- a Law of corresponding states
 - There exist alteast semi-qualititatively a universal equation of state, containing no arbitrary constant, if expressed in terms of reduced variables

$$P_r = \frac{P}{P_c}$$
; $T_r = \frac{T}{T_c} R V_r = \frac{V}{V_c}$

This law is known as law of corresponding state (3)

(b)
(i) Stability criterion in Helmholtz free
energy representation.

$$\left(\frac{3^2A}{3^{12}}\right)_{1}$$
 70

$$\Rightarrow \left(\frac{3}{3}\left(\frac{3A}{3V}\right)_{T}\right)_{T} \geqslant 0$$

(man-him) portujo

(ii) Stability Criterion in entropy representation

$$\frac{\partial^{2} S}{\partial U^{2}} \Big|_{V,N} \leq 0$$

$$\frac{\partial}{\partial U} \Big(\frac{\partial S}{\partial U} \Big)_{V,N} \Big|_{V,N} \leq 0$$

$$-\frac{1}{T^{2}} \Big(\frac{\partial T}{\partial U} \Big)_{V,N} \leq 0$$

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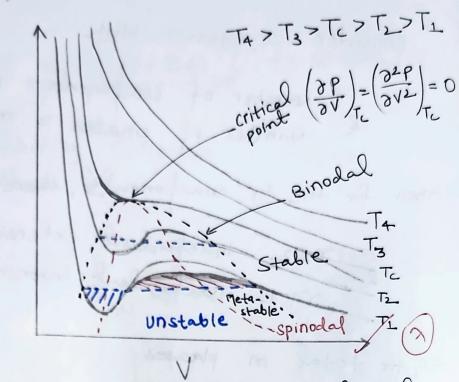
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$$\frac{\partial}{\partial U} \Big(\frac{\partial$$

Gibb's phase rule consider a system with. e number of components = x Number of phases = m Then for the 'r' components, there will be 2 r+2 independent extensive variables r+1 independent intensive variables So, for total m pleases Total number of independent intensive = variables 2+m(r-1) at equilibrium, 11,1 = 11,2 = ... = H1, m (for component 1) The above relationship for 1 component & m'phases will give (m-1) equations. NOW, for 'r' components Total number of = r(m-1)
equations = Total number Total
of independent number
Intervive of Independent
Variables Independent NOW Degree of freedom (F) (6) = 2 + m(r-1) - r(m-1) F = 2 - m + r {Gibb's phase? Rule}



(iii) Thermodynamic characteristics of critical point in terms of Helmholtz Free energy pressure, volume, tenferalure & their derivatives!

At critical point
$$\left(\frac{\partial^2 P}{\partial V^2}\right)_{T_c} = 0$$

Since,
$$\left(\frac{\partial A}{\partial V}\right)_{T_c} = -P$$

$$\left(\frac{\partial^2 A}{\partial V^2}\right)_{T_c} = -\left(\frac{\partial^2 P}{\partial V^2}\right)_{T_c} = 0$$

The stability criteria is,

$$dA = 0$$
 $d^{2}A > 0$
if $d^{2}A = 0$
 $d^{3}A > 0$
if $d^{3}A > 0$
 $d^{4}A > 0$

$$\frac{d^3A}{dV^3} = 0$$

since,

$$A = A(T,V)$$

on an isotherm.

$$d_{A} = \left(\frac{\partial_{A}}{\partial A_{3}}\right)^{L} (d_{A})^{3} = 0$$

Therefore for critical point to be stable.

However, since

$$d^{4}A = \left(\frac{\partial^{4}A}{\partial V^{4}}\right)_{T_{c}} \left(\frac{\partial^{4}A}{\partial V^{4}}$$

$$\left|\frac{\partial^4 A}{\partial V^4} > 0\right|$$

$$= - \left(\frac{33p}{3v^3}\right)_{T_c} > 0$$

$$\xi = \left(\frac{\partial T}{\partial v}\right)v$$

$$\frac{[Y,Y]}{[Y,Y]} = \frac{T[S,T] - P[Y,T]}{T[S,Y] - P[Y,Y]}$$

$$\frac{(v, v)}{T[\underline{s}, \underline{v}] - P[\underline{v}, \underline{v}]} = \frac{[\underline{s}, \underline{\tau}]}{[\underline{s}, \underline{v}]} - \frac{P[\underline{v}, \underline{\tau}]}{T[\underline{s}, \underline{v}]}$$

