the WLS fiber to the electronics, and their conversion into photoelectrons after collection by the APD. The detector is treated as uniform for this section, with the assumption that any individual differences between cells would be removed by the downstream calibration. Individual photons are not traced from emission to the electronics due to its expensive computation time, despite the ability of `Geant4` to perform this simulation. Rather, a template that parametrizes the photon transport are created and individual results are drawn from this.

The first template is a collection rate of photons by the WLS fibers as a function of time from scintillator emission and distance along the cell length from scintillator emission. The template is constructed using a ray tracing algorithm that considers a photon to be captured when it intersects with a fiber, and is shown in figure 4.4. The mean number of photons per energy deposited is a tunable parameter that is constrained using cosmic ray data. The scintillator emission spectrum and cell wall reflectivity are modeled as a function of wavelength based on data from NO ν A's quality control database. The reflectivity is also shown in figure 4.4.

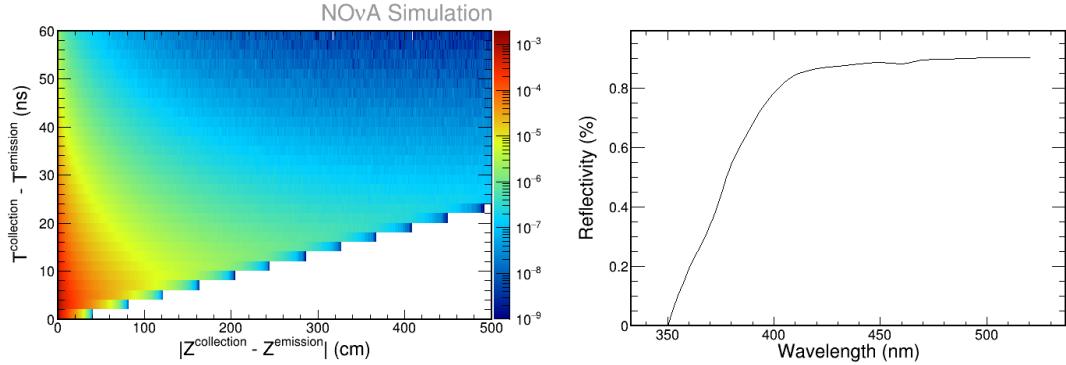


Figure 4.4: Distributions used during the simulation of photon transport. Left: collection rate as a function of time and distance travelled from scintillator emission. Right: cell wall reflectivity as a function of wavelength.

Next, the photons are transported through the fiber to the electronics. Since the fiber is looped and both ends are read out, half of the photons are transported to each end. The number of photons that reach the fiber ends is attenuated by applying an attenuation curve derived from quality control tests as a function of distance from the fiber end. Another template is constructed using a ray tracer that models transport time as a function of distance. The actual transport times are drawn from this template instead of handling the photon capture angles individually.

At this point in the simulation, there is a number of photons that reach the APD as a function of time. This is converted to photoelectrons (PEs) assuming a flat quantum efficiency of 0.85 for every APD. The exact number of PEs collected is drawn from a log-normal distribution with a mean of the expected number of PEs to account for the repeated random charge avalanche process within the APD. The final output of the photon transport simulation is a set of objects that encodes the number of PEs as a function of time, called `PhotonSignals`.

The last piece of the detector simulation models the electronics readout. This section starts with the number of PEs collected by the APDs and simulates the effects of the electronics to generate the raw signals equivalent to data. Before the PE signal from the APD is input to the circuitry simulation, an excess noise factor is applied. The excess noise PE signal is drawn from a log-normal distribution, the limit of the product of many random variables, motivated by the fact that the actual PE signal comes from many random charge multiplications. Next, three chips on the FEBs are simulated, an application-specific integrated circuit (ASIC), an analog to digital converter (ADC), and a field-programmable gate

array (FPGA).

The ASIC pre-amplifies and shapes the PE signal. A data driven effect is included at this point called APD sag, or the tendency when one APD channel observes a large hit for other channel baselines to slightly drop. The recovery from the drop can appear like real activity causing all APD pixels to register hits. Another form of noise is modeled at this point as the sum of two Gaussian Markov chains. This is done to simulate the high correlation of noise traces between adjacent samples, and is tuned to best match pedestal detector scans.

The shaped pulse is next input to the ADC, which performs two functions. The signal is first converted into an integer value, up to a maximum threshold that would saturate the ADC. These values are then sampled on a periodic basis to create a simulated time series of channel outputs. The FPGA then searches for peaks above threshold using a dual correlated sampling trace, defined from $ADC_i - ADC_{i-3}$, where ADC_j is the j th output from the ADC. A final type of noise is considered for otherwise empty cells, drawn probabilistically from templates made from histograms of ADC differences from data noise hits.

The hits above threshold are the final output of the detector simulation and are recorded as a vector of objects called RawDigits. This is the lowest level object that has the same format for both data and MC; however, the additional truth information is always stored for the MC for validation and further studies.

In real data, readout did not always occur from every cell in the full detector. Full diblocks sometimes were missing due to original commissioning or repair work. Within individual diblocks, single channels can be masked off for too high noise levels or too low data rates. The diblock configuration for data was stored by run, and the channel masks were stored by subrun. To most accurately reflect the detector running conditions, the MC was generated matching the configurations in data, or weighting the MC configurations simulated by POT.

The FD MC files generated do not include any realistic oscillation probabilities; instead, three sets of files are produced. The first set assumes that there are no oscillations, called nonswap files. These files are consist of the same percentage of ν_μ and ν_e generated naturally in the NuMI beam. The simulation does

not generate any ν_τ so these events are entirely absent from nonswap files. The second set of files, called fluxswap, completely oscillate all ν_μ to ν_e , and vice versa. Finally, the third set of files, called tauswap, completely oscillate all ν_μ and ν_e to ν_τ . All three files include the same rate of NC events. Oscillation weights as a function of true energy can be applied individually to each interaction channel at a later stage, which allows for realistic conditions that can also be tuned without generating entirely new MC.

4.3 STUDYING SIMULATION FILES

The benefits of retaining all of the information described in the previous sections is that it allows for quicker studies and validation. Truth information can be used to test the idealized proof of concept of a new analysis technique, and studies involving alternative flux models but no detector effects can be run without the time and computation expensive process of running the full simulation and reconstruction. Two packages were developed for these purposes, FluxReader and MCCheckOut.

FluxReader is a package designed to read and make distributions from Dk2Nu flux files [70]. Like Dk2Nu files, FluxReader is experiment independent and can be used for any experiment that uses the NuMI beam line. Its main design is to create sets of spectra from the Dk2Nu files, where each spectrum in the set plots the same variable(s) but shows the distribution for a unique set of neutrino flavor, parent, applied cross section, and detector location. Users are able to configure nearly everything inside macros, thus eliminating the need for recompiling code.

FluxReader makes clever use of several existing capabilities within other packages to generate its sets of spectra. Every individual neutrino ray in the flux files is re-decayed to each detector configured in a particular FluxReader macro, so each flux file can populate spectra for both the NO ν A ND and FD. The package is able to apply cross sections by reading in cross section splines from GENIE, since flux files by definition do not simulate detector interactions. The user is able to configure what cross sections are applied, with different spectra generated for each. In both of these cases, any additional histograms added to a particular set are handled automatically alleviating the need for user generated repetition of error prone code.

The strategy for evaluating beam related systematic errors was not to generate entirely new MC with

systematically shifted quantities. Instead, only shifted Dk2Nu files were created, and ratios of shifted to nominal spectra as a function of true neutrino energy were generated for each neutrino flavor (also split by sign) and provided as event weights for the fully simulated MC. These studies were performed using the FluxReader framework [71]. The specific systematics are discussed in detail in chapter 8.

While the systematic studies performed using FluxReader are the most directly related to the analysis in this dissertation, it has facilitated many other studies, from which a representative sample is briefly discussed below. The newest version of Fluka was compared against the previous version used in the NO ν A simulation, a collaboration internal study necessary to help trace major simulation differences that might arise between major MC productions [72]. A proposal to run the NuMI beam with various different horn currents was supported with studies utilizing FluxReader [73]. This test was eventually carried out, and the FluxReader distributions were used for the MC predicted energy spectra compared to the data [74]. Another study considered the effect of increasing the spacing between target fins to allow pions to more easily escape before scattering in a secondary target interaction [75].

MCCheckOut is a package designed for validation of various sections of the NO ν A simulation. It creates distributions of relevant variables for each step of the simulation to search for changes between different MC productions. Figure 4.3 shown above is one of the standard plots created by MCCheckOut. Between major productions, modules are run that compare distributions for the MC truth, FLS hits, simulated POT, and RawDigits. MCCheckOut is typically run by a single person who pushes the results to an online webpage so that a wider audience can look for and comment on simulation differences.

Ideally, only expected changes from simulation improvements are seen, but in reality, unexpected and unintended differences sometimes appear. It is vitally important to find and correct these mistakes before valuable computation time is wasted generating erroneous MC files. Several problems were discovered and fixed while this author was in charge of running simulation validation, two of which are briefly mentioned below. At one point, the window over which neutrino rays were sampled for FD interactions was shortened without taking into account the 3° angle between the beam direction and detector. This led to an absence of events in a section of the detector, as shown in figure 4.5 [76]. A similar problem

occurred later when a code change caused the location GENIE used to start to generating neutrino interactions at the ND to be pushed too far forward in Z. The result was an absence of events at the front of the detector, and specifically in the steel bookend, as shown in figure 4.6 [77]. MCCheckOut is clearly a valuable and necessary package for the MC simulation production.

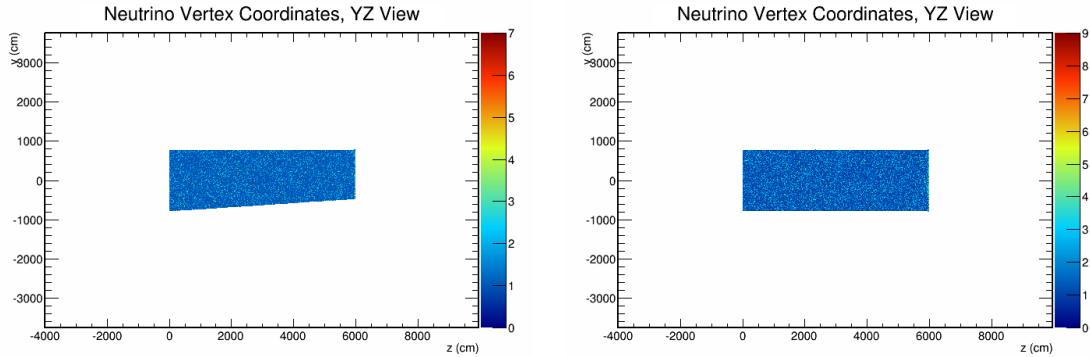


Figure 4.5: Distribution of neutrino interaction vertices in the YZ view at the FD. The left plot shows a region of detector with no events due to a shortened flux window. The right plot shows the corrected distribution after the flux window was lengthened.

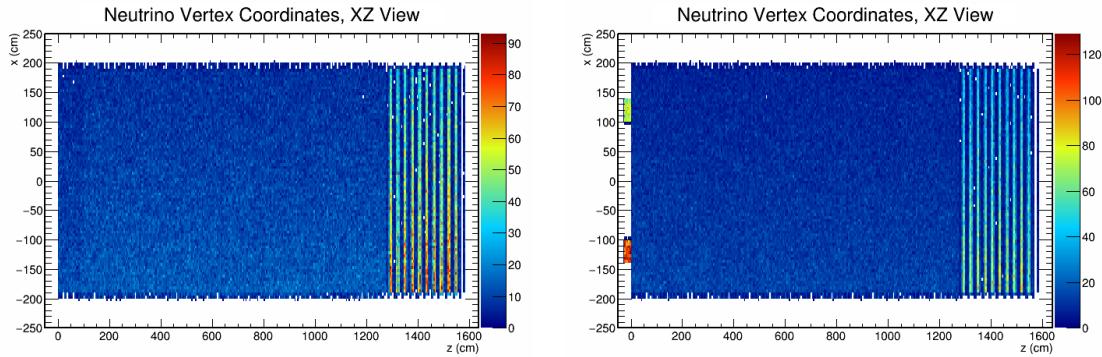


Figure 4.6: Distribution of neutrino interaction vertices in the XZ view at the ND. Rock events were removed in these distributions. There are no interaction in the steel bookend in the left plot. On the right, the simulation was corrected and events can be seen within the bookend.

5

Event Reconstruction

The most basic form of raw data collected is a vector of hits above threshold. The MC simulation described in chapter 4 also outputs events in this format. However, these objects by themselves are not very useful; instead, a certain level of reconstruction is required before real physics can be studied. The first major step in this process is to apply calibration so that the hits can be translated into a set of energy depositions, consistent throughout and across both detectors. Any number of algorithms can then be applied to create new objects or search for features, including tracks, the event vertex, or particle identifiers (PIDs). This chapter describes the calibration and elements of the reconstruction chain relevant to the NC disappearance analysis.

5.1 CALIBRATION

The purpose of calibration is to ensure a uniform detector response throughout each and across both detectors. This is done in two major steps, a relative and absolute calibration. The relative calibration

accounts for threshold effects and attenuation across a single cell. It is designed to create a uniform response throughout a cell and across a single detector. The absolute calibration creates a scale factor for each detector to convert the calibrated PE scale from the relative calibration into an energy unit. This section follows the SA notes in reference [78].

The relative calibration is designed to convert the PE signal output from the electronics into a calibrated unit, such that two equal signals from any two detector locations mean equal true energy deposited. This part of the calibration accounts for threshold effects and attenuation in the WLS fibers, outputting a corrected PE value, the PE_{Corr}.

Hits from through-going cosmic ray muons, or muons that enter and exit the detector without stopping, are used for the relative calibration. The WindowTrack algorithm, a fast algorithm that fits straight lines through hits [79], is used to produce 3D tracks from the cosmic ray events, and only those with a successful reconstruction are used. Within these events, only tricell hits are used for the calibration procedure. A tricell hit is defined as a hit in cell i within a given plane that also has hits in cells $i+1$ and $i-1$. Under special circumstances (low statistics, too many dead neighboring cells) different hits are used for a particular cell. The path length of the cosmic traveling through the cell and the distance from readout are calculated for all of the selected hits in a cell. The distance from readout is labeled W , an alias for either X or Y, such that $W = 0$ is the center of a cell and positive values of W are closer to the readout. From this information, individual histograms of the average PE/cm vs W are constructed for each cell. The relative calibration procedures apply corrections to these histograms.

The first effect handled by the relative calibration is of threshold and shielding. Thresholds refer to the issue that an energy deposition may not register for a hit at all, as opposed to simply being attenuated, if there are not enough photons that reach the APD. Shielding refers to the tendency of the detector mass to alter the average signal of a minimum ionizing particle, or MIP, as a function of distance to the readout. Both of these effects would bias the set of hits used by the calibration by preferably selecting hits with greater numbers of photons, in turn underestimating the true energy deposited in the cell. To

account for this, a correction factor is applied for each cell,

$$T = \frac{PE}{\lambda} \cdot \frac{E_{True}}{E_{MIP}} \quad (5.1)$$

where T is correction factor, PE is the number of simulated photons that the electronics register, λ is the number of photons that would be seen without fluctuations, E_{True} is the true energy deposited in the scintillator, and E_{MIP} is the energy that would be deposited based only on the particle path length through the cell. The ratio on the left accounts for the threshold correction since λ is only dependent on the simulated threshold level, and the ratio on the right accounts for the shielding correction since E_{MIP} is only dependent on the path length. Two dimensional histograms of the correction factor as a function of the cell number and distance from the readout are made for each detector and view, then fit with a polynomial to remove noise. These histograms are then used to correct the corresponding data and MC. Figure 5.1 shows examples of this correction factor used for the FD.

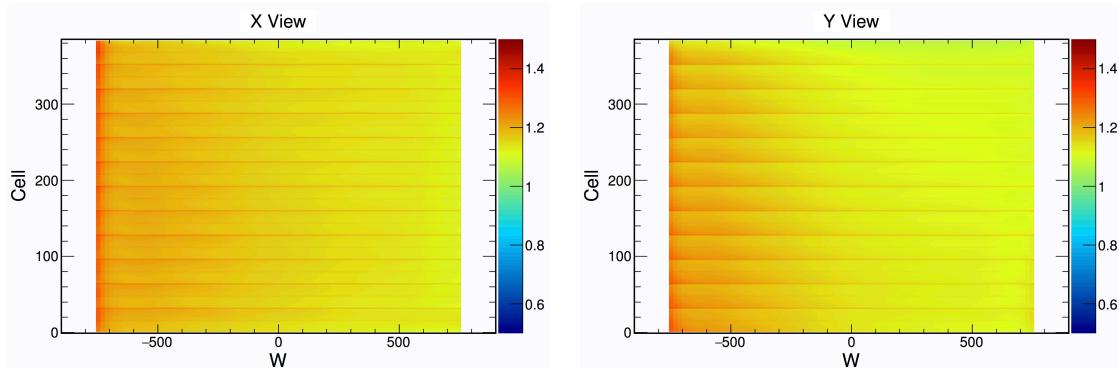


Figure 5.1: Correction factor for threshold and shielding effects at the FD as a function of cell number and distance from electronic readout. Cells in the X view are shown on the left, Y view on the right.

The next part of the relative calibration is the general form of the attenuation correction. Tricell hits are grouped by cell and fit to a double exponential to consider both short and long path light. For data, individual fits are performed for every single cell. In MC, a fit is performed on the group of all cells in a particular view and at the same location in a plane due to much lower simulated cosmic event statistics.

The double exponential has the form

$$y = C + A \left(\exp \left(\frac{W}{X} \right) + \exp \left(-\frac{L+W}{X} \right) \right) \quad (5.2)$$

where C , A , and X are the free parameters, and L is the full length of the cell. X is the cell attenuation length. The fit excludes hits in the regions closest and furthest to the readout. It includes the range $[-750, 750]$ at the FD, $[-150, 150]$ for the fully active region of the ND and for Y view muon catcher cells, and $[-150, 50]$ for X view muon catcher cells, with all numbers in cm. These central regions are chosen to exclude the significant rolloff regions at the end of the cells, which are handled differently.

