

# Coupled Laser Simulation Updates

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## Introduction and Goals

I wish to investigate the behavior of a system of lasers which are coupled to one another. In this project, I model the system as a set of Coupled Lang-Kobayashi Equations. This set of first-order, delay differential equations are solved numerically by a Python program.

**Goal:** to use this laser model to simulate the behavior of a physical compute reservoir (PRC). We will compare this (simulated) reservoir's ability to perform classification/prediction tasks against a traditional neural network.

**Motivation:** It is a slow and costly process to build an actual chip of lasers in the cleanroom, so we first verify that a system of lasers can function as a PRC in software before we actually go build it. This simulation process will also give us intuition into how the lasers should be patterned on our chip and how they should be driven/read out for best performance.

In these notes I primarily reference "Semiconductor Lasers: Stability, Instability, and Chaos" by Ohtsubo [1], which I have been reading since mid-June.

My code and figures can be found in my GitHub repo: <https://github.com/godwinduan/lasersim>

## 1 July 9-10, 2025

Began coding the Coupled Lang-Kobayashi Equations.

### 1.1 Coupled Lang-Kobayashi Equations

For a system of  $n$  lasers, the differential equations for the field in laser  $j$  are given by:

$$\begin{aligned} \frac{dE_j(t)}{dt} = & \frac{1}{2}(1 - i\alpha)(G_j - \gamma)E_j(t) + \frac{\kappa_{jj}}{\tau_{in}}E_j(t - \tau_{jj})\exp(i\omega_j\tau_{jj}) \\ & + \sum_{k=0, k \neq j}^{n-1} \frac{\kappa_{jk}}{\tau_{in}}E_k(t - \tau_{jk})\exp(-i\Delta\omega_{jk}\tau_{jk}) \end{aligned} \quad (1)$$

where  $\kappa_{jk}$  is the coupling strength from  $k$  to  $j$ ,  $\tau_{jk}$  is the time delay from  $k$  to  $j$ , and  $\Delta\omega_{jk} = \omega_k - \omega_j$ .

The first term on the right is increase in field from the laser's own field. The second term on the right is the self coupling of the laser (has a delay  $\tau_{jj}$ ). The third term on the right is the coupling from all other lasers.

To code this, we have to break up the above expression for  $E$  into a real component  $X_j(t)$  and imaginary component  $Y_j(t)$  such that  $E_j(t) = X_j(t) + iY_j(t)$ .

$$\begin{aligned} \frac{dE_j(t)}{dt} = & \frac{1}{2}(G_j - \gamma)E_j(t) - i\alpha\frac{1}{2}(G_j - \gamma)E_j(t) + \frac{\kappa_{jj}}{\tau_{in}}E_j(t - \tau_{jj})\cos(\omega_j\tau_{jj}) + i\frac{\kappa_{jj}}{\tau_{in}}E_j(t - \tau_{jj})\sin(\omega_j\tau_{jj}) \\ & + \sum_{k=0, k \neq j}^{n-1} \frac{\kappa_{jk}}{\tau_{in}}E_k(t - \tau_{jk})\cos(\Delta\omega_{jk}\tau_{jk}) - i\sum_{k=0, k \neq j}^{n-1} \frac{\kappa_{jk}}{\tau_{in}}E_k(t - \tau_{jk})\sin(\Delta\omega_{jk}\tau_{jk}) \end{aligned} \quad (2)$$

After we expand  $E$  into  $X$  and  $Y$  and distribute terms,

$$\begin{aligned} \frac{dX_j(t)}{dt} = & \frac{1}{2}(G_j - \gamma)(X_j(t) + \alpha Y_j(t)) + \frac{\kappa_{jj}}{\tau_{in}}(X_j(t - \tau_{jj})\cos(\omega_j\tau_{jj}) - Y_j(t - \tau_{jj})\sin(\omega_j\tau_{jj})) \\ & + \sum_{k=0, k \neq j}^{n-1} \frac{\kappa_{jk}}{\tau_{in}}(X_k(t - \tau_{jk})\cos(\Delta\omega_{jk}\tau_{jk}) + Y_k(t - \tau_{jk})\sin(\Delta\omega_{jk}\tau_{jk})) \end{aligned} \quad (3)$$

$$\begin{aligned} \frac{dY_j(t)}{dt} = & \frac{1}{2}(G_j - \gamma)(Y_j(t) - \alpha X_j(t)) + \frac{\kappa_{jj}}{\tau_{in}}(Y_j(t - \tau_{jj})\cos(\omega_j\tau_{jj}) + X_j(t - \tau_{jj})\sin(\omega_j\tau_{jj})) \\ & + \sum_{k=0, k \neq j}^{n-1} \frac{\kappa_{jk}}{\tau_{in}}(Y_k(t - \tau_{jk})\cos(\Delta\omega_{jk}\tau_{jk}) - X_k(t - \tau_{jk})\sin(\Delta\omega_{jk}\tau_{jk})) \end{aligned} \quad (4)$$

For the carrier density of laser  $j$ ,

$$\frac{dN_j(t)}{dt} = \frac{I_j}{e} - \gamma_e * N_j - G_j * |E|^2 \quad (5)$$

Where  $|E|^2 = X^2 + Y^2$ , and

$$G_j = \frac{g(N_j - N_0)}{1 + s|E|^2}$$

## 1.2 Comments

- Question: is it plus alpha or minus alpha in the Lang-Kobayashi Equation? In Ohtsubo's textbook, it is minus (see Ohtsubo 87), and I am following this convention, but in the original Lang-Kobayashi paper (and in Chixuan's work), it is plus alpha. As a result my equations' signs differ from Chixuan's equations on slide 11 of her generals presentation.
- I'm using individual frequencies. This is because the self-coupling term requires the laser's absolute (not relative frequency). See Ohtsubo pg. 186 for explanation of where the  $i\Delta\omega E$  term comes from. Chixuan uses this but I do not.
- We aren't missing a cavity leaking term gamma in the above equation for carrier density since the carrier density doesn't depend on that.

