

# Optimal Estimation Methods

## (Lecture 10 – Review of Dynamic Systems: Part I)

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- Typical spring-mass-damper system

$$m \ddot{x} + c \dot{x} + k x = F$$

- Free response with  $F = 0$

$$\text{let } x = Ae^{st}, \dot{x} = sAe^{st} = sx, \ddot{x} = s^2Ae^{st} = s^2x$$

- Leads to

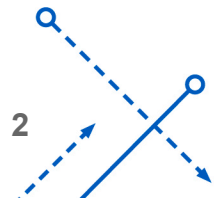
$$(ms^2 + cs + k)Ae^{st} = 0$$

- Assuming that  $Ae^{st} \neq 0$  gives the *characteristic equation*

$$ms^2 + cs + k = 0$$

- Roots are given by

$$s_{1,2} = \frac{-c \pm \sqrt{c^2 - 4mk}}{2m}$$



- Possibilities for  $s_{1,2}$ 
  - Real and unequal with  $c^2 - 4mk > 0$

$$x(t) = A_1 e^{s_1 t} + A_2 e^{s_2 t}$$

- Real and equal with  $c^2 - 4mk = 0$

$$x(t) = A_1 e^{s_1 t} + t A_2 e^{s_2 t}$$

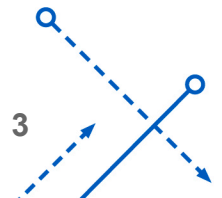
- Complex conjugates with  $c^2 - 4mk < 0$

$$x(t) = B e^{-a t} \sin(b t + \phi)$$

where

$$a = \frac{c}{2m}, \quad b = \frac{\sqrt{4mk - c^2}}{2m}$$

- Constants found through initial conditions



- Define the following

Natural Frequency  $\omega_n = \sqrt{\frac{k}{m}}$  and Damping Ratio  $\zeta = \frac{c}{2\sqrt{mk}}$

- Then

$$m s^2 + c s + k = 0 \Rightarrow s^2 + 2 \zeta \omega_n s + \omega_n^2 = 0$$

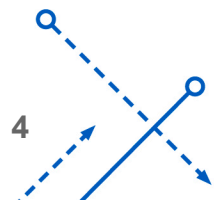
$$s_{1,2} = -\zeta \omega_n \pm \omega_n \sqrt{\zeta^2 - 1}$$

- Three cases again

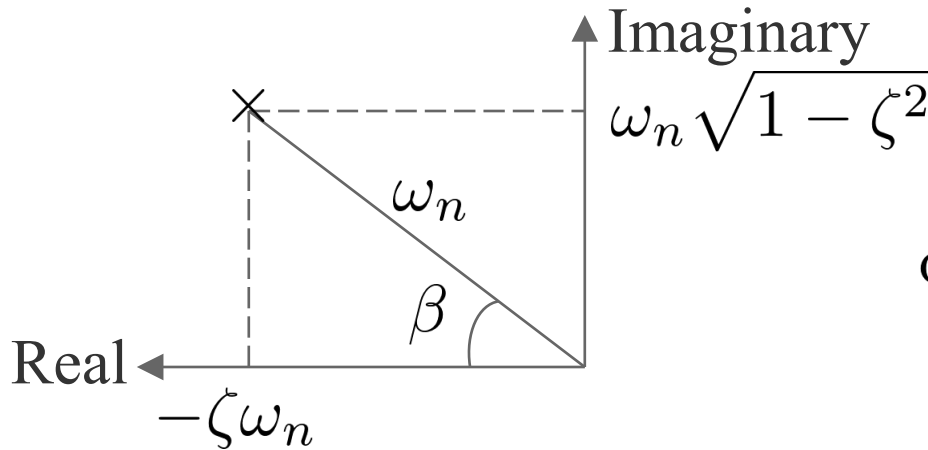
- Real and unequal with  $\zeta > 1$
- Real and equal with  $\zeta = 1$
- Complex conjugates with  $0 < \zeta < 1$

- Note if  $\zeta = 0$  then  $s_{1,2} = \pm \omega_n j$

- Also  $\omega_d \equiv \omega_n \sqrt{1 - \zeta^2}$  is the *damped natural frequency*

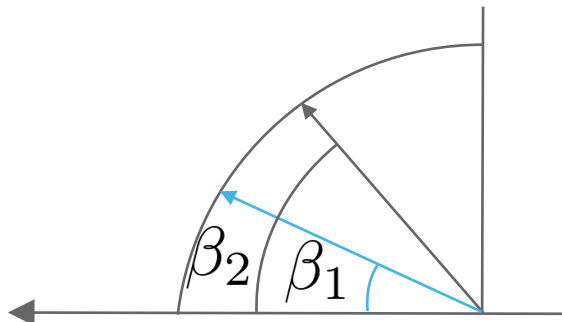
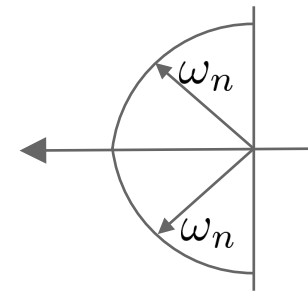