The final step in the relative calibration handles the attenuation correction at the ends of the cells and the residuals from the fit above in the central part of the cells. First, both of these regions are fit with a single LOWESS curve using a tricube weight,

$$w_i = \begin{cases} \left(1 - \left| \frac{W-W_i}{\sigma} \right|^3 \right)^3 & \text{for } |W - W_i| < \sigma \\ 0 & \text{for } |W - W_i| \geq \sigma \end{cases} \quad (5.3)$$

where W is a local point on the curve, W_i is the i th neighbor of the local point, w_i is the weight on W_i , and $\sigma = 30$ cm is the range of neighbors that affect the value of W . W is then the weighted mean of W_i . Next, a simple line of the form $y = mW + c$ is fit to a collection of the neighbor points around a point, W , to give the corrected response, y . The full LOWESS curve is approximated by linear interpolation between 20 points calculated using this procedure. The full attenuation calibration combines the results of the linear interpolation with the result of the double exponential above. Figure 5.2 shows examples of the fully fit attenuation curves used in FD data calibration.

The relative calibration combines the results of all three steps to create a uniform detector response as a function of distance to readout. Figure 5.3 shows the MC before and after calibration is applied to assess the performance of the full procedure. The regions with deviations at the end of the cells ultimately are not used for the analysis.

The absolute calibration is performed after the relative calibration to convert the corrected PE signal

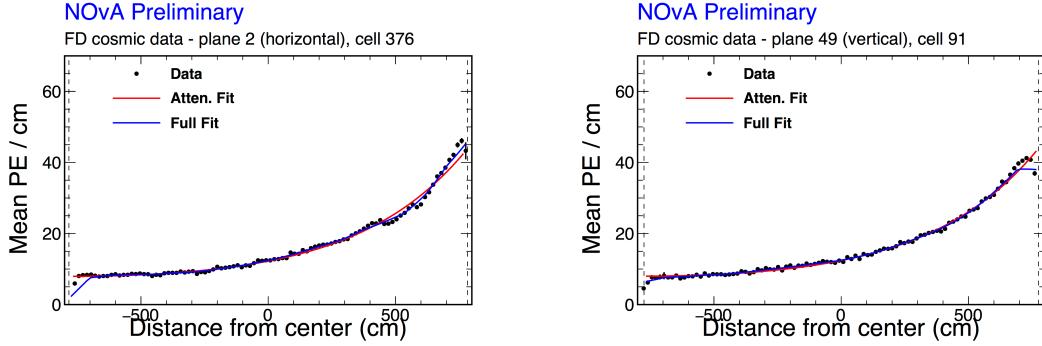


Figure 5.2: Fits to the average detector response for a cell in the FD. The red curves show just the result from the double exponential fit, and the blue curves show the full fit that also includes the LOWESS correction. The left plot is an example of a horizontal view cell, the right plot is an example of a vertical view cell.

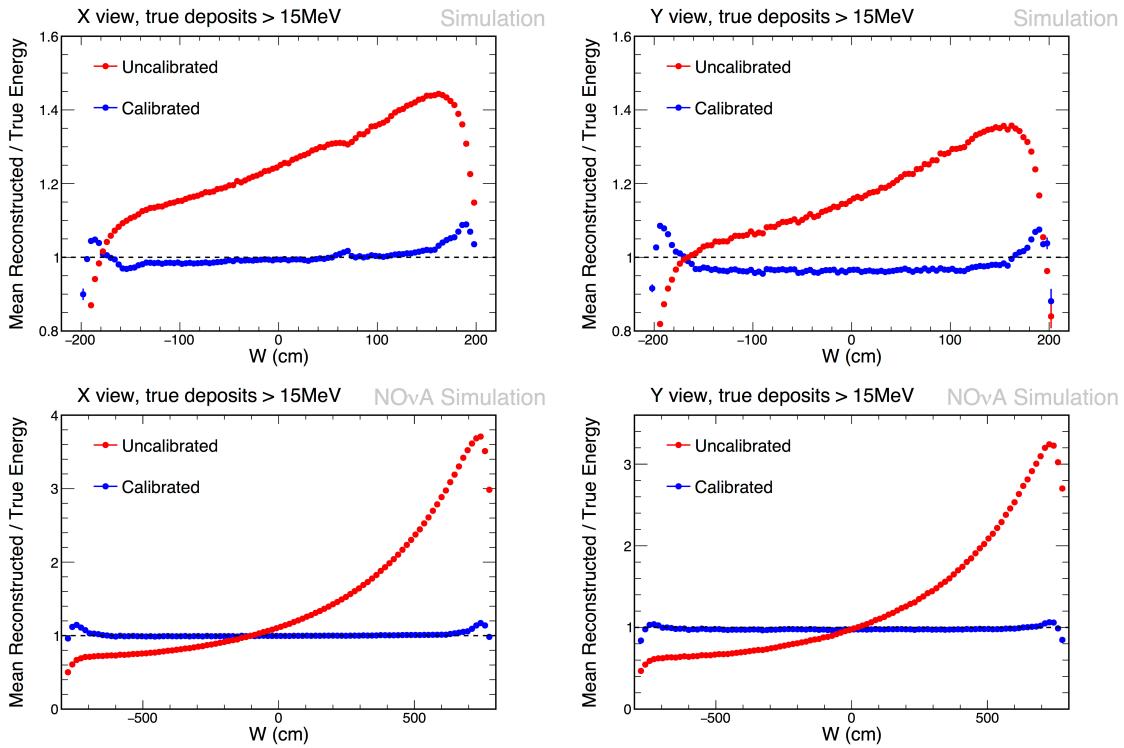


Figure 5.3: Ratios of mean reconstructed to true energy as a function of distance to readout to assess the performance of the relative calibration. The red points show the ratios before calibration; the blue points after. The left column shows calibration of the X view cells, the right column shows the Y view cells, the top row shows the ND relative calibration results, and the bottom shows the FD results.

into an energy value. Much like the relative calibration, the absolute calibration uses a sample of tricell hits from cosmic ray muons. However, in this case the energy deposited in the cell must be known in order to convert the PE signal into an energy. The energy is known either from external sources [8o] or the Bethe-Block equation [34],

$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right], \quad (5.4)$$

where $K \equiv 4\pi N_A r_e^2 m_e c^2$, r_e is the classical electron radius, $m_e c^2$ is the electron rest energy, z is the charge number of the incident particle, Z and A are the atomic number and mass of the stopping material, β and γ are calculated for the incident particle, W_{max} is the maximum energy transfer to an electron in a single collision, I is the mean excitation energy of electrons in the stopping material, and $\delta(\beta\gamma)$ is the density effect correction to ionization energy loss in the stopping material. With known energy deposition and number of PEs, determining the calibration energy scale is a simple procedure.

The particular energy used for the absolute calibration is the minimum energy deposition of muons through liquid scintillator. For NO ν A, the liquid scintillator is well approximated as chains of polyethylene, or $(C_2H_4)_n$, for which muons have a minimum average energy loss of 2.079 MeV cm²/g [8o]. The density of the scintillator was measured to be 0.8617 g/cm³ [8r], giving an average energy loss per unit length of 1.792 MeV/cm. From simulation, it was determined that this minimum energy deposition occurs between 100 – 200 cm from the end of a muon track, with only a 1.8% variation throughout that region. Figure 5.4 shows the distributions of the relative calibration corrected response as a function of distance from the muon track end for FD data and MC and the true energy deposition for the MC.

The absolute calibration energy scale is determined from distributions of tricell hits from stopping muons that occur between 100 and 200 cm from the end of the muon track. These hits are used to make one dimensional muon energy unit (MEU) distributions of the relative calibration corrected detector response for data and MC called MEU_{reco}, and the true energy deposition for MC called MEU_{truth}. The calorimetric energy scale is then taken as the mean of the MEU_{truth} distribution over the mean of the MEU_{reco} distribution. Figure 5.5 shows the distributions of the corrected detector response before the absolute calibration is applied, and the energy distributions after the absolute calibration is applied.

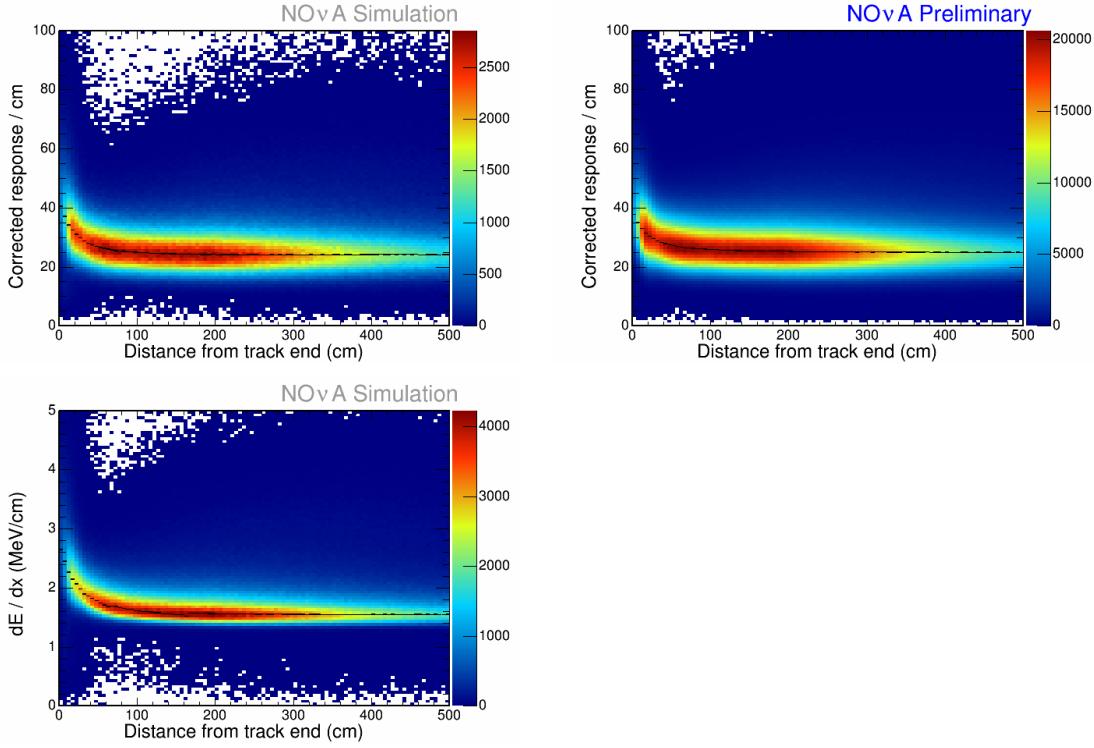


Figure 5.4: Distribution of tricell hits from cosmic ray muon tracks at the FD as a function of distance to the end of the muon track. The top row shows the response in PECorr for both data and MC. The bottom left plot shows the true energy deposited for the MC.

At this point the calibration procedure is complete. The calibration constants output by the relative and absolute calibrations are stored in a database so that the raw PE signal in any RawDigit object can be converted into an energy value.

5.2 RECONSTRUCTION CHAIN

The first step in the NC disappearance analysis is the selection of a relatively pure sample of NC events. However, the relatively raw calibrated hits are a far cry from the complete events necessary for the analysis. A set of reconstruction algorithms are thus applied to take individual hits and group them into more complete objects, often providing extra information along the way. The reconstruction procedures begin with the basic grouping of hits in spacetime, and become as complex as applying machine learning algorithms to separate events into different components. The NC disappearance analysis did not have an independent reconstruction chain from the two main NO ν A analyses. Rather, the information used

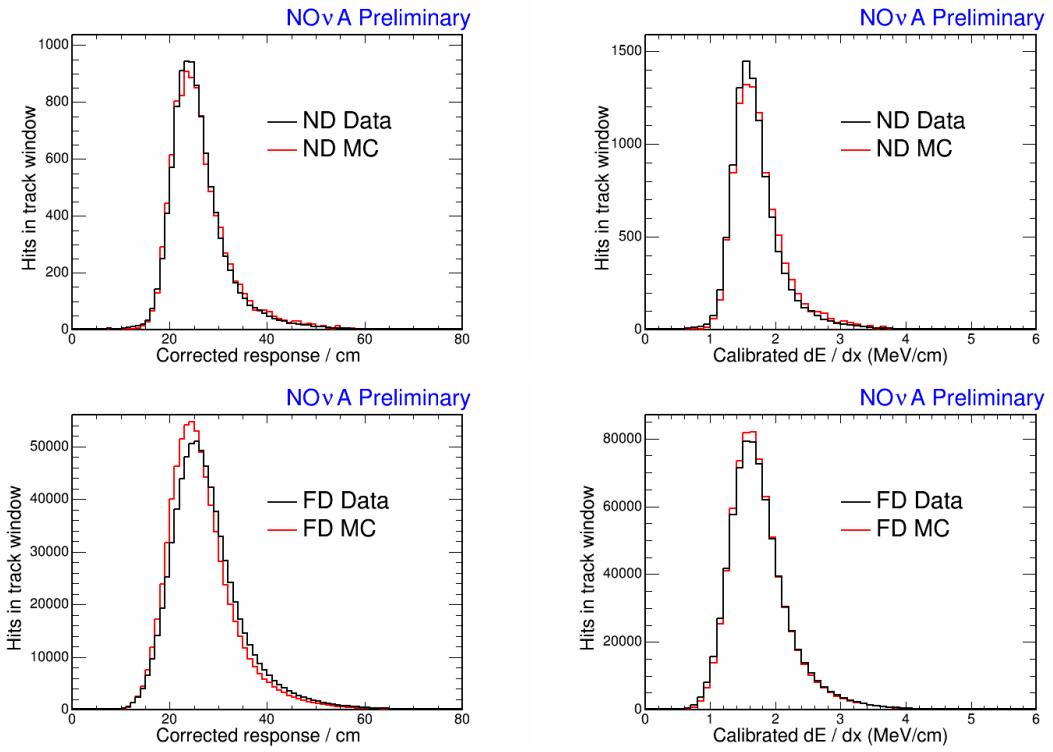


Figure 5.5: Distributions of the tricell hits from cosmic ray muons 100 – 200 cm from the end of the track. The left column shows the relative calibration corrected detector response before the absolute calibration is applied. The right column shows the absolute calibrated energy deposition. The top row shows results for the ND; the bottom shows the FD.

for the NC selection discussed in the next chapter was a mixture from the ν_e appearance and ν_μ disappearance analyses. The remainder of this section discusses the reconstruction components relevant to the analysis described in this dissertation.

The first part step in reconstruction is the clustering of hits into separate events, or slices. This is done based on the Density-Based Spatial Clustering of Application with Noise (DBSCAN) algorithm [82]. The algorithm as applied to NO ν A is described in detail in reference [83]; the main points are summarized here. The DBSCAN algorithm uses a score function to compute a distance between neighboring points. A threshold is set to determine whether two neighbors are close. Points that have more than a set number of neighbors within a that distance are labeled *core points*, and the close neighbors of the core points that are not themselves considered core points are labeled border points. Clusters are formed by iterating over the hits in an event window, computing whether a point is a core point, then expanding to the neighbors of the core point, and moving to the next cluster when the current one is entirely bounded

by border points. At the end of this procedure, slices with more than 3 hits in each view are made into a ‘physics’ slice, and any hits not assigned to a cluster are placed into a ‘noise’ slice.

The score function used between two hits in the slicing algorithm is

$$\epsilon = \left(\frac{\Delta T - |\Delta \vec{r}/c|}{T_{res}} \right)^2 + \left(\frac{\Delta Z}{D_{pen}} \right)^2 + \left(\frac{\Delta XY}{D_{pen}} \right)^2 + \left(\frac{PE_{pen}}{PE} \right)^5, \quad (5.5)$$

T_{res} is the timing resolution of the hits added in quadrature, D_{pen} is a distance penalty, PE is the number of photoelectrons in both hits added in quadrature, and PE_{pen} is a penalty on the number of photoelectrons. Hits that are in the same view vs opposite views are handled slightly differently. For hits in the same view, $\Delta \vec{r}$ is calculated in two dimensions. For hits in opposite views, $\Delta \vec{r}$ is one dimensional, ΔXY is 0, and D_{pen} is replaced with a separate, smaller, opposite view plane penalty. T_{res} was set individually for each hit during the timing calibration, discussed in reference [78]. The exponent on the PE term was set at 5 as the PE spectrum for noise falls as $PE^{-2.5}$.

The free parameters and their values are summarized in Table 5.1. The parameters were tuned and the performance of the algorithm was measured based on two metrics, completeness, the percentage of energy deposited in the scintillator from a physics interaction included in a slice, and purity, the percentage of energy in a slice that came from a single physics interaction. Values were tuned separately for the ND and FD. Notably, PE_{pen} was set to 0 for both detectors, effectively removing this term from the distance calculation.

Table 5.1: A summary of the free parameters used in the slicing algorithm, and the tuned values used for each detector.

Parameter	ND	FD
ϵ , minimum score required to be considered close neighbors	5.0	2.0
Minimum close neighbors required to be considered core point	4	4
D_{pen} , distance penalty	75.0	100.0
Opposite view plane penalty	8	4
PE_{pen} , penalty on the PE total	0	0

6

Neutral Current Event Selection

The NC disappearance analysis is performed by comparing the observed FD spectrum of NC events to the predicted number of events after extrapolation. This requires a largely pure sample of NC signal events. There are three main backgrounds for this sample, ν_e CC interactions, ν_μ CC interactions, and cosmic ray events. This chapter describes the full process for selecting a sample of NC events and eliminating backgrounds, briefly discussing data quality and event quality metrics, and fully detailing the fiducial, containment, NC selection, and cosmic rejection.

The analysis specific cuts were chosen based on the official SA files. These files were produced in tag R16-03-03, and NuS group specific DeCAF files were produced in tag S16-04-08. The cuts were trained on cosmic data from the cosmic trigger so that the actual cosmic prediction could come from the NuMI timing sideband. Beam events were oscillated assuming three flavor oscillations with the parameters listed in table 6.1. For normalization simplicity, only 14 diblock MC and data were studied. After applying this criterion, all beam spectra were scaled to 6×10^{20} POT, and the cosmic spectra were scaled to

120 s. The remainder of this note describes each level of event cut and the data/MC comparison spectra from the ND as validation.