## 2 July 11, 2025

Finished coding the Lang-Kobayashi Equations. Changed kappa values to make them dimensionless, following the convention adopted in Ohtsubo's textbook. (Chixuan had kappa include the divisor term  $\tau_{in}$ , the round trip time inside the cavity.

Ran some sample simulations:

### 2.1 One Laser Simulations

Tweaking parameters of  $\kappa$ , the self-coupling coefficient,  $\tau$ , the delay time it takes for the laser to couple to itself, and  $I_{injection}$ , both stable and chaotic behavior was observed in the single self-coupled laser.

Below in Fig. 1 the laser is biased just above threshold current and is very weakly coupled to itself. Only relaxation oscillations are observed.

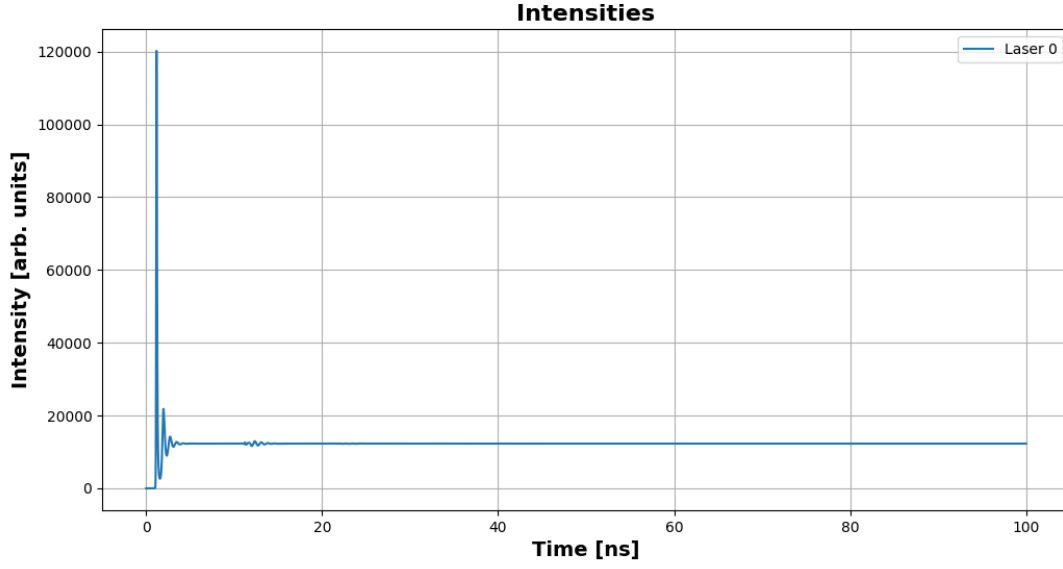


Figure 1: Intensity of relaxation oscillation of a single self coupled laser w/  $\kappa = 0.0012$ ,  $\tau = 10\text{ns}$ ,  $I_{\text{injection}} = 18\text{mA}$ .

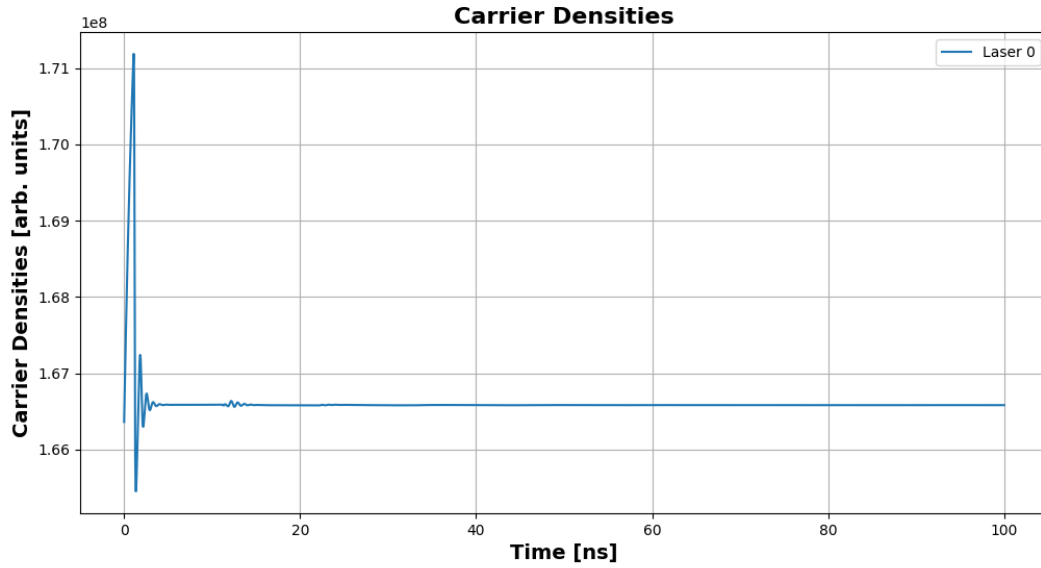


Figure 2: Carrier density during relaxation oscillation of a single self coupled laser w/  $\kappa = 0.0012$ ,  $\tau = 10\text{ns}$ ,  $I_{\text{injection}} = 18\text{mA}$ .

Also shown here is the carrier density inside the laser. This is something we can measure if needed by measuring the voltage across the laser junction. See Ohtsubo pg. 435 for how voltage relates to carrier density.

Increasing the self-coupling, we now observe chaotic oscillations, shown in Fig. 3.

