

Course: EE3013/ Semiconductor Devices and Processing
School: School of Electrical and Electronic Engineering
Week 10 – PN Junction Diodes

At the end of this lesson, you will be able to:

- Explain pn junction diode theory and operation mechanisms under different biasing conditions namely thermal equilibrium, forward bias and reverse bias
- Under different bias:
 - Draw the energy band diagram of the pn junction
 - Explain the concepts and determine the key parameters such as built-in voltage, depletion width, space charge, electric field, and junction capacitance
- Explain the concepts of injected minority carriers and current distributions, and calculation of electron and hole diffusion currents under different bias conditions
- Explain how useful capacitance versus voltage characteristics is to determine the built-in voltage and the impurity concentration in the semiconductor

- Semiconductors are mainly classified into two categories:
 - **Intrinsic** - Is chemically very pure and possesses poor conductivity, and
 - **Extrinsic** - Is an improved intrinsic semiconductor with a small amount of impurities which modifies its electrical properties and improves its conductivity.

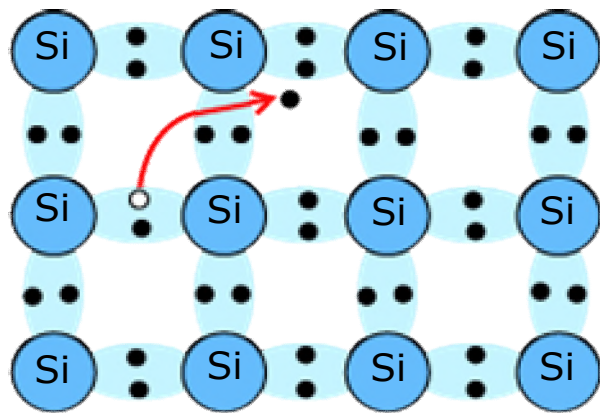


Figure 1.1 Intrinsic

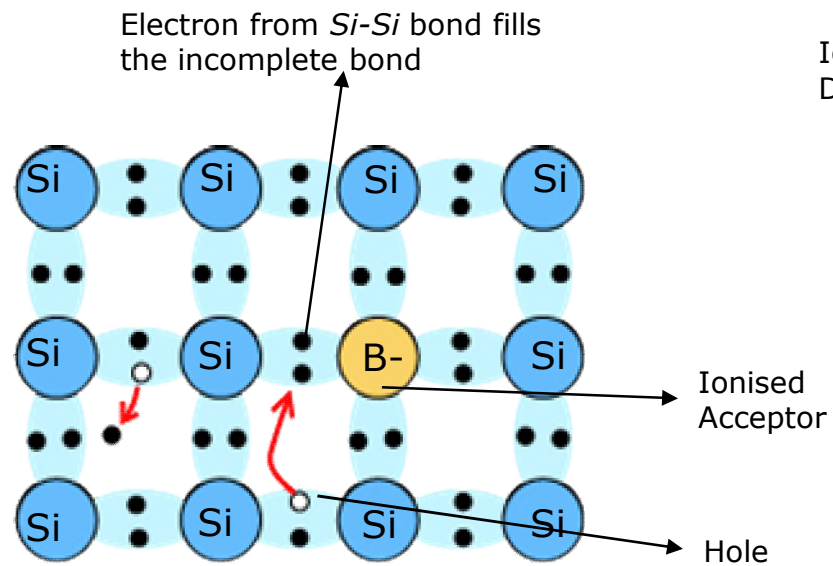


Figure 1.2 Extrinsic (p-type)

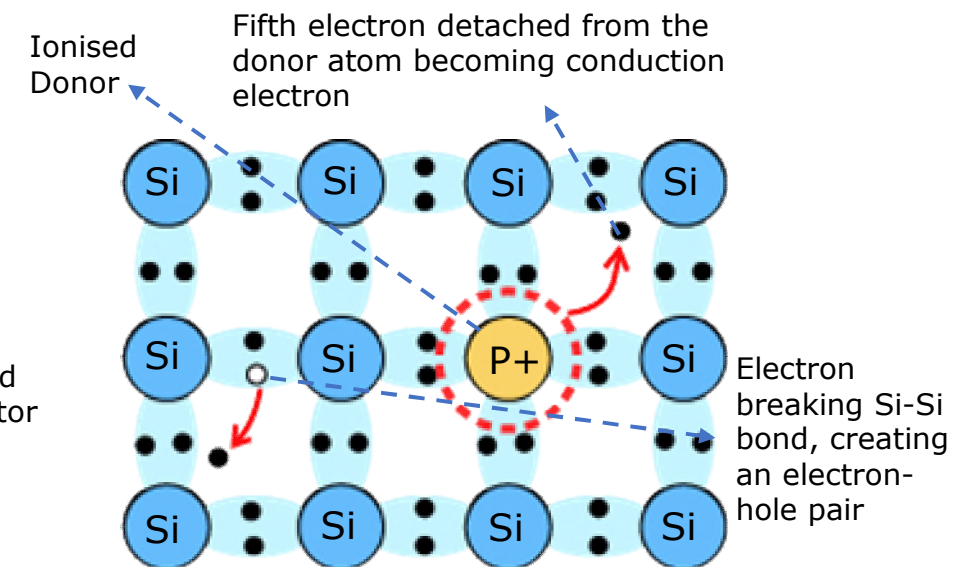


Fig. 1.3. Extrinsic (n-type)

- **P-type** - The addition of trivalent impurities (Group III elements such as Boron) to an intrinsic silicon creates deficiencies of valence electrons, called **holes**.
- **N-type** - The addition of pentavalent impurities (Group V elements such as Phosphorus) contributes free **electrons**.

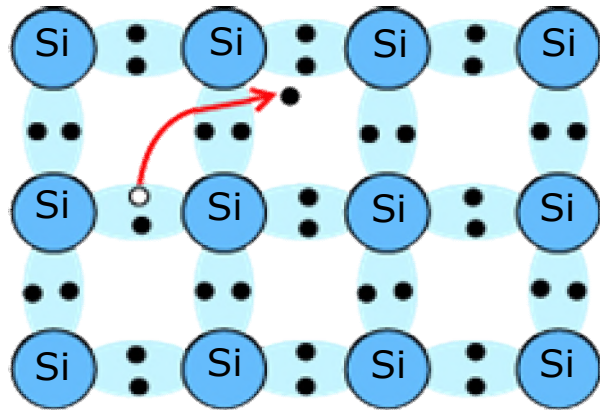


Figure 1.1 Intrinsic

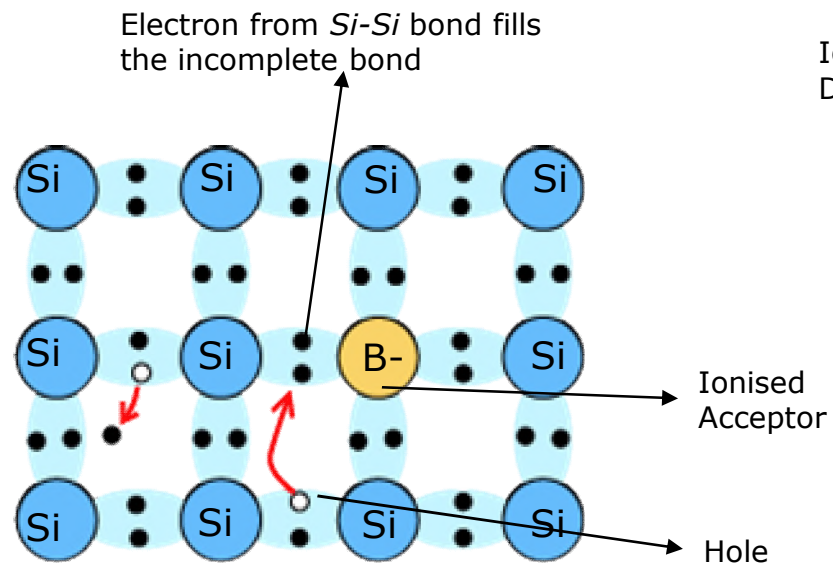


Figure 1.2 Extrinsic (p-type)

