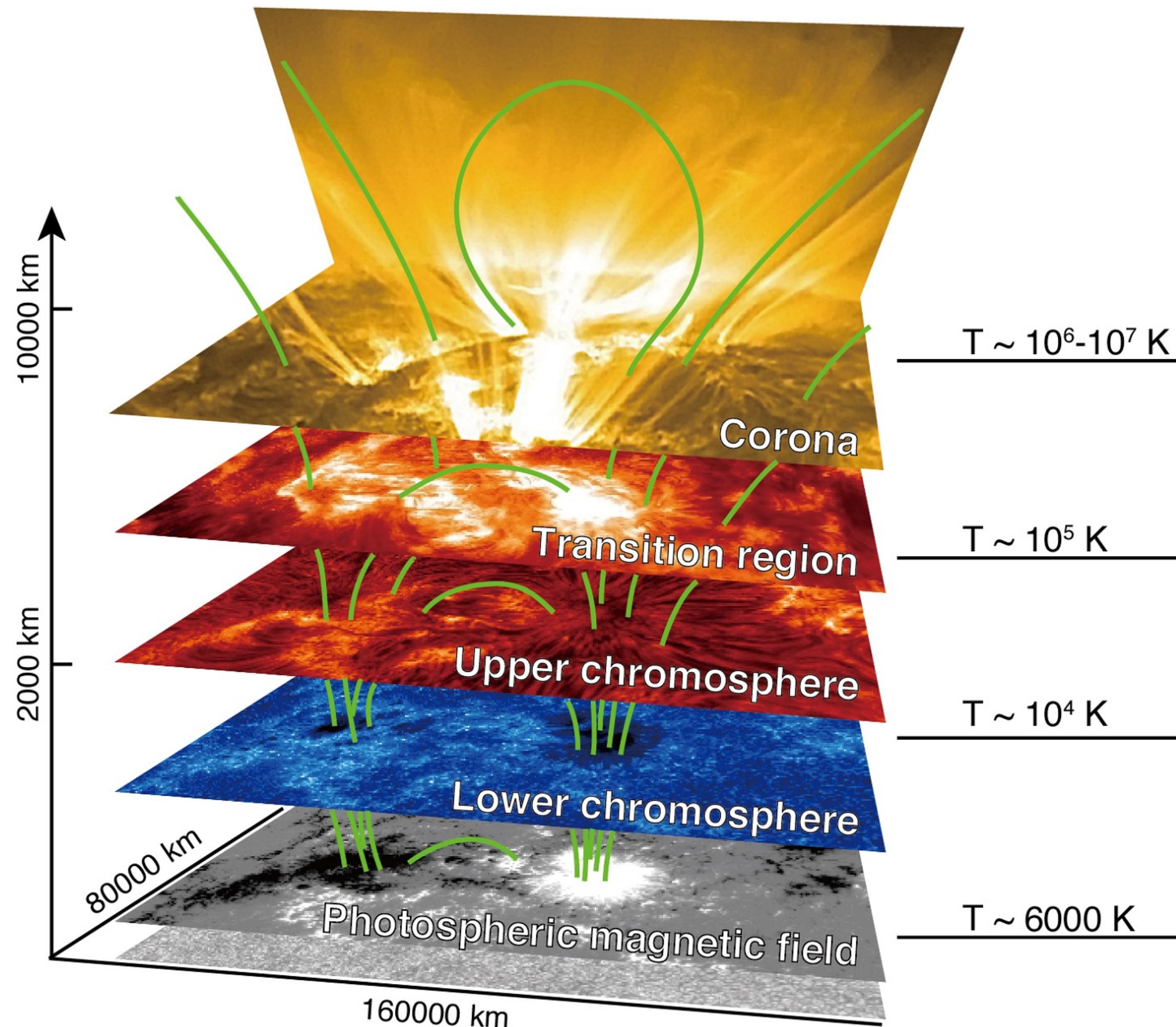


GRAHAM KERR cua & NASA/GSFC (HE/HIM/HIS)

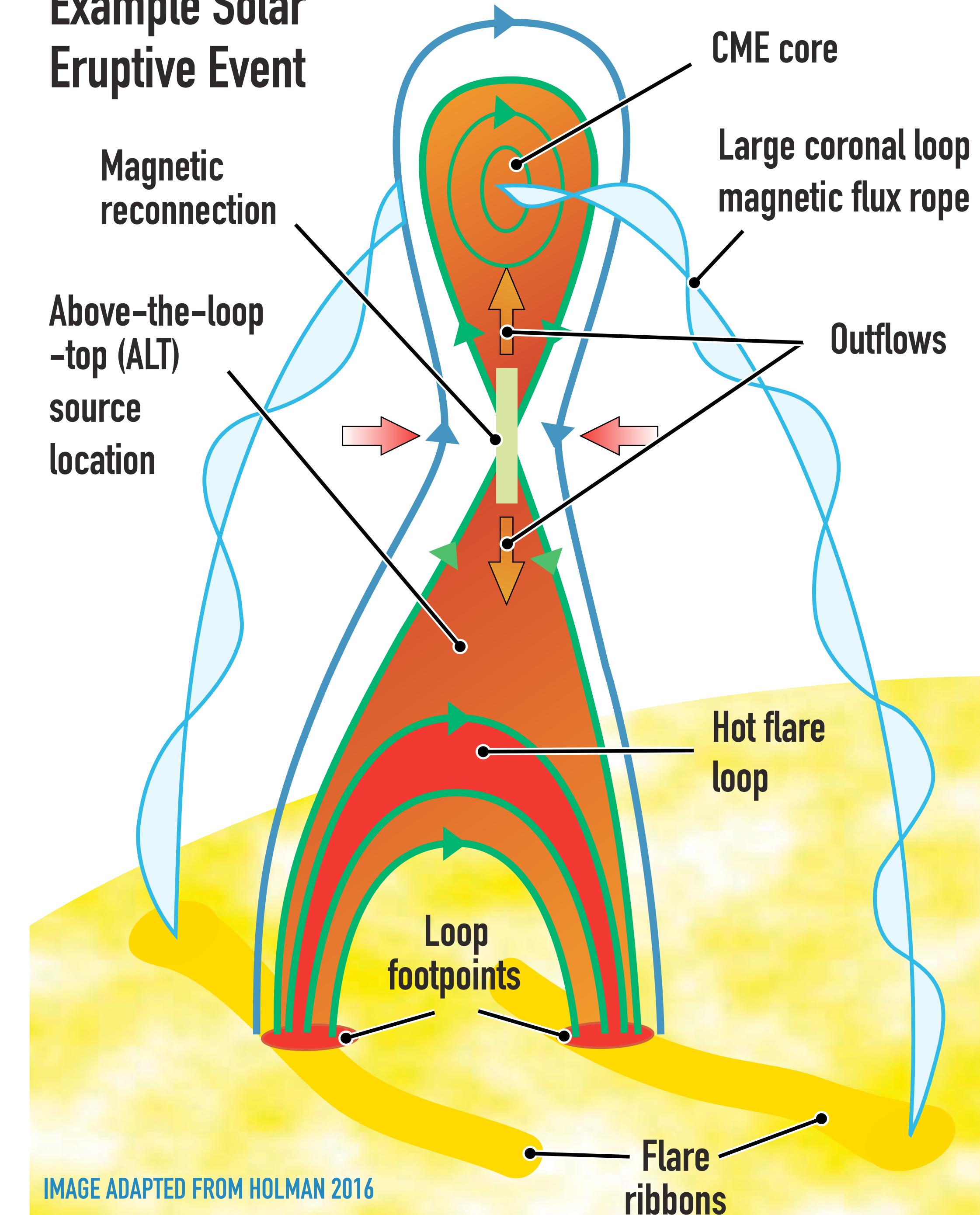
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# INTERROGATING SOLAR FLARE LOOP MODELS WITH OBSERVATIONS: RECENT SUCCESSES, CURRENT PROBLEMS, AND FUTURE DIRECTIONS