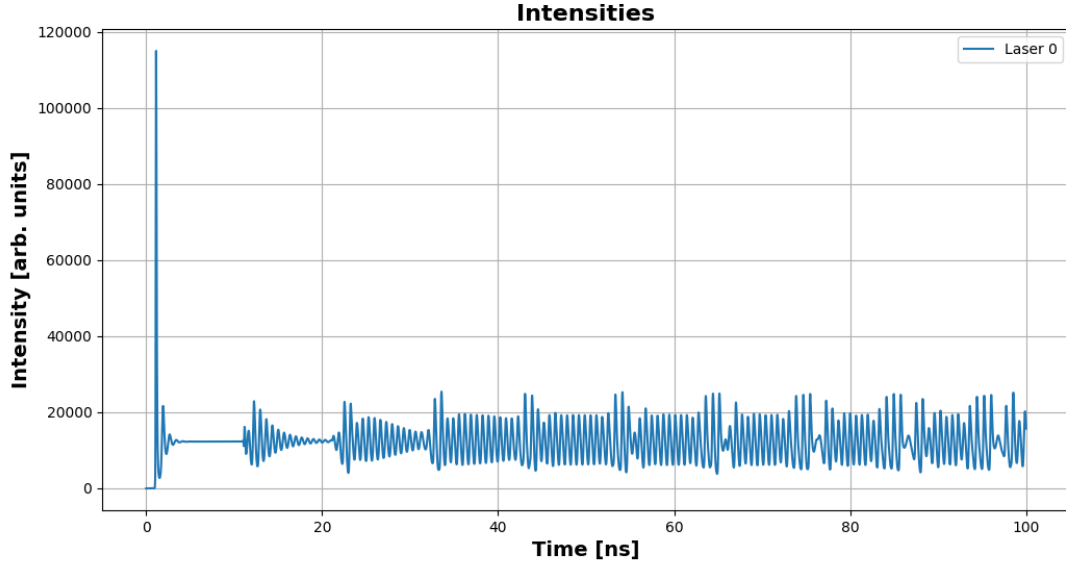


Figure 3: Relaxation oscillation of a single self coupled laser w/  $\kappa = 0.012$ ,  $\tau = 10\text{ns}$ ,  $I_{\text{injection}} = 18\text{mA}$ .

Further increasing the self-coupling, and also increasing the injection by a few mA, we can observe periodic oscillations, shown in Fig. 4.

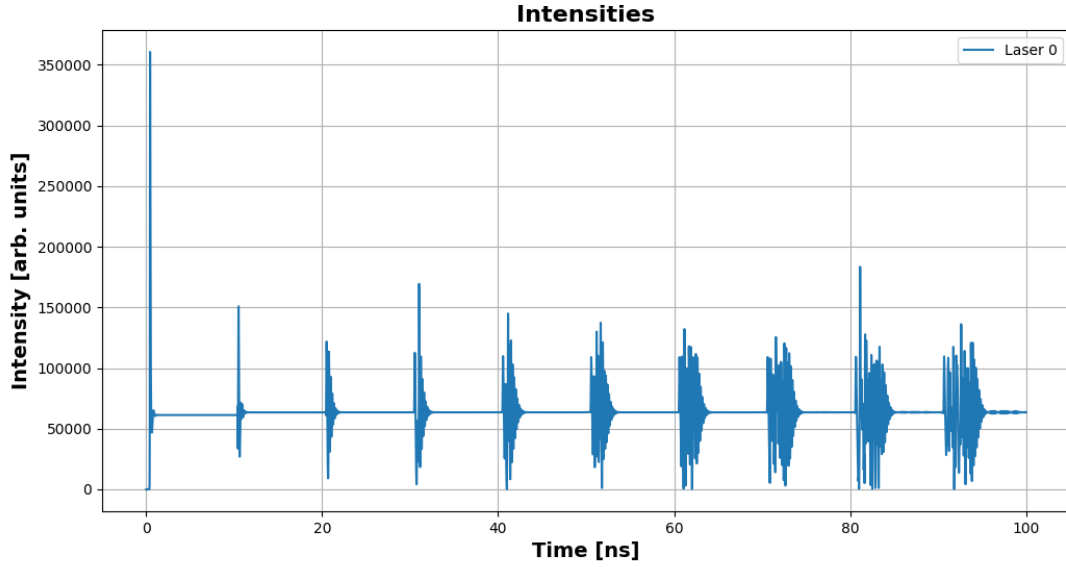


Figure 4: Relaxation oscillation of a single self coupled laser w/  $\kappa = 0.12$ ,  $\tau = 10\text{ns}$ ,  $I_{\text{injection}} = 22\text{mA}$ .

## 2.2 Two Lasers Simulations

I played around with changing the coupling parameters/delay times for two lasers coupled together. Highly chaotic behavior can be observed; here is an example where laser 0 is coupled to laser 1 as well as to itself. Both lasers are being driven at their threshold currents. Both lasers oscillate rapidly between on and fully off, shown in Fig. 5.

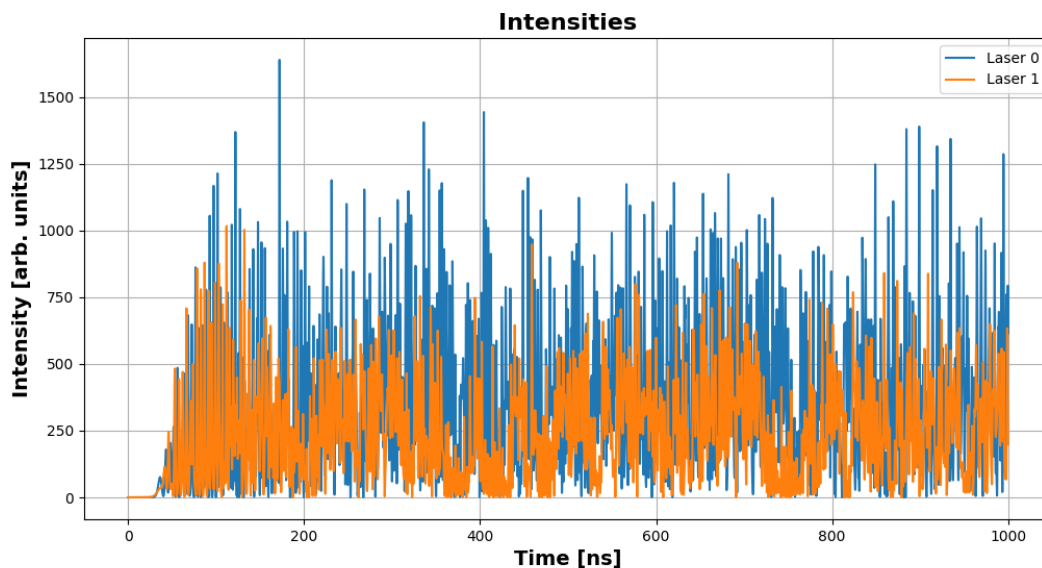


Figure 5: Highly chaotic oscillations of a two coupled lasers.