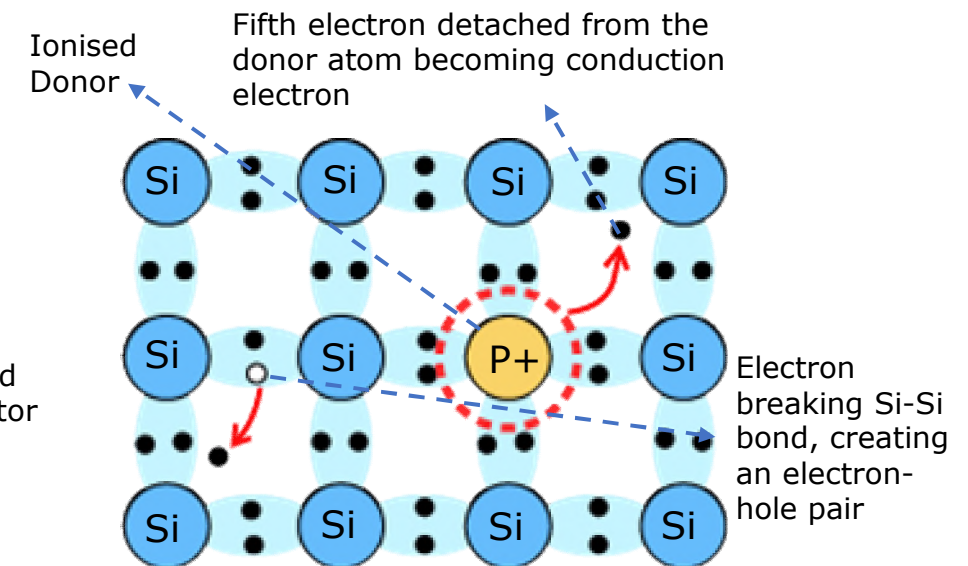


Fig. 1.3. Extrinsic (n-type)

- There are two basic transport mechanisms in a semiconductor crystal depending on the driving force:
 - Drift:** movement of charge due to electric fields
 - Diffusion:** flow of charge due to concentration gradients

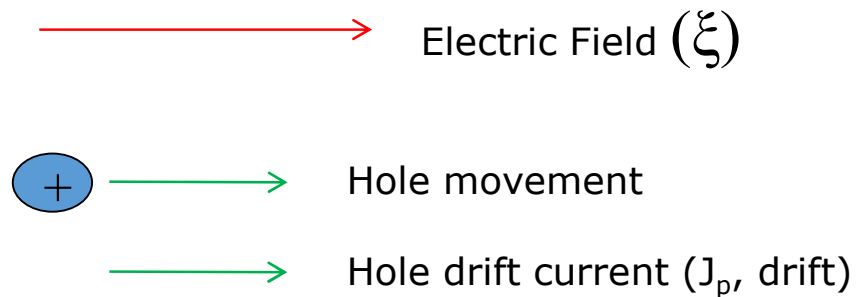


Figure 1.4 Drift current

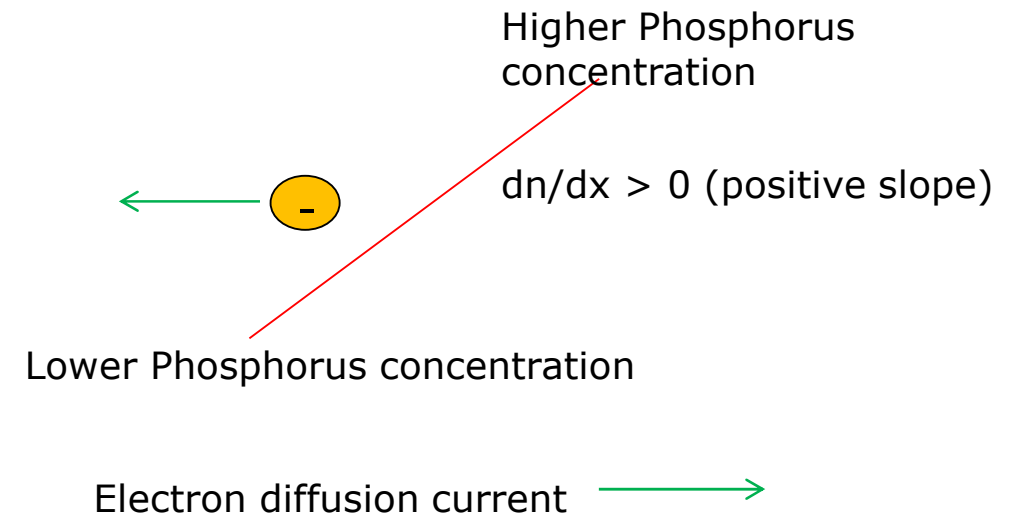
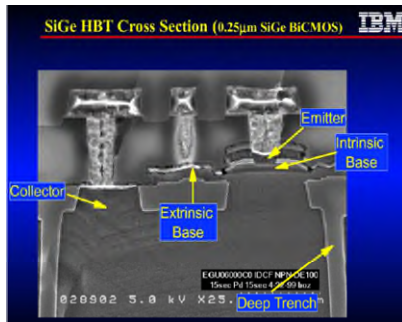
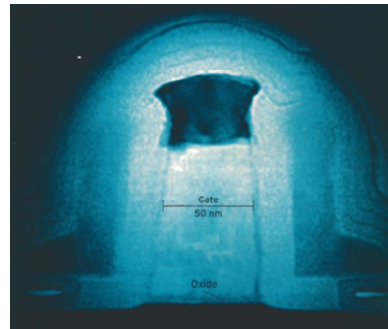


Figure 1.5 Diffusion current

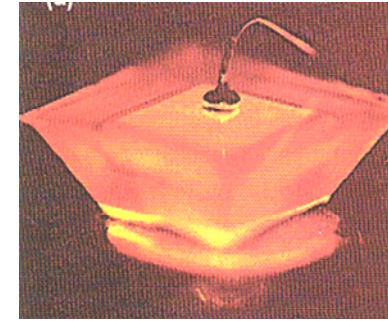
- Some basic semiconductor devices are **BJT, MOSFET, LED, laser, and photo-detectors**.
 - Electronic devices: **BJT, MOSFET**
 - Optoelectronic devices: **LED, laser, and photo-detectors**



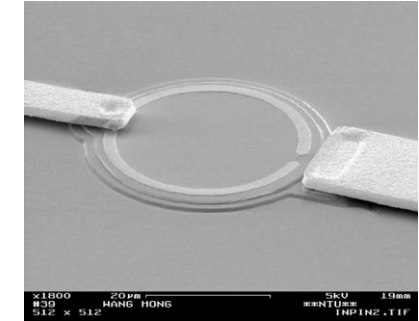
IBM's SiGe Bipolar Transistor



Intel's 65nm MOSFET



Red LED



Photodetector

Figure 1.6 Electronic devices

Figure 1.7 Optoelectronic devices

- In this course, we will mainly deal with electronic devices.

- **PN junction** is a boundary or interface between two types of semiconductor material, p-type and n-type.



Figure 1.8 PN Junction

- PN junction serves an important role in:
 - Modern electronic and optoelectronic applications (used in rectification, switching, light emission and detection, and other operations) and
 - Understanding other semiconductor devices because it is the key building block for most of the devices, such as LEDs, photodetector, MOSFETs and BJTs, etc. as mentioned in previous slide.