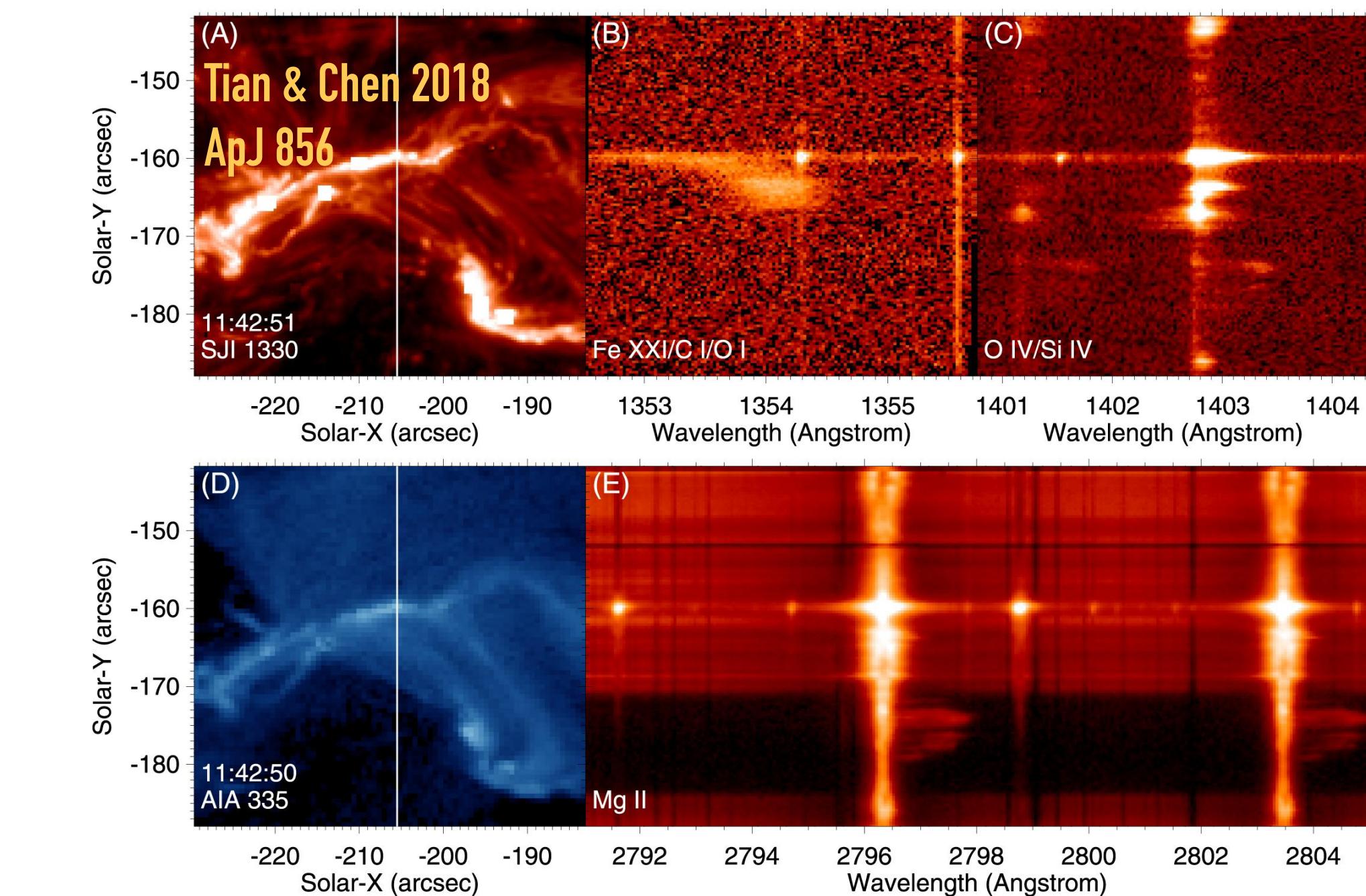
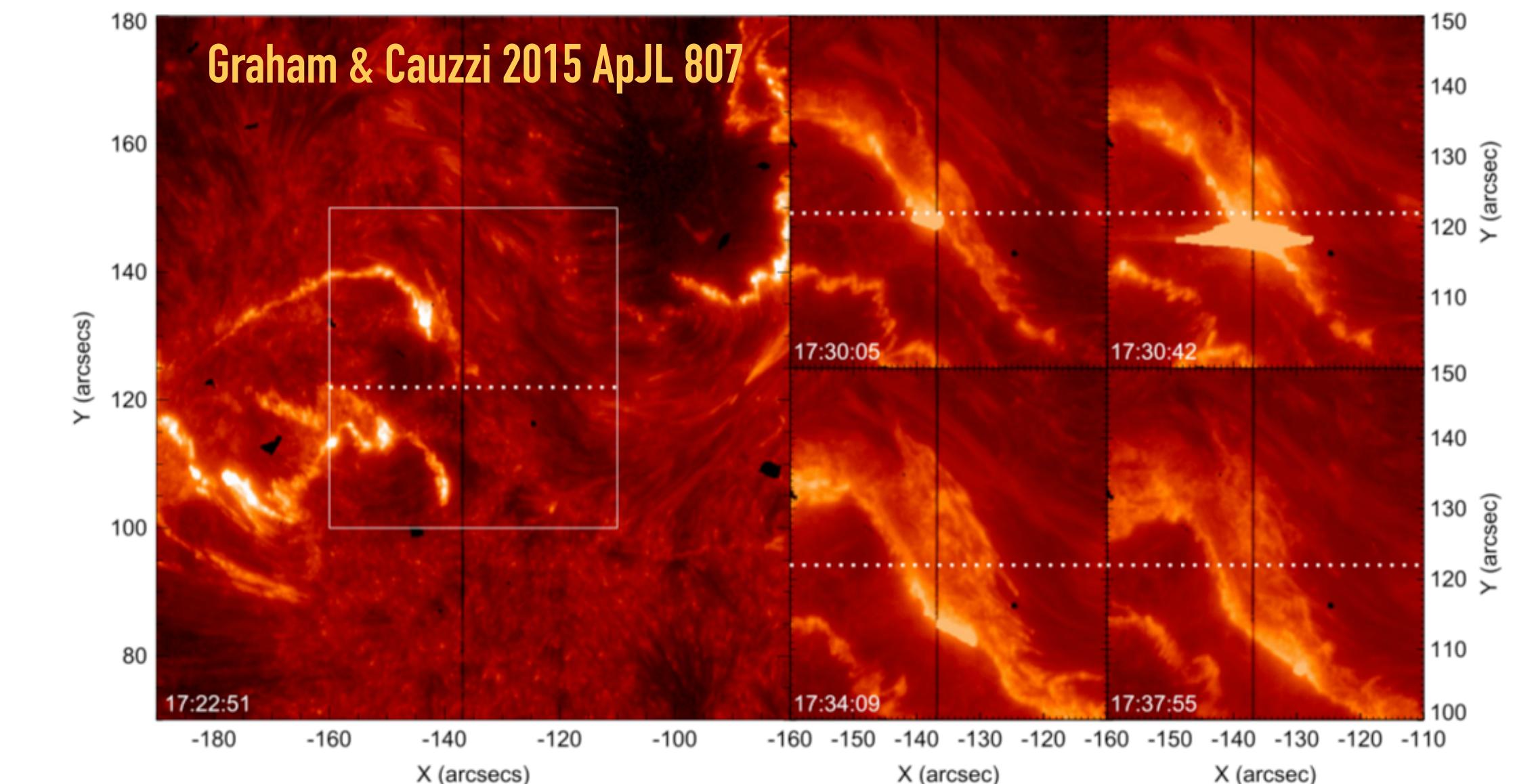
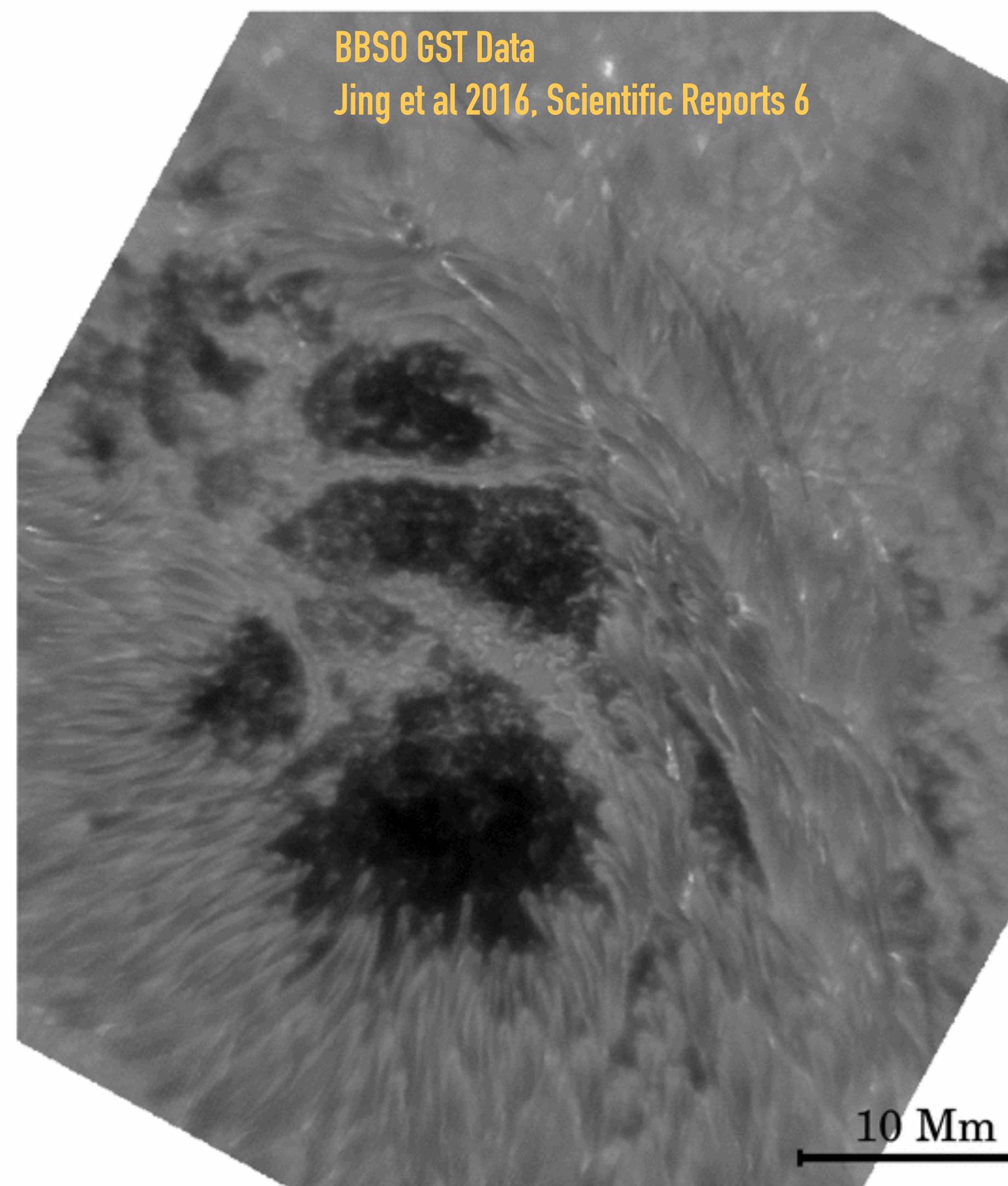
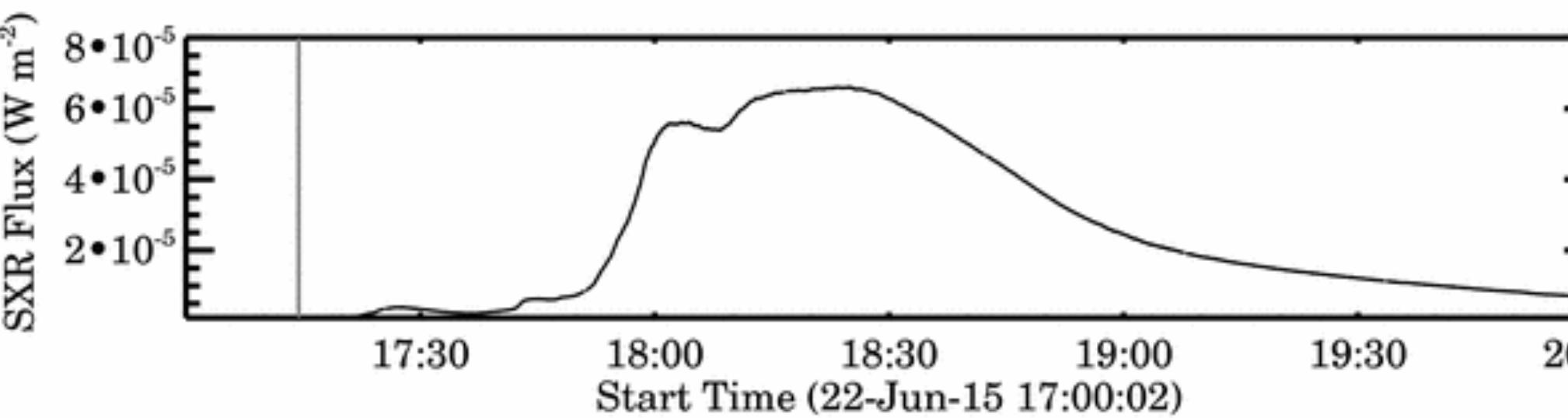
# SOLAR FLARES