## • Graphical Interpretation



$$\cos \beta = \frac{\zeta \omega_n}{\omega_n} = \zeta$$

- Roots lying on a circle centered at the origin have same natural frequency
- The closer a line is to the imaginary axis the smaller  $\zeta$  becomes



$$\cos \beta_1 > \cos \beta_2$$

$$\text{so } \zeta_1 > \zeta_2$$

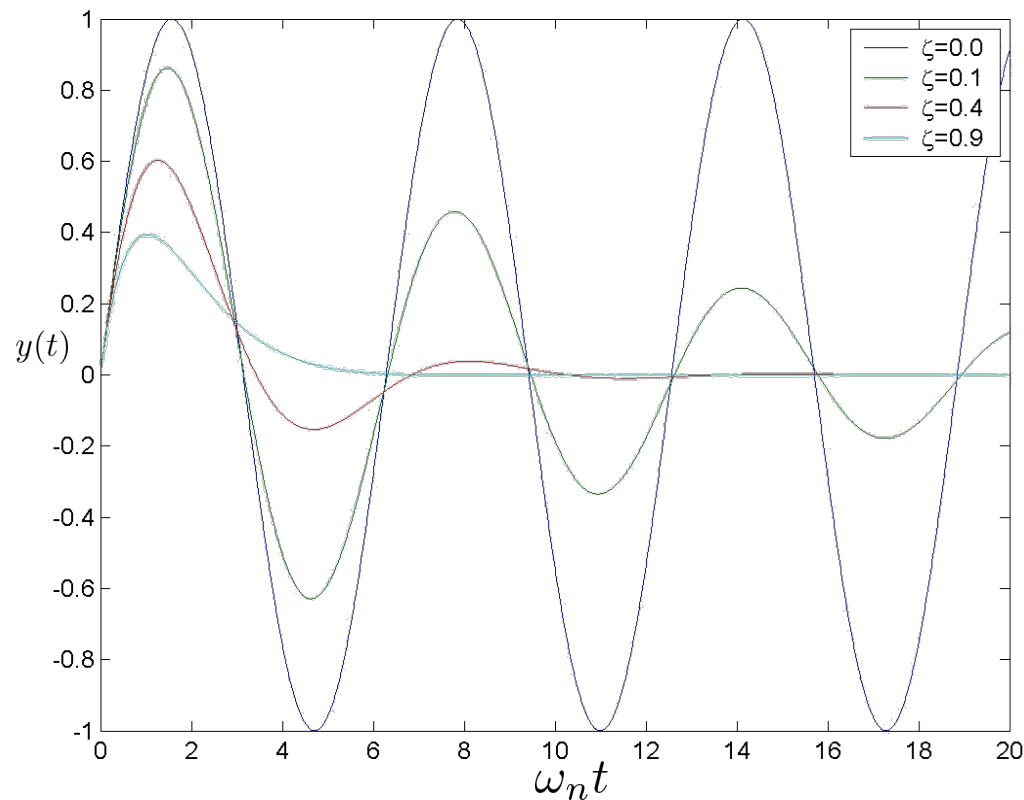
note when  $\beta = 90^\circ$  then  $\zeta = 0$

- Responses to various  $\zeta$  values

Optimal  $\zeta = \frac{\sqrt{2}}{2}$

$$Y(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}$$

$$y(t) = \frac{\omega_n}{\sqrt{1 - \zeta^2}} e^{-\zeta\omega_n t} \sin(\omega_n \sqrt{1 - \zeta^2} t), \text{ for } 0 < \zeta < 1$$



- For spring-mass system we have (no forcing)

$$m \ddot{x} + k x = 0$$

- Characteristic equation becomes

$$m s^2 + k = 0$$

- Roots are given by (purely imaginary)

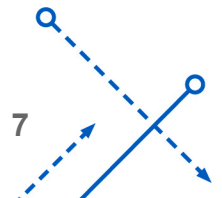
$$s_{1,2} = \pm \sqrt{\frac{k}{m}} j$$

- Assumed form of the solution

$$x(t) = B \sin(\sqrt{k/m} t + \phi)$$

- Need to find constants  $B$  and  $\phi$
- Derivative is given by

$$\dot{x}(t) = B \sqrt{k/m} \cos(\sqrt{k/m} t + \phi)$$



- Initial Conditions

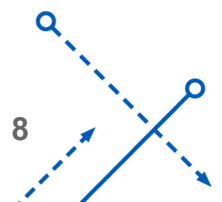
$$x(0) \equiv x_0 = B \sin \phi \text{ and } \dot{x}(0) \equiv \dot{x}_0 = B \sqrt{k/m} \cos \phi$$

- So

$$B = \frac{x_0}{\sin \phi}$$

- Also we have

$$\frac{x_0}{\dot{x}_0} = \frac{1}{\sqrt{k/m}} \tan \phi \quad \text{so} \quad \phi = \tan^{-1} \left( \frac{x_0}{\dot{x}_0} \sqrt{\frac{k}{m}} \right)$$





