

---

---

# Nuclear Reactor Thermal-Hydraulics

*NUGN520 - Homework*

---

---

By

GUILLAUME L'HER



Department of Nuclear Engineering  
COLORADO SCHOOL OF MINES

Homework submitted for the Nuclear Reactor Thermal-Hydraulics class at the Colorado School of Mines.

SPRING 2017



## TABLE OF CONTENTS

	Page
<b>1 Reactor Heat Generation</b>	<b>1</b>
1.1 . . . . .	1
1.1.1 Problem . . . . .	1
1.1.2 Solution . . . . .	1
<b>Bibliography</b>	<b>5</b>



## REACTOR HEAT GENERATION

Several exercises from the book written by M. M. El Wakil [1] are tackled in this homework. The problems in this section relate to the fifth chapter of the book, covering the subject of heat conduction in reactor elements.

## 1.1

## 1.1.1 Problem

*A 0.5-in.-diameter fuel element is made of 3 percent enriched  $UO_2$ . It is surrounded by a 0.003-in.-thick helium layer and a 0.03-in.-thick Zircaloy 2 cladding. A certain section of the element operates in boiling light water at 1000 psia. The boiling heat transfer coefficient is  $10000 \text{ Btu} \cdot \text{h}^{-1} \cdot \text{ft}^{-2} \cdot \text{F}$ , and the temperature drop in the boiling film at the section is  $30.4 \text{ F}$ . Calculate the maximum fuel temperature at that section.  $k_{He} = 0.16$ ,  $k_{clad} = 8 \text{ Btu} \cdot \text{h}^{-1} \cdot \text{ft}^{-2} \cdot \text{F}$ .*

## 1.1.2 Solution

The following assumption are made:

1.  $T_\infty$  is the temperature a certain distance from the cladding
2.  $T_\infty = 70 \text{ F}$ .
3. The radius of the fuel element is  $r_h$
4. The helium layer thickness is  $c_h$
5. The full radius of the element is  $r_z$

6. The conductivity of  $UO_2$  is taken at 800  $F$ ,  $k_f = 2.5 \text{ Btu} \cdot \text{h}^{-1} \cdot \text{ft}^{-1} \cdot \text{F}^{-1}$

Let us first compute the temperature distribution in the fuel. We know that Equation 1.1 describes the heat conduction within our usual assumptions.

$$(1.1) \quad q''' 2\pi r \Delta r L = q_{r+\Delta r} - q_r$$

We also know the relations 1.2 and 1.3.

$$(1.2) \quad q_r = -k_f A \frac{dT}{dr} = -2\pi L k_f r \frac{dT}{dr}$$

$$(1.3) \quad q_{r+\Delta r} = q_r + \frac{dq_r}{dr} \Delta r = -2\pi L k_f r \frac{dT}{dr} - 2\pi L k_f \left( r \frac{d^2 T}{dr^2} + \frac{dT}{dr} \right) \Delta r$$

Consequently, we can combine Equations 1.2 and 1.3 to obtain Equation 1.4.

$$(1.4) \quad -\frac{q'''}{k_f} = \left( \frac{d^2 T}{dr^2} + \frac{1}{r} \frac{dT}{dr} \right)$$

This equation can be written following Equation 1.5.

$$(1.5) \quad \frac{1}{r} \frac{d}{dr} \left( r \frac{dT}{dr} \right) + \frac{q'''}{k_f} = 0$$

Multiplying both sides by  $r$  and integrating, we obtain Equation 1.6.

$$(1.6) \quad r \frac{dT}{dr} = -\frac{q''' r^2}{2k_f} + C_1 = 0$$

Now dividing both sides by  $r$  and integrating again, we can compute Equation 1.7.

$$(1.7) \quad T(r) = -\frac{q''' r^2}{4k_f} + C_1 \ln(r) + C_2 = 0$$

Using the boundary conditions  $\left. \frac{dT}{dr} \right|_{r=0} = 0$  and  $T(r=0) = T_m$ , we can respectively define  $C_1 = 0$  and  $C_2 = T_m$ .