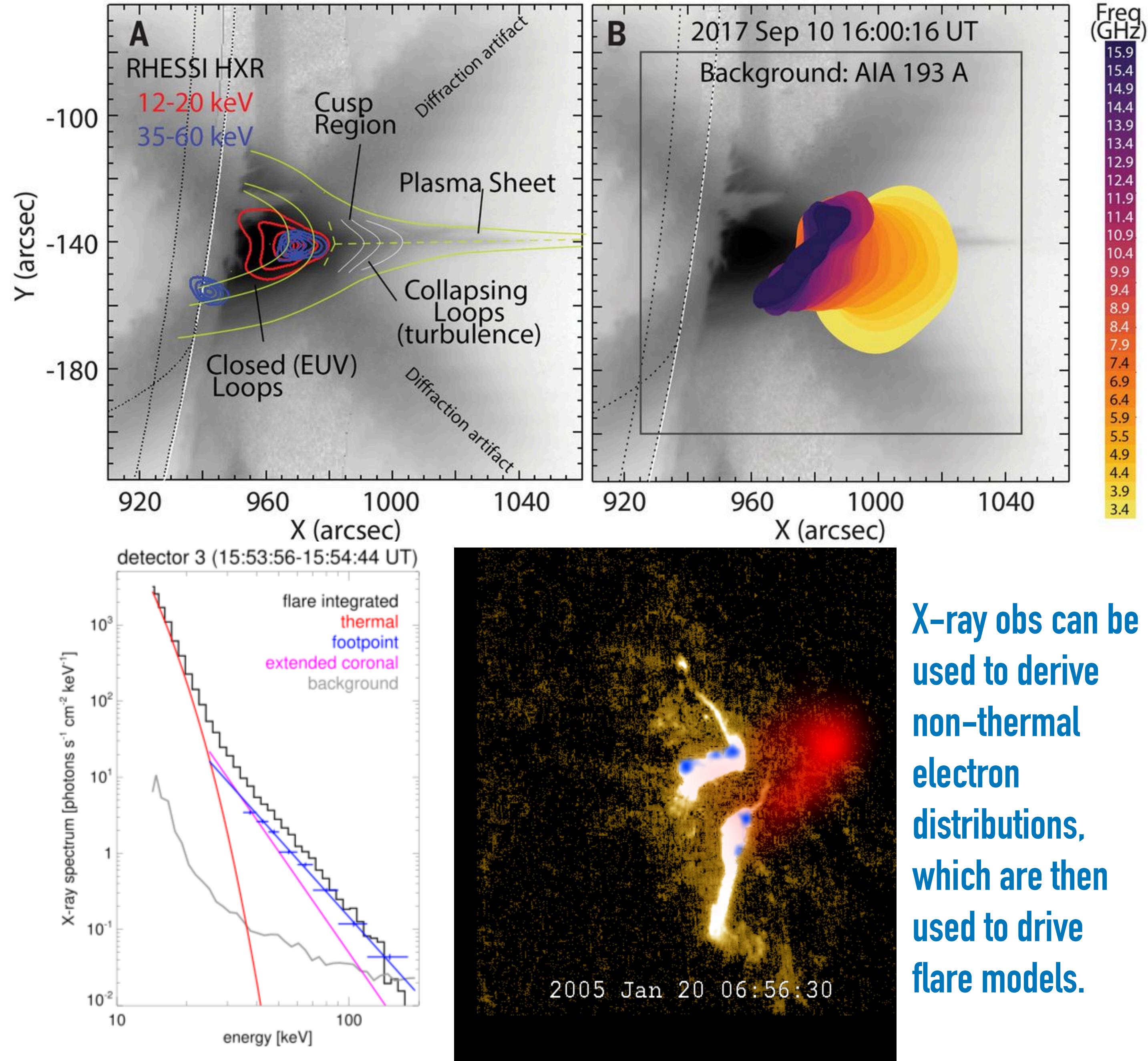


## Example Solar Eruptive Event

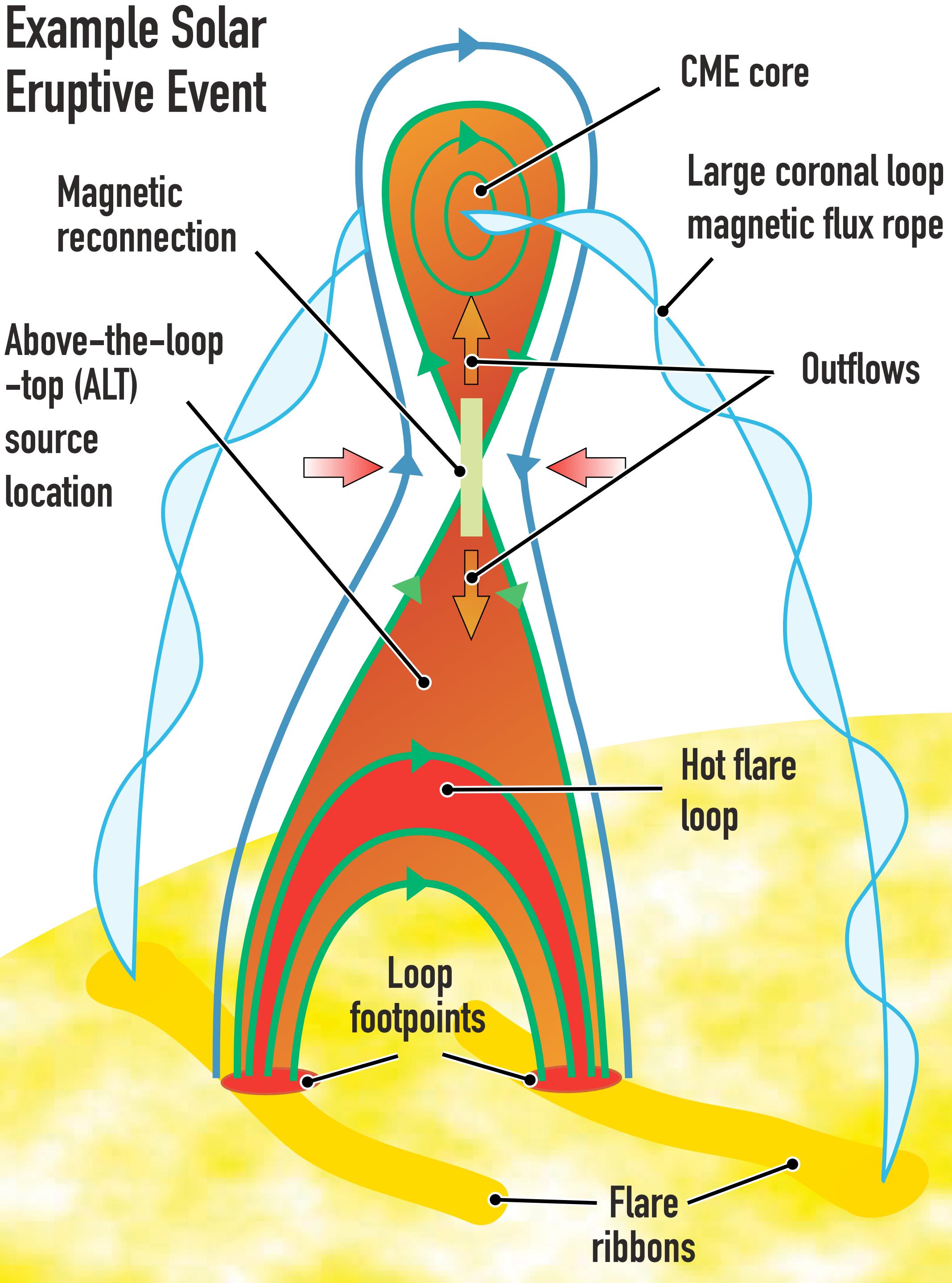


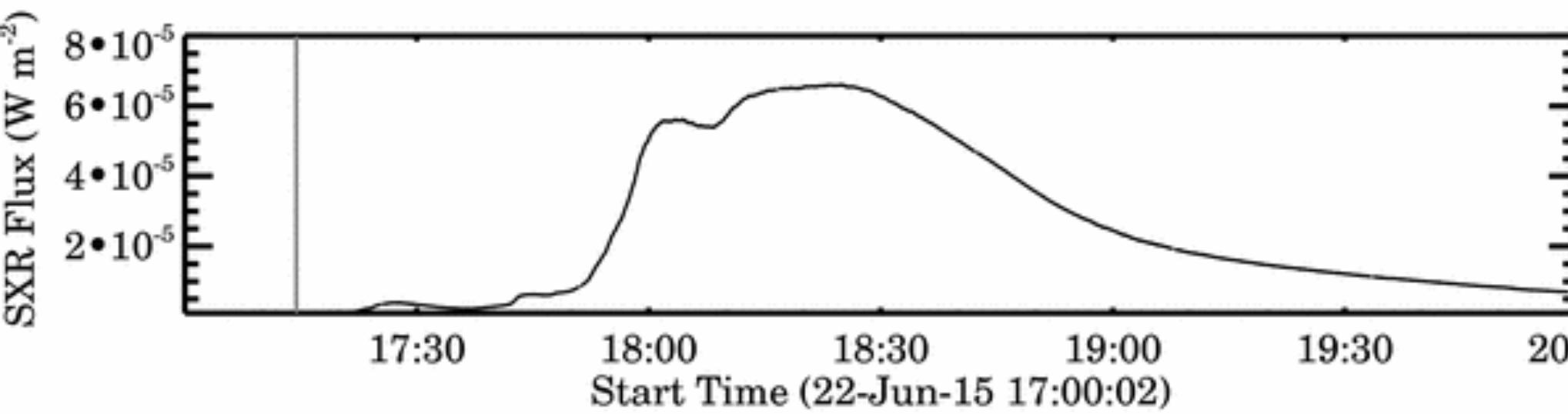
# SOLAR FLARES



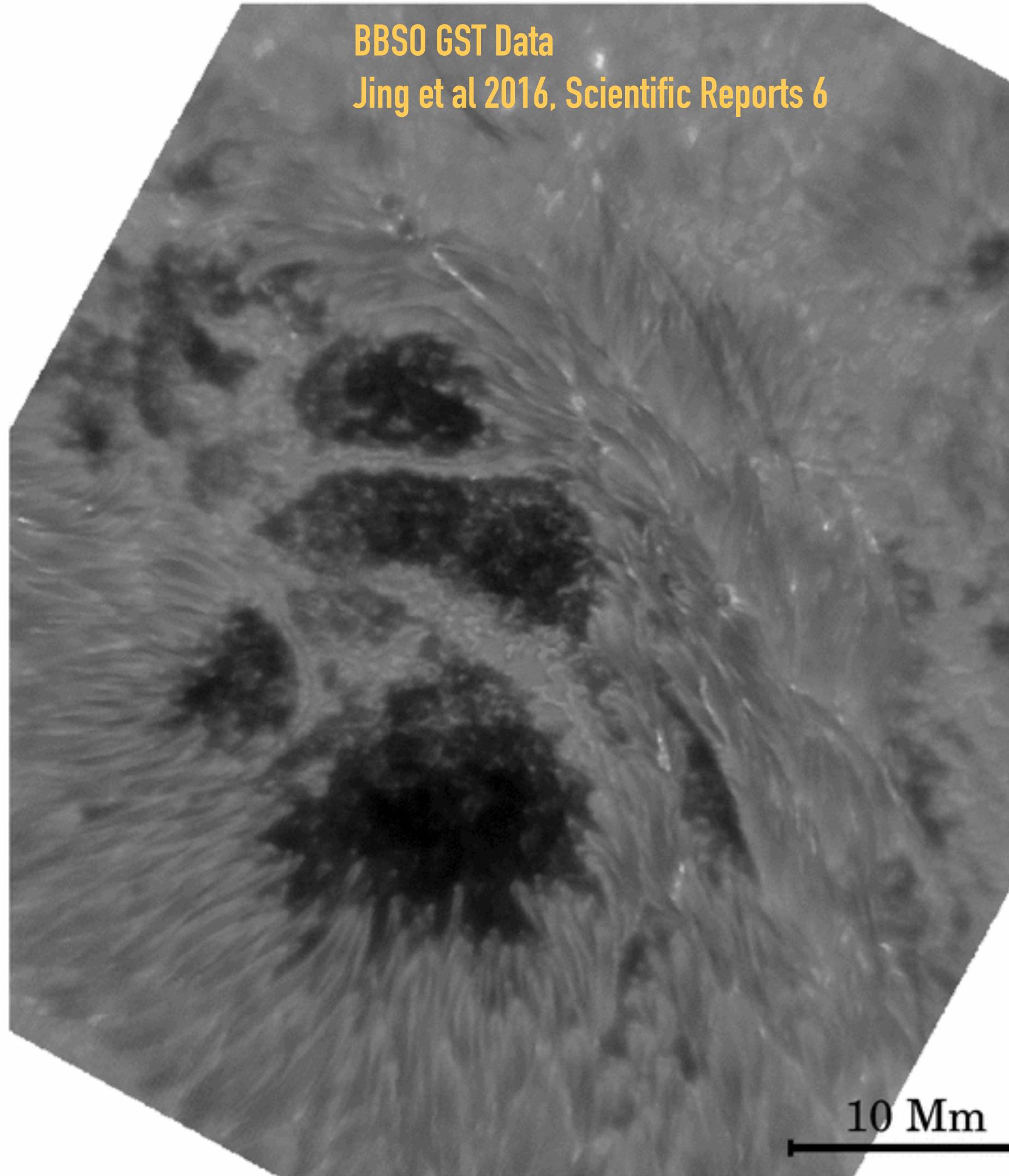


X-ray obs can be used to derive non-thermal electron distributions, which are then used to drive flare models.





