

# Homework 4

ME 590

Applied CFD and Numerical Heat Transfer

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## Problem 1

### Problem Statement

Consider a straight fin with a length of  $L_x = 7.2$  cm and a thickness of  $L_y = 8$  mm. The fin material is initially at thermal equilibrium with the surroundings at a temperature of  $T_\infty = 20^\circ\text{C}$ . At time  $t = 0$  s the base of the fin is brought to a temperature of  $T_b = 30^\circ\text{C}$  and held at that temperature. Assume the heat transfer coefficient is  $h = 40$  W/m<sup>2</sup>·K and the fin is made of aluminum alloy with the following properties.

$$\rho = 2790 \text{ kg/m}^3$$

$$C_p = 883 \text{ J/kg}\cdot\text{K}$$

$$k = 168 \text{ W/m}\cdot\text{K}$$

In Example 3.1, this system was analyzed with the assumption of 1-D behavior. You will now investigate the validity of the 1-D assumption by conducting a 2-D analysis.

- a) Use the finite volume method to derive a numerical scheme for analyzing this system. Show the symbolic form of the equations for the base nodes, interior nodes, surface nodes, and corner nodes.
- b) Using a grid spacing of  $\Delta x = \Delta y = 1$  mm, determine the maximum allowable timestep to maintain stability.
- c) Generate a graph showing the centerline temperature profile for the 2-D analysis in comparison to the 1-D analysis using an equivalent grid spacing after 30 s.
- d) Discuss the validity of the 1-D assumption used in Example 3.1 based on the results of your 2-D analysis.

## Solution

a.

### Knowns:

- **Geometry:**
  - Fin length: ( $L_x = 7.2 \{cm\} = 0.072 \{m\}$ )
  - Fin thickness: ( $L_y = 8 \text{ mm} = 0.008 \text{ m}$ )
- **Initial and Boundary Conditions:**
  - Initial fin temperature: ( $T_\infty = 20^\circ\text{C}$ )
  - Base temperature at ( $x = 0$ ): ( $T_b = 30^\circ\text{C}$ )
  - Heat transfer coefficient: ( $h = 40 \text{ W/m}^2 \cdot \text{K}$ )
- **Material Properties:**
  - Density: ( $\rho = 2790 \text{ kg/m}^3$ )
  - Specific heat capacity: ( $C_p = 883 \text{ J/kg} \cdot \text{K}$ )
  - Thermal conductivity: ( $k = 168 \text{ W/m} \cdot \text{K}$ )
- **Calculated Properties:**
  - Thermal diffusivity: ( $\alpha = \frac{k}{\rho C_p}$ )

### 1. Governing Equation:

The transient heat conduction in two dimensions is governed by the heat equation:

$$\rho C_p \frac{\partial T}{\partial t} = k \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right)$$

### 2. Interior Nodes:

For an interior node ( $i, j$ ), the energy balance over the control volume is:

$$\rho C_p \Delta x \Delta y \frac{dT_{i,j}}{dt} = k \Delta y \left( \frac{T_{i+1,j} - T_{i,j}}{\Delta x} - \frac{T_{i,j} - T_{i-1,j}}{\Delta x} \right) + k \Delta x \left( \frac{T_{i,j+1} - T_{i,j}}{\Delta y} - \frac{T_{i,j} - T_{i,j-1}}{\Delta y} \right)$$

Simplifying provides:

$$\rho C_p \Delta x \Delta y \frac{dT_{i,j}}{dt} = k \Delta y \frac{T_{i+1,j} - 2T_{i,j} + T_{i-1,j}}{\Delta x} + k \Delta x \frac{T_{i,j+1} - 2T_{i,j} + T_{i,j-1}}{\Delta y}$$

Dividing both sides by (  $\rho C p \Delta x \Delta y$  ) and defining (  $\alpha = \frac{k}{\rho C p}$  ):

$$\frac{dT_{i,j}}{dt} = \alpha \left( \frac{T_{i+1,j} - 2T_{i,j} + T_{i-1,j}}{\Delta x^2} + \frac{T_{i,j+1} - 2T_{i,j} + T_{i,j-1}}{\Delta y^2} \right)$$

**Fourier numbers:**

$$Fo_x = \frac{\alpha \Delta t}{\Delta x^2}, \quad Fo_y = \frac{\alpha \Delta t}{\Delta y^2}$$