Table 6.1: The three flavor oscillation parameters assumed for studying the NC selection cuts

Oscillation Parameter	Value
ρ	2.84 g/cm ³
Δm_{21}^2	7.53×10^{-5} eV ²
$\sin^2 2\theta_{12}$	0.846
Δm_{32}^2	2.40×10^{-3} eV ²
θ_{23}	$\pi/4$
$\sin^2 2\theta_{13}$	0.084
δ	0

6.1 DATA QUALITY

Data quality cuts were developed well before the first NO ν A results to ensure proper data taking conditions, and all analysis groups apply them as standard. These cuts are applied per spill, and spills that fail these cuts are not included in POT accounting. The cuts can be categorized into three main groups, beam quality, data quality, and timing.

The beam quality cuts were studied and set as described in reference [84]. These cuts ensure that the data was collected during acceptable conditions by cutting on metrics such as the beam position on the target, and the horn current. To be included, spills must meet all of the criteria listed in table 6.2.

Table 6.2: Beam quality cuts applied to each spill to ensure proper data taking conditions. Taken from reference [84].

Beam Quality Parameter	Minimum	Maximum
Spill POT	2.00×10^{12}	
Horn Current	-202 kA	-198 kA
Beam X and Y position on target	0.02 mm	2.00 mm
Beam X and Y width	0.57 mm	1.58 mm
Time to nearest beam spill		0.5 s

Two data quality cuts were applied to data from each detector. These cuts were motivated and set as described in [85, 86, 87], and are summarized in table 6.3. The ND observes a higher level of noise when

lights are on in the detector hall, so a cut is placed on one of the particularly noisy DCMs. Also at the ND, data is only considered good if all of the DCMs are active. A similar cut at the FD is made, but as the detector conditions were changing during commissioning, the cut is made using the LiveGeometry package that properly handles a changing detector. The other cut applied at the FD is an ‘edge metric’ that looks for tracks matching across DCM boundaries and cuts spills where the fraction of matches is too low, cutting spills where the DCMs were out of sync relatively to one another.

Table 6.3: Data quality cuts applied to each spill to ensure proper data taking conditions.

Data Quality Parameter	Detector	Metric for Spill to Pass
Number of Missing DCMs	ND	$= 0$
Lights On Effect Hit Fraction	ND	≤ 0.45
Missing DCMs from LiveGeometry	FD	$= 0$
DCM Edge Match Fraction	FD	> 0.2

Finally, a timing cut is applied to cosmic data to ensure that the data is not too close to the edge of the data taking window. For cosmic events within a given $500 \mu\text{s}$ trigger window, only events between $25 \mu\text{s} < t < 475 \mu\text{s}$ are kept.

Table 6.4 shows the number of events that pass these data quality cuts, and figure 6.1 shows the energy spectra at the ND and FD.

Table 6.4: The number of events that pass the data quality cuts, at both detectors.

Cut Level	NC	ν_μ CC	ν_e CC	Cosmic
FD:				
Data Quality	363.1	298.1	65.0	14.4×10^6
ND ($\times 10^3$):				
Data Quality	12260	86703	1085	

6.2 EVENT QUALITY

Event quality cuts are the first applied to individual events and ensure that there is enough reconstructed information to be properly analyzed and that there were no obvious reconstruction failures. As with the

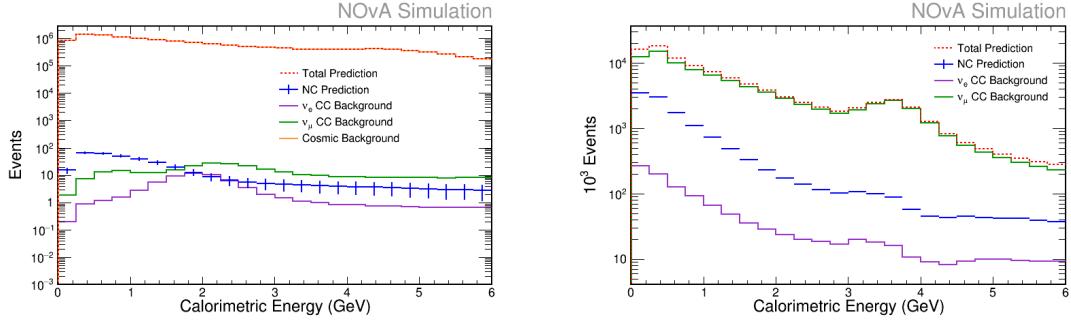


Figure 6.1: Energy spectra after Data Quality cuts for the FD (left) and ND (right).

data quality cuts above, these cuts were taken directly from the ν_e [88, 89] and ν_μ [90] analyses. Two of the cuts within this suite simply require the presence of a reconstructed vertex and a shower object. These reconstructed objects are used more extensively in later stages, so the event quality cuts make sure they are available. Other quantities considered are the number of hits per plane, distance between the reconstructed vertex and leading shower, and number of contiguous planes. Events with a high number of hits per plane are cut to remove so called ‘FEB Flashers,’ an event most often triggered by high energy cosmic rays. The events have a signature where all of the cells read out by an FEB appear to be triggered. Events with a low number of contiguous planes are most often associated with very vertical cosmic rays, so events with a low number of planes are cut. Finally, a large distance between a vertex and leading shower is considered a reconstruction failure. These cuts are summarized with the exact parameters used for the cuts in table 6.5. The number of events before and after the event quality cuts are listed in table 6.6, and figure 6.2 shows the energy spectra of the events that pass these cuts.

Table 6.5: Quality cuts applied to individual events to ensure properly reconstructed quantities.

Event Quality Metric	Metric for Event to Pass
Number of vertex objects	> 0
Number of shower objects	> 0
Number of hits per plane	< 8
Number of contiguous planes	> 2
Distance between vertex and leading shower	< 100 cm

Table 6.6: The number of events before and after application of event quality cuts, at both detectors.

Cut Level	NC	ν_μ CC	ν_e CC	Cosmic
FD:				
Data Quality	363.1	298.1	65.0	14.4×10^6
+Event Quality	226.6	121.7	59.7	0.176×10^6
ND ($\times 10^3$):				
Data Quality	12260	86703	1085	
+Event Quality	7518	45752	658	

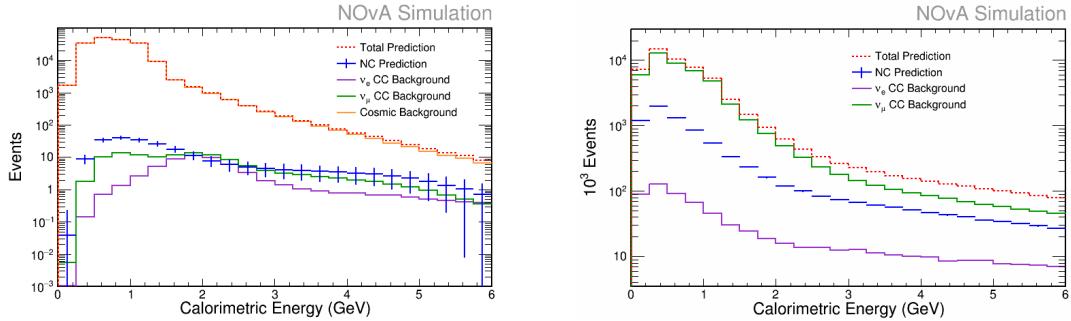


Figure 6.2: Energy spectra after Event Quality cuts for the FD (left) and ND (right).

6.3 FIDUCIAL VOLUME AND CONTAINMENT

Fiducial volume and containment cuts are applied to remove events originating outside of the detector and to ensure that events originating inside of the detector do not have activity that escapes the detector. The fiducial volume cut is a cut on the location of the reconstructed neutrino vertex. The containment cut considers the leading shower of an event and cuts the event if the start or stop point is too close to a detector wall. The specific cuts were set separately at each detector.

6.3.1 FAR DETECTOR

For the FD, the fiducial and containment cuts are the first major step in reducing the cosmic background. Since cosmic ray events originate above the detector, there are many more cosmic events with reconstructed vertices in the upper portion of the detector. The fiducial volume cut in Y is thus applied in an asymmetric way to eliminate more of the cosmic background. The fiducial volume cut on Z (the beam direction) is also asymmetric to account for cosmic rays entering the back of the detector hall where there

is a smaller amount of rock overburden to shield these events. The fiducial volume cut on X is also applied slightly asymmetrically for a slight performance gain. Events that pass these fiducial cuts must still pass the containment cuts. For the FD, the start and stop points of the leading prong must be further than 10 cm of any detector face. Figure 6.3 shows the distributions of these variables before application of the fiducial or containment cuts. Table 6.8 summarizes the event rates before and after applying fiducial and containment cuts; figure 6.4 show the energy spectra of events that pass these cuts.

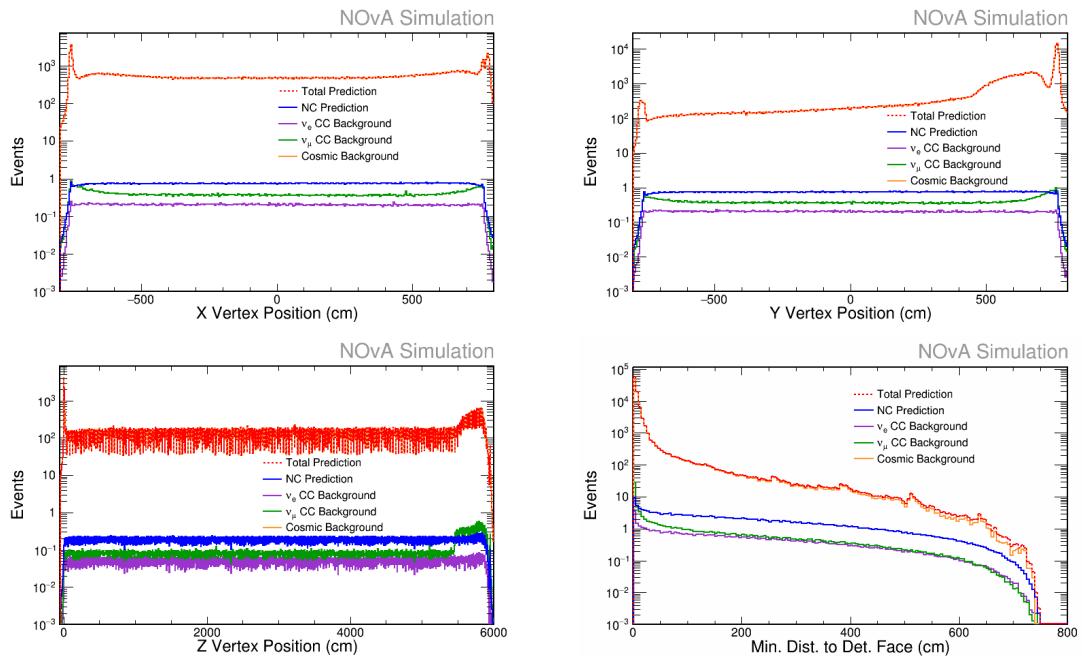


Figure 6.3: Distributions of variables used for fiducial and containment cuts at the FD. The top left, top right, and bottom left show the distributions of the reconstructed vertex. The bottom right shows the minimum distance between the start or stop of the leading prong to any detector face.

Table 6.7: Fiducial and containment cuts applied to events in the FD.

Reconstructed Quantity	Metric for Event to Pass
Reconstructed Vertex X Coordinate	$-680 \text{ cm} \leq vtxX \leq 650 \text{ cm}$
Reconstructed Vertex Y Coordinate	$-720 \text{ cm} \leq vtxY \leq 600 \text{ cm}$
Reconstructed Vertex Z Coordinate	$50 \text{ cm} \leq vtxZ \leq 5450 \text{ cm}$
Leading prong start/stop distance to detector face	$> 10 \text{ cm}$

Table 6.8: The number of events before and after application of fiducial and containment cuts at the FD.

Cut Level	NC	ν_μ CC	ν_e CC	Cosmic
...Event Quality	226.6	121.7	59.7	175900
+ Fiducial	142.2	49.2	37.2	13550
+ Containment	139.5	44.9	36.5	2295

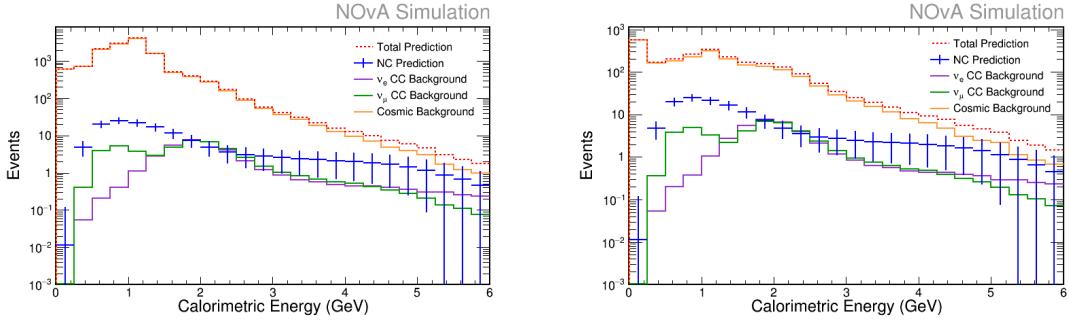


Figure 6.4: Energy spectra after fiducial and containment cuts for the FD. On the left, the cuts applied are data quality, event quality, and fiducial. The right plot also includes the containment cut.

6.3.2 NEAR DETECTOR

While the ND does not have a large cosmic background to eliminate, there are many events that originate in the rock outside of the detector that can leak into the detector. Furthermore, the small size of the ND means that many events have activity that escapes outside of the detector. The fiducial and containment cuts at the ND are designed to combat both of these effects. The fiducial cuts on the X and Y coordinates of the reconstructed vertex were applied symmetrically, with a modestly large cut to remove events that originate in the rock outside the detector. The vertex cut on Z also cuts a large portion of the detector to cut rock events that leak into the front of the detector. The containment cuts at the ND are more strict than the FD. The leading prong is required to be greater than 25 cm from each detector face, and every track in the event is also subject to this cut. For the cut on each track, the back ‘face’ of the detector is considered to be the last plane of the fully active portion of the detector, effectively cutting events with activity in the muon catcher. All of these cuts are summarized in table 6.9, table 6.10 lists the event rates before and after the cuts are applied, and figure 6.5 shows the energy spectra of events that pass these cuts.

Table 6.9: Fiducial and containment cuts applied to events in the ND.

Reconstructed Quantity	Metric for Event to Pass
Reconstructed Vertex X Coordinate	$-100 \text{ cm} \leq vtxX \leq 100 \text{ cm}$
Reconstructed Vertex Y Coordinate	$-100 \text{ cm} \leq vtxY \leq 100 \text{ cm}$
Reconstructed Vertex Z Coordinate	$200 \text{ cm} \leq vtxZ \leq 1000 \text{ cm}$
Leading prong start/stop point, distance to detector face	$> 25 \text{ cm}$
Each track start/stop point, distance to detector face	$> 25 \text{ cm}$

Table 6.10: The number of events before and after application of fiducial and containment cuts at the ND. Here, containment refers to cuts on the distance of the leading prong and each track to the closest detector face.

Cut Level	NC	ν_μ CC	ν_e CC
...Event Quality	7518	45752	658
+ Fiducial	619.3	1437.1	30.8
+ Containment	343.6	382.0	15.4

6.4 NC SELECTION

All of the cuts discussed to this point ensure that the remaining sample of events are well reconstructed, contained events. The remaining events should all be representative of their respective types, i.e., NC, ν_e CC, and ν_μ CC. The NC selection cuts are designed to select a relatively pure sample of NC events from amongst the CC backgrounds. The LID, RemID, and CVN PID distributions all have power to complete this goal, though only CVN was ultimately the only PID used. Finally, we use a cut on the number of hits for a small performance gain at the ND.

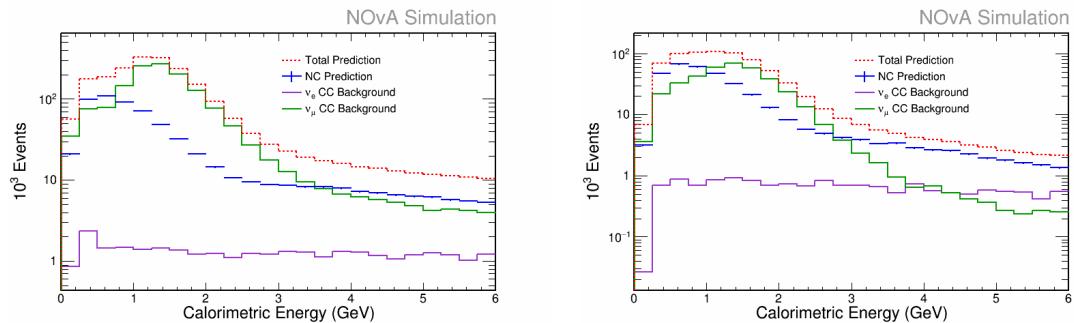


Figure 6.5: Energy spectra after fiducial and containment cuts for the ND. On the left, the cuts applied are data quality, event quality, and fiducial. The right plot also includes the containment cuts.

CVN, discussed in detail in reference [91], does not attempt to identify a single event type; rather, it gives the most likely event type among all possibilities. CVN is a convolutional neural network that only takes the raw event topology as input. The algorithm applies a series of convolutional filters to search for key event features as opposed to taking these features as input, then trains against these features. Since this PID attempts to classify all event types, cutting events that have a low NC score removes a significant majority of all background types.

RemID, discussed in detail in reference [92], is specifically designed to identify muons from charged pions, and provides separation power to remove ν_μ CC events that score highly under this PID. The PID is trained using a k -Nearest Neighbor classification on four input variables for a reconstructed track, the track energy deposition log likelihood, a scattering log likelihood, the track length, and the fraction of planes passed by the track that are uncontaminated by vertex or hadronic activity.

LID, described in reference [93], has the power to remove ν_e CC events. This PID uses an artificial neural network to test the most energetic shower against electron, muon, proton, neutron, pion (charged and neutral), and photon hypothesis. The main inputs to this PID are the longitudinal and transverse energy deposition, but it also considers the fraction of event energy contained in the primary shower, the invariant mass of all event showers, the energy contained within eight planes of the vertex, the distance between the event vertex and shower start point, and the angle of the shower away from the beam direction.

Chronologically, CVN was developed later than RemID and LID, thus these PIDs were originally used to develop the analysis. While they were replaced by CVN as the sole PID used in the analysis, they nevertheless provided a valuable check that CVN was behaving sensibly. The PID distributions for RemID and LID were studied before and after applying the cut chosen on CVN to ensure that it was removing largely the same events, and if there were any other minor performance increases to be gained. Figure 6.6 shows these distributions. Based on the results, it was decided that CVN provided all of the necessary separation power and so RemID and LID were not used in the analysis.