- In-class assignment
  - Compute the transfer function of the following system
    - Take the Laplace Transform with zero initial conditions

$$x_1 = y$$

$$x_2 = \dot{y}$$

$$\ddot{y} + \dot{y} + y = u$$

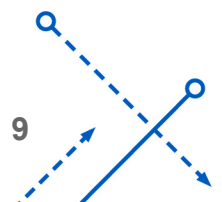
$$(s^2 + s + 1)Y(s) = U(s)$$

$$\frac{Y(s)}{U(s)} = \frac{1}{s^2 + s + 1}$$

- Then compute its state space model

$$\dot{\underline{x}} = \underbrace{\begin{pmatrix} 0 & 1 \\ -1 & -1 \end{pmatrix}}_F \underbrace{\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}}_{\underline{x}} + \underbrace{\begin{pmatrix} 0 \\ 1 \end{pmatrix}}_B \underbrace{u}_{\underline{u}}$$

$$y = \underbrace{[1 \ 0]}_H \underbrace{\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}}_{\underline{x}} + \underbrace{0}_D \underline{u}$$



- Take the Laplace Transform with zero initial conditions

$$(s^2 + s + 1)Y(s) = U(s)$$

- Transfer function is then given by

$$\frac{Y(s)}{U(s)} = \frac{1}{(s^2 + s + 1)}$$

- Define states up to the  $n - 1$  derivatives:  $\ddot{y} + \dot{y} + y = u$

$$x_1 = y,$$

$$\dot{x}_1 = x_2$$

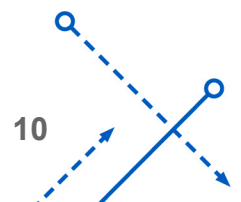
$$x_2 = \dot{y},$$

$$\dot{x}_2 = -x_2 - x_1 + u$$

- State space model is then given by

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & -1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u$$

$$y = \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 \end{bmatrix} u$$



- General form for single-input-single-output (SISO) systems

$$\dot{\mathbf{x}}(t) = F \mathbf{x}(t) + B u(t)$$

$$y(t) = H \mathbf{x}(t) + D u(t)$$

- Convert back to TF by taking the Laplace Transform

$$[sI] \mathbf{X}(s) = F \mathbf{X}(s) + B U(s)$$


$$Y(s) = H \mathbf{X}(s) + D U(s)$$

- From first equation we have

$$\mathbf{X}(s) = (sI - F)^{-1} B U(s)$$

- Substituting this into the second equation gives

$$\boxed{\frac{Y(s)}{U(s)} = H (sI - F)^{-1} B + D}$$

- Note that  $(sI - F)^{-1} = \text{adj}(sI - F) / |sI - F|$
- Then  $|sI - F|$  gives the denominator of the transfer function
- This shows that the eigenvalues of  $F$  are the roots of the characteristic equation! 

- From last example

$$F = \begin{bmatrix} 0 & 1 \\ -1 & -1 \end{bmatrix}, \quad B = \begin{bmatrix} 0 \\ 1 \end{bmatrix}, \quad H = \begin{bmatrix} 1 & 0 \end{bmatrix}, \quad D = 0$$

- First compute  $(sI - F)^{-1}$

$$(sI - F)^{-1} = \begin{bmatrix} s & -1 \\ 1 & s+1 \end{bmatrix}^{-1} = \frac{1}{s^2 + s + 1} \begin{bmatrix} s+1 & 1 \\ -1 & s \end{bmatrix}$$

- Then

$$\begin{aligned} \frac{Y(s)}{U(s)} &= H (sI - F)^{-1} B + D \\ &= \frac{1}{s^2 + s + 1} \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} s+1 & 1 \\ -1 & s \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} + 0 \\ &= \frac{1}{s^2 + s + 1} \begin{bmatrix} s+1 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \frac{1}{s^2 + s + 1} \checkmark \end{aligned}$$

- In-class assignment

$$\text{inv} \begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} D & -B \\ -C & A \end{pmatrix}$$

- Compute the transfer function of the following system

$$F = \begin{bmatrix} -0.5 & -1.5 \\ 0.5 & -0.5 \end{bmatrix}, \quad B = \frac{\sqrt{2}}{2} \begin{bmatrix} 1 \\ -1 \end{bmatrix}, \quad H = \frac{\sqrt{2}}{2} \begin{bmatrix} 1 & 1 \end{bmatrix}, \quad D = 0$$

$$sI - F = \begin{pmatrix} s+0.5 & 1.5 \\ -0.5 & s+0.5 \end{pmatrix} \Rightarrow \frac{1}{(s+0.5)^2 + (1.5)(0.5)} \begin{pmatrix} s+0.5 & -1.5 \\ 0.5 & s+0.5 \end{pmatrix}$$

$$\begin{pmatrix} \frac{s+0.5}{s^2+s+1} & \frac{-1.5}{s^2+s+1} \\ \frac{0.5}{s^2+s+1} & \frac{s+0.5}{s^2+s+1} \end{pmatrix} \underbrace{\quad}_{(sI-F)^{-1}}$$

$$H(sI-F)^{-1}B \quad \text{to } \begin{bmatrix} \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix} \begin{pmatrix} \frac{s+0.5}{s^2+s+1} & \frac{-1.5}{s^2+s+1} \\ \frac{0.5}{s^2+s+1} & \frac{s+0.5}{s^2+s+1} \end{pmatrix} \begin{pmatrix} \frac{\sqrt{2}}{2} \\ -\frac{\sqrt{2}}{2} \end{pmatrix}$$

$$= \frac{1}{s^2+s+1}$$

- First compute  $(sI - F)^{-1}$

$$(sI - F)^{-1} = \begin{bmatrix} s + 0.5 & 1.5 \\ -0.5 & s + 0.5 \end{bmatrix}^{-1} = \frac{1}{s^2 + s + 1} \begin{bmatrix} s + 0.5 & -1.5 \\ 0.5 & s + 0.5 \end{bmatrix}$$