Diode I-V Characteristics

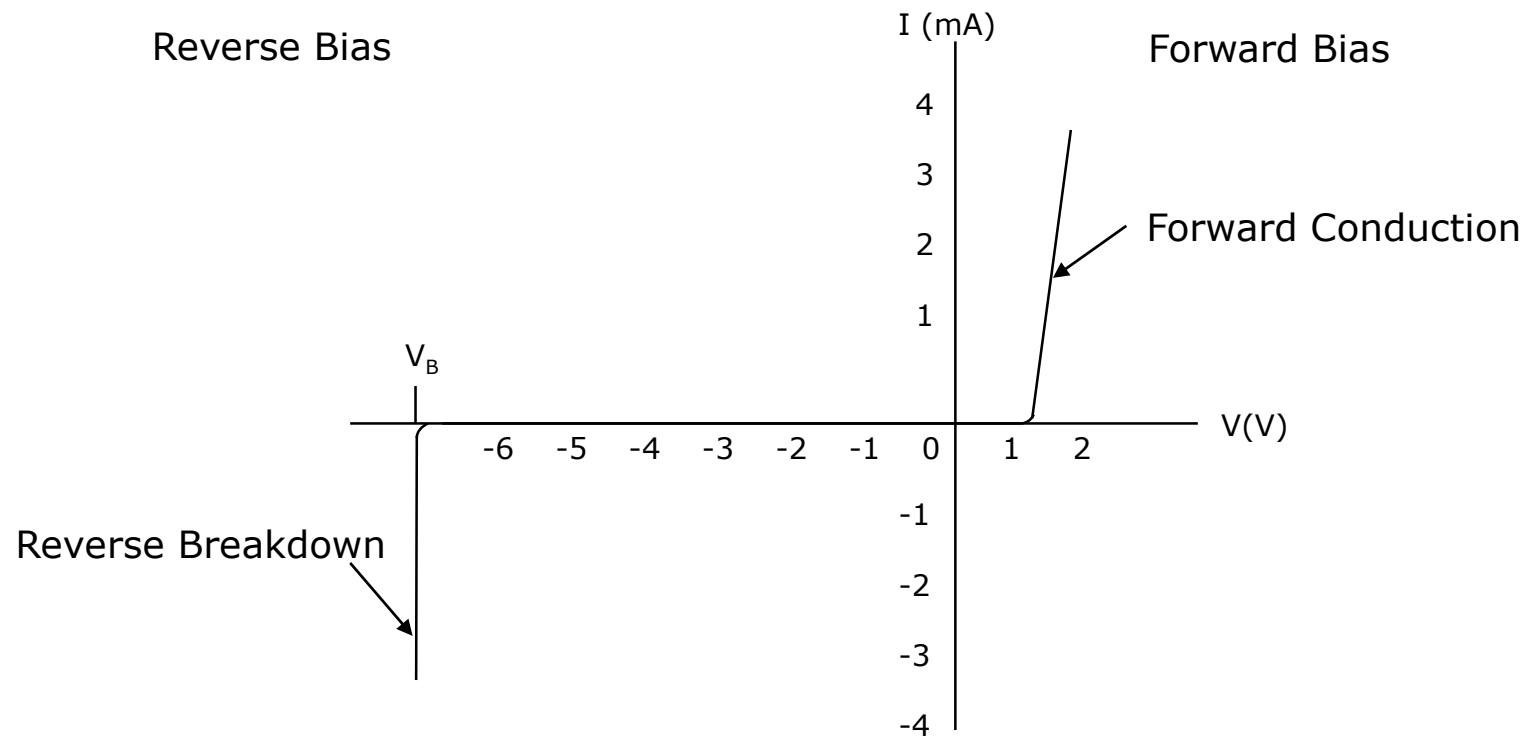


Figure 1.9 Diode I-V Characteristics

- Forward bias: p-side positive, I increases exponentially with V
- Reverse bias: p-side negative, current nearly 0 till junction breakdown
- A PN junction allows current to flow easily in only one direction

(Equation 1.1)
$$I = I_s (e^{qV/kT} - 1)$$

where I_s is the reverse saturation current and V is the bias applied

P- and N-type Semiconductors

- Assume the p-type semiconductor materials are uniformly doped.

p-type silicon

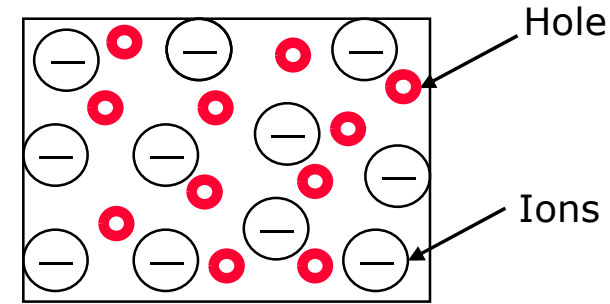
Doped by acceptors (Boron)

Majority carriers

(Equation 1.2) $p = N_A = n_i e^{\frac{E_i - E_F}{kT}}$

Minority carriers

(Equation 1.3) $n = \frac{n_i^2}{N_A}$ where $p \gg n$



p

E_c —————

E_F - - - - -

E_V 

P- and N-type Semiconductors

- Assume the n-type semiconductor materials are uniformly doped.

n-type silicon

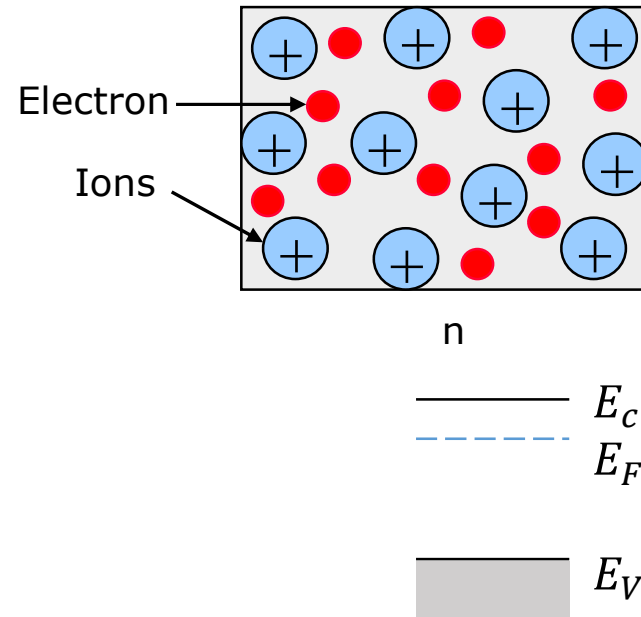
Doped by donors (Phosphorus)

Majority carriers

(Equation 1.4) $n = N_D = n_i e^{\frac{E_F - E_i}{kT}}$

Minority carriers

(Equation 1.5) $p = \frac{n_i^2}{N_D}$ where $n \gg p$



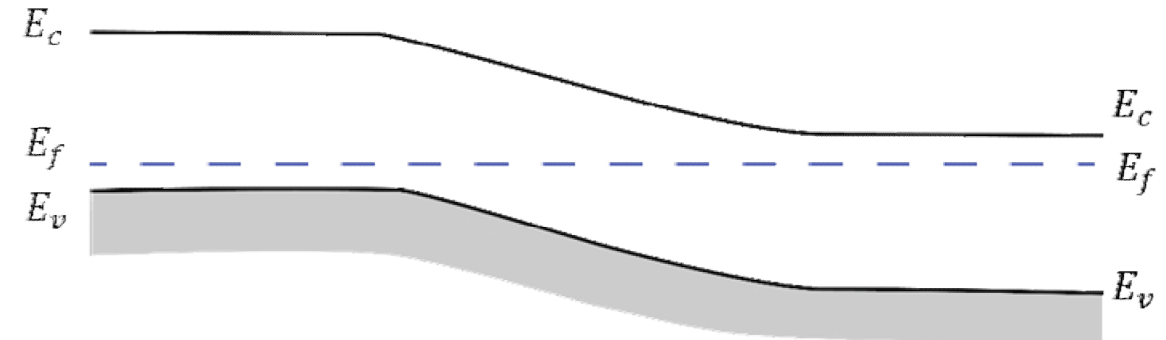
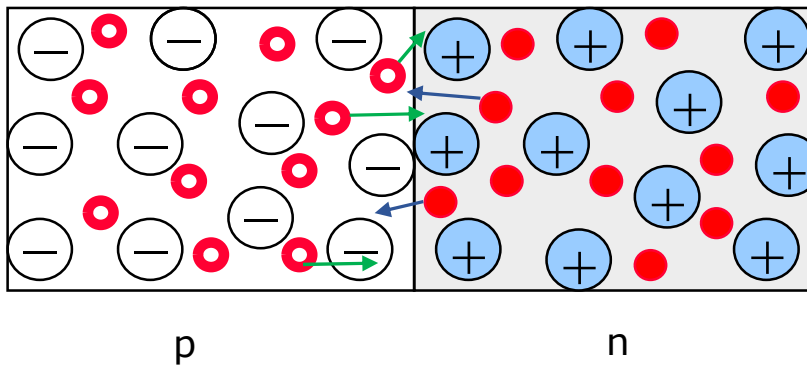
PN Junction Formation