The distributions of CVN and the number of hits at the FD are shown in figure 6.7. Table 6.11 summarizes the cuts used, table 6.12 lists the event rates before and after the selection cuts are applied, and

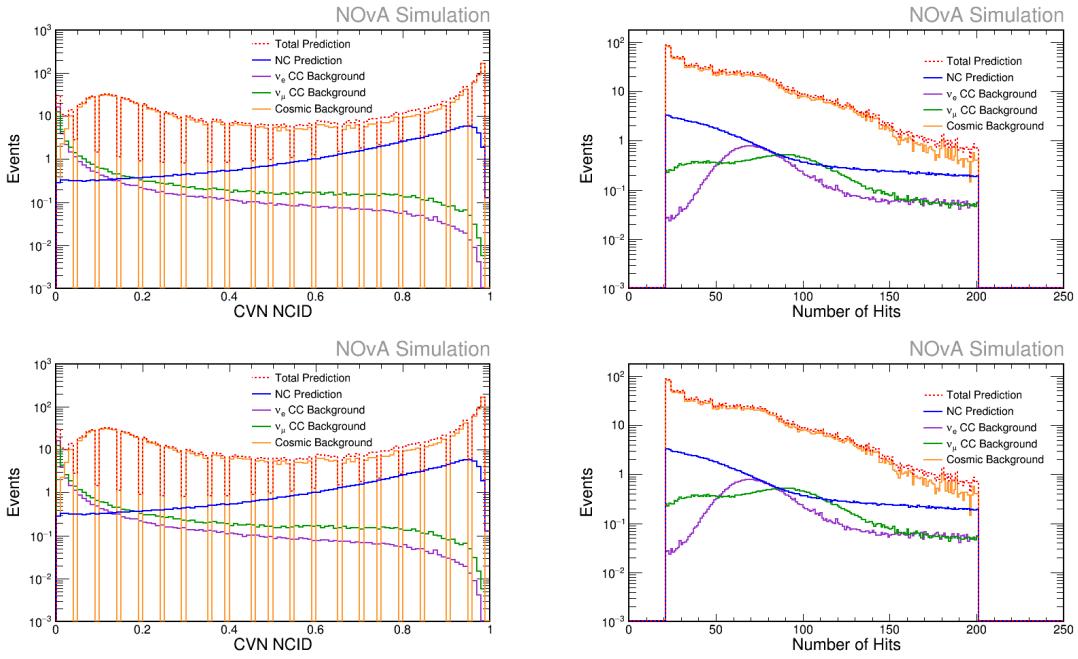


Figure 6.6: Distributions of RemID (left column) and LID (right column) at the FD before (top row) and after (bottom row) applying the cut on CVN listed in table 6.11.

figure 6.8 shows the energy distributions of events that pass these cuts.

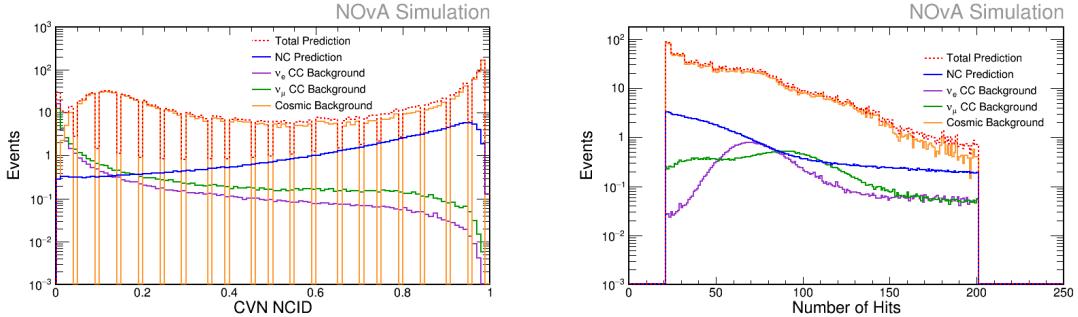


Figure 6.7: Distributions of variables used for NC selection at the FD.

6.5 COSMIC REJECTION

The final set of cuts applied to the events is cosmic rejection. While there are essentially no cosmic events at the ND, the cuts are still applied at the ND when possible to maintain similar spectra at both detectors. The main variable to discriminate against cosmics is the output of a boosted decision tree originally

Table 6.11: NC selection cuts to reject CC events, leaving a relatively pure sample of NC events.

Selection Parameter	Metric for Event to Pass
CVN NC Distribution Score	≥ 0.2
Number of Hits	≤ 200

Table 6.12: The number of events before and after application of NC selection cuts, at both detectors.

Cut Level	NC	ν_μ CC	ν_e CC	Cosmic
FD:				
...Containment	139.5	44.9	36.5	2295
+NC Selection	132.9	12.8	6.7	339.0
ND ($\times 10^3$):				
...Containment	343.6	382.0	15.4	
+NC Selection	296.9	136.0	4.2	

tuned for the ν_μ disappearance analysis [94]. This algorithm takes a number of inputs, mainly based on the reconstructed track, such as the angle from the beam, the track length, and number of hits in the track. Another powerful discriminant is the fraction of total transverse energy deposition versus total event energy deposition, originally developed for the ν_e appearance analysis [89]. The other cuts used for cosmic rejection are a cut on low average energy per hit, a harsher cut on the start/stop distance of the leading prong to the top of the detector, and removing events where the leading shower has a greatest likelihood of being a muon. The distributions of these variables at the FD are shown in figure 6.9 and table 6.13 summarizes the cosmic rejection with the values used for the actual cuts. Table 6.14 lists the event

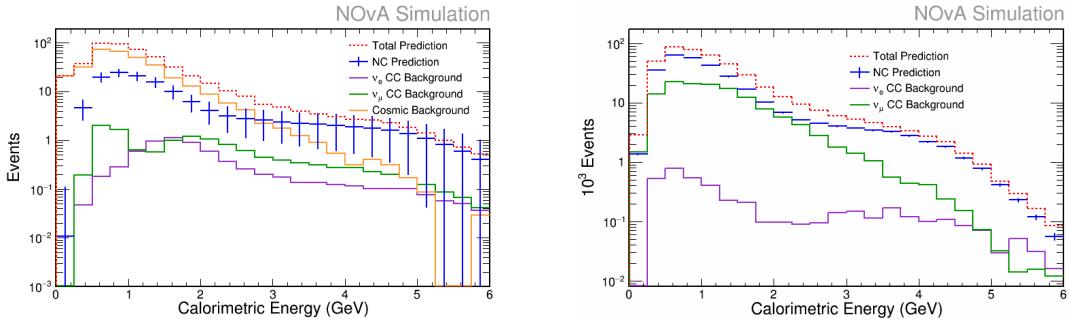


Figure 6.8: Energy spectra after NC selection cuts for the FD (left) and ND (right).

rates before and after the cosmic rejection cuts are applied, and figure 6.10 shows the energy distributions of events that pass these cuts.

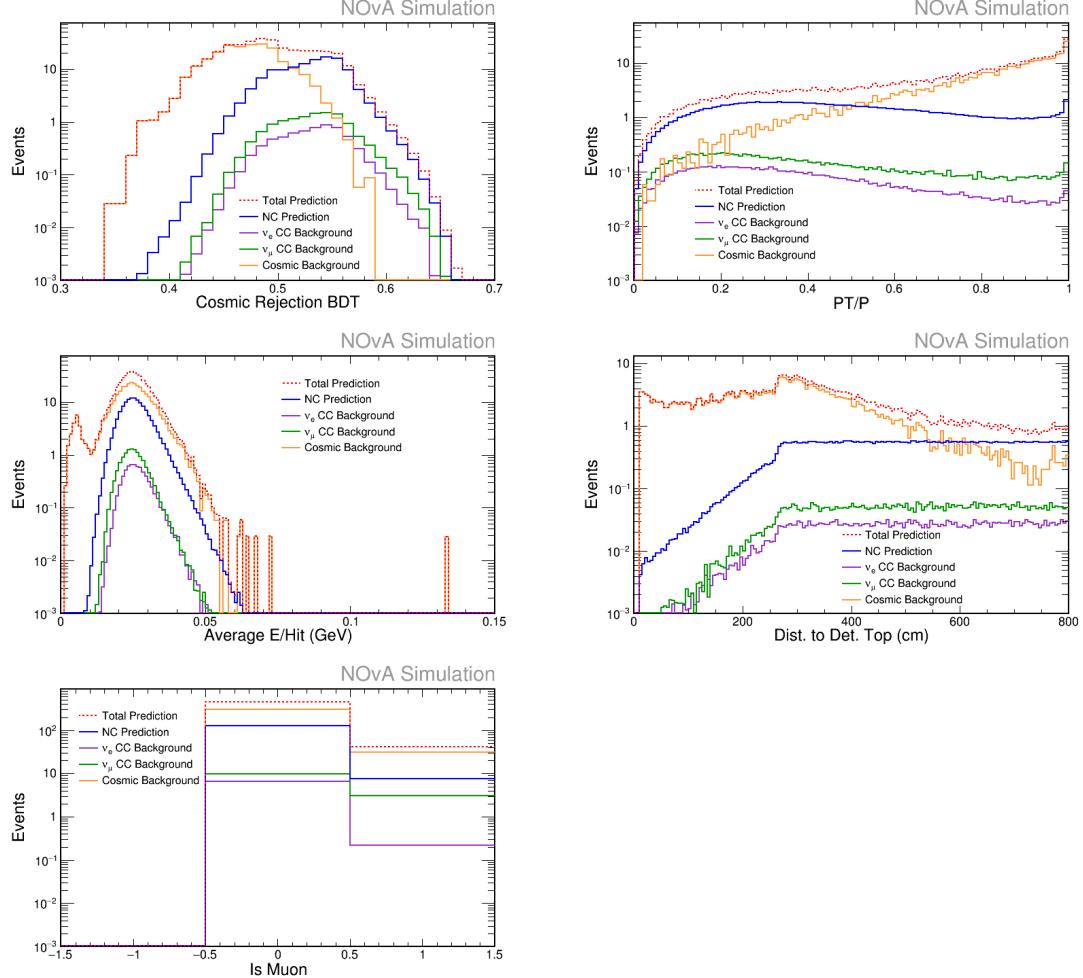


Figure 6.9: Distributions of variables used for cosmic rejection at the FD.

6.6 SUMMARY

The final selection combines all of the cuts from the previous sections. The event rates at each level of cut are summarized for the FD in table 6.15, and in table 6.16 for the ND. The final energy spectra are shown in figure 6.10.

Table 6.13: Cuts applied to reject cosmic events.

Reconstructed Quantity	Metric for Event to Pass
Cosmic Rejection BDT	> 0.5 (FD Only)
Transverse Momentum Fraction	< 0.8
Average Energy/Hit	> 18 MeV (FD) > 9 MeV (ND)
Leading Shower Best Identified as Muon	False
Leading Prong Start/Stop Distance to Top of Detector	≥ 480 cm (FD Only)

Table 6.14: The number of events after each cut level at the FD.

Cut Level	NC	ν_μ CC	ν_e CC	Cosmic
FD:				
...NC Selection	132.9	12.8	6.7	339.0
+ Cosmic Rejection	73.1	6.3	4.3	6.0
ND ($\times 10^3$):				
...NC Selection	296.9	136.0	4.2	
+ Cosmic Rejection	235.9	88.1	3.4	

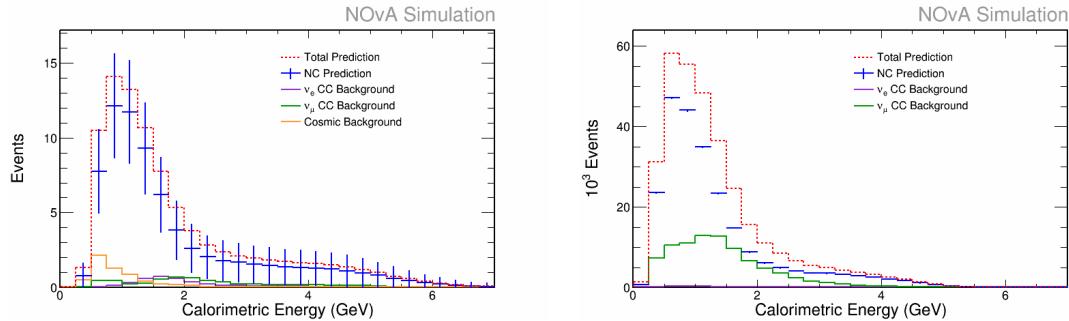


Figure 6.10: Energy spectra after all cuts for the FD (left) and ND (right).

Table 6.15: The number of events after each cut level at the FD.

Cut Level	NC	ν_μ CC	ν_e CC	Cosmic
Data Quality	363.1	298.1	65.0	14.40×10^6
+ Event Quality	226.6	121.7	59.7	175900
+ Fiducial	142.2	49.2	37.2	13550
+ Containment	139.5	44.9	36.5	2295
+ NC Selection	132.9	12.8	6.7	339.0
+ Cosmic Rejection	73.1	6.3	4.3	6.0

Table 6.16: The number of events after each cut level at the ND.

Cut Level	NC	ν_μ CC	ν_e CC
Data Quality	12260	86703	1085
+Event Quality	7518	45752	658
+ Fiducial	619.3	1437.1	30.8
+ Containment	343.6	382.0	15.4
+ NC Selection	296.9	136.0	4.2
+ Cosmic Rejection	235.9	88.1	3.4

7

Neutral Current Spectrum Prediction

Any rigorous scientific analysis must have a well defined analysis procedure. This chapter describes the bulk of that procedure for the FD NC disappearance analysis. The methods presented use the NC event selection discussed in ch. 6 to predict the FD energy spectrum.

7.1 THE CAFANA ANALYSIS CHAIN

The analysis presented in this thesis was performed in a software framework called CAFAna that analyzes Common Analysis Format files, or CAFs. The design principles of CAFAna recognize that all neutrino analyses have to perform essentially the same tasks, such as decomposing a detector spectrum into different components or applying oscillation weights to a particular event spectrum. A given task might have multiple implementations, but while the inner details are different, the number and type of end products are the same. Thus CAFAna has a fixed general analysis chain with the ability to easily swap specific implementations of any given piece of the chain.

The general analysis chain shown in fig. TODO illustrates how a near to far detector analysis can be nearly separated into smaller components that can be implemented individually. The first three components, IDecomp, IExtrap, and IPrediction, are responsible for the creation of a predicted FD spectrum. Each of these classes are pure virtual and simply declare the necessary functions for each class to hook into the next. The remainder of this chapter discusses these three classes in more detail; the classes that follow in the analysis chain are for the fitting and final results and are described in ch. 9.

Before delving into the analysis classes, it is worth briefly mentioning the basic unit of CAFAna, the Spectrum. The primary use of this class is to store the histogram of a specific variable and two methods of normalization, POT and livetime. Nearly every analysis class handles Spectrum objects in some fashion, be it storing them for internal calculations, returning them as output, or both. While the Spectrum has the ability to handle two dimensional histograms on its own, the specialized class OscillatableSpectrum stores a two dimensional histogram of a user defined variable on its X axis, and true neutrino energy on its Y axis. The function OscillatableSpectrum::Oscillated applies oscillation weights to all of the stored histogram's rows based on the true energy, then returns a one dimensional Spectrum by summing all of the contents in a column, projecting out the true energy axis.

7.2 NEAR DETECTOR DECOMPOSITION

Near detector decomposition is performed within implementations of IDecomp. The output of any decomposition is simply a set of ND spectra separated by neutrino flavor and sign; thus IDecomp declares a set of functions that each return a Spectrum object with the same X axis and contain estimates of the NC, ν_μ , $\bar{\nu}_\mu$, ν_e , and $\bar{\nu}_e$ components. Since oscillations do not necessarily affect neutrinos and anti-neutrinos equally, a split by sign must happen at some point; it was a design choice to have that occur starting with the decomposition. The most basic implementation of IDecomp, named CheatDecomp, simply fills component histograms based on MC truth information. Fig. TODO shows the inheritance structure for all of the decompositions discussed in this section.

The decomposition used for the main analysis was ProportionalDecomp, which split each bin in data based on the proportion of each event type found in the MC. Fig. TODO shows the ND decomposi-

tion based on this method. Other decompositions were used for systematic studies, including NCDecomp, NumuDecomp, NueDecompoUp, and NueDecompoDn. NCDecomp and NumuDecomp assess any data/MC difference to a single component. For NCDecomp, the ν_μ , $\bar{\nu}_\mu$, ν_e , and $\bar{\nu}_e$ components are taken directly from MC, and the difference between the sum of these components and the total data Spectrum is taken as the NC component. For NumuDecomp, it is the NC, ν_e , and $\bar{\nu}_e$ components that are taken from MC. NueDecompoUp and NueDecompoDn split the data proportionally as in ProportionalDecomp, but then fluctuate the ν_e and $\bar{\nu}_e$ components up by 100% (doubling this background) or down by 100% (removing it), respectively.

7.3 EXTRAPOLATION

Near to far detector extrapolation occurs within implementations of IExtrap.

7.4 FAR DETECTOR PREDICTION

Oscillation weights are applied within implementations of IPrediction.

8

Systematic Error Analysis

As with any analysis, the NC disappearance analysis is sensitive to a number of systematic effects. NO ν A was designed with two functionally identical detectors, so that Near Detector data can be used to constrain or correct the Far Detector prediction. Since many effects such as beam and cross section uncertainties affect the spectra at both detectors in a similar or the same way, this two detector technique leads to the reduction of these systematic errors. Other effects, such as Near Detector rock event contamination, require a data driven technique to quantify.

The general technique for analyzing systematic errors was to run the full extrapolation chain and generate a predicted spectrum with and without a systematic effect applied. As in chapter 6, the prediction assumed the same three flavor oscillation parameters listed in table 6.1, and for normalization simplicity, only 14 diblock MC and data were studied. All spectra were then scaled to 6×10^{20} POT. Each given systematic effect was used to shift the MC simulation at one or both detectors as appropriate. The resulting difference between the shifted and nominal spectra was quantified as a systematic error.

The systematic effects analyzed included uncertainties arising from the beam, GENIE simulation, Birks-Chou light yield simulation, calibration, bias from the ND containment, ND rock event contamination, ND data/MC spectrum and hadronic energy differences, noise model, MC statistics, and overall normalization. The rest of this chapter discusses each of these effects in greater detail.