- Then

$$\begin{aligned} \frac{Y(s)}{U(s)} &= H (sI - F)^{-1} B + D \\ &= \frac{1}{s^2 + s + 1} \cdot \frac{1}{2} \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} s + 0.5 & -1.5 \\ 0.5 & s + 0.5 \end{bmatrix} \begin{bmatrix} 1 \\ -1 \end{bmatrix} + 0 \\ &= \frac{1}{s^2 + s + 1} \cdot \frac{1}{2} \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} s + 2 \\ -s \end{bmatrix} \\ &= \frac{1}{s^2 + s + 1} \end{aligned}$$

- Same transfer function as before
  - There are actually an infinite amount of state space *realizations* that give the same transfer function

- Numerator dynamics; consider the following TF

$$\frac{Y(s)}{U(s)} = \frac{s+1}{(s^2+s+1)} = \frac{Y(s)}{Z(s)} \cdot \frac{Z(s)}{U(s)}$$

- Let

$$\frac{Y(s)}{Z(s)} = s+1 \quad (1) \qquad \frac{Z(s)}{U(s)} = \frac{1}{(s^2+s+1)} \quad (2)$$

- From Eq. (2) we have

$$\ddot{z} + \dot{z} + z = u, \quad \text{or} \quad \ddot{z} = -\dot{z} - z + u$$

- State space

$$\begin{aligned} x_1 &= z, & \dot{x}_1 &= x_2 \\ x_2 &= \dot{z}, & \dot{x}_2 &= -x_2 - x_1 + u \end{aligned}$$

- Same as the first example
- This is because the dominator is identical

- From Eq. (1) we have

$$y = \dot{z} + z, \quad \text{or} \quad y = x_2 + x_1$$

- State space model is then given by

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & -1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u$$

$$y = \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + [0]u$$

- Same as the first example, except that the  $H$  matrix is different
- Note that the eigenvalues of the  $F$  matrix are the same as before, because it's the same matrix (again due to the fact that the dominator is identical)



- Another example; consider the following TF

$$\frac{Y(s)}{U(s)} = \frac{s^2 + 3s + 2}{(s^2 + s + 1)} = \frac{Y(s)}{Z(s)} \cdot \frac{Z(s)}{U(s)}$$

- Let

$$\frac{Y(s)}{Z(s)} = s^2 + 3s + 2 \quad (1)$$

$$\frac{Z(s)}{U(s)} = \frac{1}{(s^2 + s + 1)} \quad (2)$$

- From Eq. (2) we have

$$\ddot{z} + \dot{z} + z = u, \quad \text{or} \quad \ddot{z} = -\dot{z} - z + u$$

- State space

$$x_1 = z,$$

$$\dot{x}_1 = x_2$$

$$x_2 = \dot{z},$$

$$\dot{x}_2 = -x_2 - x_1 + u$$

- Same as the first example
- This is again because the dominator is identical
- So the  $F$  and  $B$  matrices are the same



- From Eq. (1) we have  $y = \ddot{z} + 3\dot{z} + 2z$
- Now use the following

$$\ddot{z} = \dot{x}_2 = -x_2 - x_1 + u$$

$$\dot{z} = x_2$$

$$z = x_1$$

- Then  $y$  becomes

$$y = -x_2 - x_1 + u + 3x_2 + 2x_1 = x_1 + 2x_2 + u$$

- State space model is then

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & -1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u$$

$$y = \begin{bmatrix} 1 & 2 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + [1]u$$

Order num can't  
be > order of  
den. not realizable

- When the numerator has the same order as the denominator then the  $D$  matrix is nonzero (always happens)
  - This is called “Direct Transmission” because on output is detected directly when an input is applied



- Past example, slight return; consider the following TF

$$\frac{Y(s)}{U(s)} = \frac{s^2 + 3s + 2}{(s^2 + s + 1)} = \frac{Y(s)}{Z(s)} \cdot \frac{Z(s)}{U(s)}$$

- Let

$$\frac{Y(s)}{Z(s)} = s^2 + 3s + 2 \quad (1)$$

$$\frac{Z(s)}{U(s)} = \frac{1}{(s^2 + s + 1)} \quad (2)$$

- From Eq. (2) we have

$$\ddot{z} + \dot{z} + z = u, \quad \text{or} \quad \ddot{z} = -\dot{z} - z + u$$

- Use different states now

$$x_1 = \dot{z},$$

$$\dot{x}_1 = -x_1 - x_2 + u$$

$$x_2 = z,$$

$$\dot{x}_2 = x_1$$

- A little different than before
- So the  $F$  and  $B$  matrices are different, but the eigenvalues of  $F$  are the same



- From Eq. (1) we have  $y = \ddot{z} + 3\dot{z} + 2z$
- Now use the following

$$\ddot{z} = \dot{x}_1 = -x_1 - x_2 + u$$

$$\dot{z} = x_1$$

$$z = x_2$$

- Then  $y$  becomes

$$y = -x_1 - x_2 + u + 3x_1 + 2x_2 = 2x_1 + x_2 + u$$

- State space model is then

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} -1 & -1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} u$$

$$y = \begin{bmatrix} 2 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + [1]u$$