## 2.3 Comments

- Want to take the FFT of the chaotic oscillations to see its frequency spectrum
- What is the speed of the photodetectors that we can build on-chip? I have read they have a response of 1-10MHz. If it's too slow, we won't be able to see this phenomena.
- We could try to induce some kind of low-frequency fluctuations (LFF) (see Ohtsubo pg. 140) that would be slow enough detectable. This requires a long delay time. I wonder if we attach a single long fiber optic cable, to a system containing multiple coupled lasers, if this would offer enough long-term memory.
- Next step: add in code a way to be able to change the injection current during the simulation. Right now it is constant.

### 3 July 14, 2025

- Implemented FFT analysis of intensity vs. time, and a low pass filter on intensities to emulate the limited response speed of a photodetector.
- Refactored entire notebook for ease of use

#### 3.1 Thoughts

- We wish to measure a "steady state" laser signal. But for what configurations is there actually a steady state, and for what configurations are there only chaotic oscillations?
- Prof. Gmachl says its possible to measure the GHz frequencies by heterodyning two signals at the photodetector (Prof. Wysocki does this).
- If the laser is oscillating chaotically on to off, is its average power output lower than it would be if it was outputting at steady state? This average power is a low frequency signal and would be detectable by our photodetectors.
- Want to find the intensity vs current curve of a single, uncoupled laser according to the laser equations.
- Need an efficient way to simulate many variations on parameters.
- Far future: modify the equations to include effect of external light that does not couple into the laser yet still depletes the carrier density in the laser. This is mentioned somewhere in Ohtsubo.
- Future work: convert carrier density to voltage, to get an IV curve of a single uncoupled laser

#### 3.2 Coupling Coefficients

##### 3.2.1 Feedback Coefficient $\kappa$

Assuming the reflectivities for the front and back facets are both  $r_0$ , (see Ohtsubo pg. 87)

$$\kappa = (1 - r_0^2) \frac{r}{r_0} \tag{6}$$

where  $r$  is the external reflection.

### 3.2.2 C Parameter for single-laser self-feedback

$$C = \frac{\kappa\tau}{\tau_{in}} \sqrt{1 + \alpha^2} \quad (7)$$

When  $C < 1$ , the solitary laser is stable and no chaotic oscillations can be observed. (see Ohtsubo pg. 89)

### 3.3 Intensity vs Current Curve

I ran transient simulations with varying injection currents for a single laser with no self-coupling. I capture the value of the intensity at the end of the simulation once it has reached steady state value. As expected according to the theoretical model, the intensity is zero when current is under the threshold current, and increases linearly when current is above the threshold.

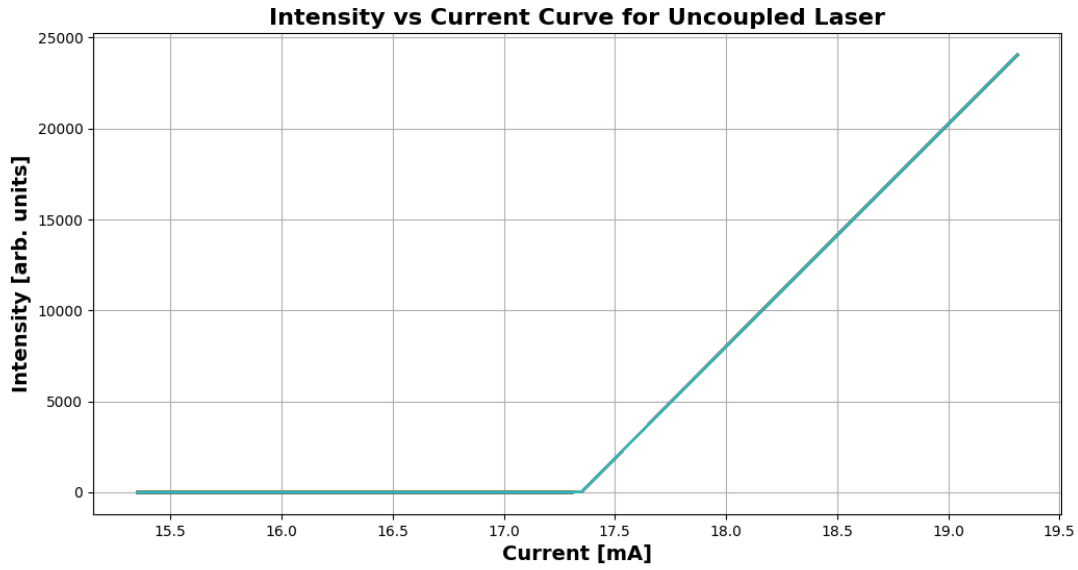


Figure 6: Intensity vs. Current Curve of a single uncoupled laser

## 4 July 15, 2025

- Implemented ability to change injection current values during simulation.
- Observed the decay of chaotic oscillations in a single oscillator for shorter self-coupling delays which changes the average intensity of the oscillator on a slower time scale (100s of ns). This might have potential to become memory.



## 4.1 Variable current during simulation

### 4.1.1 Changing parameters

I add a parameter instead of a constant current for  $I_{\text{injection}}$ , and when it is time to set a new injection current, I pause the integration, change the parameter, and resume integration.

If no discontinuity smoothing is done after the change, sometimes the integrator fails to solve the system. On the other hand, if `step_on_discontinuities()` or `adjust_diff()` is used after the change, there are huge spikes in the result as seen in 7.

