

# [Solutions] DMFT+NRG

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## Solution to Exercise (a): Complete the DMFT+NRG function

The complete version of the function is uploaded as `NRG/DMFT_Hubbard.m`. Compare with your implementation!

## Solution to Exercise (b): Obtain the DMFT+NRG solutions starting from the $U = 0$ result

*Note:* the computing time for Exercises (b) and (c) is quite huge and many log messages are generated during the computation. I suppress them in this solution.

The numerical result for this Exercise can be obtained by the following steps. First, we treat the  $U = 0$  case to obtain the initial hybridization function.

```
clear

% system parameters
U = 0;
mu = 0;
T = 1e-4;
D = 1;

% initial hybridization function
ozin = [-1;1];
RhoV2in = [1;1]/2*(D/2)^2;

% NRG parameters
Lambda = 4;
N = 30;
Nkeep = 300;

% run DMFT+NRG
[ocont,RhoV2s,iscvg] = DMFT_Hubbard (U,mu,T,D,ozin,RhoV2in,Lambda,N,Nkeep);

RhoV2init_met = RhoV2s(:,end); % initial hyb. func.
```

Now we treat finite  $U$ 's.

```
Us = (2:0.3:3.5);
mus = Us/2; % half filling

RhoV2_met = cell(1,numel(Us));

for itU = (1:numel(Us))
    [~,RhoV2_met{itU}] = DMFT_Hubbard(Us(itU),mus(itU),T,D, ...
```

```

ocont,RhoV2init_met,Lambda,N,Nkeep);
end

```

How do the results look like?

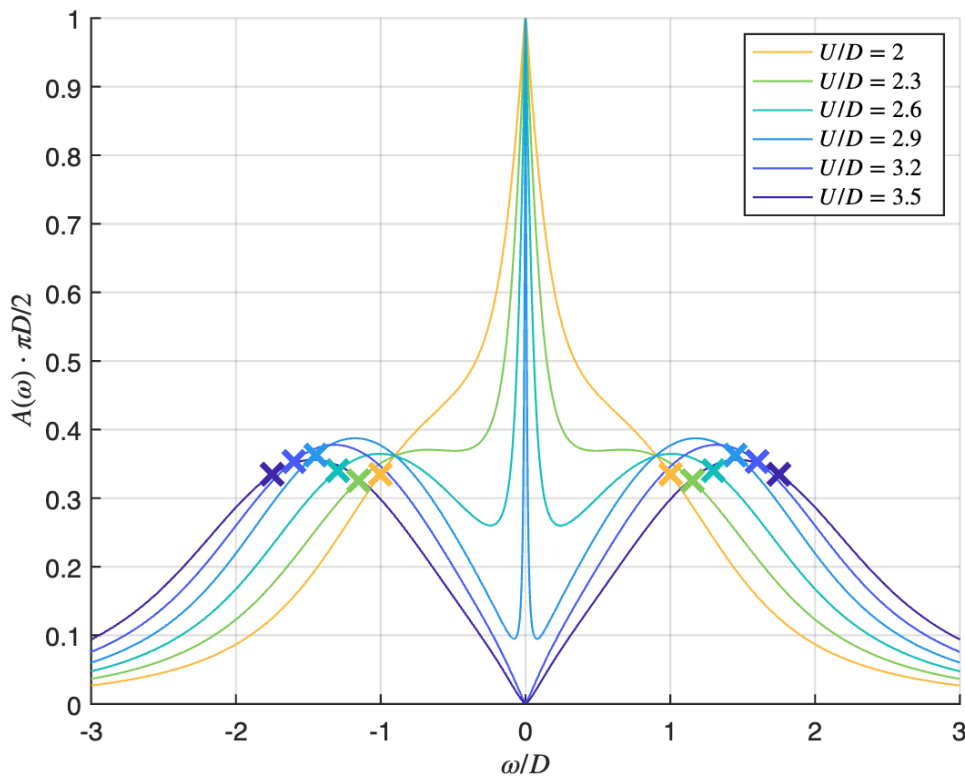
```

cmap = parula(round(numel(Us)*1.2));
cmap = flipud(cmap(1:numel(Us),:)); % to avoid too bright colors
legs = cell(1,numel(Us));

figure;
hold on;
for itU = (1:numel(Us))
    plot(ocont/D,RhoV2_met{itU}(:,end)/((D/2)^2)*(pi*D/2), ...
        'LineWidth',1,'Color',cmap(itU,:));
    legs{itU} = ['$U/D = ',sprintf('%.4g',Us(itU)/D),'$'];
end
% guide to the eye; mark \omega = \pm U/2
for itU = (1:numel(Us))
    plot([-1;1]*(Us(itU)/2/D),interp1(ocont/D, ...
        RhoV2_met{itU}(:,end)/((D/2)^2)*(pi*D/2), ...
        [-1;1]*(Us(itU)/2/D),'linear'),'x', ...
        'LineWidth',3,'MarkerSize',15,'Color',cmap(itU,:));
end
hold off;

grid on;
set(gca,'LineWidth',1,'FontSize',13);
legend(legs,'Interpreter','latex');
xlabel('$\omega / D$', 'Interpreter','latex');
ylabel('$A(\omega) \cdot \pi D/2$', 'Interpreter','latex');
xlim([-1 1]*(3*D));