- The  $D$  matrix is the same as before, but others have changed
- The transfer function is identical though
  - Just a different realization of the same system



- First compute  $(sI - F)^{-1}$

$$(sI - F)^{-1} = \begin{bmatrix} s+1 & 1 \\ -1 & s \end{bmatrix}^{-1} = \frac{1}{s^2 + s + 1} \begin{bmatrix} s & -1 \\ 1 & s+1 \end{bmatrix}$$

- Then

$$\frac{Y(s)}{U(s)} = H (sI - F)^{-1} B + D$$

$$= \frac{1}{s^2 + s + 1} \begin{bmatrix} 2 & 1 \end{bmatrix} \begin{bmatrix} s & -1 \\ 1 & s+1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} + 1$$

$$= \frac{1}{s^2 + s + 1} \begin{bmatrix} 2 & 1 \end{bmatrix} \begin{bmatrix} s \\ 1 \end{bmatrix} + 1$$

$$= \frac{2s + 1}{s^2 + s + 1} + 1$$

$$= \frac{2s + 1 + s^2 + s + 1}{s^2 + s + 1}$$

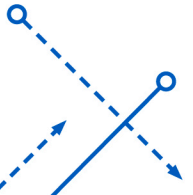
$$= \frac{s^2 + 3s + 2}{s^2 + s + 1} \quad \checkmark$$

- General ODE

$$\begin{aligned} \frac{d^n y}{dt^n} + a_{n-1} \frac{d^{n-1} y}{dt^{n-1}} + \cdots + a_1 \frac{dy}{dt} + a_0 y \\ = b_n \frac{d^n u}{dt^n} + b_{n-1} \frac{d^{n-1} u}{dt^{n-1}} + \cdots + b_1 \frac{du}{dt} + b_0 u \end{aligned}$$

- State space matrices

$$\left. \begin{aligned} x_1 &= y \\ x_2 &= \frac{dy}{dt} \\ &\vdots \\ x_n &= \frac{d^{n-1} y}{dt^{n-1}} \end{aligned} \right| \begin{aligned} F &= \begin{bmatrix} 0 & 1 & 0 & \cdots & 0 \\ 0 & 0 & 1 & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & 1 \\ -a_0 & -a_1 & -a_2 & \cdots & -a_{n-1} \end{bmatrix} \\ B &= [0 \quad 0 \quad \cdots \quad 1]^T \\ H &= [(b_0 - b_n a_0) \quad (b_1 - b_n a_1) \quad \cdots \quad (b_{n-1} - b_n a_{n-1})] \\ D &= b_n \end{aligned}$$



- Consider a multi-input-multi-output (MIMO) model

$$\dot{\mathbf{x}} = \mathbf{F} \mathbf{x} + \mathbf{B} \mathbf{u} \quad (1)$$

$$\mathbf{y} = \mathbf{H} \mathbf{x} + \mathbf{D} \mathbf{u}$$

- Now define  $\mathbf{x} = \mathbf{T} \mathbf{z}$ , where  $\mathbf{T}$  is constant and nonsingular
- Then  $\dot{\mathbf{x}} = \mathbf{T} \dot{\mathbf{z}}$
- Multiplying both sides by  $\mathbf{T}^{-1}$ , and substituting the state space model from Eq. (1) gives

$$\dot{\mathbf{z}} = \mathbf{A} \mathbf{z} + \mathbf{G} \mathbf{u}$$

$$\mathbf{y} = \mathbf{C} \mathbf{z} + \mathbf{D} \mathbf{u}$$

where

$$\mathbf{A} \equiv \mathbf{T}^{-1} \mathbf{F} \mathbf{T}$$

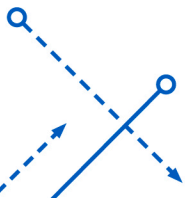
$$\mathbf{G} \equiv \mathbf{T}^{-1} \mathbf{B}$$

$$\mathbf{C} \equiv \mathbf{H} \mathbf{T}$$

- Compute the transfer function matrix from  $\mathbf{U}(s)$  to  $\mathbf{Y}(s)$

$$\begin{aligned}
 \mathbf{Y}(s) &= \left[ C (sI - A)^{-1} G + D \right] \mathbf{U}(s) \\
 &= \left[ H T (sI - T^{-1} F T)^{-1} T^{-1} B + D \right] \mathbf{U}(s) \\
 &= \left\{ H T \left[ T^{-1} (sI - F) T \right]^{-1} T^{-1} B + D \right\} \mathbf{U}(s) \\
 &= \left[ H T T^{-1} (sI - F)^{-1} T T^{-1} B + D \right] \mathbf{U}(s) \\
 &= \left[ H (sI - F)^{-1} B + D \right] \mathbf{U}(s)
 \end{aligned}$$

- Same transfer function as before
- Similarity transformation does not change transfer function!
- Many possible forms
  - Diagonal form for  $A$  (assuming distinct eigenvalues)  
can be found by using the eigenvector matrix of  $F$  for  $T$





- Cayley-Hamilton (CH) Theorem states that a matrix satisfies its own characteristic equation

- Say characteristic equation of a matrix  $A$  is given by

$$\lambda^n + a_{n-1}\lambda^{n-1} + a_{n-2}\lambda^{n-2} + \cdots + a_1\lambda + a_0 = 0$$