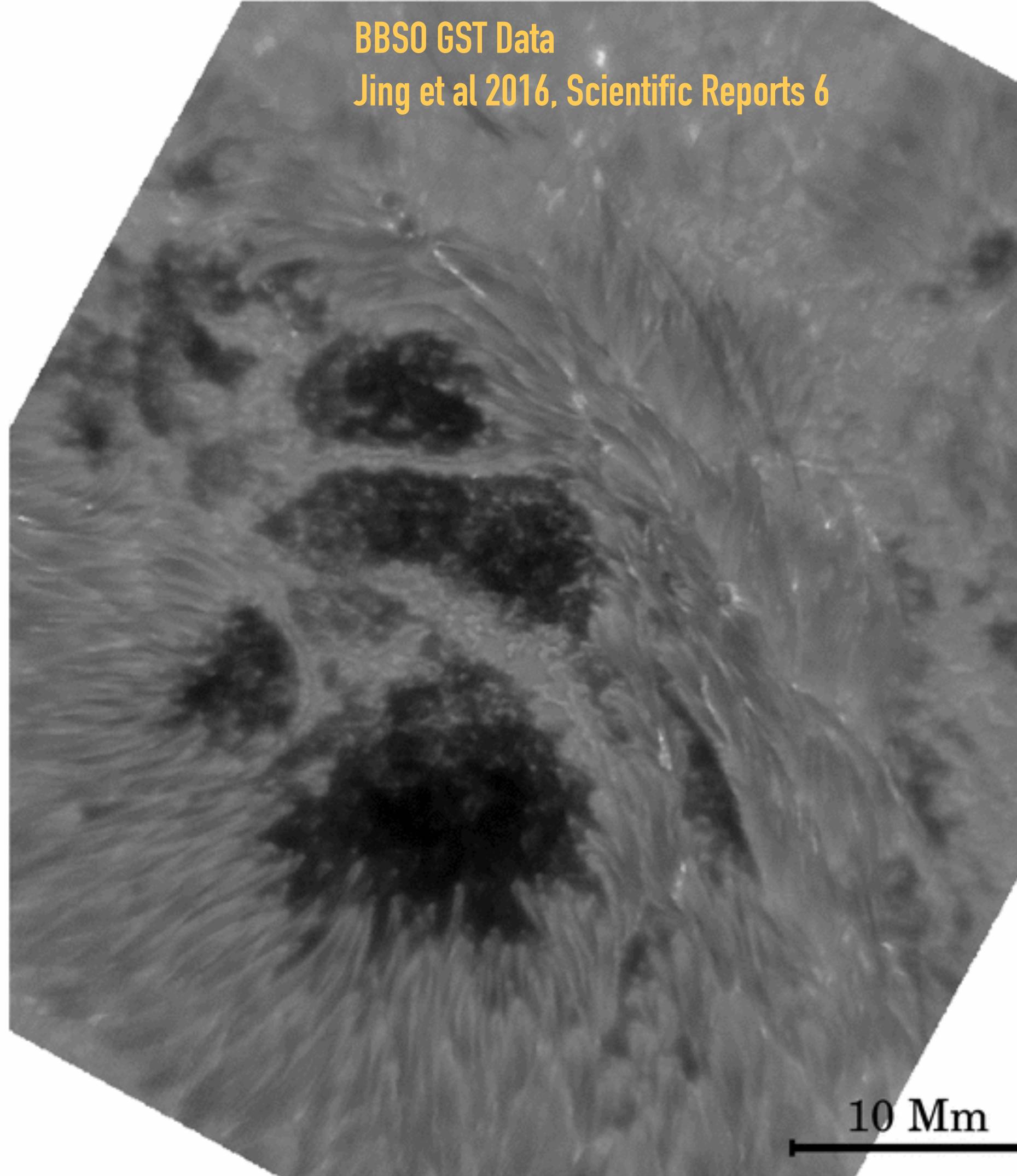
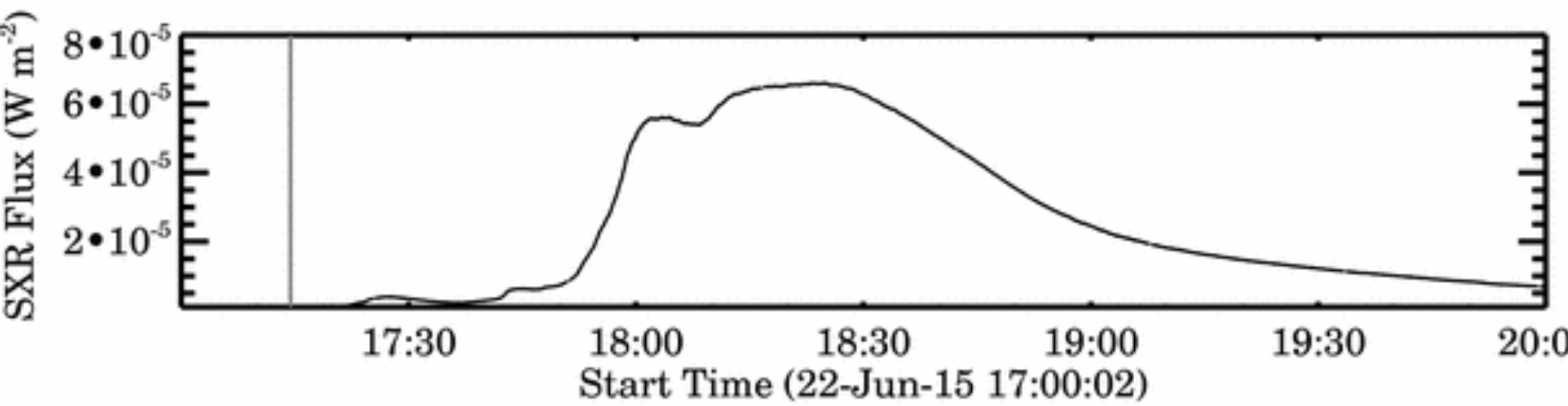
BBSO GST Data  
Jing et al 2016, Scientific Reports 6



With the fortune of increasingly high quality data comes the need for state-of-the-art numerical models to both:

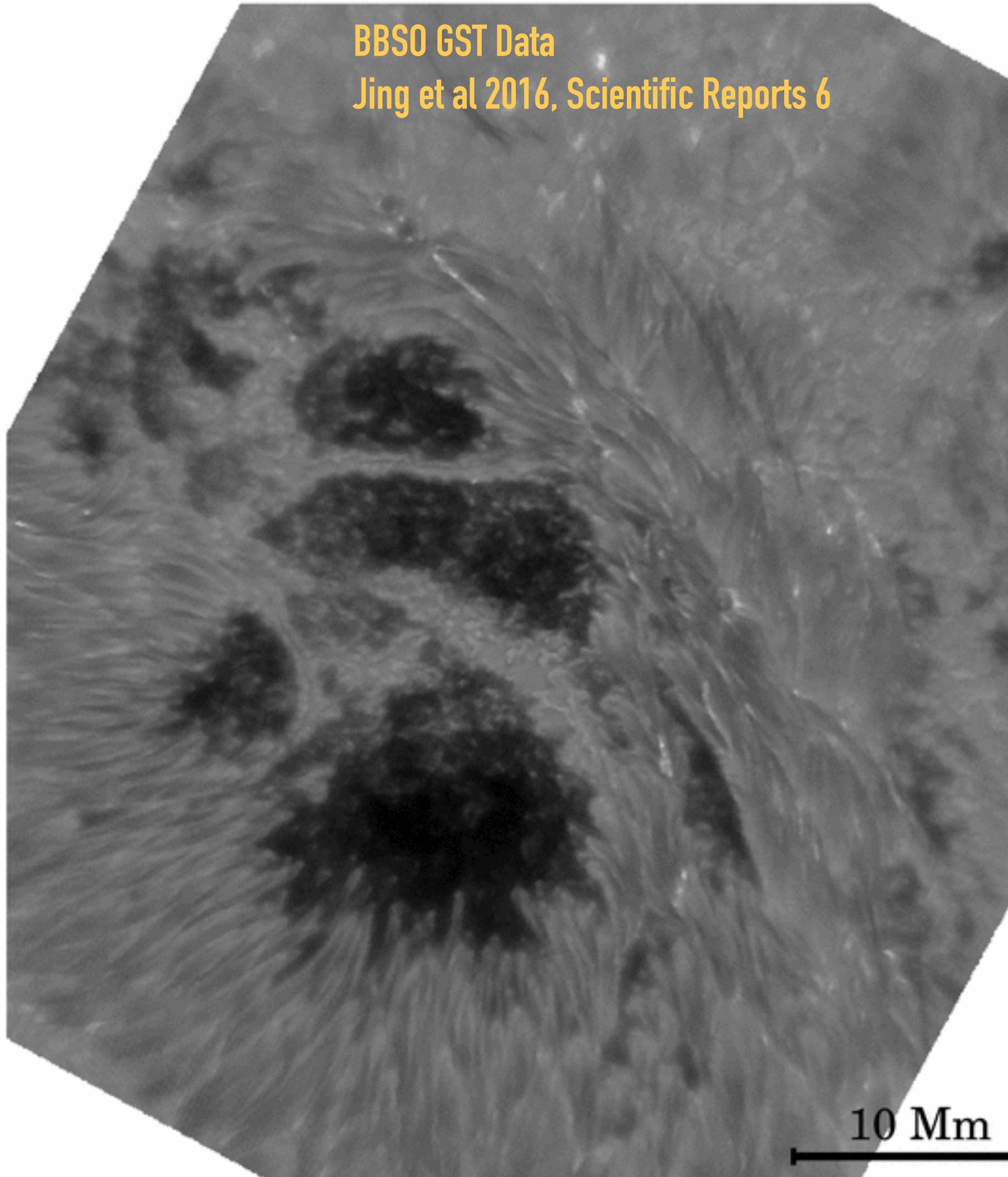
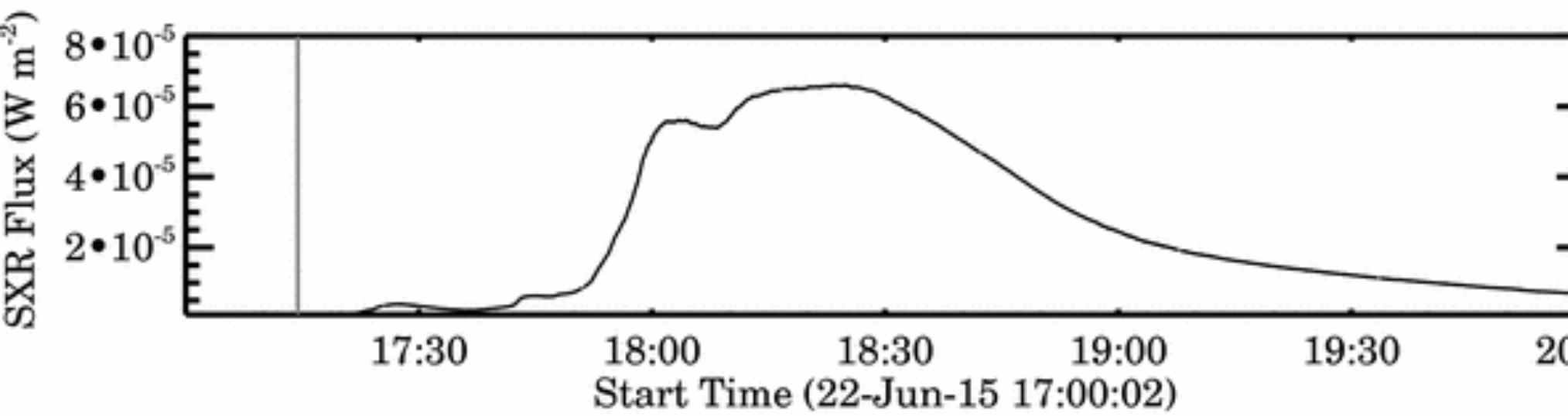
- 1) help extract and interpret what information observations contain;
- 2) determine if theory stands up to the stubborn reality of observations.