Thus, we have Equation 1.8.

$$(1.8) \quad T_m - T(r_h) = \frac{q''' r_h^2}{4k_f}$$

Now, we can solve for the temperature distribution in the Helium layer. In this layer, we can assume that no heat is generated, thus using Equation 1.9.

$$(1.9) \quad \frac{d}{dr} \left( r k_h \frac{dT}{dr} \right) = 0$$

Integrating this Equation from  $r_h$  to  $r_h + c_h$ , we can write Equation 1.10.

$$(1.10) \quad \int_{r_h}^{r_h+c_h} \frac{d}{dr} \left( r k_h \frac{dT}{dr} \right) = 0$$

Knowing Equation 1.6, we can see that  $r_h k_h d \frac{dT}{dr} = -\frac{q''' r_h^2}{2}$ , and we can consequently write Equation 1.11.

$$(1.11) \quad r k_h \frac{dT}{dr} + \frac{q''' r_h^2}{2} = 0$$

Dividing by  $r k_h$  and integrating, we obtain Equation 1.12.

$$(1.12) \quad T(r_h + c_h) - T(r_h) = \frac{q''' r_h^2}{2 k_h} \ln \left( \frac{r_h}{r_h + c_h} \right)$$

In the cladding  $Zr_2$ , the solution is identical. We can thus write Equation 1.13.

$$(1.13) \quad T(r_z) - T(r_h + c_h) = \frac{q''' (r_h + c_h)^2}{2 k_z} \ln \left( \frac{r_h + c_h}{r_z} \right)$$

And finally, we can express the heat flow to the coolant. We know that  $q''(r_z) = -k_w \frac{dT}{dr} \Big|_{r_z} = h_w (T(r_z) - T_\infty)$ . We know, using Equation 1.6 that we can write Equation 1.14.

$$(1.14) \quad \frac{q''' (r_h + c_h)^2}{2} = h_w r_z (T(r_z) - T_\infty)$$

And finally, we can write Equation 1.15.

$$(1.15) \quad T(r_z) - T_\infty = \frac{q''' (r_h + c_h)^2}{2 h_w r_z}$$

We are given the temperature drop in the coolant,  $30.4 \text{ }^\circ\text{F}$ , as well as the boiling heat transfer coefficient  $h_w = 10000 \text{ } Btu \cdot h^{-1} \cdot ft^{-2} \cdot ^\circ\text{F}^{-1}$ . Converting the radii to feet, we can calculate  $q'''$  from Equation 1.15. We obtain  $q''' = 3.309 \times 10^7 \text{ } Btu \cdot h^{-1} \cdot ft^{-3}$ . We can now plug this value back into Equations ?? ( $T(r_z) - T(r_h + c_h) = -121.8 \text{ }^\circ\text{F}$ , 1.12 ( $T(r_h + c_h) - T(r_h) = -640.6 \text{ }^\circ\text{F}$ ) and 1.8

$(T_m - T(r_h) = 1431.6 \text{ } F)$  to find the temperature  $T_m$ , which will be the maximum temperature at the section.

Consequently,  $T_m = T_\infty + 1431.6 + 640.6 + 121.8 + 30.4 = T_\infty + 2224.4 \text{ } F$ . Assuming  $T_\infty = 70$ , we obtain  $T_m = 2294.4 \text{ } F$ .

Nota bene: We did not use the fuel enrichment, nor did we use the water pressure information in our calculations. We could have used this information to compute the  $q'''$ , however the flux is unknown. In that regards, I am unsure as to the usefulness of this data.



## BIBLIOGRAPHY

- [1] M. M. EL-WAKIL, *Nuclear Heat Transport*, American Nuclear Society, 1993.
- [2] G. LEINWEBER, D. BARRY, R. BLOCK, M. RAPP, J. HOOLE, Y. DANON, R. BAHRAN, D. WILLIAMS, J. GEUTHER, AND F. SAGLIME III, *Thermal total cross sections of europium from neutron capture and transmission measurements*, Transactions of the American Nuclear Society, 107 (2012), p. 1007.