- Then the CH Theorem states

$$A^n + a_{n-1}A^{n-1} + a_{n-2}A^{n-2} + \cdots + a_1A + a_0I = 0$$

- Proof for diagonalizable case (general case can be proven too)

- Compute the following

$$M = A^n + a_{n-1}A^{n-1} + a_{n-2}A^{n-2} + \cdots + a_1A + a_0I \quad (1)$$

- Need to show that  $M = 0$

- Take an eigenvalue/eigenvector decomposition of  $A$

$$A = V \Lambda V^{-1} \quad \rightarrow \quad V^{-1}A V = \Lambda \quad (2)$$

where  $\Lambda$  is a diagonal matrix of the eigenvalues and  $V$  is the eigenvector matrix

- Left multiply Eq. (1) by  $V^{-1}$  and right multiply by  $V$

$$V^{-1} M V = V^{-1} A^n V + a_{n-1} V^{-1} A^{n-1} V + \cdots + a_0 I$$

- But

$$V^{-1} A^j V = \underbrace{(V^{-1} A V)(V^{-1} A V) \cdots (V^{-1} A V)}_{j \text{ times}} = \Lambda^j$$

where Eq. (2) was used

- So

$$V^{-1} M V = \begin{bmatrix} \lambda_1^n & & 0 \\ & \ddots & \\ 0 & & \lambda_n^n \end{bmatrix} + a_{n-1} \begin{bmatrix} \lambda_1^{n-1} & & 0 \\ & \ddots & \\ 0 & & \lambda_n^{n-1} \end{bmatrix} + \cdots + a_0 \begin{bmatrix} 1 & & 0 \\ & \ddots & \\ 0 & & 1 \end{bmatrix}$$

- Because every eigenvalue satisfies its characteristic equation then the right-hand-side of this equation is zero
- This can only be true if  $M = 0$  since  $V \neq 0$



- CH Theorem shows that every  $f(A)$  function of a square matrix can be computed as a linear combination of the first  $n - 1$  powers of  $A$

$$f(A) = \gamma_0 I + \gamma_1 A + \cdots + \gamma_{n-1} A^{n-1}$$

- Why is this true?
- From the CH Theorem we have

$$A^n = \underbrace{-a_{n-1}A^{n-1} - a_{n-2}A^{n-2} - \cdots - a_0 I}_{\text{RHS}}$$

- So then we have

$$\begin{aligned} A^{n+1} &= -a_{n-1}A^n - a_{n-2}A^{n-1} - \cdots - a_0 A \\ &= -a_{n-1}[\text{RHS}] - a_{n-2}A^{n-1} - \cdots - a_0 A \\ &= \gamma_{n-1}A^{n-1} - \gamma_{n-2}A^{n-2} + \cdots + \gamma_0 I \end{aligned}$$

where  $\gamma_i$ 's are a function of the  $a_i$ 's



- One way to compute the coefficients for distinct eigenvalues is to note that  $f(A)$  is the same linear combination of the powers of  $A$  as that  $f(\Lambda)$  is of the powers of  $\Lambda$
- For non-distinct eigenvalues first define

$$h(\lambda) \equiv \gamma_0 + \gamma_1 \lambda + \cdots + \gamma_{n-1} \lambda^{n-1}$$

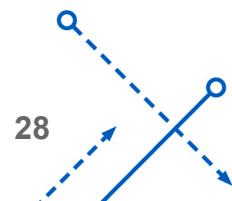
- The  $n$  unknowns  $\gamma_i$ 's can be solved using the following  $n$  equations

$$f^{(\ell)}(\lambda) = h^{(\ell)}(\lambda) \text{ for } \ell = 0, 1, \dots, n_i - 1 \text{ and } i = 1, 2, \dots, m$$

where

$$f^{(\ell)}(\lambda) \equiv \left. \frac{d^\ell f(\lambda)}{d\lambda^\ell} \right|_{\lambda=\lambda_i}$$

$$n = \sum_{i=1}^m n_i$$



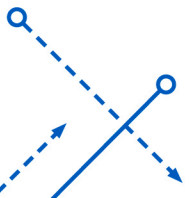
$$A = \begin{bmatrix} 1 & 0 & 4 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix}, \quad \text{compute } A^{100}$$

- Let  $f(A) = A^{100}$ , then

$$A^{100} = \gamma_0 I + \gamma_1 A + \gamma_2 A^2$$

- Eigenvalue matrix is given by  $\Lambda = \text{diag}([1 \quad +j \quad -j])$
- They are distinct so we can now use

$$\begin{aligned} F(\Lambda) &= \begin{bmatrix} 1^{100} & 0 & 0 \\ 0 & j^{100} & 0 \\ 0 & 0 & (-j)^{100} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \\ &= \gamma_0 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} + \gamma_1 \begin{bmatrix} 1 & 0 & 0 \\ 0 & j & 0 \\ 0 & 0 & -j \end{bmatrix} + \gamma_2 \begin{bmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \end{aligned}$$