8.1 BEAM

The NO ν A MC simulation involves a fully detailed model of the NuMI beam process in an attempt to create the most realistic MC possible, but systematic errors can result from any mismatch between simulation and reality. The NO ν A Beam Working Group performed studies to assess the effect that uncertainties in the simulation can have on the neutrino flux [71]. These studies included the effects of incorrectly modeling various parts of the beam transport and the effects of uncertainties in hadron production arising from fixed target experiments.

To quantify the systematic error caused by these beam uncertainties, a sample flux was generated using a systematic shift and compared to the nominal flux via a simple ratio. Results were generated separately for each neutrino flavor and for each detector. The ratios were used to modify the MC from the full simulation used as inputs to the extrapolated prediction. Finally, the shifted FD prediction was compared to the nominal prediction, with any differences being taken as the systematic error. At the end of this process, the individual errors were added in quadrature. Errors were calculated for the following systematics:

- Beam position on target varied by $\pm 0.5\text{ mm}$ in X
- Beam position on target varied by $\pm 0.5\text{ mm}$ in Y
- Beam spot size varied by $\pm 0.2\text{ mm}$ in both X and Y
- Target position varied by $+2\text{ mm}$ in Z
- Horn current varied $\pm 1\text{ kA}$
- Horn 1 position varied by $\pm 2\text{ mm}$ in both X and Y
- Horn 2 position varied by $\pm 2\text{ mm}$ in both X and Y

- Horn magnetic field changed from linear to exponential distribution
- Changes to MC beamline simulation
- Comparison between versions of FLUKA
- Comparison between FLUKA and G4NuMI
- Comparison between FLUKA and NA49

Hadron production uncertainties were combined in quadrature before being provided as weights, so this is evaluated as single systematic error. The percentage difference due to each systematic is shown in table 8.1, and the full error envelopes for NC signal and background are shown in figure 8.1. The beam systematics had an overall effect of 3.4% on the NC signal and 3.6% on the background.

Table 8.1: The percentage difference between the shifted and nominal predictions for the number of FD events due to beam systematics.

Systematic	NC Difference (%)	Background Difference (%)
Beam Position, X	0.5	0.4
Beam Position, Y	0.2	0.2
Beam Spot Size	0.2	0.3
Target Position	0.04	0.1
Horn Current	0.1	0.1
Horn 1 Position	0.6	0.8
Horn 2 Position	0.2	0.2
Horn B Field	0.05	0.05
Geometry Simulation Change	0.4	1.2
FLUKA Versions	0.04	0.4
FLUKA v G4NuMI	0.4	1.2
FLUKA v NA49	3.0	2.6
Combined	3.4	3.6

8.2 BIRKS-CHOU LIGHT YIELD SIMULATION

As discussed in section 4.2, the MC simulation employs the Birks-Chou Law (4.1) to model the relationship between scintillator light yield and energy deposition. This empirical law depends on several constants dependent on the particular scintillator material. These constants were estimated via

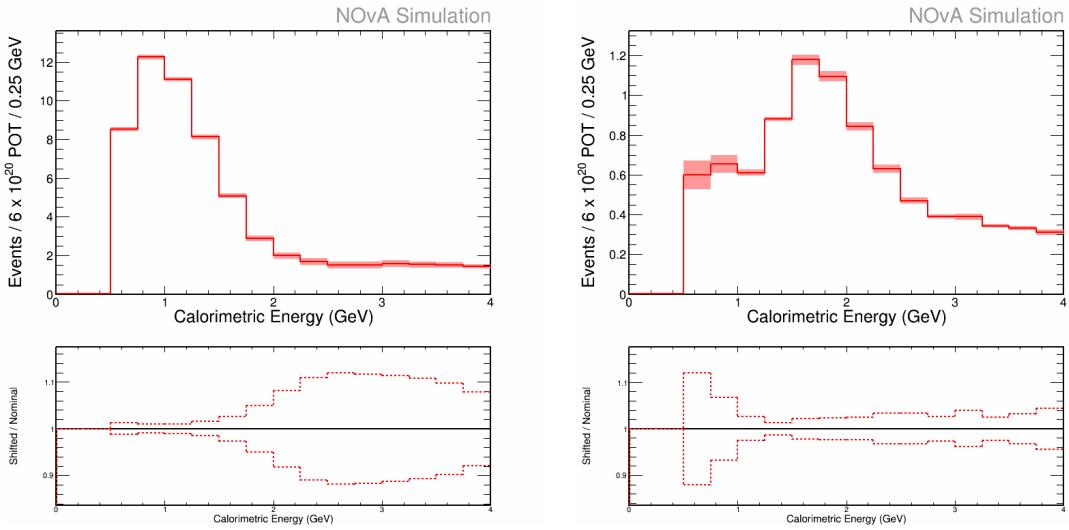


Figure 8.1: Beam systematics error envelope on the NC signal (left) and background (right) event spectra, after extrapolation. The envelope was calculated by adding in quadrature the larger of $| + 1\sigma |$ and $| - 1\sigma |$ for each individual systematic.

a data driven study for NO ν A [69]. The results of the study were $k_B = 0.04 \text{ cm/MeV}$ and $k_C = -0.0005 \text{ (cm/MeV)}^2$.

The systematic error based on the Birks-Chou light yield simulation was quantified by comparing the nominal FD prediction to predicted spectra using alternative Birks-Chou parameter constants. The values reported in the study from [69] were much larger than other typical measurements, so two MC samples were generated with more traditional values, one with $k_B = 0.01 \text{ cm/MeV}$ called BirksB, the other with $k_B = 0.02 \text{ cm/MeV}$ called BirksC, both with $k_C = 0$. Shifted FD event spectra were predicted by extrapolating the same set of ND data as the nominal prediction, but using the MC with alternative Birks-Chou model parameters. The error was taken as the percentage difference between the nominal and shifted predictions. Table 8.2 shows the percentage differences from both MC samples; figure 8.2 shows the shifted event spectra compared to nominal. Instead of combining the individual errors in quadrature, the shifted sample with the larger overall difference from the BirksC model taken as the systematic error, which had a 2.4% effect on the NC signal, and a 1.8% effect on the background.

Table 8.2: The percentage difference between the shifted and nominal predictions for the number of FD events due to extrapolation using MC with alternative Birks-Chou model parameters.

Model	NC Difference (%)	Background Difference (%)
BirksB	0.5	1.2
BirksC	2.4	1.8

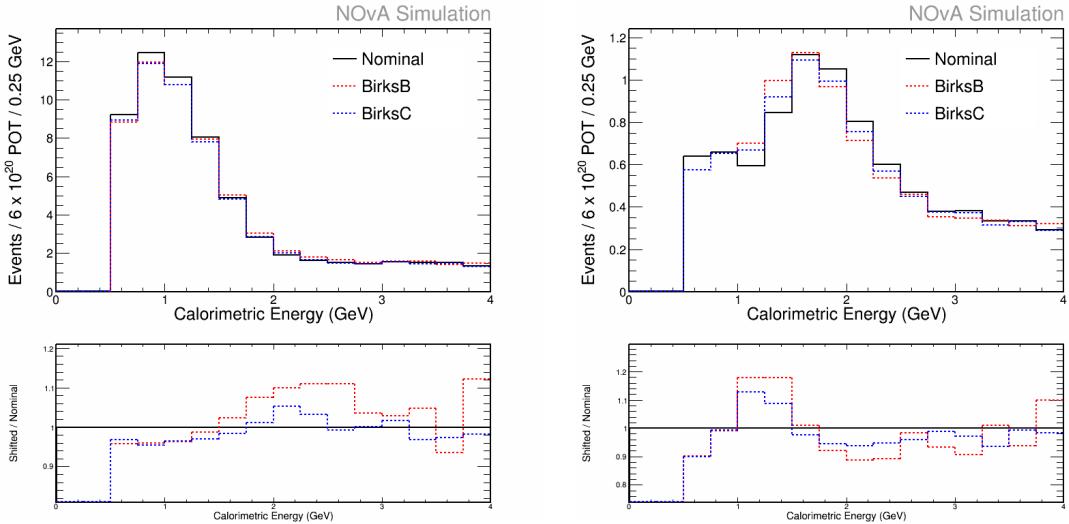


Figure 8.2: Shifted FD predictions due to extrapolation of ND data with MC using different Birks-Chou parameter values. The NC signal spectrum is on the left; the background spectrum is on the right.

8.3 CALIBRATION

The calibration procedure is designed to make a constant energy response both across each detector and between the two, but any problem can introduce a systematic error. This error was evaluated by studying various MC samples with an engineered miscalibration. The effects studied included a miscalibration that varied as a function of the cell length and an overall scale miscalibration. The functional miscalibration was studied separately between the X and Y views of the detector. The scale miscalibration was applied as a 5% effect both up and down at each detector. These effects are motivated by studying data/MC comparisons for various samples [78, 95, 96] and were updated based on the most recent round of data [97].

The systematic error was evaluated in two regimes. The first applied the same miscalibration to the MC at both detectors, and a shifted prediction was then generated using these shifted as inputs into the

extrapolation. This procedure was followed for each type of miscalibration, and all of these systematics were added in quadrature at the end. The overall scale miscalibration was also applied as a miscalibration to a single detector to measure the effect of a missed absolute calibration. A shifted prediction was then generated as above, but with this procedure only the maximum overall effect was used as the systematic error that was then added in quadrature with the systematics from above. Table 8.3 shows the percentage difference for each of the calibration systematics. The effects of the calibration systematics on the NC signal and CC background spectra are shown in figures 8.3 and 8.4 when applied to both detectors and one detector, respectively. The overall error was 5.8% on the NC signal and 6.0% on the background.

Table 8.3: The percentage difference between the shifted and nominal predictions for the number of FD events due to deliberately applied miscalibrations. For the flat scale miscalibration applied to 1 detector, the largest overall effect from the flat scale down at the ND was taken as the systematic.

Miscalibration	NC Difference (%)	Background Difference (%)
Miscalibration applied to both detectors:		
Sloped X	0.5	0.5
Sloped Y	1.6	1.1
Flat Scale Up	1.6	1.2
Flat Scale Down	3.5	2.3
Miscalibration applied to one detector:		
Flat Scale Up, ND	2.2	3.8
Flat Scale Down, ND	4.3	5.5
Flat Scale Up, FD	0.6	5.6
Flat Scale Down, FD	0.6	2.4

8.4 GENIE SIMULATION

Neutrino interactions in the NO ν A simulation are generated using GENIE [63], a generator that involves a detailed physics modeling of cross sections, hadronization, and final state interactions. GENIE includes a plethora of parameters that alter individual physics input quantities, with the parameters themselves acting as systematic uncertainties, either dialing up or down a particular quantity by a standard deviation recommended by the GENIE authors.

The systematic uncertainties from physics modeling were evaluated by using the GENIE parameters

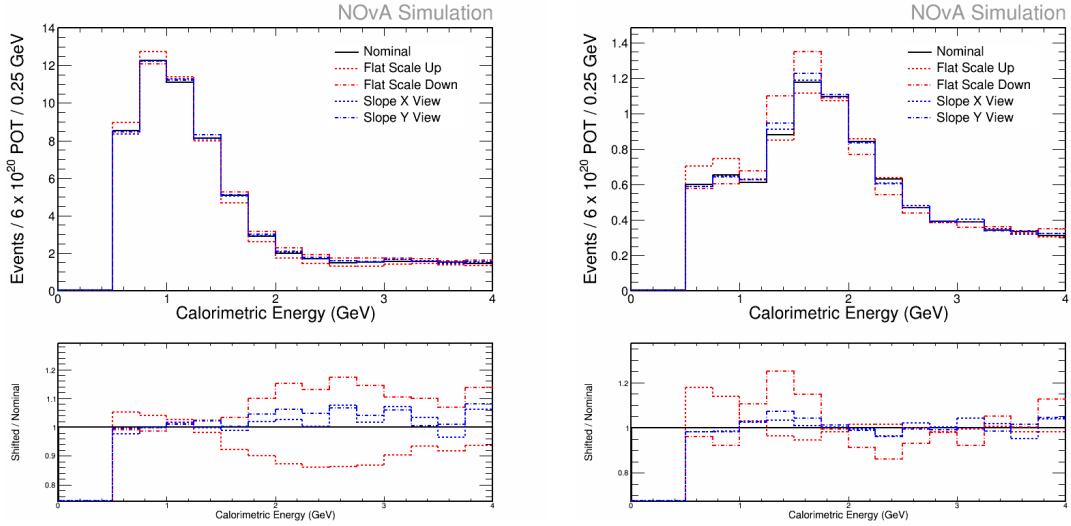


Figure 8.3: Shifted FD predictions due to miscalibration applied to both detector MC samples. The NC signal spectrum is on the left; the background spectrum is on the right.

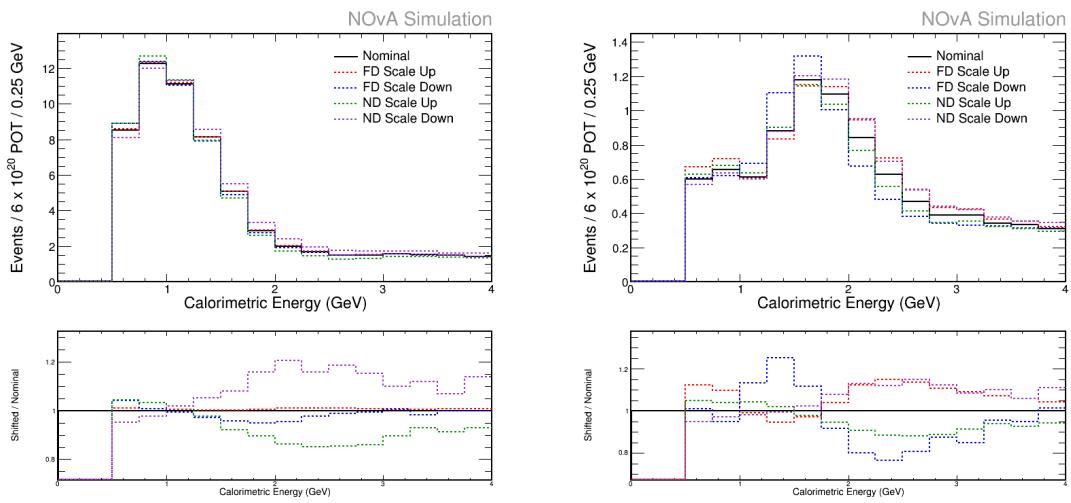


Figure 8.4: Shifted FD predictions due to miscalibration applied to one detector MC sample. The NC signal spectrum is on the left; the background spectrum is on the right.

as event reweights. Nominal MC was produced using the parameters and weights as provided by the GENIE authors, but also included a table of weights to modify an events ‘worth’ based on how much the different GENIE parameters shifted. Using the reweight table, a nominal FD prediction was produced from the default MC, and a shifted prediction was produced for each of the parameters provided by GENIE. The percentage difference was taken as the systematic error for the particular parameter. Table 8.4 lists all of the parameters considered for the systematic study, and the associated error. Figure 8.5 shows the final systematic error envelope. The combined systematic error, adding the individual contributions in quadrature, was 1.6% for the NC signal and 4.8% for the background.

Table 8.4: The systematic error, in percentage difference, for each GENIE systematic parameter. The description and standard deviations come from reference [63, 65]. Parameters with a standard deviation marked as ‘D’ are discrete and thus do not have a standard deviation estimate.

Parameter Description and Standard Deviation		NC Diff. (%)	Bkg. Diff. (%)
Axial mass for NC elastic	$\pm 25\%$	0.4	0.00
Strange axial form factor η for NC elastic	$\pm 30\%$	0.02	0.00
Normalization Factor for CCQE	$+25\%$ -15%	0.00	1.3
Axial mass CCQE distribution shape	$+25\%$ -15%	0.0	0.4
CCQE Pauli suppression (via changes in Fermi level k_F)	$\pm 35\%$	0.00	0.6
Choice of CCQE vector form factors (BBA05 \leftrightarrow Dipole)	D	0.00	0.3
Axial mass for CC resonance neutrino production	$\pm 20\%$	0.00	2.2
Vector mass for CC resonance neutrino production	$\pm 10\%$	0.00	1.1
Axial mass for NC resonance neutrino production	$\pm 20\%$	0.7	0.00
Vector mass for NC resonance neutrino production	$\pm 10\%$	0.2	0.00
Axial mass for CC and NC coherent pion production	$\pm 50\%$	0.4	0.3
Nuclear size param controlling π absorption in RS model	$\pm 10\%$	0.4	0.3
Non-resonance bkg in νp CC1 π reactions	$\pm 50\%$	0.00	0.4
Non-resonance bkg in νp CC2 π reactions	$\pm 50\%$	0.00	0.9
Non-resonance bkg in νn CC1 π reactions	$\pm 50\%$	0.00	1.9
Non-resonance bkg in νn CC2 π reactions	$\pm 50\%$	0.00	0.9
Non-resonance bkg in νp NC1 π reactions	$\pm 50\%$	0.2	0.00
Non-resonance bkg in νp NC2 π reactions	$\pm 50\%$	0.2	0.00
Non-resonance bkg in νn NC1 π reactions	$\pm 50\%$	0.2	0.00
Non-resonance bkg in νn NC2 π reactions	$\pm 50\%$	0.1	0.00
Non-resonance bkg in $\bar{\nu} p$ CC1 π reactions	$\pm 50\%$	0.00	0.1
Non-resonance bkg in $\bar{\nu} p$ CC2 π reactions	$\pm 50\%$	0.00	0.02
Non-resonance bkg in $\bar{\nu} n$ CC1 π reactions	$\pm 50\%$	0.00	0.01