Field-aligned loop models (1D) of solar flares allow us to study the detailed physics and complex feedback between the radiation and hydrodynamics at the **high spatial and temporal resolution demanded during flares**.



- Highlight just a few recent(-ish) results.\*\*
- Discuss where models are currently not consistent with observations, and potential future directions.
- Present what I and some colleagues think is feasible way forward to make progress towards an end-to-end model of solar eruptions.

\*\* Sorry if I don't cover your particular study or favorite result... I do likely cover them in a recent review article though!



- Highlight just a few recent(-ish) results.\*\*
- Discuss where models are currently not consistent with observations, and potential future directions.
- Present what I and some colleagues think is feasible way forward to make progress towards an end-to-end model of solar eruptions.

#### REVIEW article

Front. Astron. Space Sci., 21 December 2022  
Sec. Stellar and Solar Physics  
Volume 9 - 2022 |  
<https://doi.org/10.3389/fspas.2022.1060856>

This article is part of the Research Topic  
Flare Observations in the IRIS Era: What Have We Learned, and What's Next?  
[View all 8 Articles >](#)

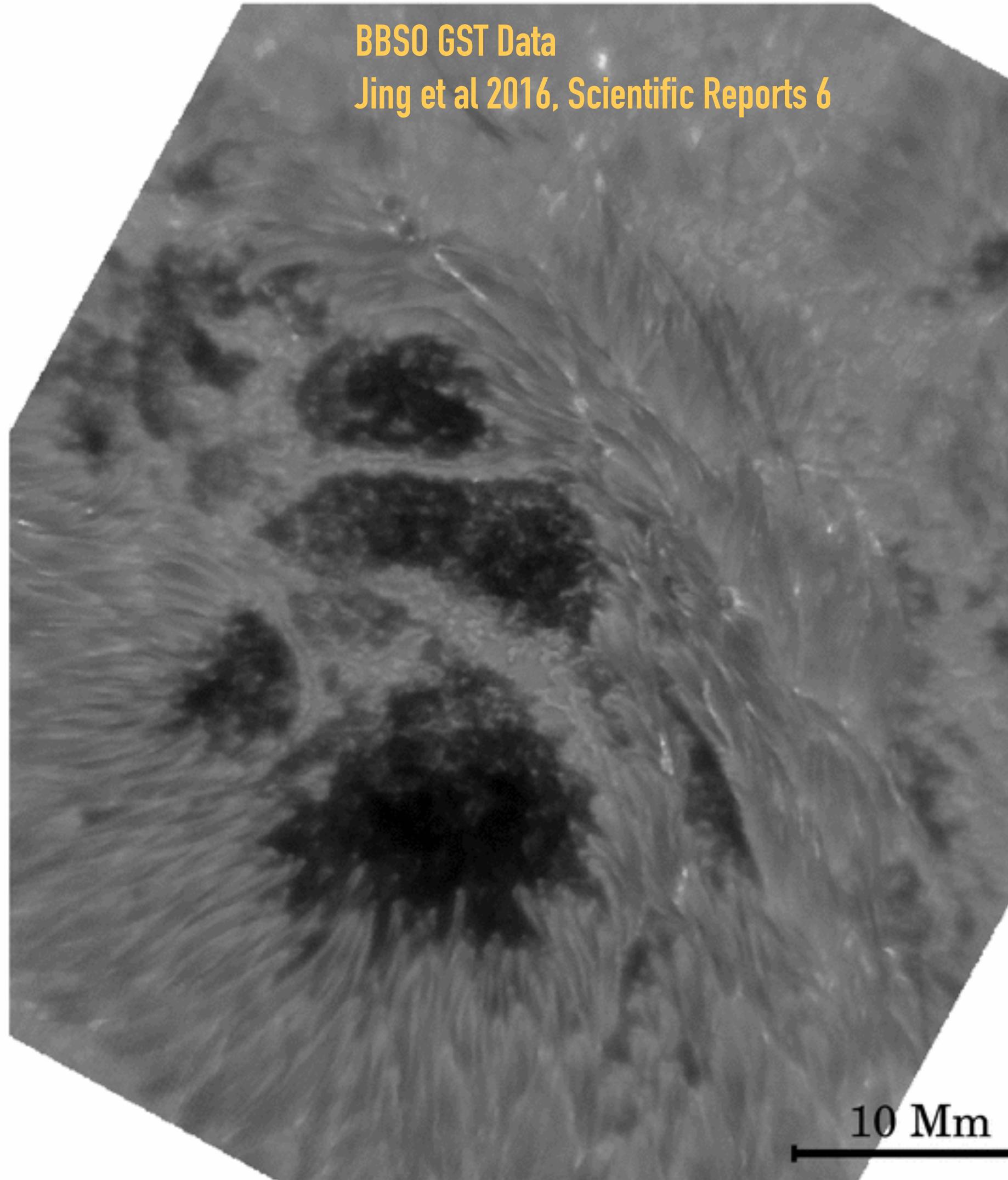
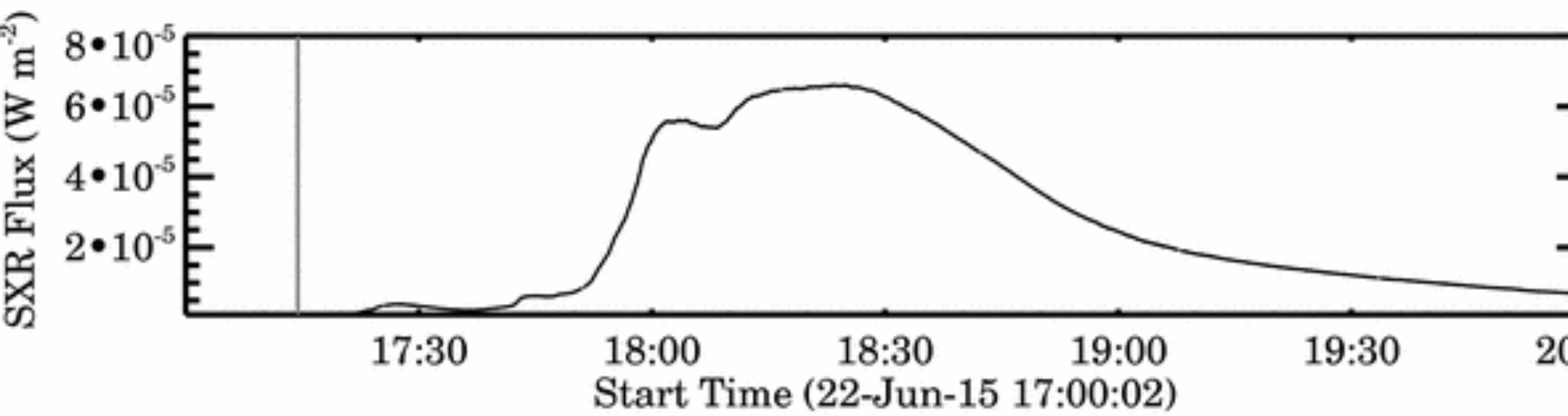
### Interrogating solar flare loop models with IRIS observations 1: Overview of the models, and mass flows

#### REVIEW article

Front. Astron. Space Sci., 06 January 2023  
Sec. Stellar and Solar Physics  
Volume 9 - 2022 |  
<https://doi.org/10.3389/fspas.2022.1060862>

This article is part of the Research Topic  
Flare Observations in the IRIS Era: What Have We Learned, and What's Next?  
[View all 8 Articles >](#)

### Interrogating solar flare loop models with IRIS observations 2: Plasma properties, energy transport, and future directions



- Highlight just a few recent(-ish) results.\*\*
- Discuss where models are currently not consistent with observations, and potential future directions.
- Present what I and some colleagues think is feasible way forward to make progress towards an end-to-end model of solar eruptions.