- When p- and n-type semiconductors contact together, large carrier concentration gradients at the junction cause carrier diffusion (diffusion currents J_p and J_n).

Holes flow from left to right

Hole current density J_p **to right**

$$J_p = -qD_p \frac{dp}{dx} \quad (\text{Equation 1.6})$$



(Equation 1.7) $J_n = -qD_n \frac{dn}{dx}$

Electrons flow to left

But electron current density J_n to right

PN Junction Formation (Cont'd.)

- As holes leave the p-side, some negative acceptor ions (N_A^-) near the junction are left uncompensated (since the acceptors are fixed in the semiconductor lattice while the holes are mobile).
- Similarly, some positive donor ions (N_D^+) are left as electrons leave n-side.

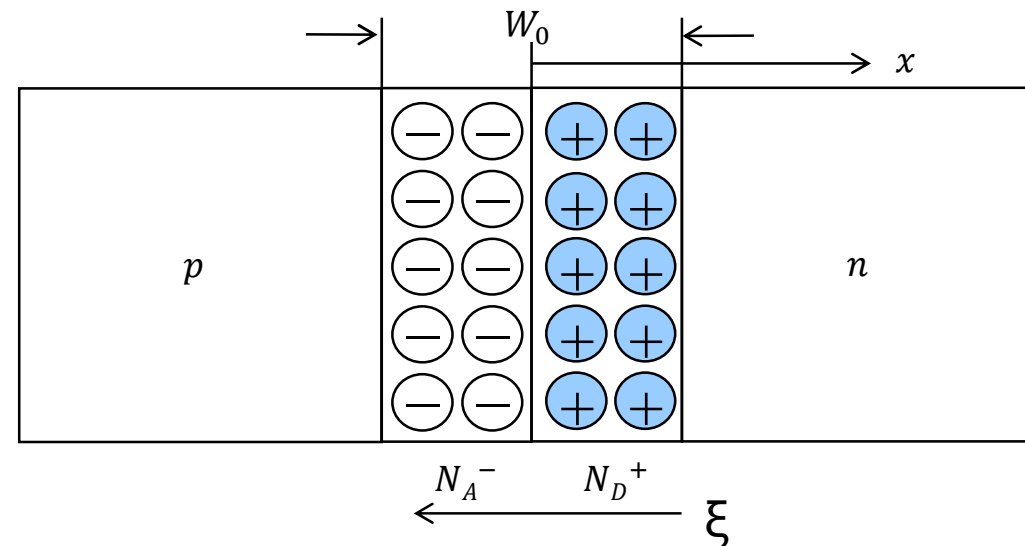
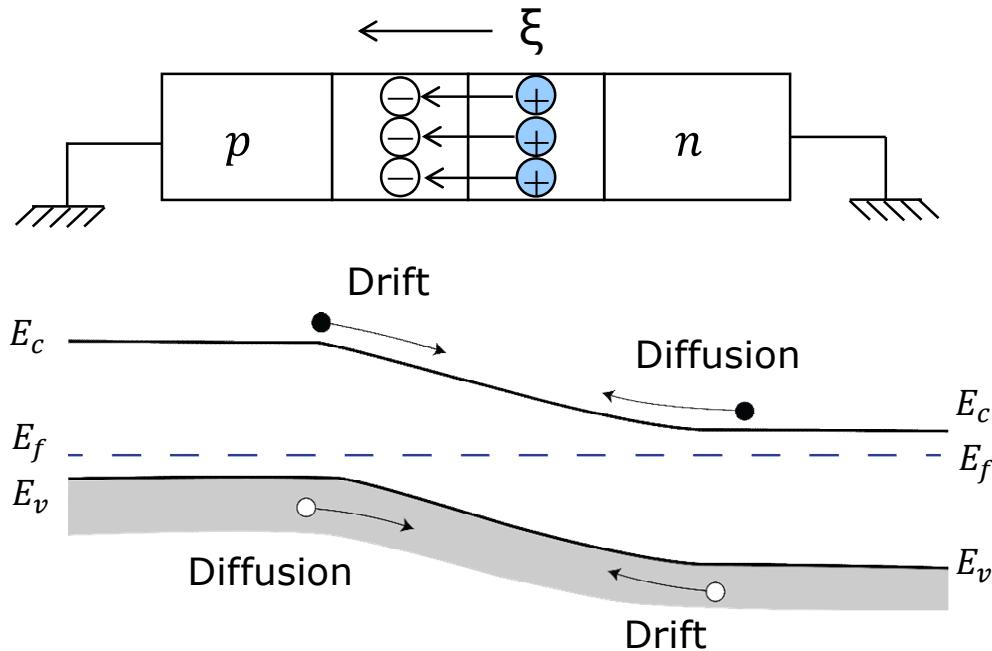


Figure 1.10

- Negative and positive space charges form a depletion region.
- They create an **electrical field**, ξ , directed from positive charges towards negative charges.

PN Junction Formation (Cont'd.)

- The electrical field creates two drift current components, $J_p(\text{drift})$ and $J_n(\text{drift})$, and in the directions opposite to the diffusion current components.



$$J_n(\text{drift}) = J_n \text{ diffusion} \quad (\text{Equation 1.8})$$

$$J_p(\text{drift}) = J_p \text{ diffusion} \quad (\text{Equation 1.9})$$

- At thermal equilibrium, (that is, the steady-state condition at a given temperature without any external excitation), the net electron and hole currents flowing across the junction are zero.

PN Junction Formation (Cont'd.)

- That is, each type of drift current must exactly cancel the diffusion current.

