

## Exploration of Dynamics

Date \_\_\_\_\_  
Page \_\_\_\_\_

The dynamics of the map

$$x_{n+1} = f_1(x_n) = 4x_n(1-x_n) \text{ if } x_n < x_{n-1}$$

$$x_{n+1} = f_2(x_n) = rx_n(1-x_n) \text{ if } x_n > x_{n-1}$$

without memory dependence, the logistic map  $x_{n+1} = rx_n(1-x_n)$  is chaotic with Lyapunov Exponent of  $\ln(2)$  and an invariant density distribution which is continuous in  $0 \leq x \leq 1$ .

The map  $x_{n+1} = rx_n(1-x_n)$ , by itself shows fixed points for  $n \leq 3$  followed by cycles of period  $2^n$  and so on.

In the presence of memory the above map does not begin with a fixed point but with a 2-cycle.

We have a possible 2-cycle  $x_1, x_2$  with

$$x_2 = 4x_1(1-x_1) \quad \text{--- (3.3)}$$

$$x_1 = rx_2(1-x_2) \quad \text{--- (3.4)}$$

These change to

$$x_2 = 4rx_2(1-x_2)(1-rx_2(1-x_2)) \quad \text{and} \quad \text{--- (3.5)}$$

$$x_1 = 4rx_1(1-x_1)(1-rx_1(1-x_1)) \quad \text{--- (3.6)}$$

$x_1, x_{m_2}$  are fixed points of two iterated functions  
 $F(n) \& G(n)$

$$F(n) = 4rn(1-n)(1-4n(1-n)) \quad - (3.7)$$

$$G(n) = 4rn(1-n)(1-rx(1-n)) \quad - (3.8)$$

stability of  $x < \frac{1}{4}$ :

For  $r < \frac{1}{4}$ , the fixed point  $n=0$  is stable for both  $F(n) \& G(n)$ . To verify this.

① Fixed points :-  $F(n) = n$  and  $G(n) = n$  & solve for  $n'$

② Derivative of fixed point: calculate the derivative at  $n=0$  for  $F(n)$

$F(0) = 0 \rightarrow$  So '0' is fixed point.

lets calculate  $F'(n)$

$$F(n) = 4rx(1-n)(1-4n(1-n))$$

$$F'(n) = \frac{d}{dn} (4rn(1-n)(1-4n(1-n)))$$

$$\frac{d}{dn} (4rn(1-n)(1-4n+4n^2))$$

$$\frac{d}{dn} ((4rn - 4rn^2)(1 - 4n + 4n^2))$$

$$\frac{d}{dn} [(4rn - 16rn^2 + 16rn^3 - 4rn^2 + 16rx^3 + 16rx^4)]$$

$$F'(0) = 4r$$

Since  $|F'(0)| = 4r < 1$  for  $r < \frac{1}{4}$   $n=0$  is stable.

Two - cycle Elements for  $r > r_y$

For  $r > r_y$ ,  $n=0$  becomes unstable, and a two-cycle emerges

The Two-cycle points  $n_1$  and  $n_2$  such that

$$F(F(n)) = n$$

$$G(G(n)) = n$$

This implies Solving the second iteration of the function to find these points.

Destabilization of pitchfork Bifurcation occurs at (using  $(F_n)$ )

$$-1 = 4r_c [1 - 8r_c(1 - m_c)](1 - m_c)$$

$$(-2r_c) = 1 - 2r_c + \sqrt{4r_c^2 + 4r_c - 9}$$
$$\frac{2(4r_c - 5)}{2(4r_c - 5)}$$

$$-1 = 4r_c \left[ 1 - 8 \left( \frac{1 - 2r_c + \sqrt{4r_c^2 + 4r_c - 9}}{2(4r_c - 5)} \right) \right] \left[ 1 - \left( \frac{1 - 2r_c + \sqrt{4r_c^2 + 4r_c - 9}}{2(4r_c - 5)} \right) \right]$$
$$\left[ \left( 1 - \frac{1 - 2r_c + \sqrt{4r_c^2 + 4r_c - 9}}{2(4r_c - 5)} \right) \right]$$

By solving this complex Expression we get.

$$r_c \approx 1.113.$$

## Cycle Bifurcation

Above  $r = r_c$  the system undergo from 2-cycle to a 4-cycle to an 8-cycle and eventually into chaos.

These Bifurcations are characterised by the Period-doubling route to chaos.

### Fixed points for $r \approx 1$

Near  $r=1$ , another fixed point solution exists. For a restricted basin of attraction, we can find a stable solution to  $f^2(n) = x_1$ .

