

Physics 606  
Quantum Mechanics I  
Professor Aleksei Zheltikov

Homework #7

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## 1 Problem #1

To find the expectation value,  $\langle \hat{L}^2 \rangle$ , in the state with a wave function

$$\psi(\theta, \varphi) = A \sin \theta \cos \varphi$$

we note that  $\psi(\theta, \varphi)$  is a spherical harmonic with  $l = 1$ . This implies that

$$\begin{aligned} \langle \hat{L}^2 \rangle &= \int_{d\Omega} \psi^*(\theta, \varphi) \hat{L}^2 \psi(\theta, \varphi) d\Omega \\ &= \int_{d\Omega} \psi^*(\theta, \varphi) \hbar^2(1)(1+1) \psi(\theta, \varphi) d\Omega \\ &= 2\hbar^2 \end{aligned}$$

## 2 Problem #2

To find the eigenvalue of the  $\hat{L}^2$  operator corresponding to the eigenfunction

$$Y(\theta, \varphi) = A(3 \cos^2 \theta - 1 + \sin(2\theta) \cos \varphi)$$

we note that  $Y(\theta, \varphi)$  is the eigenfunction that corresponds to the sum of two spherical harmonics

$$Y(\theta, \varphi) = A(Y_{20} + 2Y_{21})$$

note both are in the  $l = 2$  state. This implies

$$\begin{aligned} \hat{L}^2 Y(\theta, \varphi) &= A(\hat{L}^2 Y_{20} + 2\hat{L}^2 Y_{21}) \\ &= A(\hbar^2(2)(2+1)Y_{20} + \hbar^2(2)(2+1)2Y_{21}) \\ &= 6\hbar^2 Y(\theta, \varphi) \end{aligned}$$

### 3 Problem #3

For a  $s$ -state ( $l = 0$ ) particle with mass,  $m_0$ , in an spherically symmetric infinitely deep rectangular potential well of radius  $a$  is described by the wave function

$$\psi_{klm} = A j_l(kr) Y_{lm}(\theta, \varphi)$$

for  $l = 0$  we have the state

$$\psi_{k00} = A \frac{\sin(kr)}{kr}$$

Where we can find the normalization constant by noting the continuity condition that  $\psi_{k00}(r = 0) = \psi_{k00}(r = a) = 0$  which implies that  $k = n\pi/a$ . Where  $n$  becomes the principle quantum number

$$\begin{aligned} 1 &= \int_V \psi_{k00}^* \psi_{k00} dV \\ &= 4\pi \int_0^a A^2 \left( \frac{\sin(kr)}{kr} \right)^2 r^2 dr \\ &= \frac{4\pi a^2}{n^2 \pi^2} A^2 \int_0^a \sin^2(kr) dr \\ &= \frac{4a^2}{n^2 \pi} A^2 \frac{a}{2} \\ &\Downarrow \\ A &= n \sqrt{\frac{\pi}{2a^3}} \end{aligned}$$

we can find the expectation value of  $r$  of a particle in this state by

$$\begin{aligned} \langle r \rangle &= \int_V \psi_{k00}^* r \psi_{k00} dV \\ &= 4\pi \int_0^a A \frac{\sin(kr)}{kr} r A \frac{\sin(kr)}{kr} r^2 dr \\ &= 4\pi \frac{n^2 \pi}{2a^3} \frac{a^2}{n^2 \pi^2} \int_0^a r \sin^2(kr) dr \\ &= \frac{2a^2}{a} \frac{1}{4} = \frac{a}{2} \end{aligned}$$

which is the same for one dimensional infinite square well which is what we expect. The same follows for  $\langle r^2 \rangle$

$$\begin{aligned} \langle r^2 \rangle &= \int_V \psi_{k00}^* r^2 \psi_{k00} dV \\ &= \frac{2}{a} \int_0^a r^2 \sin^2(kr) dr \\ &= \frac{a^2}{6} \left( 2 - \frac{3}{(n\pi)^2} \right) \end{aligned}$$

Which means we can find the variance as

$$\langle (\Delta r)^2 \rangle = \langle r^2 \rangle - \langle r \rangle^2 = \frac{a^2}{12} \left( 1 - \frac{6}{(n\pi)^2} \right)$$

## 4 Problem #4

For a particle of mass,  $m_0$ , in the ground state of a spherically symmetric infinitely deep square potential we a wave function given by

$$\psi_{100}(r) = \frac{1}{\sqrt{2a\pi}} \frac{\sin(\pi r/a)}{r}$$

which in order to find the momentum probability distribution we find  $\psi$  in momentum representation as

$$\begin{aligned}\phi_{100}(p) &= \frac{1}{\sqrt{2\pi\hbar}} \int_V \psi_{100}(r) \exp\left(\frac{-ipr}{\hbar}\right) dV \\ &= \frac{4\pi}{2\pi\sqrt{\hbar}} \int_0^a \sin\left(\frac{\pi r}{a}\right) \exp\left(\frac{-ipr}{\hbar}\right) dr \\ &= \frac{2}{\sqrt{\hbar a}} a\pi \frac{1 + e^{-ipa/\hbar}}{\pi^2 - (pa/\hbar)^2} \\ &= 2\pi\sqrt{\frac{a}{\hbar}} \frac{1 + e^{-ipa/\hbar}}{\pi^2 - (pa/\hbar)^2}\end{aligned}$$

So the momentum probability distribution is given by

$$\begin{aligned}|\phi(p)|^2 &= \frac{4\pi^2 a}{\hbar} \frac{1 + e^{ipa/\hbar}}{\pi^2 - (pa/\hbar)^2} \frac{1 + e^{-ipa/\hbar}}{\pi^2 - (pa/\hbar)^2} \\ &= \frac{4\pi^2 a}{\hbar} \frac{2 + e^{-ipa/\hbar} + e^{ipa/\hbar}}{(\pi^2 - (pa/\hbar)^2)^2} \\ &= \frac{8\pi^2 a}{\hbar} \frac{1 + \cos(pa/\hbar)}{(\pi^2 - (pa/\hbar)^2)^2}\end{aligned}$$