$$= \frac{(T, S)}{(Y, S)} + \frac{f(T, Y)}{f(S, Y)}$$

g, Maxwell relationship

 $\begin{cases} \frac{1}{2} \sqrt{2} & \frac{1}{2} \end{cases} = 0$

Y: Q = T(S, V)

$$\frac{(T,U)}{(Y,U)} = -(P,V) + \frac{1}{2}(\frac{X}{4})$$

$$= \frac{(Y,P)}{(Y,U)} + \frac{1}{4}(\frac{X}{4})$$

$$= \frac{(Y,P)}{(Y,U)} + \frac{1}{4}(\frac{X}{4})$$

$$\dot{\xi} = \frac{1}{c_V} \left\{ -\frac{T_X}{K_T} + P \right\}$$

$$\begin{cases} \frac{1}{k_T} = \begin{bmatrix} Y, P \\ Y, T \end{bmatrix} \end{cases}$$

As we know;

$$c_{v} = q - \frac{V x^{2} T}{k T}$$

$$\frac{1}{\sqrt{2}} = \frac{1}{\sqrt{2}} \left[\frac$$

$$P = \frac{RT}{V-b} - \frac{9}{V^2} - \frac{1}{V}$$

(i)
$$K_T = -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_T$$
 (2)

$$\left(\frac{\partial P}{\partial Y}\right)_{T} = \left(\frac{\partial}{\partial Y} \left\{\frac{RT}{Y-b} - \frac{\alpha}{Y^{2}}\right\}\right)_{T}$$

$$\left(\frac{\partial P}{\partial Y}\right)_{T} = \frac{-RT}{(Y-b)^{2}} + \frac{2a}{Y^{3}}$$

$$= \frac{1}{2a} \left(\frac{\partial Y}{\partial P} \right)_{T} = \frac{1}{2a} \left(\frac{PT}{Y^{3}} - \frac{PT}{(Y-b)^{2}} \right)$$

$$k_{T} = -\frac{1}{V} \left\{ \frac{1}{\frac{2a}{V^{3}} - \frac{RT}{(V-b)^{2}}} \right\}$$

$$k_{\overline{z}} = \frac{1}{\frac{RT\underline{y}}{(Y-b)^2} - \frac{2a}{Y^2}}$$

$$K_{T} = \frac{Y^{2}(Y-b)^{2}}{RTY^{3}-2a(Y-b)^{2}}$$

$$x = \frac{1}{\sqrt{2L}} \left(\frac{3L}{2L} \right)^{b}$$

$$P = \frac{RT}{V-b} - \frac{a}{V^2}$$

$$\left(\frac{\partial P}{\partial T}\right)_{p} = \left(\frac{\partial}{\partial T} \left\{ \frac{RT}{Y-b} - \frac{Q}{Y^{2}} \right\} \right)_{p}$$

$$0 = \frac{R}{V-b} \times \frac{3T}{5T} + RT \left(\frac{3}{5T} \left(\frac{1}{V-b}\right)\right)_{P}$$

$$- \alpha \left(\frac{3}{5T} \left(\frac{1}{V^{2}}\right)\right)_{P}$$

$$0 = \frac{RT}{V-b} - \frac{RT}{(V-b)^2} \left(\frac{\partial V}{\partial T} \right)_p + \frac{2a}{V^3} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left\{\frac{2a}{V^3} - \frac{RT}{(V-b)^2}\right\} \left(\frac{3V}{3T}\right)_p = -\frac{R}{V-b}$$

$$\left(\frac{PT}{(V-b)^2} - \frac{2a}{V^3}\right)\left(\frac{\partial V}{\partial T}\right)_p = \frac{R}{V-b}$$

$$\left\{ \frac{RT}{(Y-b)} - \frac{2a(Y-b)}{Y^3} \right\} \left(\frac{2Y}{2T} \right)_p = R$$

$$\left(\frac{\partial V}{\partial T}\right)_{p} = \frac{RT}{V-b} - \frac{2a(V-b)}{V^{3}}$$

$$X = \frac{1}{y} \left(\frac{\partial y}{\partial T} \right)_{p}$$

$$X = \frac{1}{y} \left(\frac{\partial y}{\partial T} \right$$

4) Given,
$$\frac{1}{Ta} = \frac{3}{2} Na \frac{R}{Va} - CI$$

Na = 3 kmole $N_b = 2$ kmole

Initially, Ti = 400 K

Ti = 700 K

from equation (1) & (2)