### Basin of Attraction and Crisis

For  $r < 3$ , this solution is stable. At  $r_o = 2.7460$  the Basin of attraction for this fixed point collides with that of the period cycle - 8-cycle, leading to crisis. After the crisis, the only the fixed point attractor is observed.

## Period Adding Bifurcations

Date \_\_\_\_\_  
Page \_\_\_\_\_

### Behaviour $r=2$ :

- For  $r=2$ , the map has unique feature: The behaviour of the initial conditions relative to  $x_0 = \frac{1}{2}$
- If  $x_n < \frac{1}{2}$  then  $x_{n+1} > x_n$  but will not exceed  $x_{n+1} = \frac{1}{2}$ . This leads to  $x$  converging to the fixed point  $x = \frac{1}{2}$ .
- If  $x_n > \frac{1}{2}$ , then  $x_{n+1} < x_n$ , which triggers a switch to a different map, and the system does not necessarily reach a fixed point.

### For $r > 2$ :

- 1) periodic windows: - The system starts showing periodic behaviours with very high periods.
- 2) periodic decrements: - As  $r$  increases slightly, the period of oscillations decreases rapidly, described by inverse period adding bifurcations.
- 3) Ribbon structure: - The bifurcation diagram shows a complex ribbon-like ~~structure~~ pattern due to branch merging and splitting.

This scenario involved Border collision bifurcation, where the fixed points of two smooth maps intersect at a critical value,  $n_c$ .

This Dynamic model system can be modelled by a set of piecewise smooth maps. One-dimensional piecewise smooth maps can be defined in the following way:-

We take two maps  $g(n) \{ h(n; r_0)$

$$g(n) = 4n(1-n) \text{ for } n < n_c$$

$$h(n; r_0) = (2r_0 - 4)n(1-n) \text{ for } n \geq n_c$$

$$n_{n+1} = \left[ r_0 - (4 - r_0) \tanh\left(\frac{n_n - n_c}{c}\right) \right] n(1-n)$$

Border Collision and Fixed points.

A Border Collision occurs when the fixed point of these maps cross the border at  $n = n_c$ . we need to analyze the fixed points  $n_L^*$  and  $n_R^*$  of  $g(n)$  and  $h(n; r_0)$  respectively, determine the conditions under they collide at  $n = n_c$ .

Fixed points of  $g(n)$

The fixed points of  $g(n)$  is obtained by solving:

$$n = g(n) = 4n(1-n)$$

Solving this Equation:-

$$n = 4n(1-n)$$

$$n = 4n - 4n^2$$

$$y^2 - 3n = 0$$

$$n(y-n) = 0$$

So, the fixed points are

$$n=0, \text{ or } n=3/4$$

We denote the fixed point within the relevant interval

$n < n_c$  or  $n_L^*$ . For the given condition, we assume  $n_c = 1/2$ .

$$\text{Since } n_L^* = 3/4$$

Since  $n_L^* = 3/4 > 1/2$  the fixed point  $n_L^*$  cannot cross the borders.

### Fixed points for $h(n; \infty)$

The fixed points of  $h(n; r_0)$  is obtained by solving:

$$n = h(n; r_0) = (2r_0 - 4)n(1-n)$$

Solving the equation.

$$n = (2r_0 - 4)n(1-n)$$

$$n - (2r_0 - 4)n - (2r_0 - 4)n^2 = 0$$

$$(2r_0 - 4)n^2 - (2r_0 - 4)n = 0$$

$$n(2r_0 - 4)n - (2r_0 - 4) = 0$$

$$\text{So, the fixed points are } n=0 \text{ or } n = \frac{2r_0 - 4}{2r_0 - 4}$$

We denote the fixed point within the relevant interval

$$\pi \geq \pi_c \text{ or } \pi_R^*$$

$$\pi_R^* = \frac{2\gamma_0 - 5}{2\gamma_0 - 4}$$

### Collision Condition

For a border collision to occur, the fixed points  $\pi_R^*$  must collide with the border  $\pi_c = \frac{1}{2}$ .

$$\pi_R^* = \pi_c$$

$$\frac{2\gamma_0 - 5}{2\gamma_0 - 4} = \frac{1}{2}$$

Solving for  $\gamma_0$ :

$$2(2\gamma_0 - 5) = 2\gamma_0 - 4$$

$$4\gamma_0 - 10 = 2\gamma_0 - 4$$

$$4\gamma_0 - 2\gamma_0 = 6$$

$$2\gamma_0 = 6$$

$$\boxed{\gamma_0 = 3}$$

→ At  $\gamma_0 = 3$ ,  $\pi_R^*$  collides with  $\pi_c = \frac{1}{2}$ , resulting in a border collision bifurcation.