```



For  $U/D \leq 2.9$ , the spectral function  $A(\omega)$  has a peak centered at  $\omega = 0$ , and the peak gets narrowed as  $U$  increases. This peak is called the quasiparticle peak, and its presence means that the system is metallic. (Hence `_met` in the variable names `RhoV2init_met` and `RhoV2_met`.) The height of the quasiparticle peak at  $T = 0$  is pinned to  $2/\pi D$ , independent of  $U$ . Indeed, the height is dictated by the Luttinger theorem, which relates the height with the filling and the non-interacting density of states.

Moreover, for  $U/D \leq 2.9$  there are two side peaks centered around  $\omega = \pm U/2$ . These side peaks are called the Hubbard bands. For  $U/D \leq 2.3$ , the Hubbard bands look like mere shoulders, but for larger  $U$  they are better separated from the quasiparticle peak.

For  $U/D \geq 3.2$ , the quasiparticle peak disappears and only the Hubbard bands remain. There is a gap at  $\omega = 0$ , which is the Mott gap. So the system is insulating.

We see that for  $U/D \leq 2.9$  the DMFT loop takes more iterations as  $U$  increases.

```
disp(Us/D);
```

```
2.0000 2.3000 2.6000 2.9000 3.2000 3.5000
```

```
disp(cellfun(@(x) size(x,2), RhoV2_met));
```

```
5 7 10 23 21 14
```

Why? Let's plot how the curves change as the iteration proceeds for  $U/D = 2.9$ .