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## NEXT-GENERATION COMPREHENSIVE DATA-DRIVEN MODELS OF SOLAR ERUPTIVE EVENTS

A white paper in response to the Solar and Space Physics (Heliophysics) Decadal Survey

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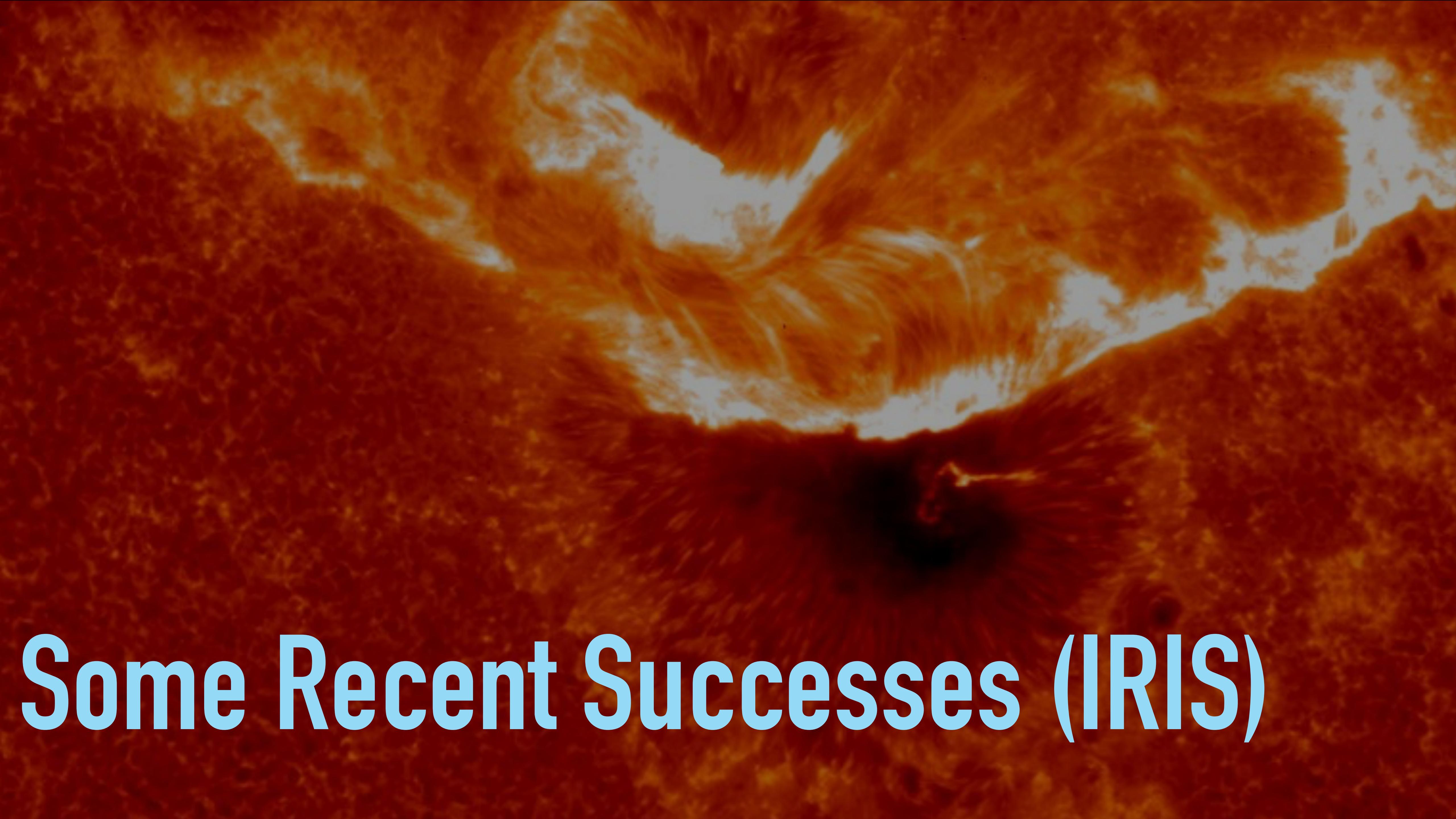
### Principal Author

Joel C. Allred<sup>1</sup> (joel.c.allred@nasa.gov)

### Co-authors

Graham S. Kerr<sup>1,2</sup>, Meriem Alaoui<sup>1,3</sup>, Juan Camilo Buitrago-Casas<sup>4</sup>, Amir Caspi<sup>5</sup>, Bin Chen<sup>6</sup>, Thomas Y. Chen<sup>7</sup>, Lindsay Glesener<sup>8</sup>, Silvina E. Guidoni<sup>9</sup>, Fan Guo<sup>10</sup>, Judith T. Karpen<sup>1</sup>, Sophie Musset<sup>11</sup>, Katharine K. Reeves<sup>12</sup>, Albert Y. Shih<sup>1</sup>