Parameter Description and Standard Deviation	NC Diff. (%)	Bkg. Diff. (%)
Non-resonance bkg in $\bar{\nu}n$ CC2 π reactions	$\pm 50\%$	0.00
Non-resonance bkg in $\bar{\nu}p$ NC1 π reactions	$\pm 50\%$	0.02
Non-resonance bkg in $\bar{\nu}p$ NC2 π reactions	$\pm 50\%$	0.02
Non-resonance bkg in $\bar{\nu}n$ NC1 π reactions	$\pm 50\%$	0.02
Non-resonance bkg in $\bar{\nu}n$ NC2 π reactions	$\pm 50\%$	0.02
A_{HT} higher-twist param in BY model scaling variable ξ_w	$\pm 25\%$	0.01
B_{HT} higher-twist param in BY model scaling variable ξ_w	$\pm 25\%$	0.3
C_{V1u} u valence GRV98 PDF correction param in BY model	$\pm 30\%$	0.02
C_{V2u} u valence GRV98 PDF correction param in BY model	$\pm 40\%$	0.01
Pion transverse momentum (p_T) for $N\pi$ states in AGKY	D	0.01
Pion Feynman x (x_F) for $N\pi$ states in AGKY	D	0.01
Pion angular distribution in $\Delta \rightarrow \pi N$ (isotropic \leftrightarrow RS)	D	0.1
Branching ratio for radiative resonance decays	$\pm 50\%$	0.02
Branching ratio for single- η resonance decays	$\pm 50\%$	0.1
Nucleon mean free path (total rescattering probability)	$\pm 20\%$	0.2
Nucleon charge exchange probability	$\pm 50\%$	0.1
Nucleon elastic reaction probability	$\pm 30\%$	0.1
Nucleon inelastic reaction probability	$\pm 40\%$	0.1
Nucleon absorption probability	$\pm 20\%$	0.05
Nucleon π -production probability	$\pm 20\%$	0.04
π mean free path (total rescattering probability)	$\pm 20\%$	0.4
π charge exchange probability	$\pm 50\%$	0.1
π elastic reaction probability	$\pm 10\%$	0.1
π inelastic reaction probability	$\pm 40\%$	0.6
π absorption probability	$\pm 20\%$	0.6
π π -production probability	$\pm 20\%$	0.03
MEC event scale	$+50\%$	0.00
RPA event weight	Varied	0.00
Combined		1.6
		4.8

8.5 ND CONTAINMENT

The ND does not see an effective point source of neutrinos like the FD due to its proximity to the beam source. As a result, the neutrino flux is not uniform across the detector, and so the energy spectrum of neutrinos seen by the two NO ν A detectors is slightly different. To study the effect this has on the ex-

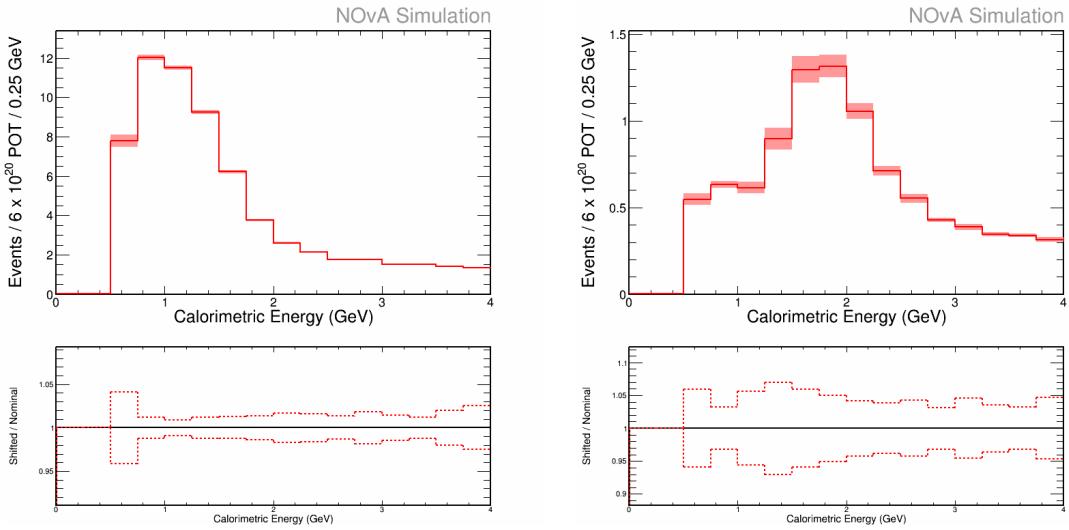


Figure 8.5: GENIE systematics error envelope on the NC signal (left) and background (right) event spectra, after extrapolation. The envelope was calculated by adding in quadrature the larger of $| + 1\sigma |$ and $| - 1\sigma |$ for each individual systematic.

trapolated prediction, multiple predicted FD spectra were generated using subsamples of the ND. The fiducial volume of the ND was split in half along each axis, and split into an inner and outer half (the overall containment criteria was left in tact). The extrapolation was performed using each of these ‘half detectors.’ The error was taken as the percentage difference from the shifted prediction to the nominal. Table 8.5 shows the results from each of these extrapolations and figure 8.6 show the shifted spectra. Like the Birks-Chou systematic, the largest overall difference was taken as the systematic error for a 1.0% effect on the NC signal and a 0.6% effect on the background.

Table 8.5: The percentage difference between the shifted and nominal predictions for the number of FD events after extrapolation using half of the fiducial volume at the ND.

ND Half	NC Difference (%)	Background Difference (%)
West (+X)	0.3	0.01
East (-X)	0.2	0.1
Top (+Y)	0.3	0.01
Bottom (-Y)	0.3	0.01
Front (Low Z)	0.1	0.3
Back (High Z)	0.1	0.4
Inner (Low $ X , Y $)	0.8	0.02
Outer (High $ X , Y $)	1.0	0.6

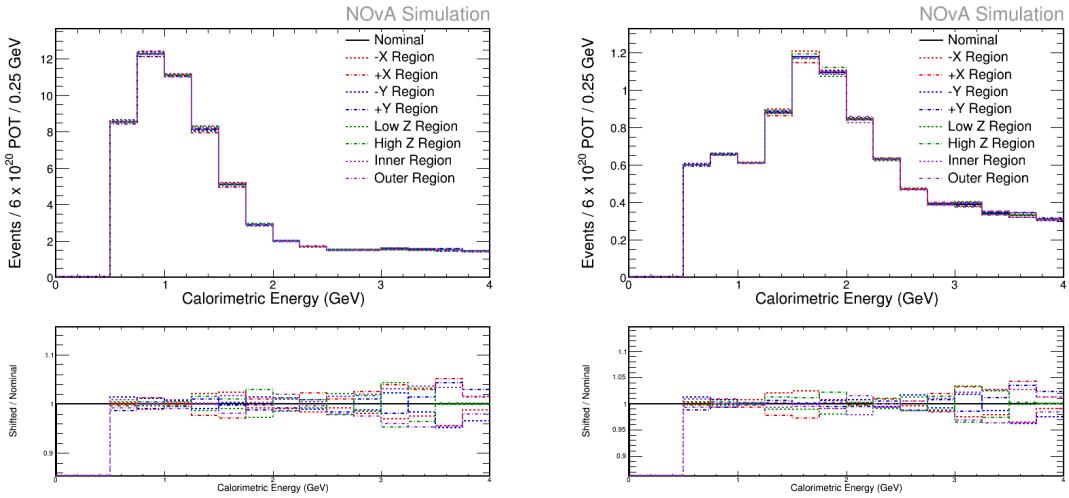


Figure 8.6: Shifted FD predictions after extrapolation using only half of the ND fiducial volume. The NC signal spectrum is on the left; the background spectrum is on the right.

8.6 ND Rock Event Contamination

The MC simulation does include neutrino interactions that occur in the rock that surrounds the ND. These events often leak into the detector volume, and while most of them are cut away by fiducial and containment cuts, there are some that remain. Those events that do remain cannot be reconstructed properly as their origins are outside of the detector.

The systematic error that is incurred due to the rock event contamination was estimated by predicting the FD event spectrum from an extrapolation with rock events removed by MC truth and comparing to the nominal predicted spectrum. The events were only removed from the ND MC sample, requiring that the true neutrino vertex was inside the detector to remain. The shifted spectra are shown in figure 8.7. This systematic amounted to an overall 4.1% shift on the NC signal spectrum and a 1.7% shift on the background spectrum.

8.7 ND DATA/MC DIFFERENCE AND CC BACKGROUND

To assess the results over the MC simulation, data was compared to MC at the ND, and the event spectra were found to be an imperfect match; see figure 8.8. The FD prediction is made by decomposing ND data into NC, ν_μ CC and ν_e CC components proportional to the amounts in the MC, so a data/MC

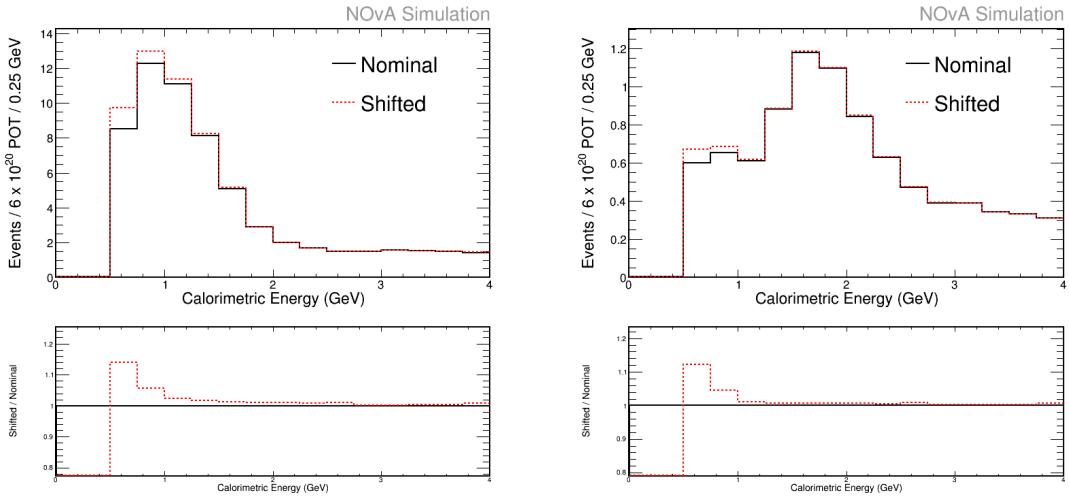


Figure 8.7: Shifted vs nominal spectra for the NC signal (left) and background (right). The shifted spectra are extrapolated after removal of rock events in the ND MC by truth.

difference can cause a systematic effect. This error was treated as a hybrid decomposition and CC background error, and as such, was quantified using a hybrid method.

For the first NO ν A ν_μ disappearance and ν_e appearance analyses, the two groups used different methods to estimate their background uncertainties. The ν_μ analysis noted that their backgrounds were small and thus placed a 100% uncertainty on each [43]. The ν_e analysis assigned any discrepancy between ND data and MC to each of the three ND background components, used each shifted decomposition to perform an extrapolation, compared the shifted predictions to nominal, and took the shift with the largest difference as the systematic error [42].

The NC disappearance analysis used a combination of the above techniques to evaluate the ND data/MC and CC background systematic error. The background from the beam ν_e component was calculated to be small, so this component was allowed to vary by $\pm 100\%$. The shifted spectrum was used for extrapolation, and the percentage difference between the nominal and shifted predictions was taken as the error. For both the NC and ν_μ CC components, the error was taken following the same procedure as the ν_e analysis. Results from both these two studies were then combined in quadrature. Shifting the beam ν_e component up and down produced the same net change in event numbers, albeit in opposite directions. Assigning the data/MC difference to the ν_μ CC component caused the larger

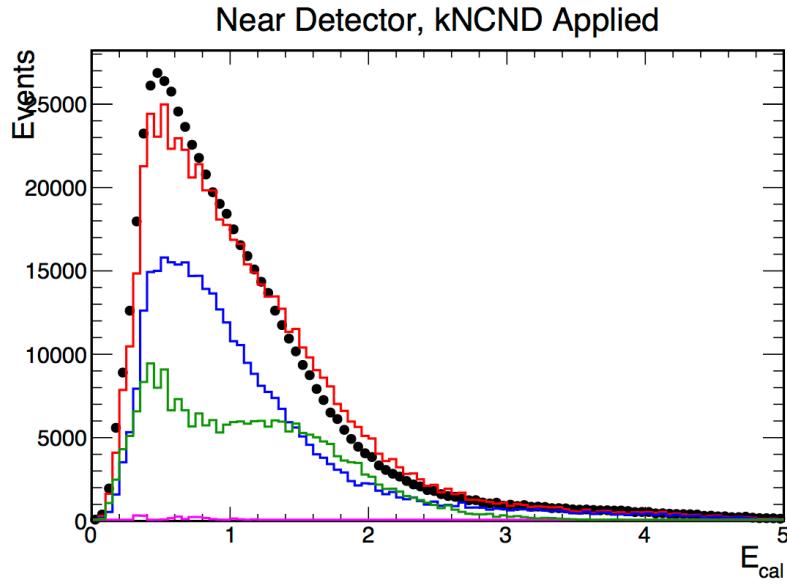


Figure 8.8: Energy spectra for data and MC at the ND, from reference [98]

overall difference, compared to assignment to the NC component. The overall error was thus calculated by combining in quadrature the errors from assigning the data/MC difference to the ν_μ CC component and shifting the beam ν_e component by $|100\%|$, for a result of a 7.0% effect on the NC signal and a 10.4% effect on the CC background. Table 8.6 summarizes the results from these studies, figure 8.9 shows the changes in the predicted spectrum due to each systematic shift.

Table 8.6: Systematic error covering ND data/MC discrepancy and CC backgrounds. The percentage differences were calculated by shifting the ND decomposition using the listed error method, performing an extrapolation, and comparing the predicted results to nominal.

Error Method	NC Difference (%)	Background Difference (%)
Beam $\nu_e +100\%$	0.00	8.2
Beam $\nu_e -100\%$	0.00	8.2
ND Data/MC difference assigned to NC component	5.2	7.7
ND Data/MC difference assigned to ν_μ CC component	7.0	6.4
Overall	7.0	10.4

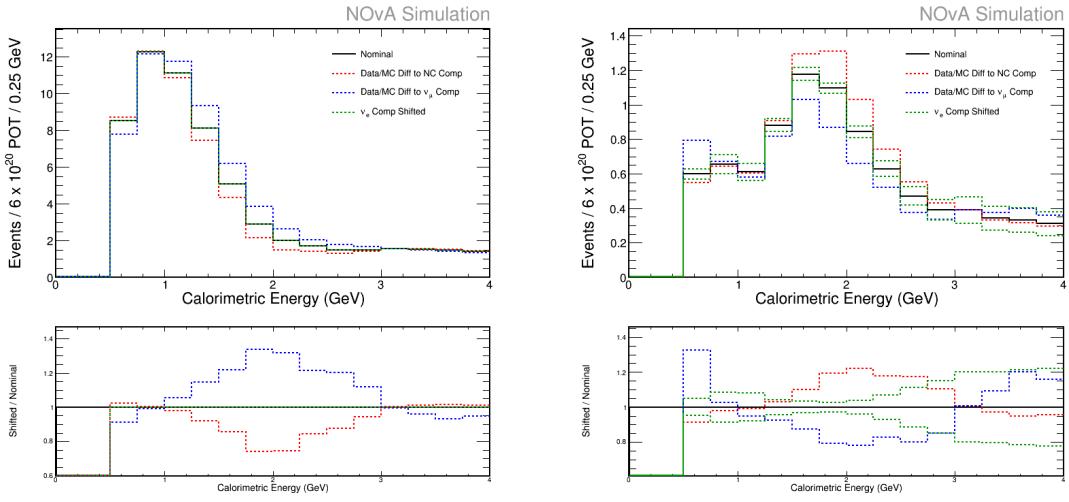


Figure 8.9: Shifted FD predictions from studies to quantify ND data/MC difference and CC background uncertainties.

8.8 MC STATISTICS

In a perfect world, there would be enough MC statistics that this section would be unnecessary, but alas, a perfect world this is not. To estimate the systematic error due to MC statistics, the MC was split into five uniformly sized samples, and each was used to extrapolate the same set of ND data. One of the resultant predicted FD spectra was labeled as nominal, and the other four were compared to this. The bin by bin differences were taken as the error between the nominal spectrum and the ‘shifted’ spectrum. To come up with an overall uncertainty, all of the errors were added in quadrature and the result was divided by the square root of the number of samples, or $\sqrt{4} = 2$. The result was a 2.0% error on the NC signal spectrum and a 4.8% error on the background spectrum.

8.9 NOISE MODEL

A sample of MC was generated for the FD using a different model of noise to study the systematic error caused by the noise model. To quantify the error, the predicted FD event spectra from the nominal and alternative MC samples were compared against each other after applying normal three flavor oscillation weights. The difference seen was 0.9% for the NC signal, and 0.7% for the background. As the shifted sample was only generated for the FD MC and the error was comparatively small, this systematic was

added into the overall normalization.

8.10 OVERALL NORMALIZATION

Several independent effects contributed to an overall normalization systematic error. A 0.5% error on the POT counting came from a small difference in the two toroids that determine the POT in a spill [71]. Uncertainties in the masses of the various parts of the NO ν A detectors contributed another error of 0.7% [99]. Finally, a study of the reconstruction efficiency between ND data and MC showed a 2.9% difference between the number of events that get fully reconstructed [100], which was taken directly as a contribution to the normalization error. These three effects and the noise systematic were combined in quadrature and constituted a systematic error of 3.2% on the NC signal and 3.1% on the background.

8.11 SYSTEMATIC ERROR SUMMARY

Table 8.7 shows a summary of all of the systematics, as well as an overall error. The overall error was calculated by summing the error from each row in quadrature. The final systematic error on the NC signal is 11.6% and the error on the background is 14.8%.

Table 8.7: A summary of the individual systematic errors for the NC disappearance analysis. The errors are percentage differences between the nominal and shifted predicted spectra.

Systematic	NC Difference (%)	Background Difference (%)
Beam	3.4	3.6
Birks-Chou	2.4	1.8
Calibration	5.8	6.0
GENIE	1.6	4.8
ND Containment	1.0	0.6
ND Rock Contamination	4.1	1.7
ND Data/MC	7.0	10.4
MC Statistics	2.0	4.8
Overall Normalization	3.2	3.1
Combined	11.6	14.8

All of the systematics discussed to this point were also calculated at the ND. Many of these, including the Beam, Birks-Chou, and GENIE, are much larger without mitigation through extrapolation. The

combined effects were studied to be certain that any data/MC discrepancy was covered under the full ND systematic error band. The data/MC spectra are shown in figure 8.10. Note that the systematic error that assesses any data/MC difference to a single component does not contribute to the full band. The technique used for this systematic did not change the total number of events since both the nominal and shifted spectra respected the normalization from the data. As the systematic error band was calculated using all events, by construction this data driven systematic effect does not have an effect on the total error band.