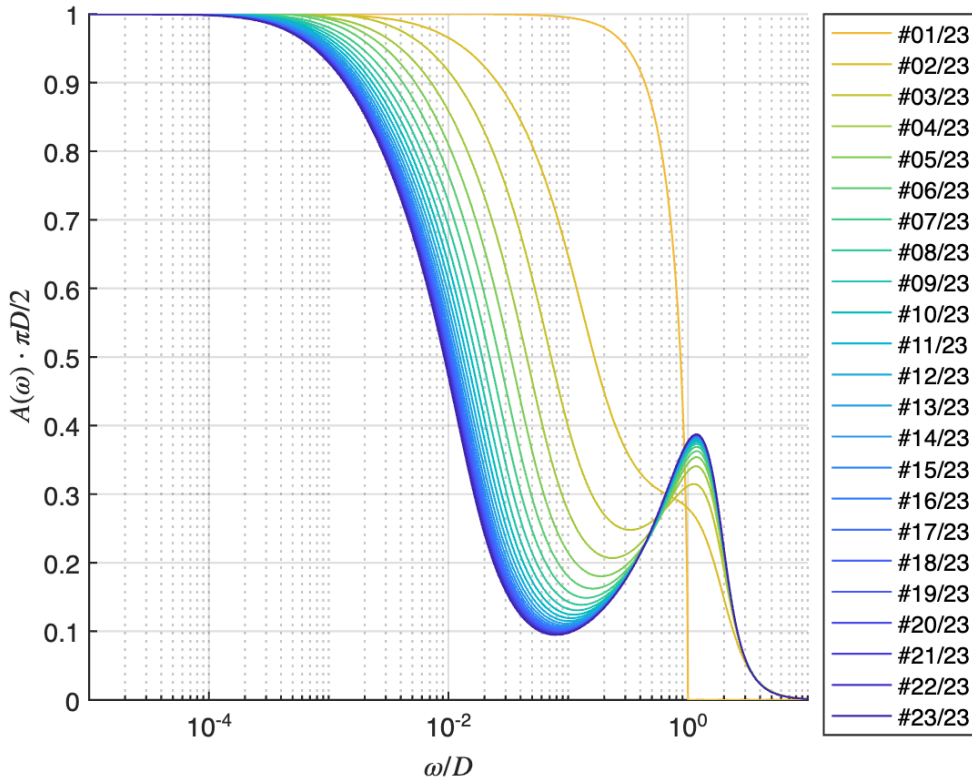
```

itU = 4;
cmap = parula(round(size(RhoV2_met{itU},2)*1.2));
cmap = flipud(cmap(1:size(RhoV2_met{itU},2),:)); % to avoid too bright colors
legs = cell(1,size(RhoV2_met{itU},2));

figure;
hold on;
for itd = (1:size(RhoV2_met{itU},2))
    plot(ocont/D,RhoV2_met{itU}(:,itd)/((D/2)^2)*(pi*D/2), ...
        'LineWidth',1,'Color',cmap(itd,:));
    legs{itd} = sprintf('#%02i/%02i',[itd size(RhoV2_met{itU},2)]);
end
hold off;

grid on;
set(gca,'LineWidth',1,'FontSize',13,'XScale','log');
legend(legs,'Location','eastoutside');
xlabel('$\omega / D$', 'Interpreter','latex');
ylabel('$A(\omega) \cdot \pi D/2$', 'Interpreter','latex');
xlim([0.1*T/D 10]);

```



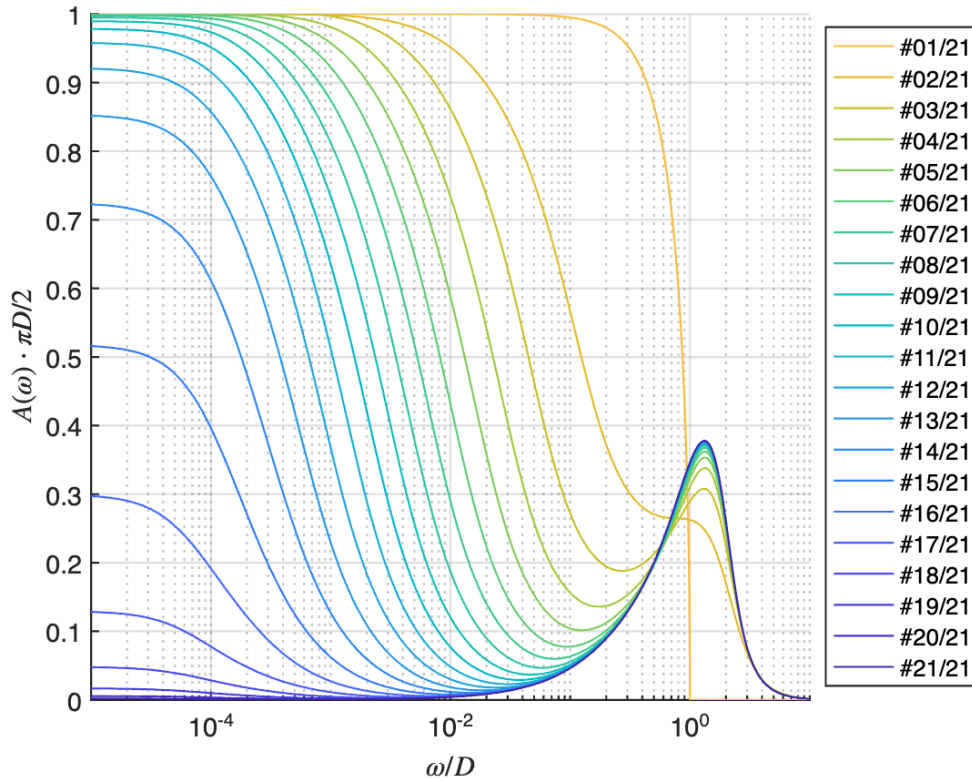
The peak centered at  $\omega = 0$  gets narrower as the iteration proceeds, until it reaches the converged shape of the quasiparticle peak. For larger  $U$  (while being in the metallic phase), the quasiparticle peak is narrower, so it takes more iterations to converge.