Va = 3 NaRTa

 $= \frac{3}{2} \times 3 \times 8^{-314} \times 4^{00}$

Va = 14965.2 KJ

Ub = 5 Nb RTb

 $=\frac{5}{2}X2X8-314X700$

Ub = 29099 KJ

Total internal energy of the system (U)

 $U = U_a^i + U_b^i$ = 14965.2 + 29099

Finally

wall is made diathermal & both the systems are allowed to reach the systems are allowed to reach the state of thermal equilibrium.

$$\Rightarrow \frac{1}{T_a} = \frac{1}{T_b}$$

from (1) & (2)

$$\frac{3}{2}N_{a} = \frac{5}{2}N_{b} \frac{R}{V_{b}}$$

$$V_{a} = \frac{3}{5} \frac{N_{a}N_{b}}{N_{b}} \frac{N_{b}}{V_{b}}$$

$$V_{a} = \frac{3}{5} \times \frac{3}{2} \times V_{b}$$

$$V_{a} = 0.9 V_{b}$$

$$U_a + U_b = U = 0.9U_b + U_b$$

$$U_{b} = 23191.68 \text{ FJ}$$

$$U_{a} = 20872.52 \text{ FJ}$$

$$T_a = \frac{2}{3} \times \frac{1}{3} \times \frac{20872.52}{8.314}$$

from eq 2

$$T_b = \frac{2}{5} \times \frac{1}{N_b} \times \frac{U_b}{R}$$

$$= \frac{2}{5} \times \frac{1}{2} \times \frac{23191.68}{8.314}$$

$$T_b = T_a = 557.895 \text{ K}$$

(3) for van der waalk fluid

$$P = \frac{NRT}{V-Nb} - \frac{aN^2}{V^2} - 0$$

$$\left(\frac{\partial P}{\partial N}\right)_{T,N} = \frac{RT}{V-Nb} + \frac{NbRT}{(V-Nb)^2} - \frac{2aN}{V^2}$$
for
$$\left(\frac{\partial P}{\partial N}\right)_{T,N} < 0$$

$$\frac{RT}{V-Nb} + \frac{NbRT}{(V-Nb)^2} - \frac{2aN}{V^2} < 0$$

$$\frac{RTV^2}{aN(V-Nb)} + \frac{NbV^2RT}{aN(V-Nb)^2} - 2 < 0$$

$$\frac{RTV^2}{aN(V-Nb)^2} + \frac{Nb(V+Nb)^2R^{ATR}}{aN(V+Nb-NV)^2}$$

$$\frac{AR}{aN^2b^2} + \frac{Nb(V+Nb)^2R^{ATR}}{aN^2b^2} + \frac{Nb(V+Nb)^2R^{ATR}}{aN$$

$$\frac{T^{*} \vee^{*2}}{\vee^{*-1}} + \frac{T^{*} \vee^{*2}}{(\vee^{*-1})^{2}} < 2$$

$$\frac{T^{*} \vee^{*2}}{\vee^{*-1}} \left(1 + \frac{1}{\sqrt{2}} \right) < 2$$

$$\frac{T^{*} \vee^{*2}}{(\vee^{*-1})} \left(\frac{2}{\sqrt{2}} \right) < 2$$

$$\frac{T^{*} \vee^{*2}}{(\vee^{*-1})} \left(\frac{2}{\sqrt{2}} \right) < 2$$

$$\frac{T^{*} \vee^{*3}}{(\vee^{*-1})} \left(\frac{2}{\sqrt{2}} \right) < 2$$

$$\frac{T^{*}}{(\vee^{*-1})} \left(\frac{2}{\sqrt{2}} \right) < 2$$

$$\frac{T^{*}$$