- Gives the following three equations

$$\gamma_0 + \gamma_1 + \gamma_2 = 1$$

$$\gamma_0 + \gamma_1 j - \gamma_2 = 1$$

$$\gamma_0 - \gamma_1 j + \gamma_2 = 1$$

- By inspection the solutions are

$$\gamma_0 = 1, \quad \gamma_1 = 0, \quad \gamma_2 = 0$$

- Then

$$A = \begin{bmatrix} 1 & 0 & 4 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix}^{100} = \gamma_0 I + \gamma_1 A + \gamma_2 A^2$$

$$= 1 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} + 0 \begin{bmatrix} 1 & 0 & 4 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix} + 0 \begin{bmatrix} 1 & 0 & 4 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix}^2$$

$$= I$$



$$A = \begin{bmatrix} 0 & 0 & -2 \\ 0 & 1 & 0 \\ 1 & 0 & 3 \end{bmatrix}, \quad \text{compute } e^{At}$$

- Eigenvalue matrix is given by  $\Lambda = \text{diag}([1 \ 1 \ 2])$ 
  - Note that there are repeated eigenvalues
- Let  $f(\lambda) = e^{\lambda t}$ , then  $f'(\lambda) = t e^{\lambda t}$  (note: derivative w.r.t.  $\lambda$  to not  $t$ )

$$h(\lambda) \equiv \gamma_0 + \gamma_1 \lambda + \gamma_2 \lambda^2, \quad h'(\lambda) = \gamma_1 + 2\gamma_2 \lambda$$

- Need to solve the following equations

$$f(1) = h(1) : \quad e^t = \gamma_0 + \gamma_1 + \gamma_2$$

$$f'(1) = h'(1) : \quad t e^t = \gamma_1 + 2\gamma_2$$

$$f(2) = h(2) : \quad e^{2t} = \gamma_0 + 2\gamma_1 + 4\gamma_2$$

- Solving these equations gives

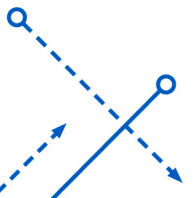
$$\gamma_0 = -2te^t + e^{2t}$$

$$\gamma_1 = 3te^t + 2e^t - 2e^{2t}$$

$$\gamma_2 = e^{2t} - e^t - te^t$$

- Then

$$\begin{aligned} e^{At} &= \gamma_0 I + \gamma_1 A + \gamma_2 A^2 \\ &= \gamma_0 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} + \gamma_1 \begin{bmatrix} 0 & 0 & -2 \\ 0 & 1 & 0 \\ 1 & 0 & 3 \end{bmatrix} + \gamma_2 \begin{bmatrix} -2 & 0 & -6 \\ 0 & 1 & 0 \\ 3 & 0 & 7 \end{bmatrix} \\ &= \begin{bmatrix} 2e^t - e^{2t} & 0 & 2e^t - 2e^{2t} \\ 0 & e^t & 0 \\ e^{2t} - e^t & 0 & 2e^{2t} - e^t \end{bmatrix} \end{aligned}$$





- Look at an expansion of  $e^{\lambda t}$

$$e^{\lambda t} = 1 + \lambda t + \frac{\lambda^2 t^2}{2!} + \dots$$

- Then

$$e^{At} = I + At + A^2 \frac{t^2}{2!} + \dots = \sum_{k=0}^{\infty} \frac{1}{k!} t^k A^k \quad (1)$$

- Some useful properties

$$e^0 = I$$

$$e^{A(t_1+t_2)} = e^{At_1} e^{At_2}$$

$$[e^{At}]^{-1} = e^{-At}$$

- The last one is especially useful because you don't need to take the inverse of an entire matrix!

- Take time derivative of Eq. (1)

$$\frac{d}{dt}e^{At} = \sum_{k=1}^{\infty} \frac{1}{(k-1)!} t^{k-1} A^k \quad (2)$$

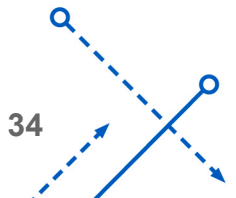
- This can be rewritten as

$$\frac{d}{dt}e^{At} = A \left( \sum_{k=0}^{\infty} \frac{1}{k!} t^k A^k \right) = \left( \sum_{k=0}^{\infty} \frac{1}{k!} t^k A^k \right) A$$

- Note that because of the  $A^k$  term in Eq. (2), when the sum was rewritten an extra  $A$  matrix is needed, which can be pulled out
- Also, note that the  $A$  can be pulled out on either the right or left side
- Then from Eq. (1)

$$\frac{d}{dt}e^{At} = A e^{At} = e^{At} A \quad (3)$$

- Note that  $e^{(A+B)t} \neq e^{At} e^{Bt}$  in general unless  $AB = BA$ , i.e. the matrices  $A$  and  $B$  commute



- Laplace Transform of the time derivative of a function

$$\mathcal{L}[f'(t)] = s\mathcal{L}[f(t)] - f(0)$$

- Using this concept on the time derivative of  $e^{At}$  gives

$$A \mathcal{L}[e^{At}] = s\mathcal{L}[e^{At}] - e^0$$

where Eq. (3) was used

- Note that the first form of Eq. (3) was specifically used so that  $A$  appears on the left (can be pulled on the right too since  $s$  is a scalar)
- Noting that  $e^0 = I$  and collecting terms yields

$$(sI - A) \mathcal{L}[e^{At}] = I$$

- Then

$$\mathcal{L}[e^{At}] = (sI - A)^{-1}$$

- Taking the inverse Laplace Transform of both sides gives

$$e^{At} = \mathcal{L}^{-1} [(sI - A)^{-1}]$$

$$A = \begin{bmatrix} 0 & 0 & -2 \\ 0 & 1 & 0 \\ 1 & 0 & 3 \end{bmatrix}, \quad \text{compute } e^{At}, \quad (sI - A) = \begin{bmatrix} s & 0 & 2 \\ 0 & s-1 & 0 \\ -1 & 0 & s-3 \end{bmatrix}$$