Substituting the Fourier numbers provides the **Equation for Interior Nodes**:

$$T_{i,j}^{n+1} = T_{i,j}^n + Fo_x (T_{i+1,j}^n - 2T_{i,j}^n + T_{i-1,j}^n) + Fo_y (T_{i,j+1}^n - 2T_{i,j}^n + T_{i,j-1}^n)$$

### 3. Top Surface Nodes:

At the top surface ( $y = L_y$ ), the convection boundary condition is:

$$-k \left( \frac{\partial T}{\partial y} \right)_{\text{boundary}} = h(T_{i,j} - T_{\infty})$$

Using a finite difference approximation for the temperature gradient:

$$\left( \frac{\partial T}{\partial y} \right)_{\text{boundary}} \approx \frac{T_{i,j+1} - T_{i,j}}{\Delta y}$$

Substituting into the boundary condition and solving for ( $T_{i,j+1}$ ) provides:

$$T_{i,j+1} = T_{i,j} - \frac{h \Delta y}{k} (T_{i,j} - T_{\infty})$$

**Biot number:**

$$Bi = \frac{h\Delta y}{k}$$

Substitute  $(T_{i,j+1})$  into the energy balance equation for  $(T_{i,j})$  provides:

$$\frac{dT_{i,j}}{dt} = \alpha \left( \frac{T_{i+1,j} - 2T_{i,j} + T_{i-1,j}}{\Delta x^2} + \frac{T_{i,j-1} - T_{i,j}}{\Delta y^2} - \frac{Bi}{\Delta y^2} (T_{i,j} - T_{\infty}) \right)$$

Discretizing in time:

$$T_{i,j}^{n+1} = T_{i,j}^n + Fo_x (T_{i+1,j}^n - 2T_{i,j}^n + T_{i-1,j}^n) + Fo_y (T_{i,j-1}^n - T_{i,j}^n) - Fo_{Bi} (T_{i,j}^n - T_{\infty})$$

Where:

$$Fo_{Bi} = \alpha \Delta t \frac{Bi}{\Delta y^2}$$

**Equation for Top Surface Nodes:**

$$T_{i,j}^{n+1} = T_{i,j}^n + Fo_x (T_{i+1,j}^n - 2T_{i,j}^n + T_{i-1,j}^n) + Fo_y (T_{i,j-1}^n - T_{i,j}^n) - Fo_{Bi} (T_{i,j}^n - T_{\infty})$$

#### **4. Top Right Corner Node:**

At corner nodes, convection occurs on two faces. Applying the same reasoning as the surface nodes, we simplify using the convection boundary conditions:

- Top face:

$$T_{i,j+1} = T_{i,j} - Bi_y (T_{i,j} - T_{\infty})$$

- Right face:

$$T_{i+1,j} = T_{i,j} - Bi_x (T_{i,j} - T_{\infty})$$

Where:

$$Bi_x = \frac{h\Delta x}{k}, \quad Bi_y = \frac{h\Delta y}{k}$$

Substitute these into the energy balance equation:

$$\frac{dT_{i,j}}{dt} = \alpha \left( \frac{T_{i-1,j} - T_{i,j}}{\Delta x^2} + \frac{T_{i,j-1} - T_{i,j}}{\Delta y^2} - \left( \frac{Bi_x}{\Delta x^2} + \frac{Bi_y}{\Delta y^2} \right) (T_{i,j} - T_\infty) \right)$$

Discretizing provides:

$$T_{i,j}^{n+1} = T_{i,j}^n + Fo_x(T_{i-1,j}^n - T_{i,j}^n) + Fo_y(T_{i,j-1}^n - T_{i,j}^n) - Fo_{Bi}^{\text{corner}}(T_{i,j}^n - T_\infty)$$

Where:

$$Fo_{Bi}^{\text{corner}} = \alpha \Delta t \left( \frac{Bi_x}{\Delta x^2} + \frac{Bi_y}{\Delta y^2} \right)$$

**Equation for Corner Nodes:**

$$T_{i,j}^{n+1} = T_{i,j}^n + Fo_x(T_{i-1,j}^n - T_{i,j}^n) + Fo_y(T_{i,j-1}^n - T_{i,j}^n) - Fo_{Bi}^{\text{corner}}(T_{i,j}^n - T_\infty)$$