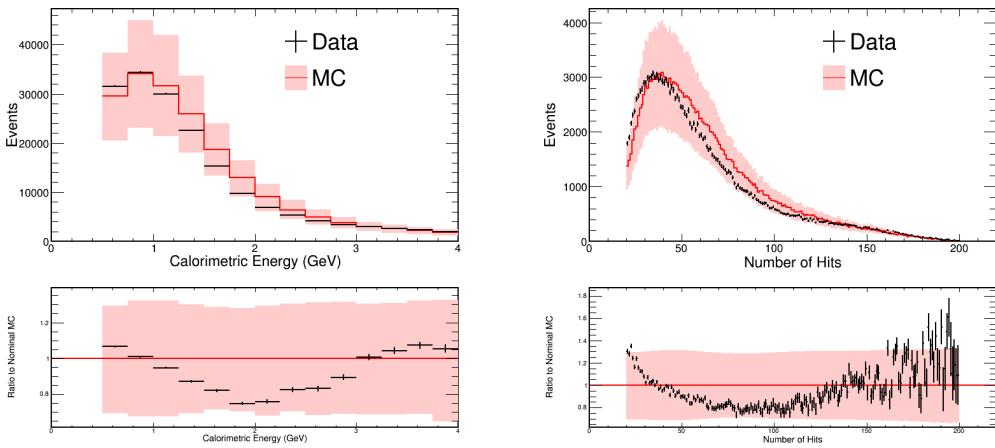


Figure 8.10: ND Data/MC spectra with full systematic error band. The calorimetric energy spectrum is shown on the left, the number of hits spectrum is shown on the right.

9

Analysis Results

The analysis presented in this dissertation was designed to make several measurements. As this dissertation was written alongside the first NO ν A NC analysis, one primary goal of the analysis was to demonstrate the successful ability of NO ν A to identify NC events and measure the ratio of observed to predicted NC events in a 3 flavor hypothesis, R_{NC} .

$$R_{NC} \equiv \frac{N_{Obs} - N_{Pred}^{Bkg}}{N_{Pred}^{NC}} \quad (9.1)$$

The other (arguably more important) goal was to start contributing to the global data on the sterile mixing parameters by extracting measurements on the mixing angles θ_{24} , θ_{34} and matrix elements $|U_{\mu 4}|^2$ and $|U_{\tau 4}|^2$.

Results in this chapter were collected between February 2014 and May 2016. For the FD, this corresponds to 6.69×10^{20} POT, or 6.05×10^{20} full detector equivalent POT. For the ND, 3.71×10^{20} POT was collected.

The first section in this chapter describes the fitting methods used to make the mixing angle and matrix element measurements. Before presenting the results, the next two sections discuss preliminary tests to demonstrate satisfactory performance of the analysis. The final section in this chapter presents the ultimate results.

9.1 FITTING METHOD

It was decided to make rate only measurements for the first NO ν A NC disappearance analysis. Thus, even though the extrapolation and prediction were performed in 250 MeV bins of calorimetric energy, only the integrated event totals were input to the fitting framework.

Fitting was performed allowing some parameters to float, some held to best fit values, and others set to 0 due to insensitivity. As mentioned above, a major analysis goal was to measure θ_{24} and θ_{34} , so these angles were allowed to float between 0° and 45° . Values outside of this range are either equivalent to this range through redefinition of the angles, or already highly disfavored through previous experiments and in a region difficult for fitting due to many local minima. θ_{23} was allowed to float with a Gaussian constraint with mean 45.8° and standard deviation 3.2° , the best fit measurement by T2K [101]. All of the remaining parameters were held fixed to the values listed in table 9.1.

Table 9.1: The parameters held fixed and the values they were set at for fitting.

Parameter	Value
Δm_{21}^2	$7.53 \times [10^{-5} \text{ eV}^2]$
Δm_{32}^2	$2.37 \times [10^{-3} \text{ eV}^2]$
$\sin^2 2\theta_{12}$	0.846
$\sin^2 2\theta_{13}$	0.085
θ_{14}	0
δ_{13}	0
δ_{24}	0
δ_{14}	0

The best fit values were found by calculating the minimum χ^2 value and allowing individual system-

atics to float with a penalty,

$$\chi^2 = 2 \left(N_{Pred} - N_{Obs} + N_{Obs} \ln \frac{N_{Obs}}{N_{Pred}} \right) + \left(\frac{\theta_{23} - \mu_{23}}{\sigma_{23}} \right)^2 + \sum_{i=1}^N \left(\frac{\sigma_i^{BF}}{\sigma_i} \right)^2, \quad (9.2)$$

where μ_{23} and σ_{23} are the mean and standard deviation used in the Gaussian constraint on the value of θ_{23} , σ_i^{BF} is the number of standard deviations the i th systematic error is shifted in the best fit, and σ_i is one standard deviation for the i th systematic error. Equation 9.2 is the standard formula for calculating χ^2 with penalty terms, where the first term is the basic value for χ^2 , and the second and third terms add a unit for each shift of θ_{23} and the systematic errors by their respective variances. For one and two dimensional χ^2 surfaces, the remaining angles not shown were profiled.

9.2 ND DATA/MC COMPARISONS

Before diving head first into the FD data, the first necessary analysis benchmark was ND data/MC comparisons. Fig 9.1 shows the event energy distributions and fig.s 9.2, 9.3, and 9.4 show the distributions for all of the variables used for selection discussed in chapter 6.

While most of the distributions differ by at most a small normalization, a few have an apparent shift, notably the energy and number of hits distributions. These particular discrepancies were actually expected due to a known mismodeling of NC interactions.

Very recently there was experimental evidence for a new ‘mode’ of interaction. A systematic scale for this event type was added to and evaluated with the GENIE systematics discussed in section 8.4. The inclusion of these effects drastically improved the data/MC agreement for the NO ν A CC analyses. Unfortunately, while the experimental data was available allowing for study of these effects on CC events, the data for NC events is largely nonexistent. Consequently the NO ν A simulation does not include these effects for NC events despite widespread belief that similar effects should be seen. This would largely explain the data/MC differences seen in the energy and number of hits distributions.

Two sanity checks were performed to ensure that the data/MC difference would not negatively impact the analysis. First, the ND data/MC energy distribution was shown with a full systematic error

NOvA Preliminary

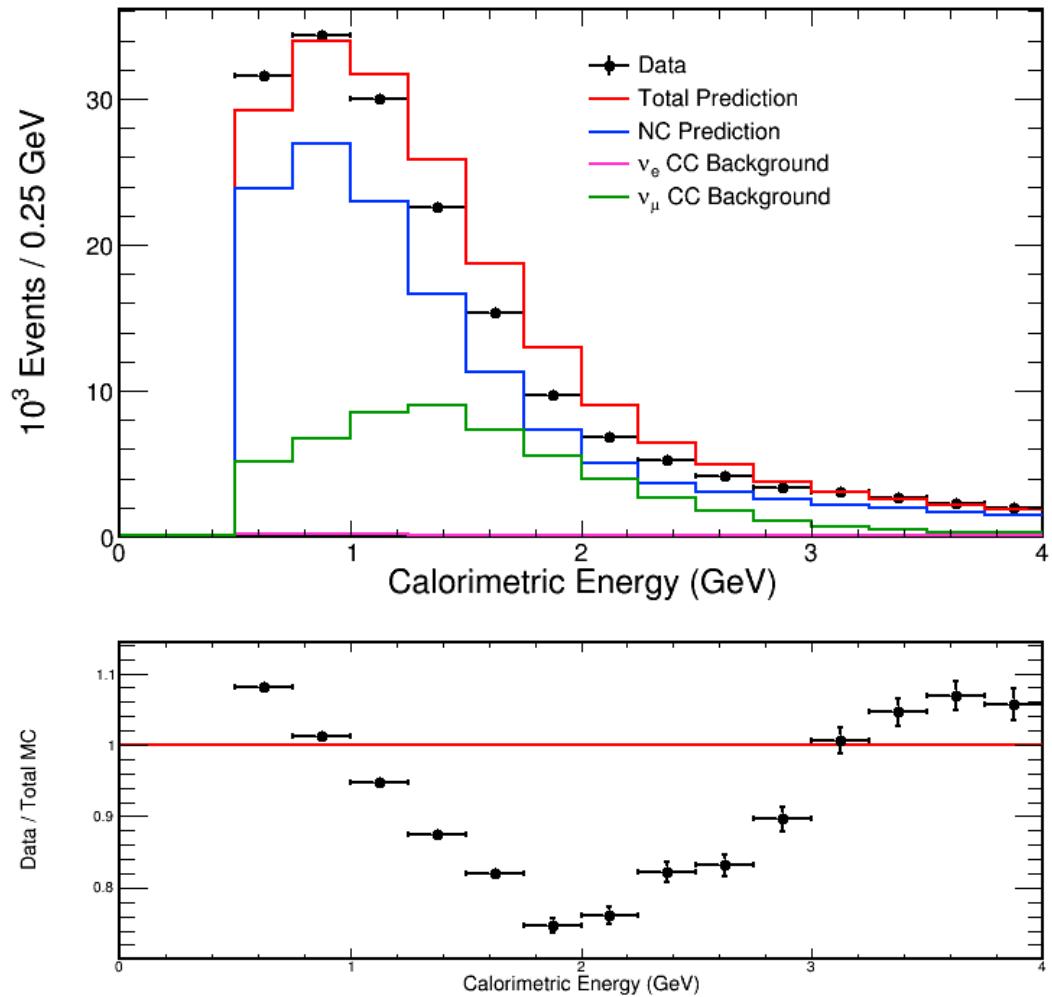


Figure 9.1: ND data/MC comparison of the calorimetric energy.

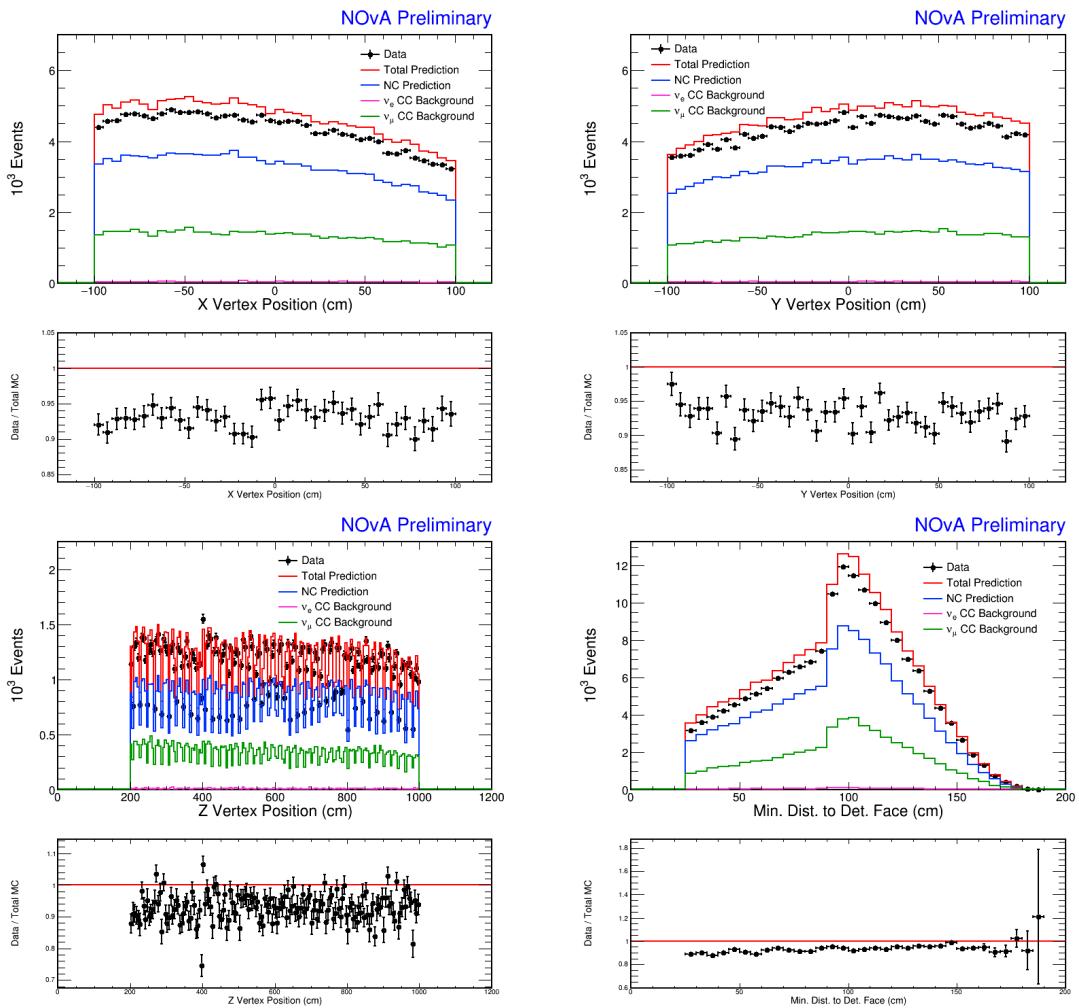


Figure 9.2: ND data/MC comparison of the reconstructed vertex position and containment distributions.

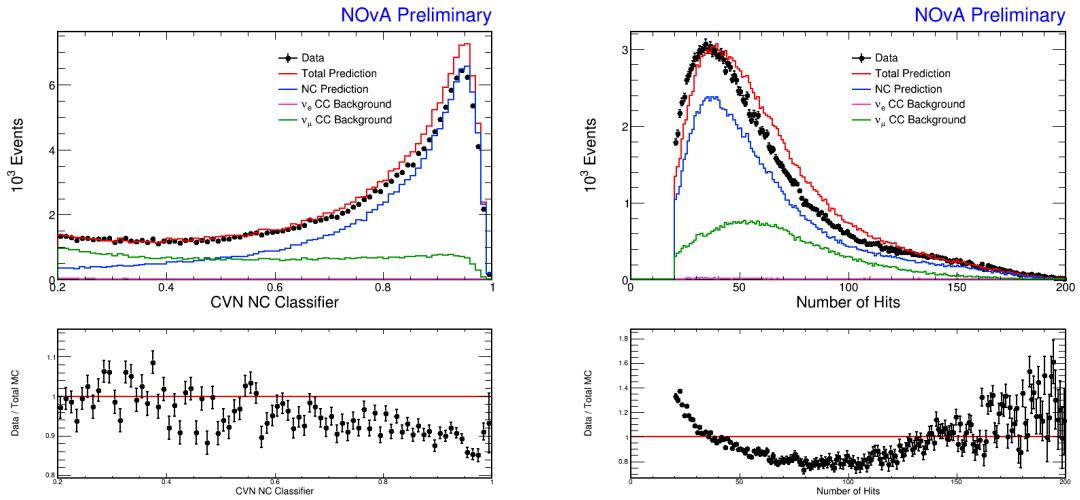


Figure 9.3: ND data/MC comparison of the CVN and number of hits distributions.

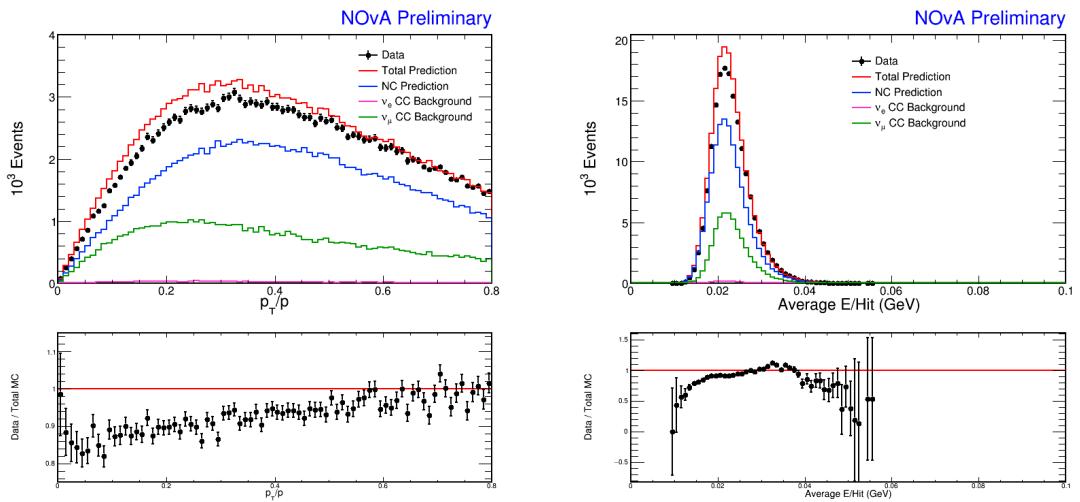


Figure 9.4: ND data/MC comparison of particle transverse momentum and energy per hit distributions.

band to show the difference is covered by systematics. This was shown in figure 8.10. Secondly, the effect of shifting either the data or MC in relation to the other on the predicted event rate and 1D angle sensitivities was studied. This was done by shifting the data up or the MC down by the ratio of the means of the energy distributions for MC and data, $1.39 \text{ GeV} / 1.33 \text{ GeV} \approx 4.5\%$ and also for a much larger shift of 10%. The event counts are shown in table 9.2 and the angle sensitivities are shown in figure 9.5. Due to the nature of a counting experiment, since the overall event rates are essentially unchanged, the effect of these rather large shifts is negligible. As a result, it was decided to take the data and MC as is and push for model improvements for future analyses.

Table 9.2: The number of predicted events at the FD after applying an energy shift to data or MC.