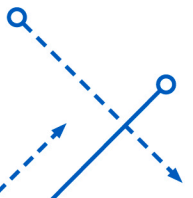
- $(sI - A)^{-1}$  can be shown to be given by

$$(sI - A)^{-1} = \begin{bmatrix} \frac{2}{s-1} - \frac{1}{s-2} & 0 & \frac{2}{s-1} - \frac{2}{s-2} \\ 0 & \frac{1}{s-1} & 0 \\ \frac{1}{s-2} - \frac{1}{s-1} & 0 & \frac{2}{s-2} - \frac{1}{s-1} \end{bmatrix}$$

- Check  $(sI - A)^{-1} (sI - A) = I$  to make sure it's correct
- Taking the inverse Laplace Transform gives

$$e^{At} = \begin{bmatrix} 2e^t - e^{2t} & 0 & 2e^t - 2e^{2t} \\ 0 & e^t & 0 \\ e^{2t} - e^t & 0 & 2e^{2t} - e^t \end{bmatrix}$$

- Same result as before



- Consider the following time-varying system

$$\dot{\mathbf{x}}(t) = F(t) \mathbf{x}(t) + B(t) \mathbf{u}(t), \quad \mathbf{x}(t_0) = \mathbf{x}_0$$

- Claim the solution for  $\mathbf{x}(t)$  is given by

$$\mathbf{x}(t) = \Phi(t, t_0) \mathbf{x}(t_0) + \int_{t_0}^t \Phi(t, \tau) B(\tau) \mathbf{u}(\tau) d\tau$$

where  $\Phi(t, \tau)$  is the *state transition matrix*

- Some useful properties

$$\frac{d}{dt} \Phi(t, \tau) = F(t) \Phi(t, \tau)$$

$$\Phi(\tau, \tau) = I$$

$$\Phi(\tau, t) = \Phi^{-1}(t, \tau)$$

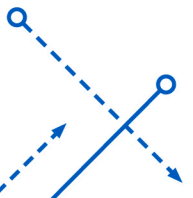
$$\Phi(t_2, t_0) = \Phi(t_2, t_1) \Phi(t_1, t_0)$$

- Take derivative to prove it is indeed the solution

$$\begin{aligned}
 \dot{\mathbf{x}}(t) &= \frac{d}{dt} \Phi(t, t_0) \mathbf{x}(t_0) + \frac{d}{dt} \int_{t_0}^t \Phi(t, \tau) B(\tau) \mathbf{u}(\tau) d\tau \\
 &= F(t) \Phi(t, t_0) \mathbf{x}(t_0) + \overset{I}{\Phi(t, t)} B(t) \mathbf{u}(t) + \int_{t_0}^t F(t) \Phi(t, \tau) B(\tau) \mathbf{u}(\tau) d\tau \\
 &= F(t) \underbrace{\left[ \Phi(t, t_0) \mathbf{x}(t_0) + \int_{t_0}^t \Phi(t, \tau) B(\tau) \mathbf{u}(\tau) d\tau \right]}_{\mathbf{x}(t)} + B(t) \mathbf{u}(t) \\
 &= F(t) \mathbf{x}(t) + B(t) \mathbf{u}(t) \quad \checkmark
 \end{aligned}$$

- Check initial conditions

$$\begin{aligned}
 \mathbf{x}(t_0) &= \overset{I}{\Phi(t_0, t_0)} \mathbf{x}(t_0) + \overset{0}{\int_{t_0}^{t_0} \Phi(t_0, \tau) B(\tau) \mathbf{u}(\tau) d\tau} \\
 &= \mathbf{x}_0 \quad \checkmark
 \end{aligned}$$



- State transition matrix reduces down to  $\Phi(t, \tau) = e^{F(t-\tau)}$ 
  - This is *matrix exponential* of  $F$
  - Can be found using the inverse Laplace Transform

$$e^{Ft} = \mathcal{L}^{-1} [(sI - F)^{-1}]$$

- Example

$$F = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

- Then

$$(sI - F)^{-1} = \frac{1}{s^2} \begin{bmatrix} s & 1 \\ 0 & s \end{bmatrix} = \begin{bmatrix} \frac{1}{s} & \frac{1}{s^2} \\ 0 & \frac{1}{s} \end{bmatrix}$$

- Inverse Laplace Transform gives

$$e^{Ft} = \begin{bmatrix} 1 & t \\ 0 & 1 \end{bmatrix}$$

- Sometimes using the series expansion is easier though

$$e^{Ft} = I + Ft + F^2 \frac{t^2}{2!} + \dots$$

- Look at this matrix again

$$F = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

- Since  $F^2 = 0$  then

$$\begin{aligned} e^{Ft} &= \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 0 & t \\ 0 & 0 \end{bmatrix} + 0 + 0 + \dots \\ &= \begin{bmatrix} 1 & t \\ 0 & 1 \end{bmatrix} \end{aligned}$$

- Same result as before