**5. Base Nodes at ( $x = 0$ ):**

At the base, the temperature is maintained at ( $T_b = 30^\circ\text{C}$ ). Thus:

$$T_{i,j}^{n+1} = T_b$$

## 6. Tip Nodes at ( $x = L_x$ ):

Applying the convection boundary condition at ( $x = L_x$ ) provides:

$$T_{i+1,j} = T_{i,j} - Bi_x(T_{i,j} - T_\infty)$$

## Equation for Tip Nodes:

$$T_{i,j}^{n+1} = T_{i,j}^n + Fo_x(T_{i-1,j}^n - T_{i,j}^n) + Fo_y(T_{i,j+1}^n - 2T_{i,j}^n + T_{i,j-1}^n) - Fo_{Bi}(T_{i,j}^n - T_\infty)$$

**b.**

Grid spacing:

$$\Delta x = \Delta y = 1 \text{ [mm]} = 0.001 \text{ [m]}$$

## 1. Calculate the Thermal Diffusivity ( $\alpha$ )

$$\rho C_p = (2790, \text{kg/m}^3)(883, \text{J/kg} \cdot \text{K}) = 2,463,570, \text{J/m}^3 \cdot \text{K}$$

Then, compute ( $\alpha$ ):

$$\alpha = \frac{168, \text{W/m} \cdot \text{K}}{2,463,570 \text{ J/m}^3 \cdot \text{K}} \approx 6.821 \times 10^{-5}, \text{m}^2/\text{s}$$

## 2. Stability for Explicit Method:

For stability in an explicit 2D transient heat conduction scheme, the combined Fourier numbers in the ( $x$ ) and ( $y$ ) directions must satisfy:

$$Fo_x + Fo_y \leq \frac{1}{2}$$

Where the Fourier numbers are defined as:

$$Fo_x = \frac{\alpha \Delta t}{\Delta x^2}, \quad Fo_y = \frac{\alpha \Delta t}{\Delta y^2}$$

Since (  $\Delta x = \Delta y$  ), the range simplifies to:

$$2Fo_x \leq \frac{1}{2} \Rightarrow Fo_x \leq \frac{1}{4}$$

### 3. Solve for Maximum Allowable Timestep ( $\Delta t_{\max}$ )

Express (  $Fo_x$  ) in terms of (  $\Delta t$  ):

$$Fo_x = \frac{\alpha \Delta t}{\Delta x^2}$$

Substitute (  $Fo_x \leq \frac{1}{4}$  ) to find (  $\Delta t_{\max}$  ):

$$\frac{\alpha \Delta t_{\max}}{\Delta x^2} \leq \frac{1}{4}$$

Solve for (  $\Delta t_{\max}$  ):

$$\Delta t_{\max} \leq \frac{1}{4} \cdot \frac{\Delta x^2}{\alpha}$$

### 4. Compute ( $\Delta t_{\max}$ ):

Plugging in values provides:

$$\Delta x = 0.001, \text{m}, \quad \alpha = 6.821 \times 10^{-5} \text{ m}^2/\text{s}$$

Calculate (  $\Delta x^2$  ):

$$\Delta x^2 = (0.001, \text{m})^2 = 1 \times 10^{-6}, \text{m}^2$$

Compute (  $\Delta t_{\max}$  ):

$$\Delta t_{\max} \leq \frac{1}{4} \cdot \frac{1 \times 10^{-6} \text{ m}^2}{6.821 \times 10^{-5} \text{ m}^2/\text{s}} = \frac{2.5 \times 10^{-7} \text{ m}^2}{6.821 \times 10^{-5} \text{ m}^2/\text{s}}$$

Simplify:

$$\Delta t_{\max} \leq \frac{2.5 \times 10^{-7}}{6.821 \times 10^{-5}} \text{ [s]}$$

The maximum allowable timestep to maintain stability with the given grid spacing is approximately:

$$\Delta t_{\max} = 0.003664 \text{ seconds}$$

c.

Below are two graphs displaying the temperature field and profile at  $t = 30$  [s] exported from the submitted MATLAB script ME590\_HW4\_P1. The one-dimensional solution is a minor modification of the source code provided ME590\_Ex3\_1 and the two-dimensional solution follows the Point Jacobi Numerical Method for solving temperature profiles.