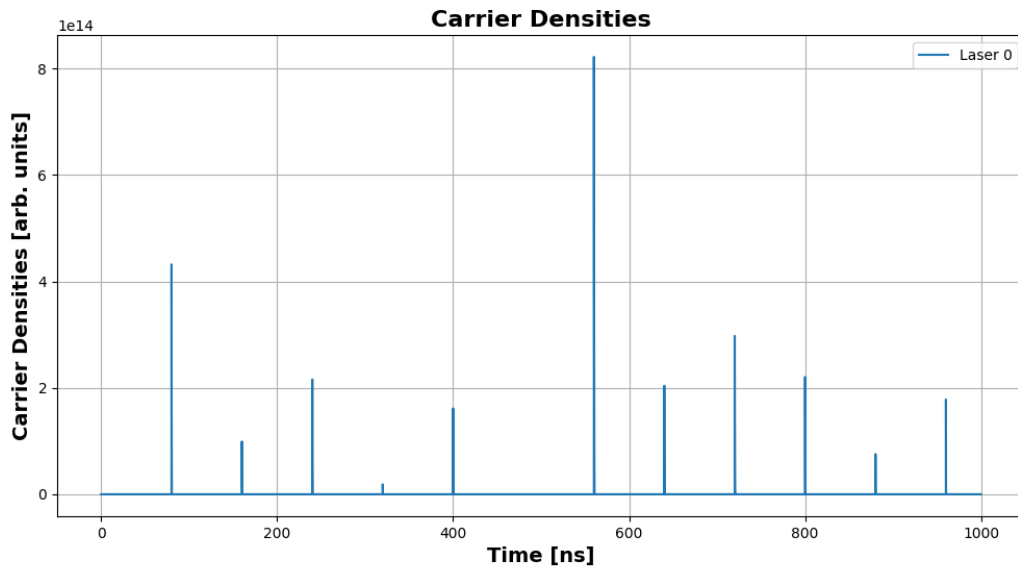


Figure 7: Huge spikes due to discontinuous injection current changes.

### 4.1.2 Smoothed current changes

I coded an alternate simulation method, not using the parameter changing method, but instead using the continuous function “conditional”. It resulted in smaller spikes but was extraordinarily slow. This was only for 1 laser, so having this for a system of say 100 lasers would be impossible.

So, I went back to the parameter-changing method and fixed the “shift\_ratio” parameter to a much larger value (now equal to 1), which smooths out the discontinuity over a larger area. Now the spikes in carrier density are much more tolerable, almost not noticeable. Increasing “shift\_ratio” did not have a noticeable effect on the intensity.

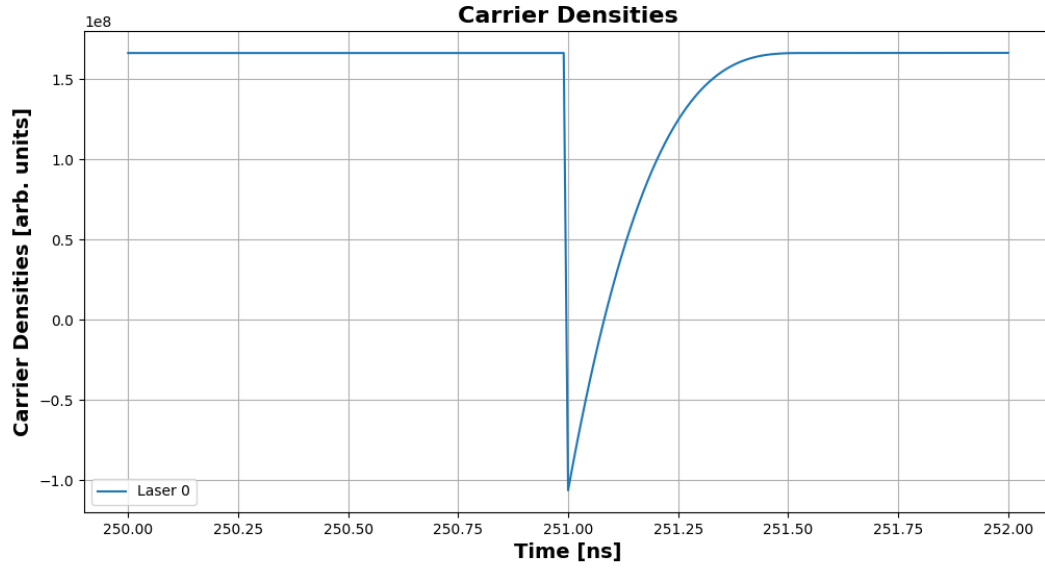


Figure 8: Carrier density after a injection current change with “shift\_ratio” = 0.01

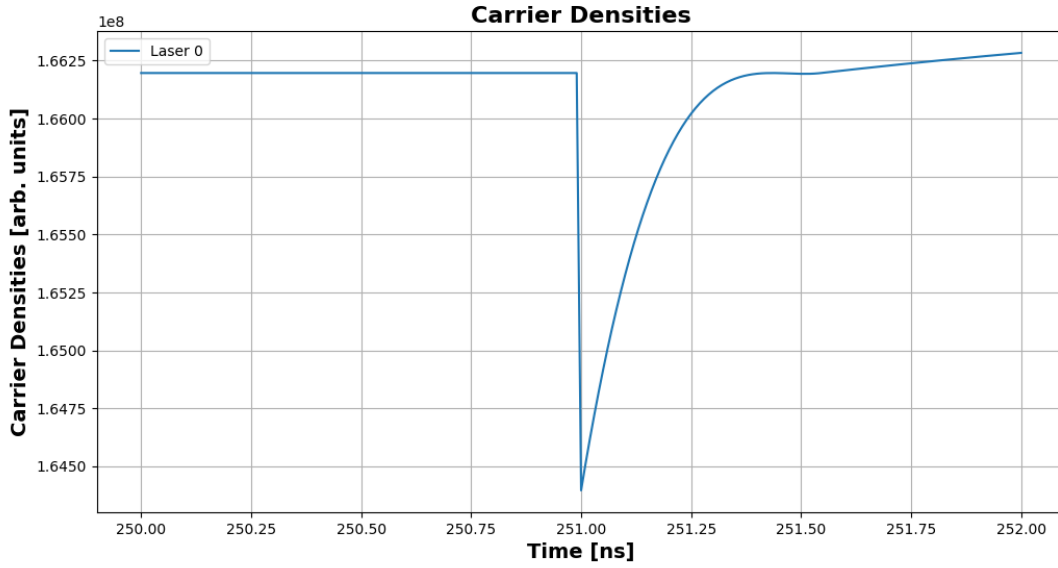


Figure 9: Carrier density after a injection current change with “shift\_ratio” = 0.1