$$J_p = q \mu_p p \xi - q D_p \frac{dp}{dx} = 0 \quad (\text{Equation 1.10})$$

$$J_n = q \mu_n n \xi + q D_n \frac{dn}{dx} = 0 \quad (\text{Equation 1.11})$$

Where,

$D_p(D_n)$ is the diffusion coefficient of holes (electrons),

$\mu_p(\mu_n)$ is hole (electron) mobility,

$P(n)$ is the hole (electron) concentration,

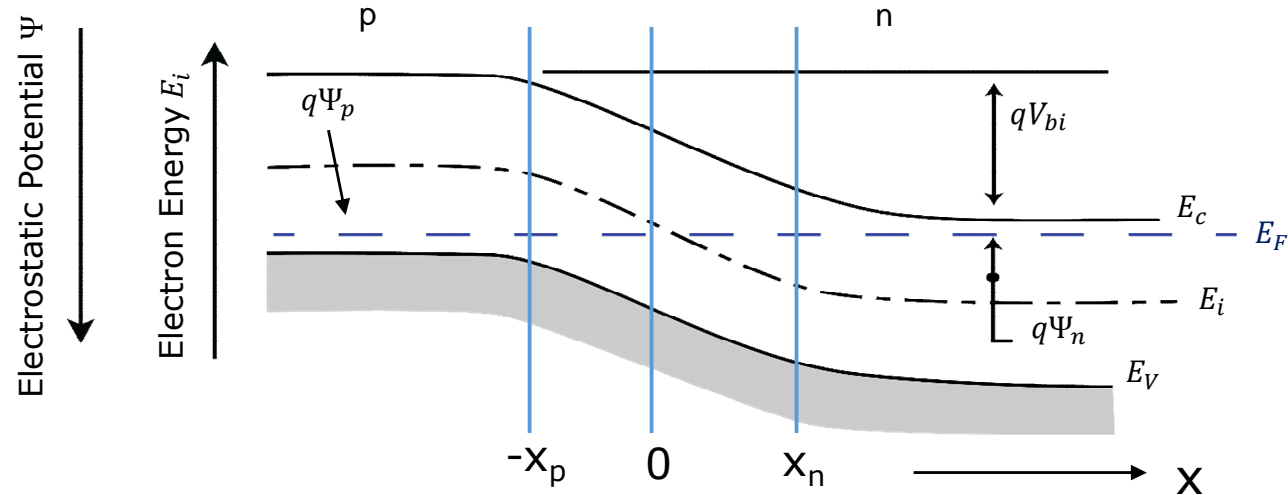
k is the Boltzmann constant, and

ξ is the electrical field.

- $D = (kT/q)\mu$ is the Einstein relation.
- Hence in thermal equilibrium, $J = J_n + J_p = 0$.

Built-in Potential V_{bi}

Figure 1.11



$$\Psi_p \equiv -\frac{1}{q}(E_i - E_F)| x \leq -x_p = -\frac{kT}{q} \ln \frac{N_A}{n_i} \quad (\text{Equation 1.12}) \quad \Psi_n \equiv -\frac{1}{q}(E_i - E_F)| x \geq x_n = \frac{kT}{q} \ln \frac{N_D}{n_i} \quad (\text{Equation 1.13})$$

- The total electrostatic potential difference between the p-side and n-side neutral regions at thermal equilibrium is called the **built-in potential** V_{bi} .

$$V_{bi} = \Psi_n - \Psi_p = \frac{kT}{q} \ln \frac{N_A N_D}{n_i^2} \quad (\text{Equation 1.14})$$

- Note that there is an energy band bending at junction due to V_{bi} . this “potential barrier” that prevents any further flow of electrons and holes under thermal equilibrium.

Depletion Width and Maximum Electric Field

- Overall space charge neutrality requires:

$$N_A x_p = N_D x_n \quad (\text{Equation 1.15})$$

- The total depletion layer width is given by:

$$W = x_p + x_n = \sqrt{\frac{2\epsilon_s}{q} \left[\frac{N_A + N_D}{N_A N_D} \right] V_{bi}} \quad (\text{Equation 1.16})$$

- The electric field can be found by solving the Poisson's equation: $d^2\phi(x)/dx^2 = -\rho(x)/\epsilon_s = -d\xi(x)/dx$ (Equation 1.17)
- This is equivalent to integrating the space charge density plot (a) which results in the electric field plot in (b).
- The maximum **electric field** ξ_m occurs at $x = 0$ is

$$\text{given by: } \xi_m = \frac{qN_D x_n}{\epsilon_s} = \frac{qN_A x_p}{\epsilon_s} \quad (\text{Equation 1.18})$$

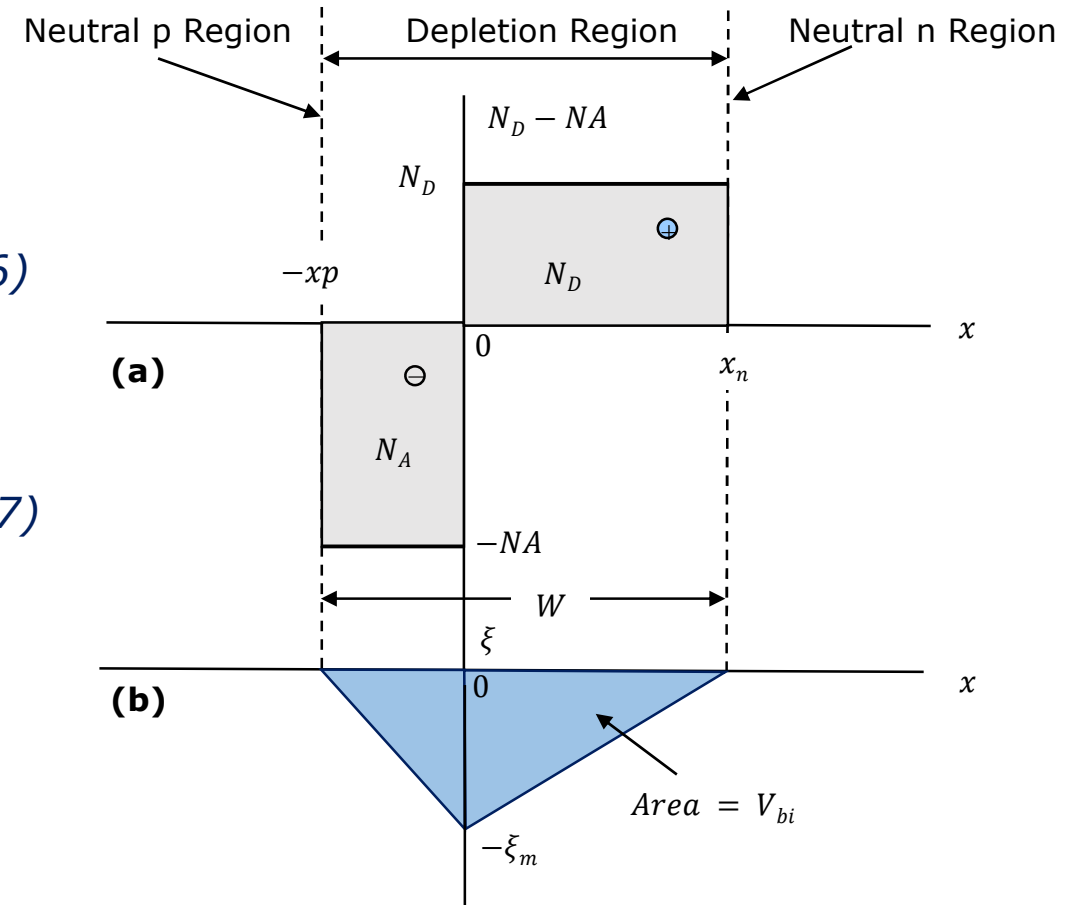


Figure 1.12

One-sided Abrupt Junction

- When the impurity concentration on one side of an abrupt junction is much higher than that of the other side, the junction is called a **one-sided abrupt junction**, e.g., a p^+ - n junction.
- $x_p \ll x_n$, and the W can be simplified to:

$$W \approx x_n = \sqrt{\frac{2\epsilon_s V_{bi}}{qN_D}} \quad (\text{Equation 1.19})$$

- The maximum electric field is $\xi_m = \frac{qN_B W}{\epsilon_s}$ (Equation 1.20) where N_B is the lightly doped bulk concentration.

