

Physics 607  
Statistical Physics and Thermodynamics  
Professor Valery Pokrovsky

Homework #8

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# 1 Problem #1

- (1) To find the Debye screening length,  $\lambda_D$ , of a semiconductor with a dielectric constant  $\varepsilon = 20$  and density of electrons  $n_0 = 10^{18} \text{ cm}^{-3}$  at room temperature,  $T = 300 \text{ K}$  we can use the reciprocal length,  $\kappa$ , given as

$$\kappa^2 = \frac{4\pi e^2}{T} \sum_a n_{a0} z_a^2$$

Noting that for a semiconductor the ions are fixed so we only consider the electrons in the sum. This implies that

$$\lambda_D = \frac{1}{\kappa} = \sqrt{\frac{\varepsilon \varepsilon_0 k_B T}{4\pi n_0 e^2}}$$

note that we added the constants  $k_B$  and  $\varepsilon_0$  to get the correct dimensionality. So for the semiconductor we calculate

$$\begin{aligned} \lambda_D &= \sqrt{\frac{\varepsilon \varepsilon_0 k_B T}{4\pi n_0 e^2}} = \sqrt{\frac{(20)(8.85 \times 10^{-12} \text{ C}^2 \text{ N}^{-1} \text{ m}^{-2})(1.38 \times 10^{-23} \text{ J K}^{-1})(300 \text{ K})}{4\pi(10^{24} \text{ m}^{-3})(1.6 \times 10^{-19} \text{ C})^2}} \\ &= 1.51 \text{ nm} \end{aligned}$$

- (2) For a semiconductor in the form of a rectangular slab with thickness  $L$ , where  $L$  is much smaller than the other two dimensions, we attach two electrodes to the faces of the slab and apply a bias voltage,  $V$ . To solve for the electric field we solve the *Poisson Equation* given by

$$\nabla^2 \phi - \kappa^2 \phi = 0 \tag{1.1}$$

where  $\kappa$  is defined above. Taking the directions that create the face as infinite we assume that  $\phi$  is only dependent on  $z$  which yields a solution of the form

$$\phi(z) = C_1 e^{\kappa z} + C_2 e^{-\kappa z}$$

applying the boundary conditions

$$\phi(z=0) = 0 \quad \phi(z=L) = V$$

we find that

$$\begin{aligned} \phi(z=0) = 0 &= C_1 e^{\kappa 0} + C_2 e^{-\kappa 0} \\ &\Downarrow \\ C_1 &= -C_2 \\ &\Downarrow \\ \phi(z) &= 0 = C (e^{\kappa z} - e^{-\kappa z}) = C \sinh(\kappa z) \end{aligned}$$

and

$$\begin{aligned} \phi(z=L) &= V = C \sinh(\kappa L) \\ &\Downarrow \\ C &= \frac{V}{\sinh(\kappa L)} \end{aligned}$$

So the potential as a function of  $z$  is

$$\phi(z) = \frac{V}{\sinh(\kappa L)} \sinh(\kappa z)$$

(3) Note if we take the limit as the thickness becomes infinite we have the following condition

$$\kappa L \rightarrow \infty$$

we see that in this limit we have

$$\lim_{\kappa L \rightarrow \infty} \frac{V}{\sinh(\kappa L)} = 0$$

So we see that as we take the thickness to be infinite the screening will completely cancel the bias voltage.

## 2 Problem #2

- (a) Given the physical characterization of the triple point of water found experimentally as

$$T_t = 273.16 \text{ K} \quad P_t = 612 \text{ Pa} \quad \rho_L = 1 \text{ g cm}^{-3} \quad \rho_S = 0.894 \text{ g cm}^{-3}$$

where  $\rho_L$  is the density of water in the liquid state and  $\rho_S$  is the density of water in the solid state. Noting the data for the freezing point of water at atmospheric pressure is

$$T = 273.15 \text{ K} \quad P_t = 101 \text{ kPa} \quad \rho_L = 1 \text{ g cm}^{-3} \quad \rho_S = 0.894 \text{ g cm}^{-3}$$

we can find the heat of ice melting assuming that it does not change between triple point and the freezing point at normal pressure by taking the *Clapeyron-Clausius formula*

$$\frac{dP}{dT} = \frac{q}{T(v_2 - v_1)} \quad (2.1)$$

and solving for the heat,  $q$ . Note we can approximate the derivative by taking the two different data points from the freezing point to the triple point. So equation 2.1 becomes

$$\begin{aligned} q &= \frac{dP}{dT} T(v_2 - v_1) \\ &= \frac{P_t - P}{T_t - T} T \left( \frac{1}{\rho_L} - \frac{1}{\rho_S} \right) \\ &= \frac{612 \text{ Pa} - 101000 \text{ Pa}}{273.16 \text{ K} - 273.15 \text{ K}} (273.15 \text{ K}) \left( \frac{1}{1 \times 10^6 \text{ g m}^{-3}} - \frac{1}{0.894 \times 10^6 \text{ g m}^{-3}} \right) \\ &= 325 \text{ J g}^{-1} \end{aligned}$$

- (b) Using equation 2.1 and the given value for the latent heat of vaporization

$$q = 2264.76 \text{ J g}^{-1}$$

we can calculate the derivative of boiling temperature

$$\begin{aligned} \frac{dT_b}{dP} &= \frac{T(v_2 - v_1)}{q} \\ &= \frac{T(\rho_V^{-1} - \rho_L^{-1})}{q} \end{aligned}$$

and the same follows for ice vaporization

$$\begin{aligned} \frac{dT_v}{dP} &= \frac{T(v_2 - v_1)}{q} \\ &= \frac{T(\rho_V^{-1} - \rho_S^{-1})}{q} \end{aligned}$$

where we take  $\rho_V$  as the density of water vapor.

- (c) When we take water at the freezing temperature,  $T = 273.15 \text{ K}$ , and atmospheric pressure,  $P = 103 \text{ kPa}$ , and decrease the pressure until the pressure reaches  $P' = 300 \text{ Pa}$  We see that this patch will pass through the triple point described in part (a). This implies that the water will skip the liquid phase and sublime from solid to vapor.