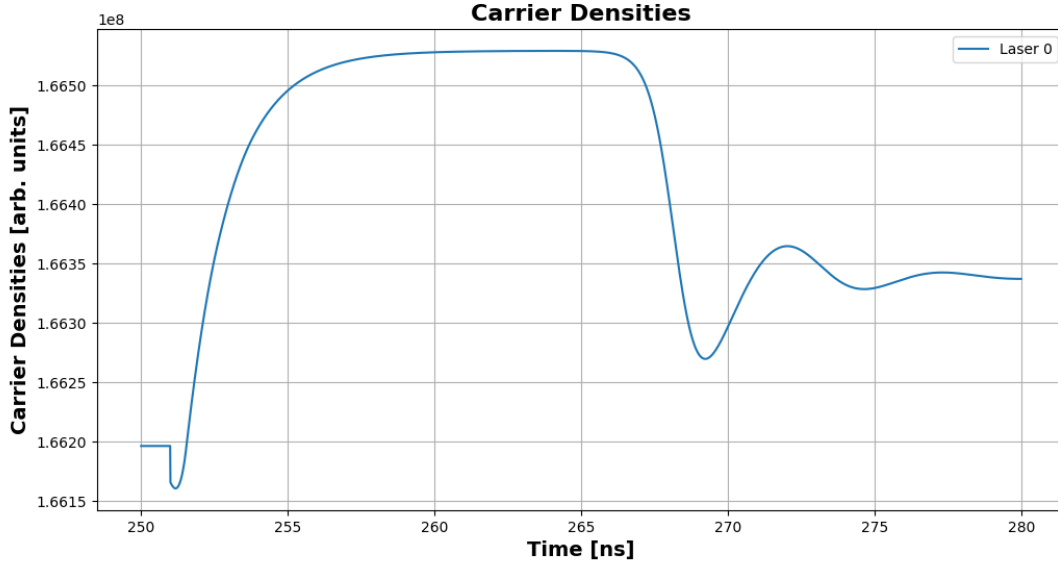


Figure 10: Carrier density after a injection current change with “shift\_ratio” = 1

#### 4.2 One Laser Experiment 5: Effect of coupling/chaos on average power

- Gotta ask Chixuan where she got the laser parameters so I can get an accurate value of the length of the laser. This way I can calculate the parameter C accurately (right now I’m using a random length and oscillations stop for some C values greater than 1)
- Once C is accurate, we can experimentally determine the criteria of the minimum coupling and coupling delay needed for oscillations.
- Just realized: even though the intensity of the field inside the laser increases, the total light that escapes the laser is not necessarily increasing? Because increasing kappa here increases the
- What I did: measures mean value of intensity period to determine the avg. power. Calculates stddev of intensity over a time period to determine if there are chaotic oscillations or not.

Found that chaotic oscillations do vary the slow-varying average intensity of the laser, which is measurable by photodiodes. See below two figures.

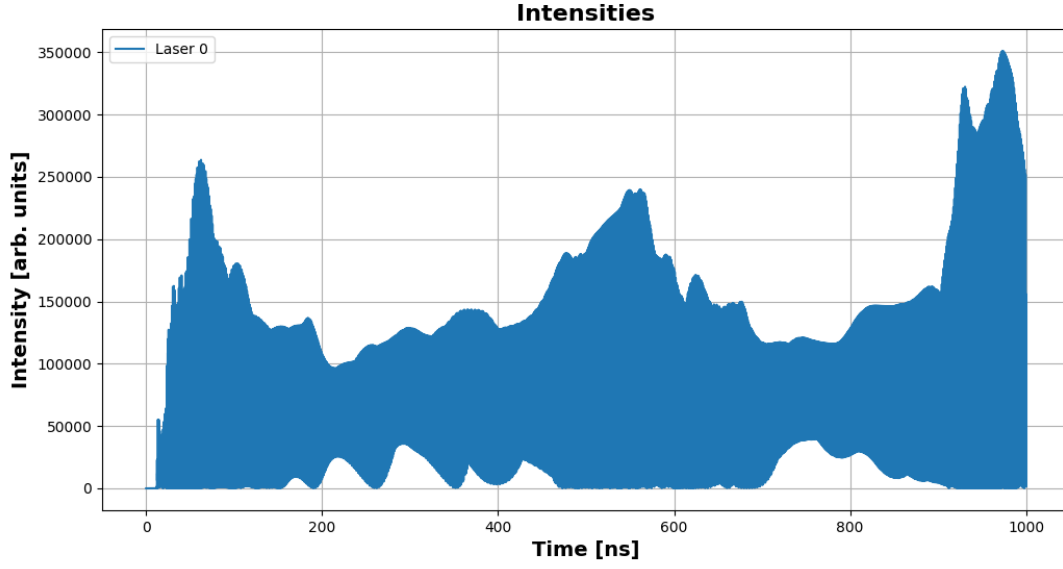


Figure 11: Intensity for a self-coupled laser w/  $\kappa = 1$ ,  $\tau = 1\text{ns}$ , driven at threshold current.