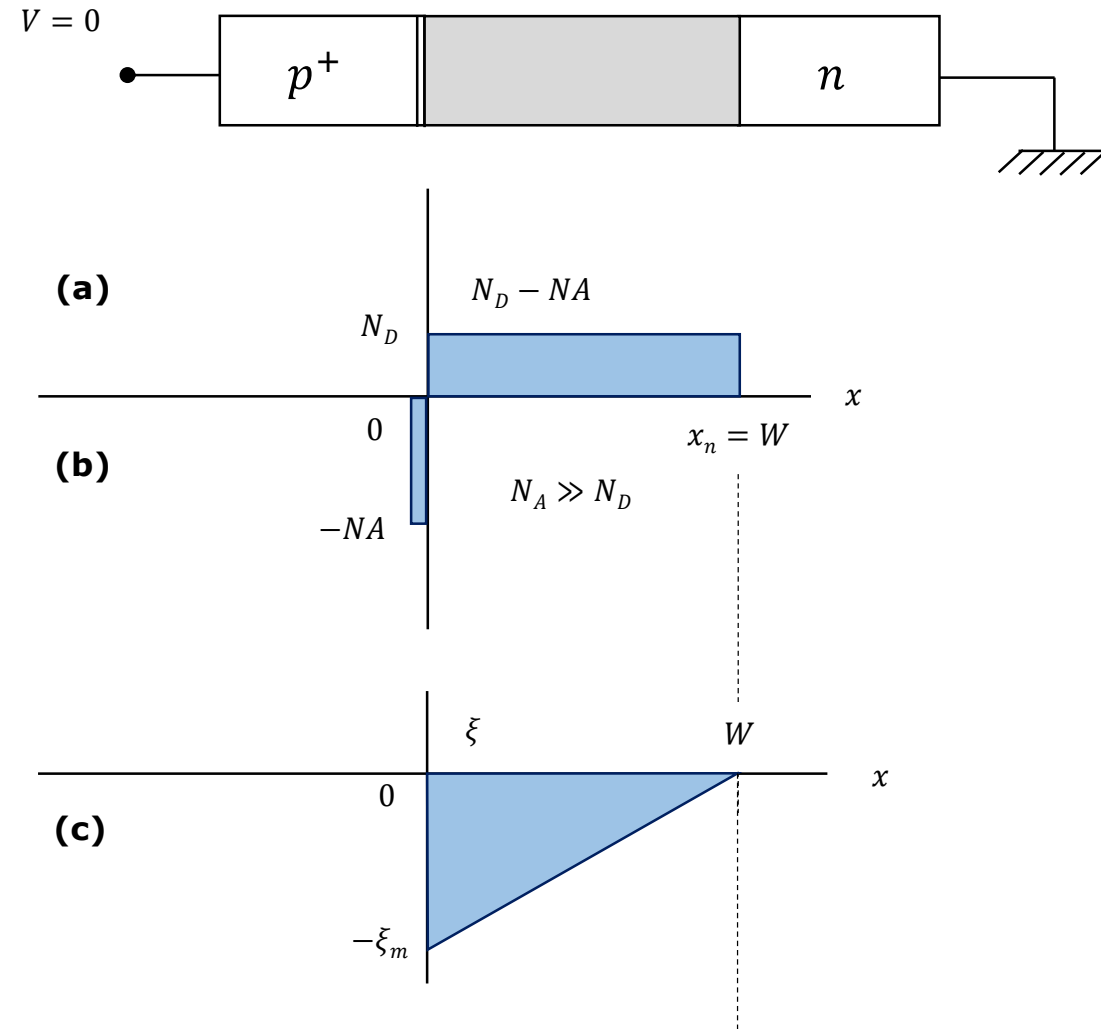
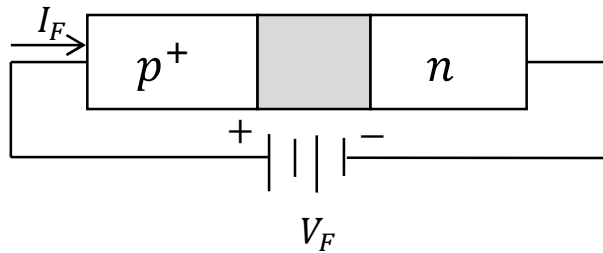


Figure 1.13

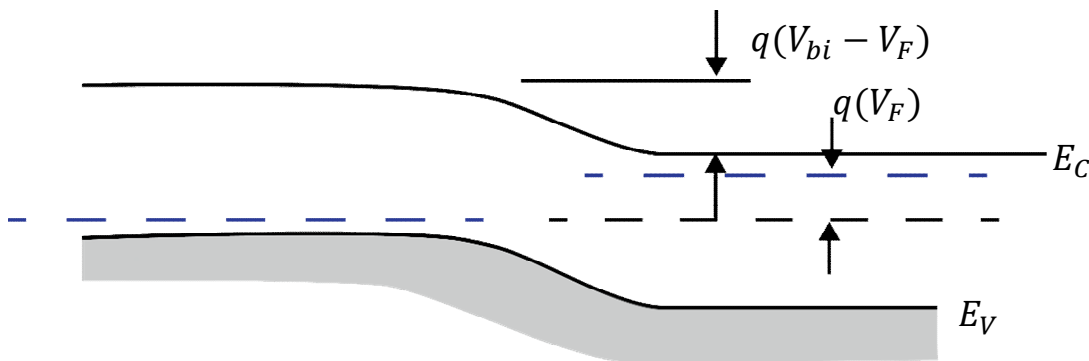
P-N Junction under Bias

Forward bias (V_F):

- Depletion layer width becomes smaller.



- The total electrostatic potential becomes $V_{bi} - V_F$.

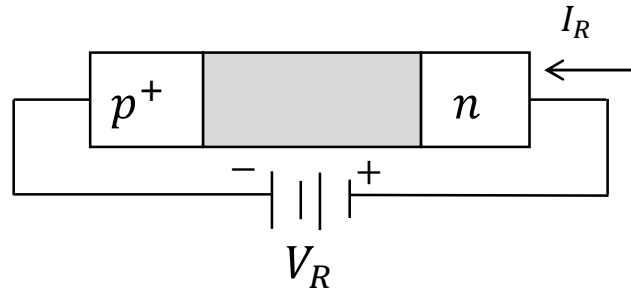


$$W_{FB} = \sqrt{\frac{2\epsilon_s}{q} \left[\frac{N_A + N_D}{N_A N_D} \right] (V_{bi} - V_F)} \quad (\text{Equation 1.21})$$

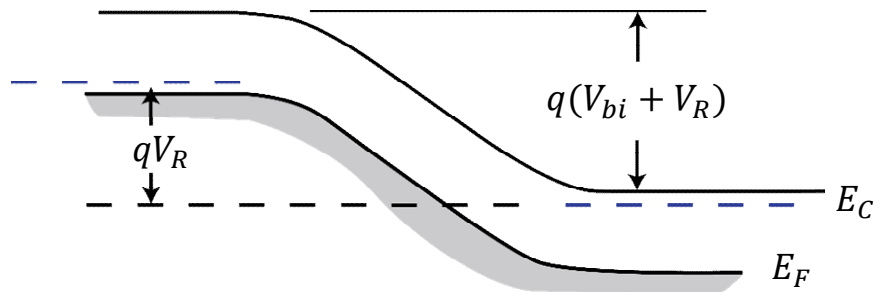
P-N Junction under Bias

Reverse bias (V_R):

- Depletion layer width increased.