How about the case of  $U/D = 3.2$ ?

```
itU = 5;
cmap = parula(round(size(RhoV2_met{itU},2)*1.2));
cmap = flipud(cmap(1:size(RhoV2_met{itU},2),:)); % to avoid too bright colors
legs = cell(1,size(RhoV2_met{itU},2));

figure;
hold on;
for itd = (1:size(RhoV2_met{itU},2))
    plot(ocont/D,RhoV2_met{itU}(:,itd)/((D/2)^2)*(pi*D/2), ...
        'LineWidth',1,'Color',cmap(itd,:));
    legs{itd} = sprintf('#%02i/%02i',[itd size(RhoV2_met{itU},2)]);
end
hold off;

grid on;
set(gca,'LineWidth',1,'FontSize',13,'XScale','log');
legend(legs,'Location','eastoutside');
xlabel('$\omega / D$', 'Interpreter', 'latex');
ylabel('$A(\omega) \cdot \pi D/2$', 'Interpreter', 'latex');
xlim([0.1*T/D 10]);
```



The central peak narrows as the iteration goes, and once its width becomes comparable with  $T = 10^{-4}D$ , the peak height lowers. Then finally the peak disappears, opening the Mott gap.

## Solution to Exercise (c): Obtain the DMFT+NRG solutions starting from the $U = 3.5D$ result

We continue from the script executed for Exercise (b) above. We take the converged hybridization function for  $U = 3.5D$  as the initial hybridization function.

```
RhoV2init_ins = RhoV2_met{end}(:,end);

RhoV2_ins = cell(1,numel(Us));

for itU = (1:numel(Us))
    [~,RhoV2_ins{itU}] = DMFT_Hubbard(Us(itU),mus(itU),T,D, ...
        ocont,RhoV2init_ins,Lambda,N,Nkeep);
end
```