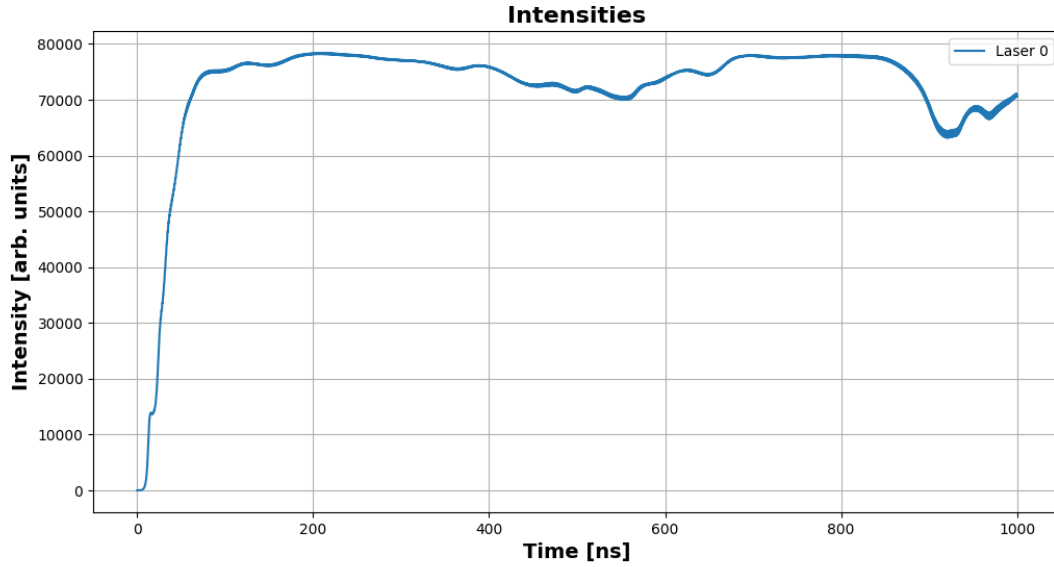


Figure 12: Low-Pass Filtered Intensity for a self-coupled laser w/  $\kappa = 1$ ,  $\tau = 1\text{ns}$ , driven at threshold current.

### 4.3 Decay of Chaos for short coupling delays

I had been simulating stuff at 10ns coupling delay but when I reduced it to 1ns I found that the initial chaotic oscillations die out after 100ns. See Fig. 13 and Fig. 14

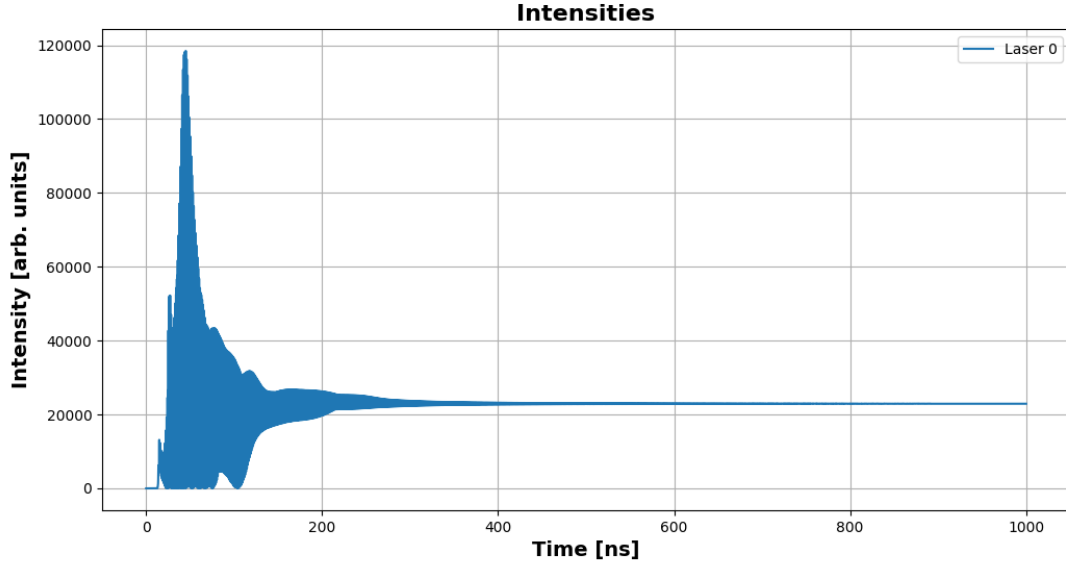


Figure 13: Intensity for a self-coupled laser w/  $\kappa = 0.5$ ,  $\tau = 1\text{ns}$ , driven at threshold current.

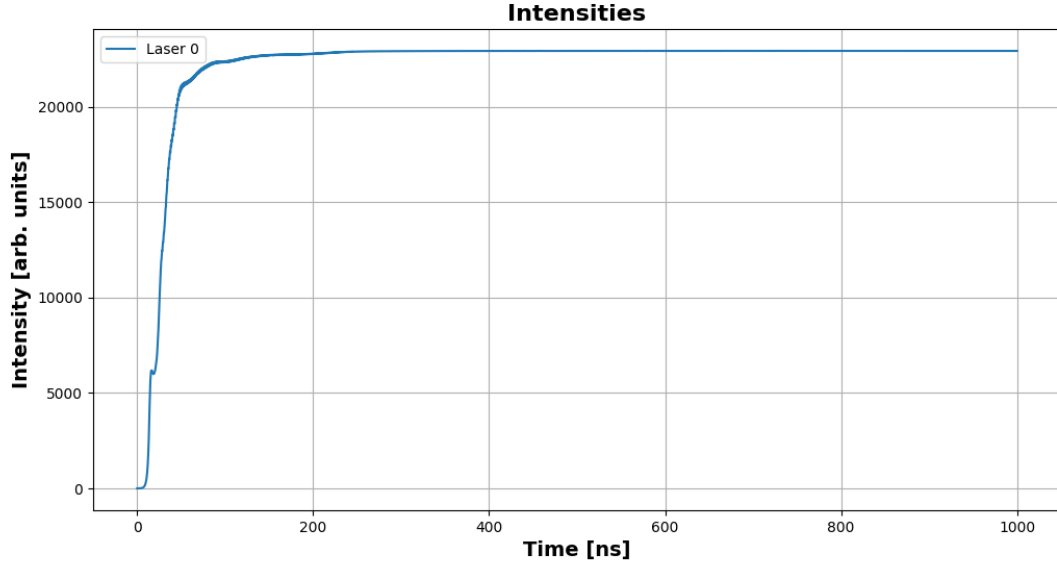


Figure 14: Low-Pass Filtered Intensity for a self-coupled laser w/  $\kappa = 0.5$ ,  $\tau = 1\text{ns}$ , driven at threshold current.

Compare Fig. 14 with the much quicker rise time of the weakly self-coupled laser shown in Fig. 15. This shows that the very high frequency chaotic oscillations which die out give rise to longer time scale variations in average intensity (which we can measure with our photodiodes).