- The total electrostatic potential becomes $V_{bi} - V_R$.



$$W_{RB} = \sqrt{\frac{2\epsilon_s}{q} \left[\frac{N_A + N_D}{N_A N_D} \right] (V_{bi} + |V_R|)} \quad (\text{Equation 1.22})$$

Injected Minority Carriers and Current Distributions

(a) Forward bias

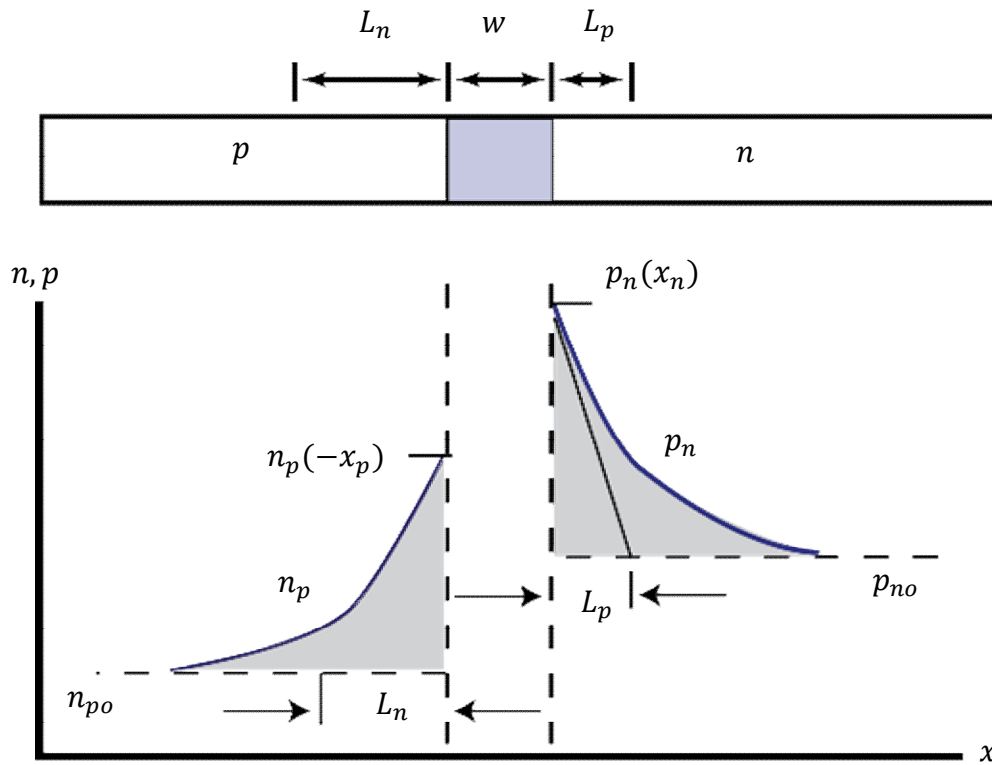


Figure 1.14

Minority carriers concentration at depletion edges:

$$p_n(x=x_n) = P_{no} e^{qV_F/kT} \quad (\text{Equation 1.23a})$$

$$n_p(x=-x_p) = n_{po} e^{qV_F/kT} \quad (\text{Equation 1.23b})$$

Diffusion currents:

$$J_p(x) = -qD_p dp(x)/dx \quad (\text{Equation 1.24a})$$

$$J_n(x) = -qD_n dn(x)/dx \quad (\text{Equation 1.24b})$$

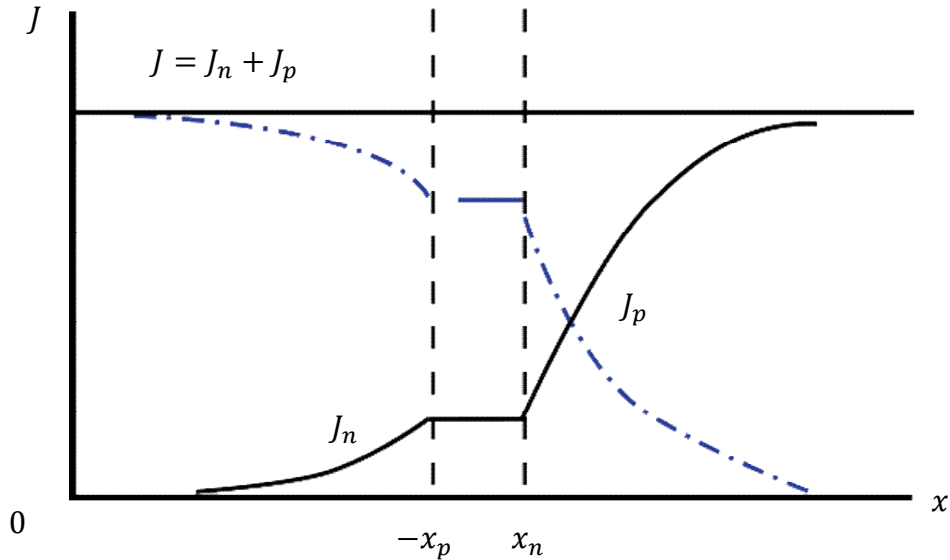
Hole and electron diffusion current densities at depletion edges:

$$I_p(x) = \left(\frac{qD_p p_{no}}{L_p} \right) \left[e^{qV_F/kT} - 1 \right] \quad (\text{Equation 1.25a})$$

$$I_n(x) = \left(\frac{qD_n n_{po}}{L_n} \right) \left[e^{qV_F/kT} - 1 \right] \quad (\text{Equation 1.25b})$$

Injected Minority Carriers and Current Distributions

(a) Forward bias



Total current density:

$$J = J_n(-x_p) + J_p(x_n) = \left(\frac{qD_p p_{no}}{L_p} + \frac{qD_p n_{po}}{L_n} \right) \left[e^{qV_F/kT} - 1 \right] \quad (\text{Equation 1.26})$$

Under Forward bias:

$$J_F \sim \left(\frac{qD_p p_{no}}{L_p} + \frac{qD_p n_{po}}{L_n} \right) \left[e^{qV_F/kT} \right] \quad \text{since } e^{qV_F/kT} \gg 1 \quad (\text{Equation 1.27})$$

Figure 1.15

(b) Reverse bias

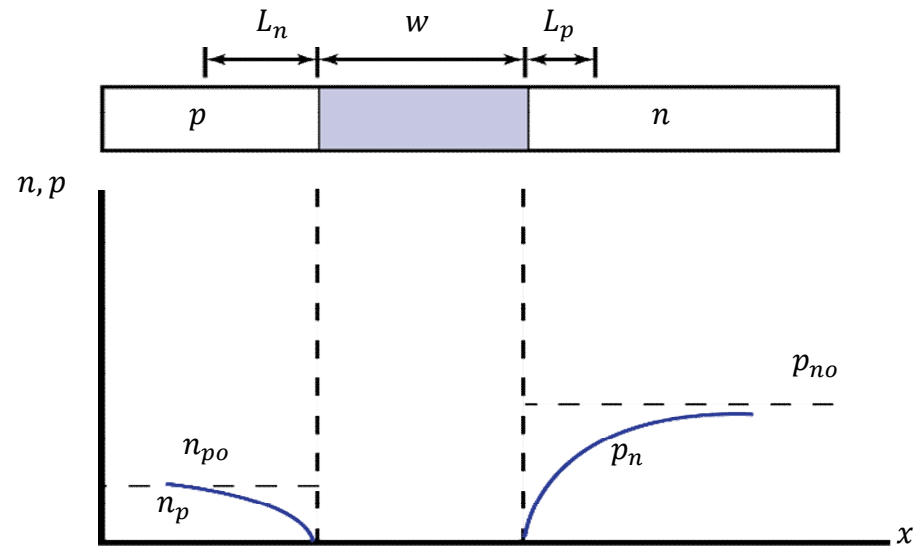


Figure 1.16

Minority carriers concentration at depletion edges:

$$p_n(x=x_n) = P_{no} e^{-qV_{|R|}/kT} \ll p_{no} \quad (\text{Equation 1.28a})$$

$$n_p(x=-x_p) = n_{po} e^{-qV_{|R|}/kT} \ll n_{po} \quad (\text{Equation 1.28b})$$

Diffusion currents:

$$J_p(x) = -qD_p dp(x)/dx \quad (\text{Equation 1.29a})$$

$$J_n(x) = -qD_n dn(x)/dx \quad (\text{Equation 1.29b})$$

Hole and electron diffusion current densities at depletion edges:

$$I_p(x) = \left(\frac{qD_p p_{no}}{L_p} \right) \left[e^{-qV_{|R|}/kT} - 1 \right] \sim 0 \quad (\text{Equation 1.30a})$$

$$I_n(x) = \left(\frac{qD_n n_{po}}{L_n} \right) \left[e^{-qV_{|R|}/kT} - 1 \right] \sim 0 \quad (\text{Equation 1.30b})$$