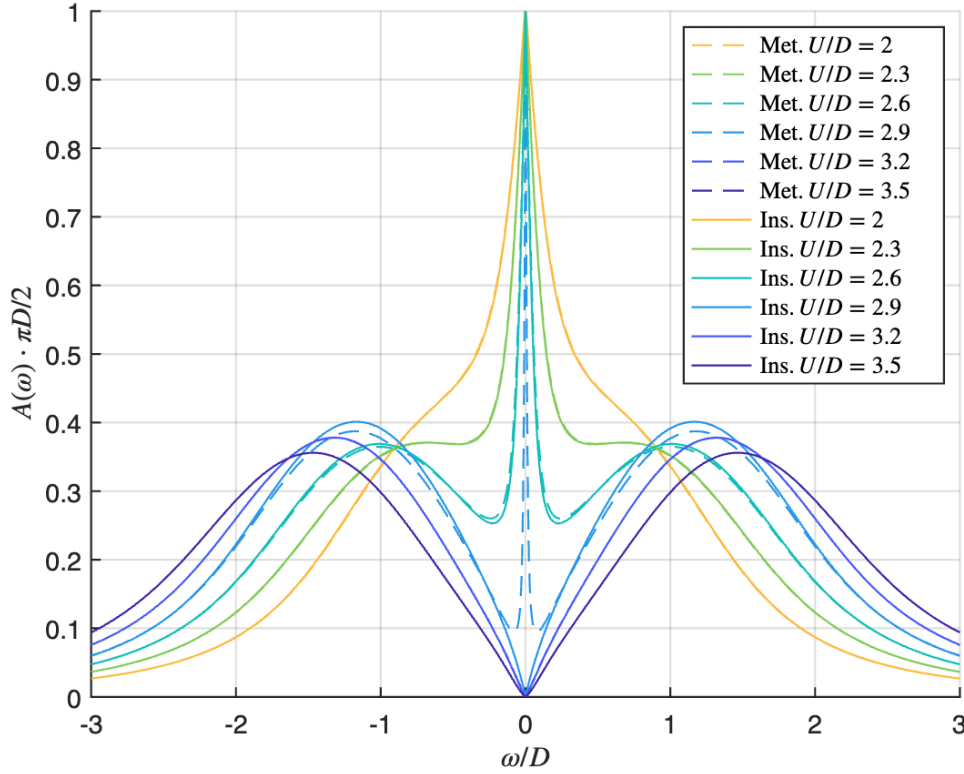
Here we use the suffix `_ins` in the variable names, since we start calculations from the insulating solution for  $U = 3.5D$ .

Plot the converged spectral functions. To compare with the results from Exercise (b), we plot them altogether.

```
cmap = parula(round(numel(Us)*1.2));
cmap = flipud(cmap(1:numel(Us),:)); % to avoid too bright colors
legs = cell(numel(Us),2);

figure;
hold on;
for itU = (1:numel(Us))
    plot(ocont/D,RhoV2_met{itU}(:,end)/((D/2)^2)*(pi*D/2), ...
        'LineWidth',1,'LineStyle','--','Color',cmap(itU,:));
    legs{itU,1} = ['Met.  $U/D =$ ',sprintf('%.4g',Us(itU)/D),'$'];
end
for itU = (1:numel(Us))
    plot(ocont/D,RhoV2_ins{itU}(:,end)/((D/2)^2)*(pi*D/2), ...
        'LineWidth',1,'LineStyle','-','Color',cmap(itU,:));
    legs{itU,2} = ['Ins.  $U/D =$ ',sprintf('%.4g',Us(itU)/D),'$'];
end
hold off;

grid on;
set(gca,'LineWidth',1,'FontSize',13);
legend(legs(:).','Interpreter','latex');
xlabel('$\omega / D$', 'Interpreter','latex');
ylabel('$A(\omega) \cdot \pi D/2$', 'Interpreter','latex');
xlim([-1 1]*(3*D));
```



We find that the results for  $U/D = 2.9$  differ depending on the choice of the initial hybridization. If the initial hybridization is metallic (insulating), the result is metallic (insulating). The presence of the two self-consistent solutions means that the system is in the co-existence regime for  $U/D = 2.9$ . In experiments, the co-existence can be confirmed by observing hysteresis.

According to the state-of-the-art DMFT+NRG calculation [[S.-S. B. Lee, J. von Delft, and A. Weichselbaum, Phys. Rev. Lett. \*\*119\*\*, 236402 \(2017\)](#)], the system can be only metallic for  $U < U_{c1} = 2.37(2)D$  and only insulating for  $U > U_{c2} = 2.91(1)D$ . Between  $U_{c1}$  and  $U_{c2}$ , it is the co-existence regime.