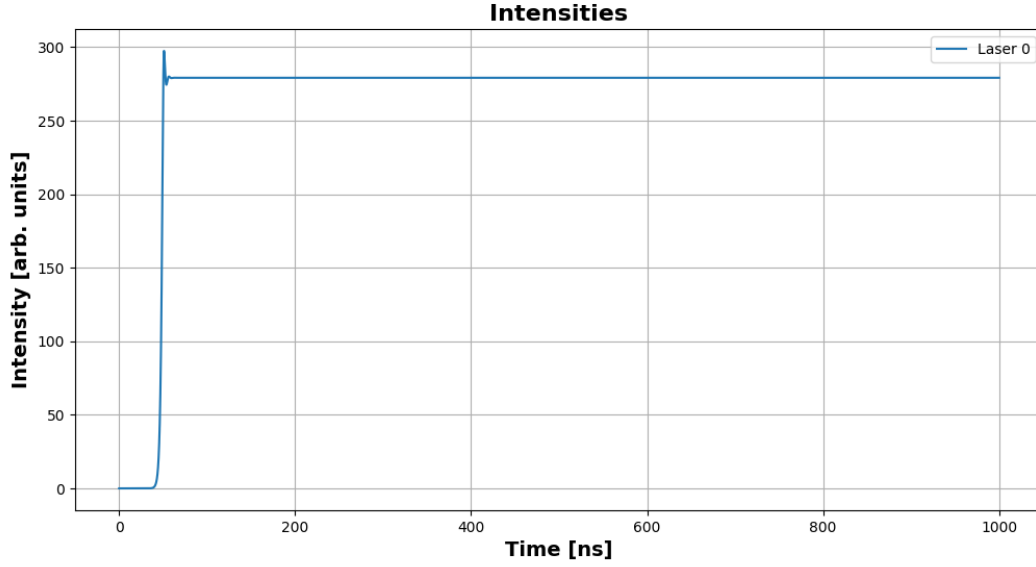


Figure 15: Low-Pass Filtered Intensity for a weakly self-coupled laser w/  $\kappa = 0.01$ ,  $\tau = 1\text{ns}$ , driven at threshold current.

How long is chaos sustained before it dies out? What does it depend on? If we introduce noise from spontaneous emissions, will the oscillations never die out? If we change the injection current, will the chaos come back?

If this chaos decay is a predictable phenomena it may be useful as a memory in a physical compute reservoir.

#### 4.4 Thoughts

- What exactly does the gain saturation term  $(1 + E^2)$  do? Do we only observe it for large outputs?
- Implement a sliding window which measures the stddev of intensity over that window, dividing it by the average intensity over that window, as a measure of how chaotic it is.
- I worry that the delays on our small chip will be so small, there will not be chaotic oscillations. The C parameter of stability is proportional to kappa (the amount of coupling) and also to the coupling delay. If the coupling delay is too small, even complete coupling will not be enough to trigger chaos.
- On one hand, we could connect lasers to one another by waveguides to increase coupling. We could also add waveguide delays in a chip to add small delays.

- On the other hand, once we add in more than one laser, frequency mismatches will cause chaos in otherwise stable lasers. (this is in Ohtsubo)

## 5 July 16, 2025

Investigated how light of one laser changes the threshold current of the other.

### 5.1 Threshold Current Changes of Two Coupled Lasers

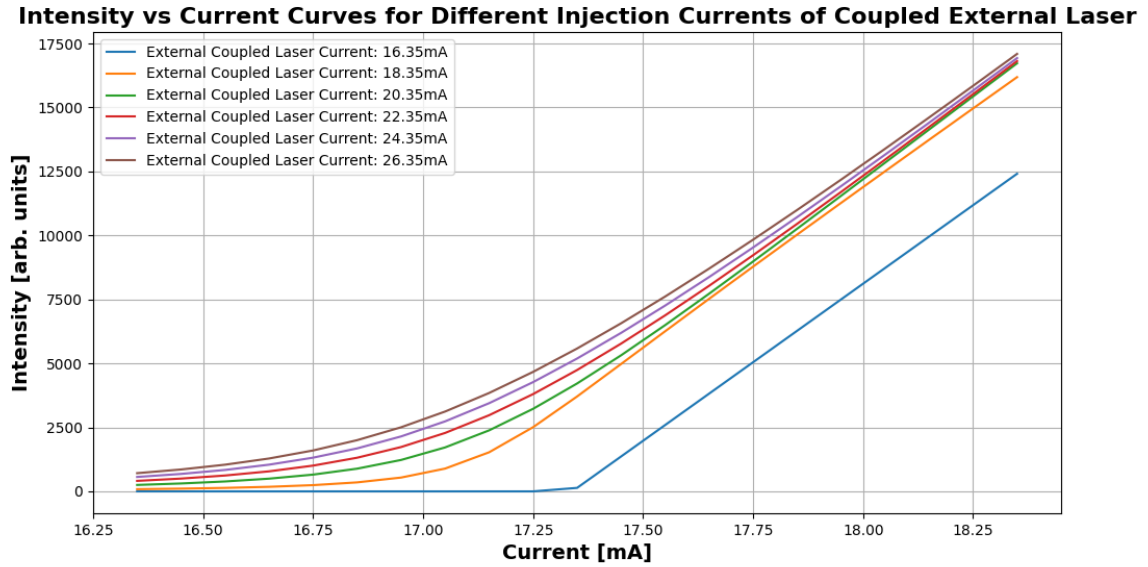


Figure 16: Intensity vs. Current curves for a laser coupled to an external laser. Different injection values of the external laser are simulated, which change the curve's characteristics.

These curves stop looking like a ReLU function and look more like a softmax  $\ln(1 + e^x)$  function. Linear regression is no longer a great way to find the threshold because the threshold has become. We could try modeling it instead as a softmax  $\ln(a + e^{b(x-c)})$ .

### 5.2 Meeting notes

## References

- [1] J. Ohtsubo, Semiconductor lasers and theory, in: Semiconductor Lasers: Stability, Instability and Chaos, Springer, 2017, pp. 33–82.