→ As  $\gamma_0$  increases beyond '3', the fixed points  $\pi_R^*$  moves past  $\pi_c$ , leading to stabilized fixed point & the emergence of a chaotic band.

$$i.e. (r_0 = 3)$$

As we move towards  $r=2$  from above, the distance between two consecutive nodes decreases. If we define the quantity.

$$S = \lim_{n \rightarrow \infty} \frac{r_n - r_{n+1}}{r_{n+1} - r_n}$$

$\downarrow$   
f.eigenvalue  
constant

where  $r_n$  is the value of ' $r$ ' at  $(n+1)$  cycle goes to  $n$ -cycle, then  $S$  converges to 1.414.

In ribbon-like structure, Every node is a fixed point of the return maps  $f_2^{n-6}(f(f_2(n_m)))$  and  $f_2^{n-6}(f_2(f(n_m)))$

where ' $n$ ' denotes the no. of cycles executed at that particular ' $r$ ' and  $f(n_m)$  is the composite function of  $f_1(n_m)$  and  $f_2(n_m)$ .

Convergence of the fixed point from left to right at nodes ensures that these nodes are also fixed points of  $f_2(n_m)$ .

so for each node:

$$f_2^{n-6}(f_2(f(x_m^*))) = f_2(n_m^*) = x_m^*$$

$$i.e. n_m^* = 1 - \frac{1}{\sqrt{5}}$$

$$f_2 \left( f_2 \left( f \left( 1 - \frac{1}{r} \right) \right) \right) = 1 - \frac{1}{r}$$

↓ r-value  $\rightarrow r_n$  at the  $n^{\text{th}}$  cycle.

This can be decomposed as:

$$\begin{aligned} & f_2^{n-7} \left( f_2 \left( f \left( 1 - \frac{1}{r} \right) \right) \right) - \left( 1 - \frac{1}{r} \right) \\ & \times \left[ r f_2^{n-7} \left( f_2 \left( f \left( 1 - \frac{1}{r} \right) \right) \right) - 1 \right] = 0. \end{aligned}$$

$$r f_2^{n-7} \left( f_2 \left( f \left( 1 - \frac{1}{r} \right) \right) \right) = 1$$

if for  $n=N$  the period adding phenomenon stops, then

$$r f_2^{N-6} \left( f \left( 1 - \frac{1}{r} \right) \right) = 1$$

$$r f_2^{N-5} \left( f \left( 1 - \frac{1}{r} \right) \right) = 1$$

when period stops  $r_N \rightarrow r_{N+1}$  the above gives

$$r_N f_2 \left( \frac{1}{r_N} \right) = 1$$

which gives.

$$\frac{1}{r_n} = r_n \left( \frac{1}{r_n} \right) \left( 1 - \frac{1}{r_n} \right)$$

$$\frac{1}{r_n} = 1 - \frac{1}{r_n}$$

$$\text{Q} \quad \frac{2}{r_n} = 1$$

$$r_n = 2$$

So, at  $r=2$ , the period adding phenomenon stops.

Let find  $s_n$

Near  $r=2$ , let  $r_n = 2 + \Delta_n$  then

$$r_n^n \Delta_n^2 = 1$$

$$r_{n+1}^n \Delta_{n+1}^2 = 1$$

$$r_n \approx k_N$$

$$s_n = \frac{\Delta_n - \Delta_{n+1}}{\Delta_{n+1} - \Delta_n}$$

$$s_n = \sqrt{r_n} = \sqrt{2} \rightarrow \text{for large value of } r_n.$$

→ Beyond  $r = 2.5427$ , there is a chaotic band with the emergence of a periodic window at

$$r = 2.7244$$

For  $2.7564 \leq r \leq 2.7876$ , we see a cycle of period 11,

Each element of the '11' cycle exhibits a sequence of period doubling bifurcation for  $r < 2.7564$  and  $r > 2.7876$

~~Focusing particular Element of the 11-cycle the two sets of~~

As we further increase the value of  $r'$ , there are two large periodic windows. At  $r = 2.9499$ , a 6-cycle is formed,

period - 6 fixed points bifurcate and take period doubling route to chaos.

Like a logistic map a period 3-cycle is formed and takes an Intermittent route to chaos.

At  $r = 3.8384$  period becomes Stable and goes through period doubling bifurcation