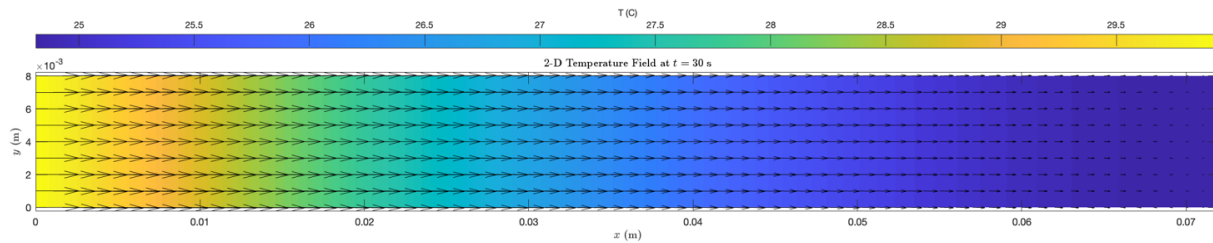


Figure 1 – Plot of 2D Temperature Field at  $T = 30s$

Figure 1 depicts the temperature field in 2 dimensions of the rectangular fin using the provided MATLAB function `plotTemperatureField2D(T, X, Y, k)`. This plot displays temperature values with a color gradient. This provides information of the uniform temperature gradient at all y-values and shows the change in temperature along the x-axis. The arrow size provides information of the direction and magnitude of the heat transfer, showing the flow of heat to the right with the largest magnitude near  $x = 0$  [m] and smallest at  $x = 0.072$  [m]



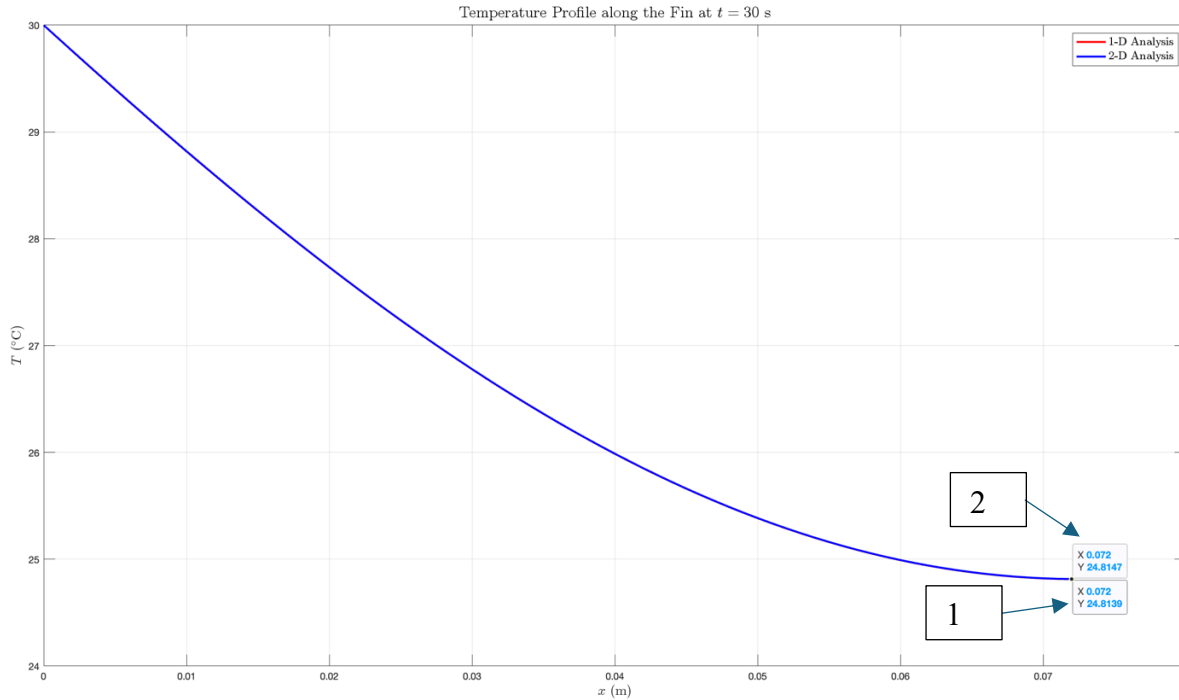


Figure 2 – Plot of 1D and 2D Numerical Solutions of the Temperature Profile at  $T = 30s$

Figure 2 depicts the plot of a one-dimensional and two-dimensional Point Jacobi Numerical Method to solve the temperature profile along the fin. The deviation of the 2 curves is small, resulting in only one visible line on the graph. The two final solutions are shown in the Figure where the two-dimensional solution at  $x = 0.072$  [m] and  $t = 30$  [s] is larger (24.8147 [C]) than the one-dimensional solution (24.8139 [c]) by only 0.0008 [c].

d.

As discussed in the Figure 2 description, the deviation of the one-dimensional and two-dimensional numerical solutions are so small they are almost negligible at 0.0008 [c]. For this reason, the one-dimensional numerical solution is a valid approximation of the temperature profile assuming the two-dimensional case is valid.