(b) Reverse bias

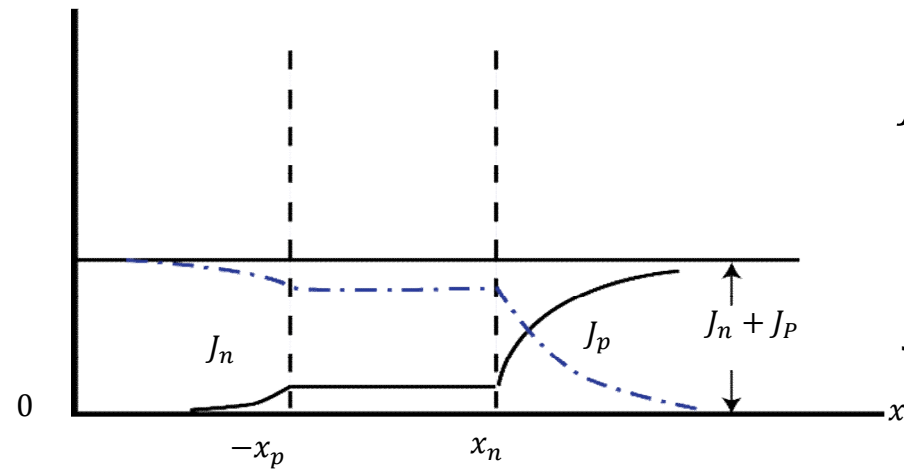


Figure 1.17

Total current density:

$$J = J_n(-x_p) + J_p(x_n) = \left(\frac{qD_p p_{no}}{L_p} + \frac{qD_p n_{po}}{L_n} \right) \left[e^{-qV_R/kT} - 1 \right] \sim 0 \quad (\text{Equation 1.31})$$

Under Reverse bias:

$$J_R \sim - \left(\frac{qD_p p_{no}}{L_p} + \frac{qD_p n_{po}}{L_n} \right) = J_{\text{sat}} \sim 0 \quad \text{since } e^{-qV_R/kT} \ll 1 \quad (\text{Equation 1.32})$$

Depletion Capacitance

- For one-sided abrupt junction, the depletion capacitance per unit area:

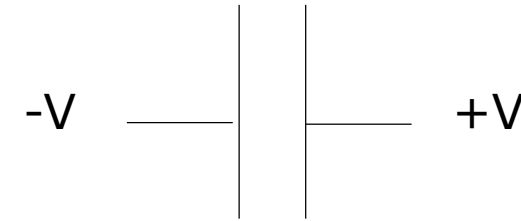
$$C_j = \frac{\epsilon_s}{W} = \sqrt{\frac{q\epsilon_s N_B}{2(V_{bi} - V)}} \quad (\text{Equation 1.33})$$

ϵ_s is the dielectric constant of the semiconductor material and W is the total depletion width.

- Or

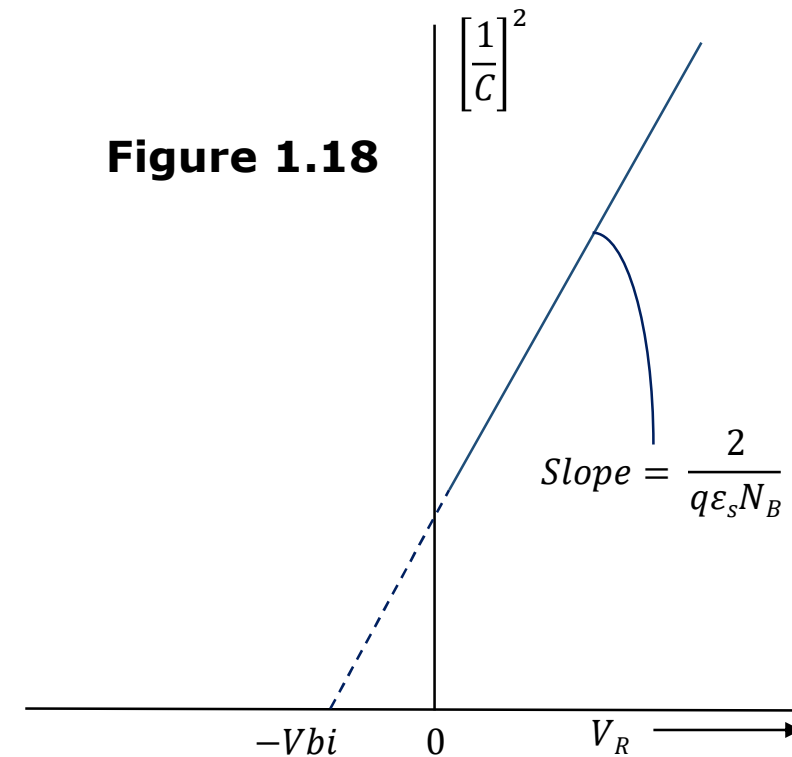
$$\frac{1}{C_j^2} = \frac{2(V_{bi} - V)}{q\epsilon_s N_B} \quad (\text{Equation 1.34})$$

- A plot of $1/C_j^2$ versus V produces a straight line for a one-sided abrupt junction.
- The slope links to the impurity concentration N_B and the intercept at $(1/C_j^2 = 0)$ gives V_{bi} .



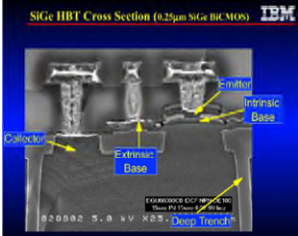
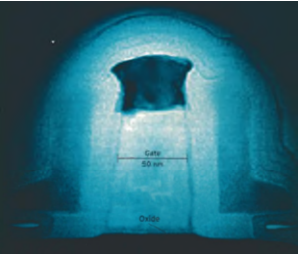
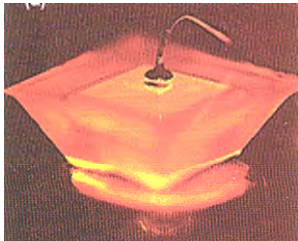
$C = \epsilon/d$ (capacitance per unit area)

Figure 1.18



- It is important to understand the operation principles of PN junction as it is the basic building block for many other semiconductor devices.
- The key parameters of a PN junction are: built-in voltage (V_{bi}), depletion width (W), space charge densities ($N_A x_p = N_D x_n$), electric field (ξ) and depletion capacitance (C_j).
- Under Forward Bias $\rightarrow V = V_{bi} - V_F$ W decreases, $N_A x_p$ and $N_D x_n$ decreases, ξ decreases and C_j increases.
- Under Reverse Bias $\rightarrow V = V_{bi} + |V_R|$, W increases, $N_A x_p$ and $N_D x_n$ increases, ξ increases and C_j decreases.

- Under forward bias, minority carriers injection gives rise to diffusion currents I_p and I_n . The total current $I_T = I_p + I_n$ increases.
- Under reverse bias, minority carriers get depleted and the total current (or reverse saturation current) is very small.
- The capacitance versus voltage characteristics can be used to determine the built-in voltage (V_{bi}) and impurity concentration (N_B) in the semiconductor.

No.	Slide No.	Image	Reference
1.	6		<i>D. (n.d.). Status and Trends of SiGe BiCMOS Technology. Retrieved March 8, 2017, from http://asia.stanford.edu/events/spring01/slides/harameSlides.pdf</i>
2.	6		<i>90nm_transistor. (n.d.). Retrieved March 08, 2017, from http://www.ixbt.com/</i>
3.	6		<i>Optoelectronic Devices-Types, Applications, Threshold Frequency Define. (2014, July 09). Retrieved March 08, 2017, from http://www.circuitstoday.com/optoelectronic-devices</i>

No.	Slide No.	Image	Reference
4.	6		<i>Silicon-Germanium multi-quantum . (n.d.). Retrieved March 08, 2017, from http://imagebank.osa.org/</i>