Shift	All	NC	ν_μ CC	ν_e CC	Cosmic	FOM
Nominal	83.5	60.6	4.6	3.6	14.3	6.633
Data 4.5% Up	84.6	61.0	4.7	3.7	14.9	6.634
Data 10% Up	86.2	61.4	4.9	3.7	15.8	6.612
MC 4.5% Down	85.1	61.8	4.8	3.7	14.3	6.702
MC 10% Down	87.2	63.4	5.2	3.8	14.3	6.792

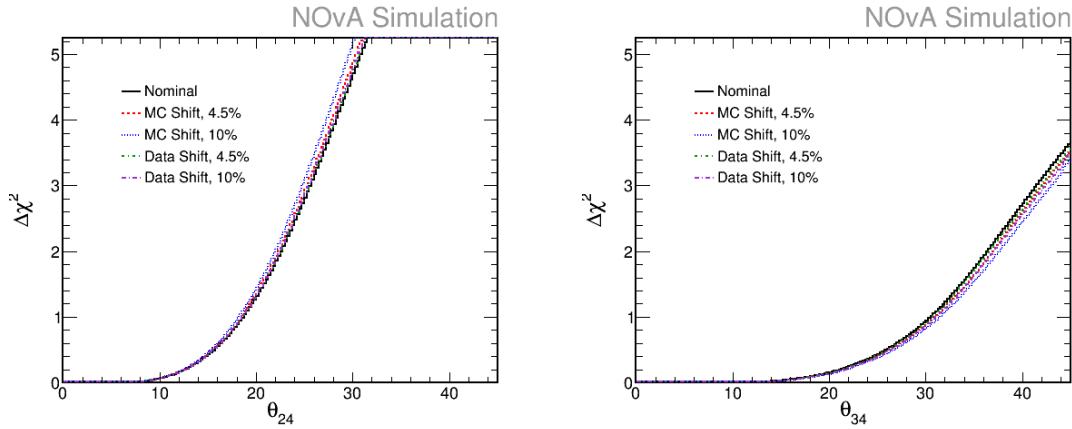


Figure 9.5: Angle sensitivities after applying an energy shift to either MC or data. Left: θ_{24} , Right: θ_{34} . The fits are statistics only and do not include any systematic errors.

9.3 SIDEBAND STUDIES

Three different sidebands were studied to test the performance of the analysis, a high energy sideband, a low CVN sideband, and mid cosmic BDT sideband. The predicted and observed event rates for each sideband are shown in table 9.3.

Table 9.3: The observed and predicted events at the FD for each sideband.

Sideband	Data	Total MC	NC	ν_μ CC	ν_e CC	Cosmic
High Energy	15	8.2	6.5	0.8	0.5	0.1
Low CVN	35	32.7	3.7	5.5	19.3	4.0
Mid BDT	17	14.5	6.0	0.5	0.4	7.5

The high energy sideband applied the standard selection and considered events between 4 and 6 GeV, a region chosen due to its high purity of NC events. The predicted and observed event distributions are shown in figure 9.6, 8.2 events were predicted and 15 ± 4 were observed in data. While the observed rate is slightly high, the low statistics means that the observation is within 2σ of the prediction. Furthermore, the discrepancy is largely driven by a single bin rather than a systematic offset. This result was thus interpreted as validation for the general analysis procedure.

The low CVN sideband considered events that fail the CVN cut, or those with a CVN NC score below 0.2. This region was used as validation for the CVN selector. The predicted and observed event distributions are shown in figure 9.7, 32.7 events were predicted and 35 ± 6 were observed in data. The overall agreement in both shape and rate for this sideband provided great confidence in the performance of CVN.

The mid cosmic BDT sideband considers events with a BDT score between 0.42 and 0.5, a region which fails the standard selection cuts but still has NC events. The predicted and observed event distributions are shown in figure 9.8, 14.5 events were predicted and 17 ± 4 were observed in data. This sideband also showed excellent agreement in both shape and rate, providing further validation for the general analysis procedure, and specifically for the main cosmic rejection variable.

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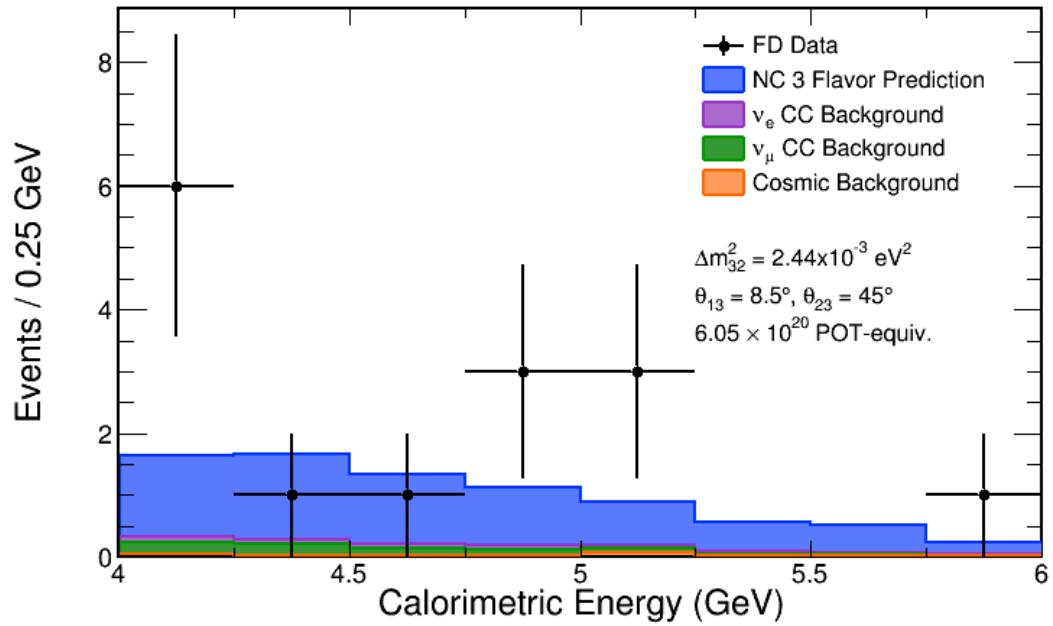


Figure 9.6: The observed and predicted FD event rates for the high energy sideband.

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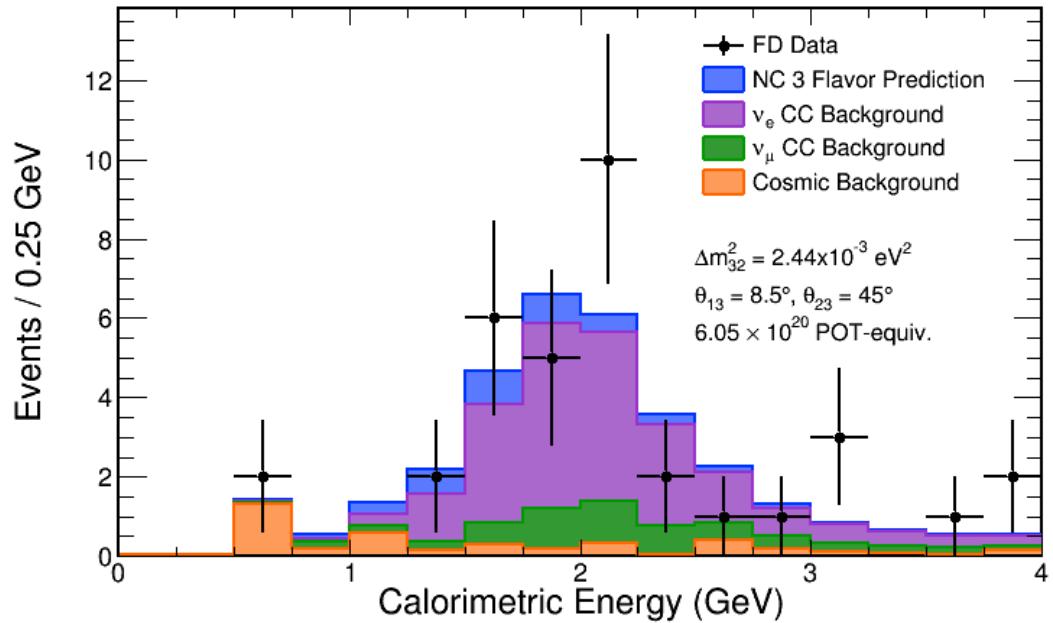


Figure 9.7: The observed and predicted FD event rates for the low CVN sideband.

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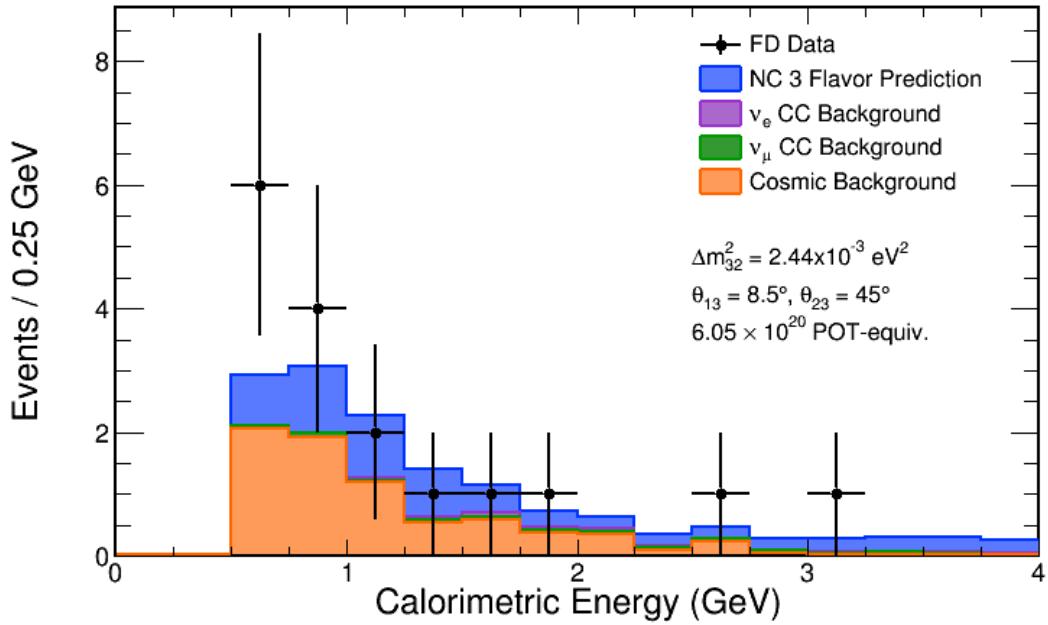


Figure 9.8: The observed and predicted FD event rates for the mid cosmic BDT sideband.

9.4 RESULTS

This analysis predicted that there would be $83.8 \pm a.b(\text{stat.})^{+c.d}_{-e.f}(\text{syst.})$ NC-like events selected in the FD data. 95 ± 9.7 were observed, corresponding to a measurement of $R_{NC} = 1.186 \pm 0.161(\text{stat.})^{+0.081}_{-0.129}(\text{syst.})$. If active neutrinos mixed with sterile neutrinos, the number of observed events would be depleted, causing R_{NC} to be less than one. Thus, the measurement of R_{NC} is consistent with the no sterile mixing hypothesis.

The energy distribution of the predicted and observed events is shown in figure 9.9. The χ^2 of this distribution is 23.31 for 14 bins, or 1.665/bin. To better assess the understanding of the observed data, data/MC comparisons distributions were studied for many of the variables used for selection, including regions cut for the analysis. These can be seen in fig.s 9.10, 9.11, and 9.12.

NOvA Preliminary

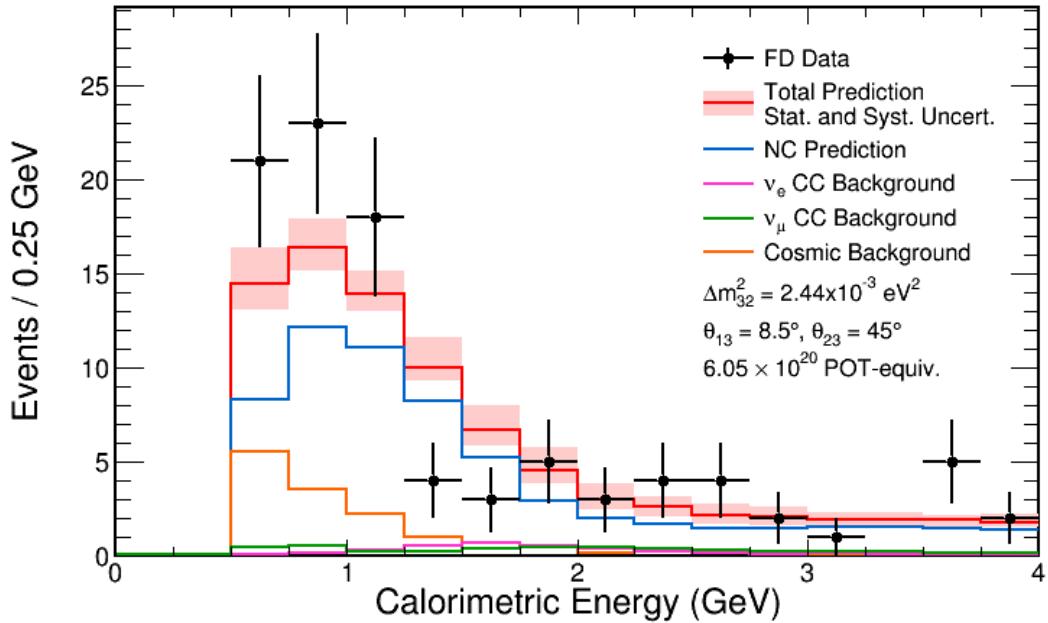


Figure 9.9: Calorimetric energy distribution of events predicted and observed in the FD.

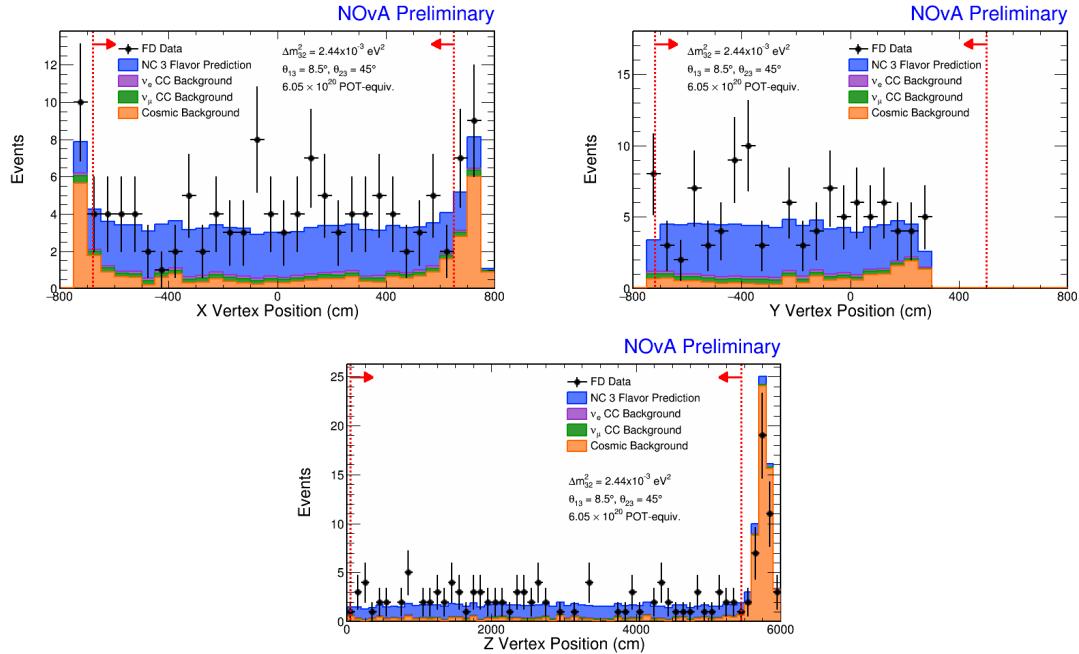


Figure 9.10: FD data/MC comparison of the reconstructed vertex position distributions. The dashed lines and arrows indicate the regions kept for the analysis.

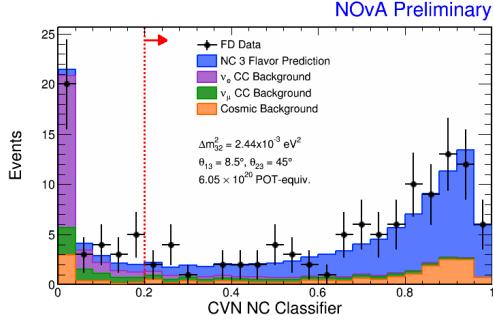


Figure 9.11: FD data/MC comparison of the CVN distribution. The dashed line and arrow indicate the region kept for the analysis.

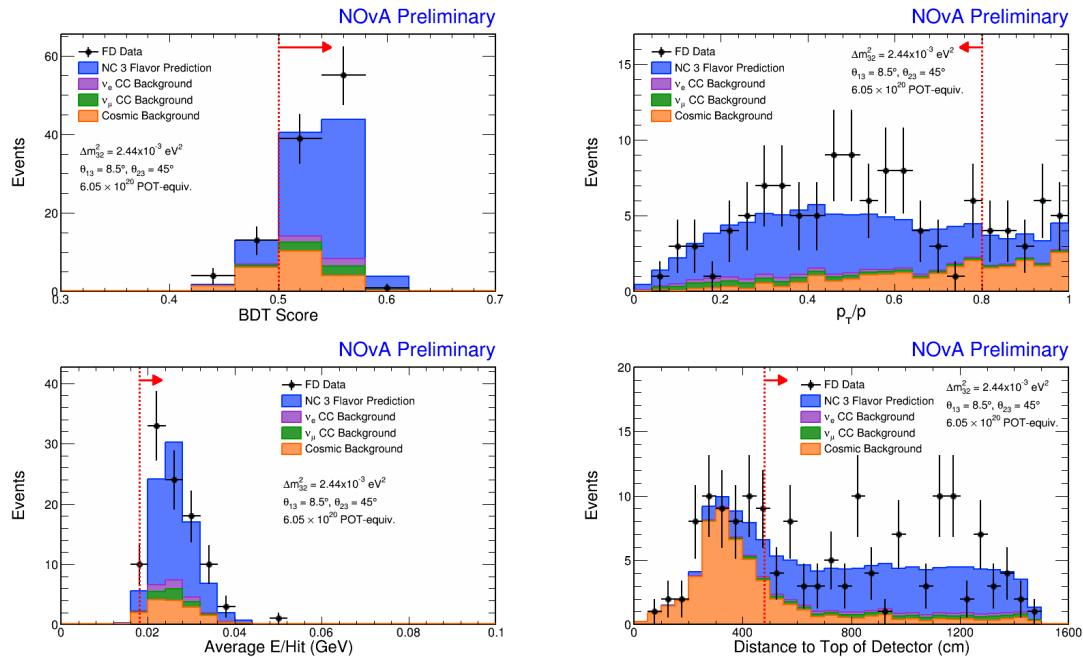


Figure 9.12: FD data/MC comparisons of the variables used for cosmic rejection. The dashed lines and arrows indicate the regions kept for the analysis.

10

Conclusions and Future Improvements

10.1 CONCLUSIONS

The results of this analysis are consistent with no sterile neutrinos.

10.2 FUTURE IMPROVEMENTS

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