Physics 615

Methods of Theoretical Physics I Professor Katrin Becker

Homework #12

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1 Problem #1

To prove the identity

$$\exp(i\vec{\sigma}\cdot\vec{n}\omega) = \cos\omega\cdot\mathbf{1} + i\vec{\sigma}\vec{n}\sin\omega \tag{1.1}$$

where σ_i are the Pauli matrices, 1 is the 2×2 identity matrix, and \vec{n} is a unit vector in \mathbb{R}^3 , we take the exponential of a matrix, A, to be defined as

$$e^{A} = \sum_{n=0}^{\infty} \frac{1}{n!} A^{n} \tag{1.2}$$

Therefore using equation 1.2 on equation 1.1 where $\vec{\sigma} \cdot \vec{n}$ is taken to be the matrix, A

$$\exp(i\vec{\sigma} \cdot \vec{n}\omega) = \sum_{n=0}^{\infty} \frac{1}{n!} (i\omega\vec{\sigma} \cdot \vec{n})^n$$

$$= \sum_{n=0}^{\infty} \frac{i^n}{n!} (\vec{\sigma} \cdot \vec{n})^n$$

$$= \sum_{n=0}^{\infty} \frac{(-1)^{2n} \omega^{2n}}{(2n)!} (\vec{\sigma} \cdot \vec{n})^{2n} + i \sum_{n=0}^{\infty} \frac{(-1)^{2n+1} \omega^{2n+1}}{(2n+1)!} (\vec{\sigma} \cdot \vec{n})^{2n+1}$$

We note that the Pauli matrices are given as

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \qquad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \qquad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Which implies that

$$\vec{\sigma} \cdot \vec{n} = \frac{1}{\sqrt{3}} \left(\begin{array}{cc} 1 & 1-i \\ 1+i & -1 \end{array} \right)$$

Note the factor of $\sqrt{3}$ comes from the normalization of n in \mathbb{R}^3 . Using this we can see that

$$(\vec{\sigma} \cdot \vec{n})^2 = \frac{1}{3} \begin{pmatrix} 1 & 1-i \\ 1+i & -1 \end{pmatrix} \begin{pmatrix} 1 & 1-i \\ 1+i & -1 \end{pmatrix}$$

$$= \frac{1}{3} \begin{pmatrix} 1+(1-i)(1+i) & (1-i)-(1-i) \\ (1+i)-(1+i) & (1+i)(1-i)+1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \mathbf{1}$$

And we can generalize to all even powers to say

$$(\vec{\sigma} \cdot \vec{n})^{2n} = \mathbf{1}$$

Therefore

$$\exp(i\vec{\sigma}\cdot\vec{n}\omega) = \sum_{n=0}^{\infty} \frac{(-1)^{2n}\omega^{2n}}{(2n)!} \mathbf{1} + i\sum_{n=0}^{\infty} \frac{(-1)^{2n+1}\omega^{2n+1}}{(2n+1)!} (\vec{\sigma}\cdot\vec{n}) \mathbf{1}$$
$$= \cos\omega \cdot \mathbf{1} + i\vec{\sigma}\vec{n}\sin\omega$$

2 Problem #2

Given Hermitian matrices A, B and unitary matrices C, D. Which implies that

$$A=A^{\dagger}, \qquad B=B^{\dagger}, \qquad C^{\dagger}C=CC^{\dagger}=1 \qquad , D^{\dagger}D=DD^{\dagger}=1$$

1) We can show that

$$(C^{-1}AC)^{\dagger} = C^{\dagger}A^{\dagger} (C^{-1})^{\dagger}$$
$$= C^{-1}A^{\dagger}C$$
$$= C^{-1}AC$$

Therefore $C^{-1}AC$ is Hermitian.

2) We can show that

$$\begin{split} C^{-1}DC(C^{-1}DC)^{\dagger} &= C^{-1}D\mathcal{C}C^{\dagger}D^{\dagger}(C^{-1})^{\dagger} \\ &= C^{-1}DD^{\dagger}(C^{-1})^{\dagger} \\ &= C^{-1}(C^{-1})^{\dagger} \\ &= C^{\dagger}C = 1 \end{split}$$

Therefore C^1DC is a unitary matrix.

3) We can show that

$$(i(AB - BA))^{\dagger} = -i((AB)^{\dagger} - (BA)^{\dagger})$$
$$= -i(B^{\dagger}A^{\dagger} - A^{\dagger}B^{\dagger})$$
$$= -i(BA - AB) = i(AB - BA)$$

Therefore i(AB - BA) is Hermitian.

3 Problem #3

Given the unitary matrix 1 we can find the eigenvalues λ of this matrix by enforcing the condition

$$\det(\mathbf{A} - \lambda \mathbf{1}) = 0$$

where we take A to be the unitary matrix. This implies that

$$\det(\mathbf{1} - \lambda \mathbf{1}) = \det((1 - \lambda)\mathbf{1}) = (1 - \lambda)^n$$

where n is the rank of the unitary matrix. So the equation

$$(1-\lambda)^n = 0$$

implies that $|\lambda| = 1$ or that λ is unimodular.

4 Problem #4

For the center, Z, of a group, G, defined as the set of elements $z \in G$ which commute with all elements within G

$$Z = \{ z \in G | zg = gz, \ \forall g \in G \}$$

We note that the group Z is an Abelian group by definition. This is due to the fact for elements $f, g \in Z$ by definition we have fg = gf. Therefore we have a commutative operator (multiplication).