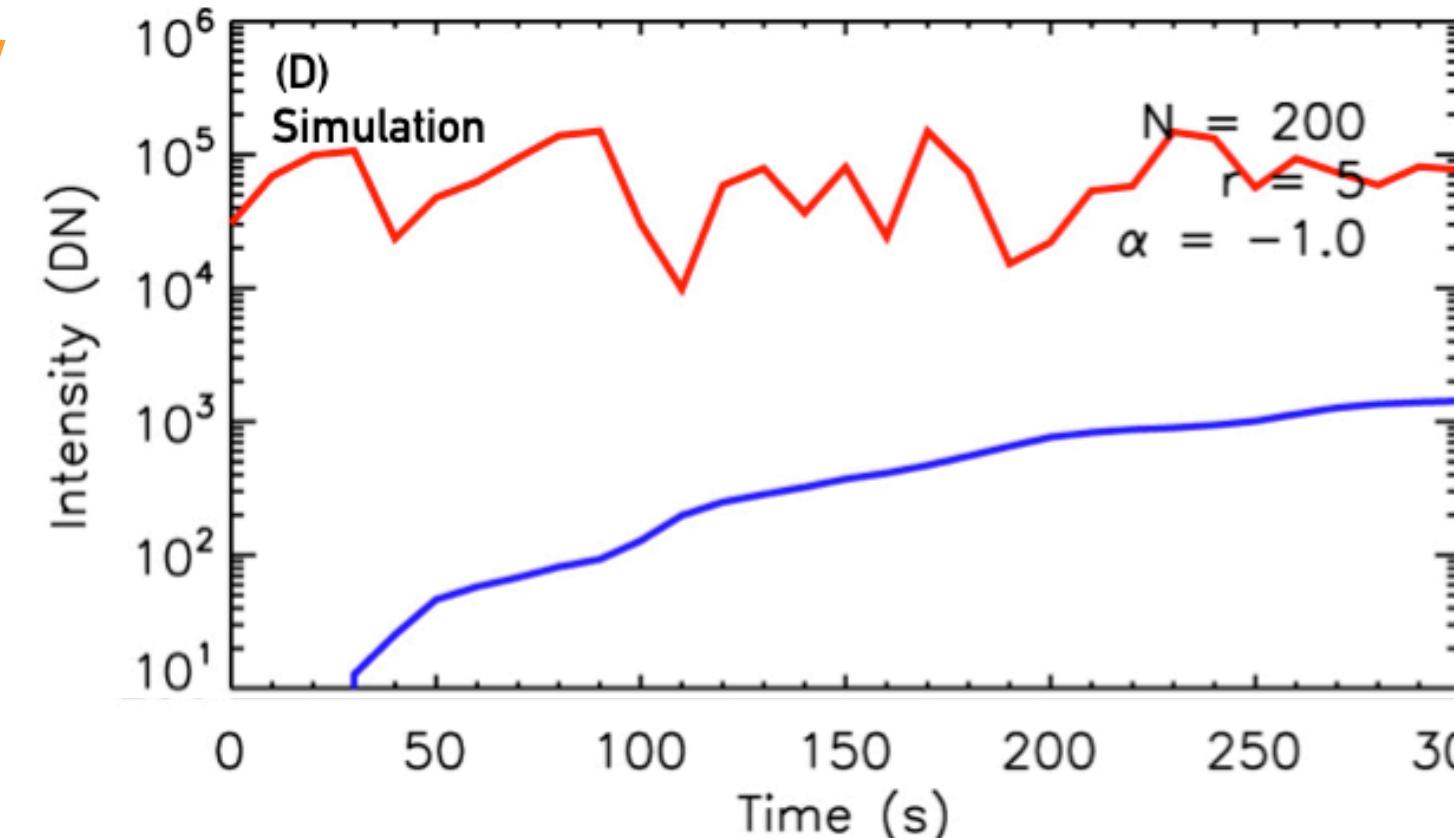
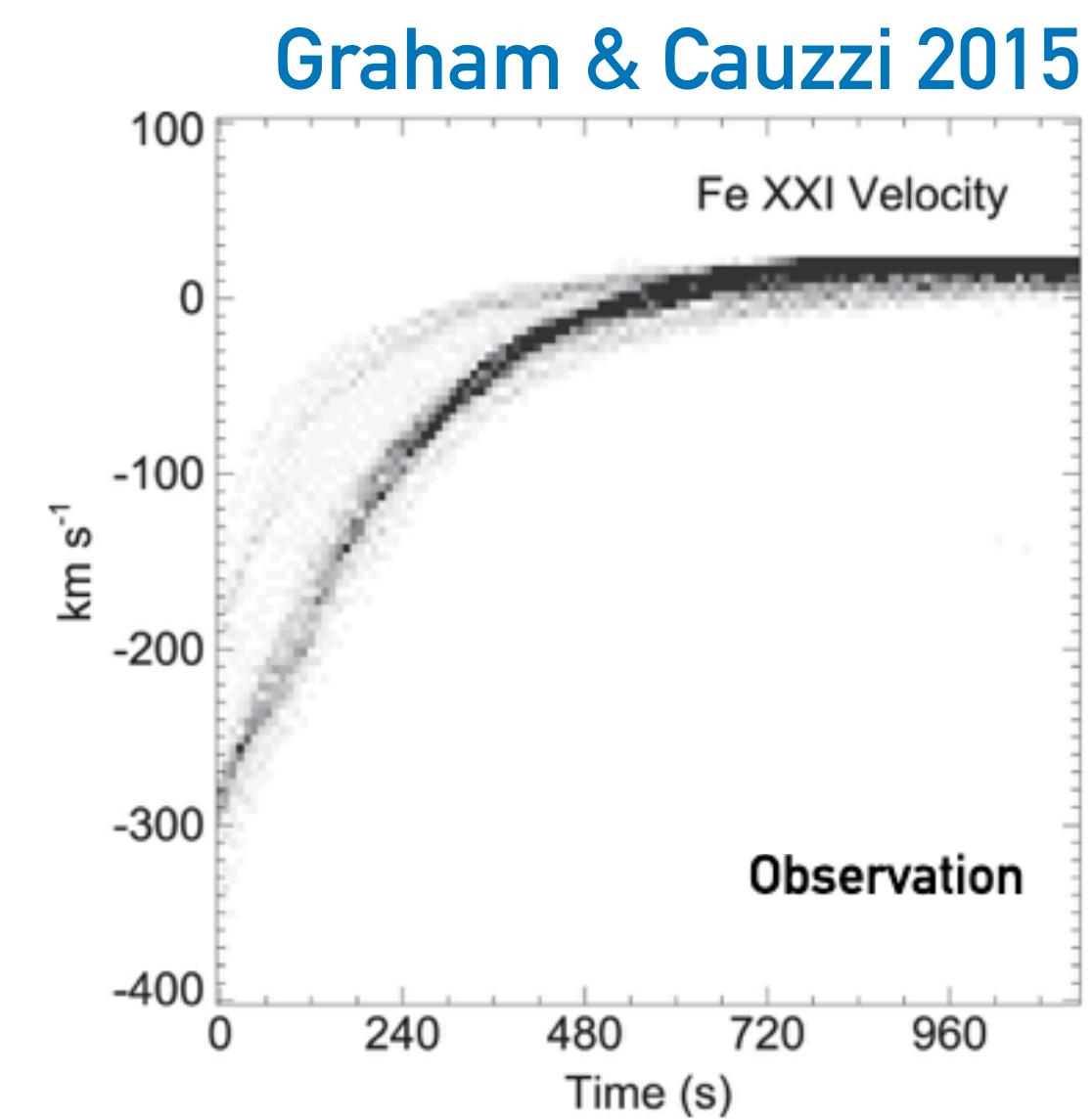
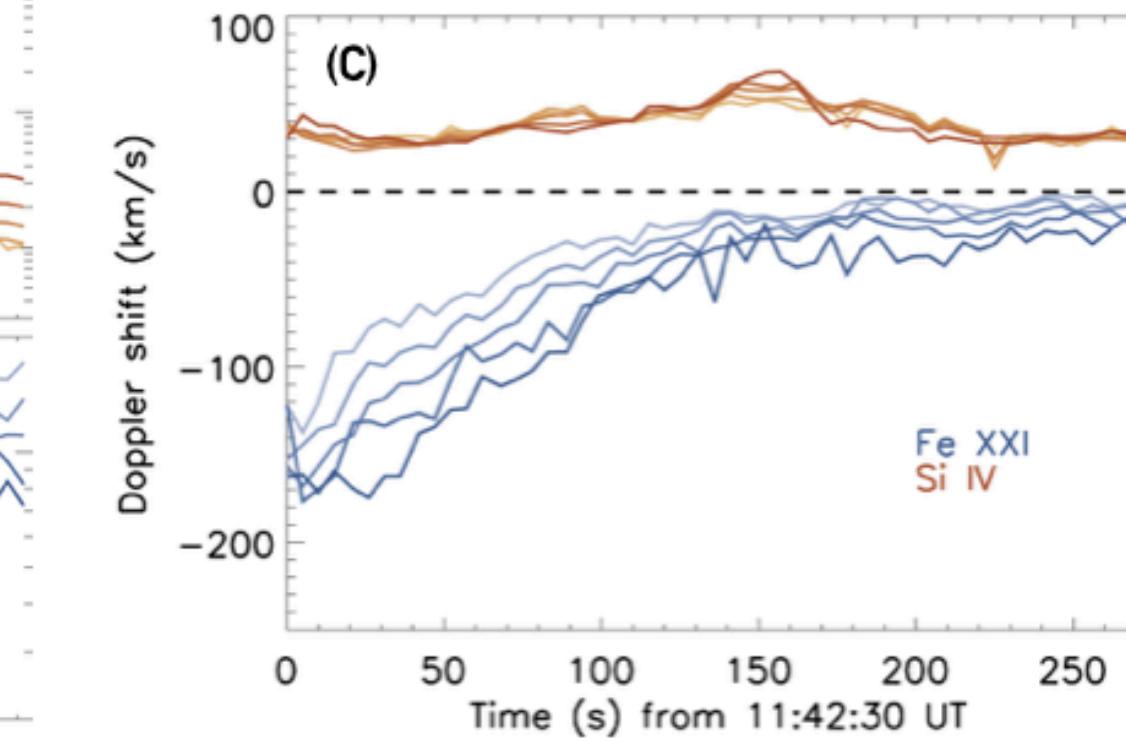
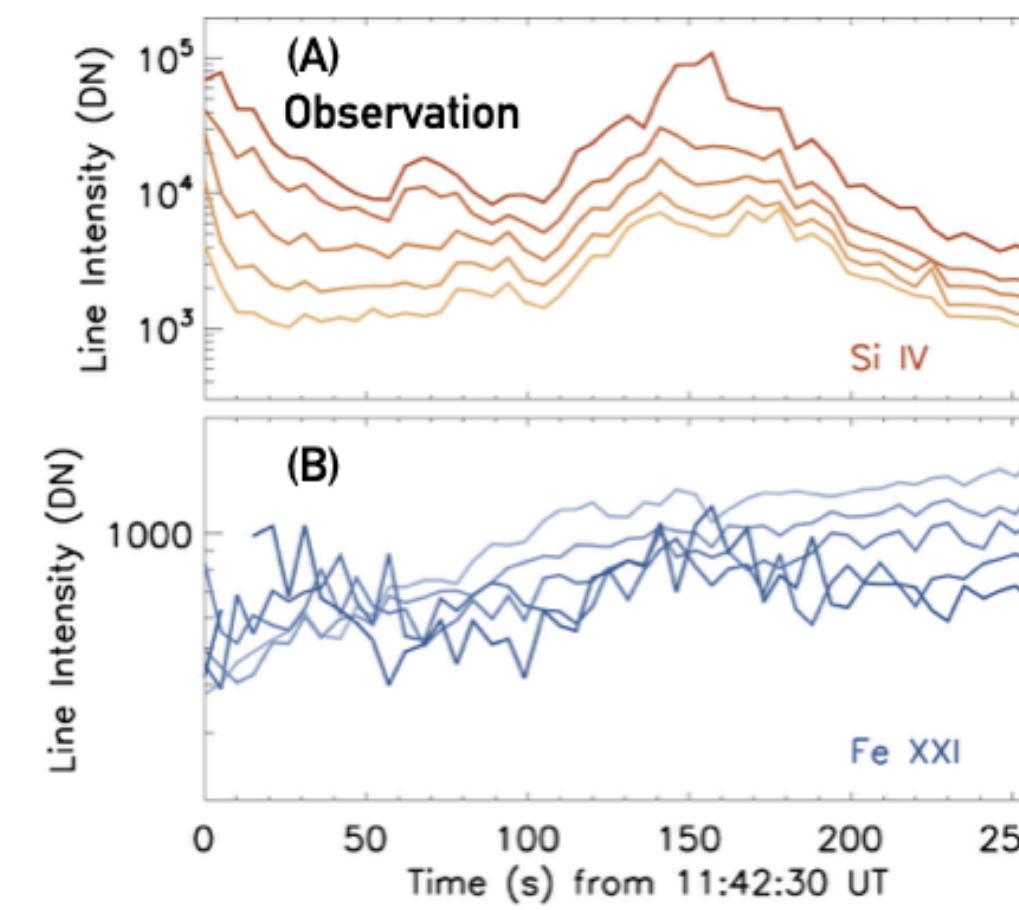
Code	Key Features	Geometry	Heating	Chromosphere	References
<b>Flarix (RHD)</b>	Accurate chromosphere and radiative transfer; test particle code to calculate heating (incl return currents, non-thermal collisions).	Half loop, fixed but optimised grid.	Electron beam (test particle code); Conduction	Full NLTE for H, Ca II (and sometimes Mg II)	Kasparova et al 2009; Varady et al 2010; Heinzel et al. 2016
<b>RADYN (RHD)</b>	Accurate chromosphere and radiative transfer; well developed flare physics (incl return currents, non-thermal collisions)	Half loop, adaptive grid with fixed # points to resolve where needed.	Particle beam (Fokker-Planck, both electron and ions); Alfvénic waves (approximated); Conduction (incl. suppression)	Full NLTE, non-equilibrium, for H, He I, He II, Ca II (and sometimes Mg II)	Carlsson & Stein 1995, 1997, 2002; Allred et al. 2005, 2015, 2020, 2022; Kerr et al 2016
<b>HYDRAD (HD)</b>	Non-equilibrium ionization for minor species; flexible geometry; two fluid plasma	Full loop, full AMR, magnetic field expansion.	Electron Beam (Hawley & Fisher 1994); Alfvénic waves; Conduction	Approximation to NLTE for H (Sollum); Carlsson & Leenaarts radiative losses	Bradshaw & Cargill 2013; Reep et al. 2019



Some Recent Successes (IRIS)

# MASS FLOWS IN SOLAR FLARES: LONG DURATION UPFLOWS

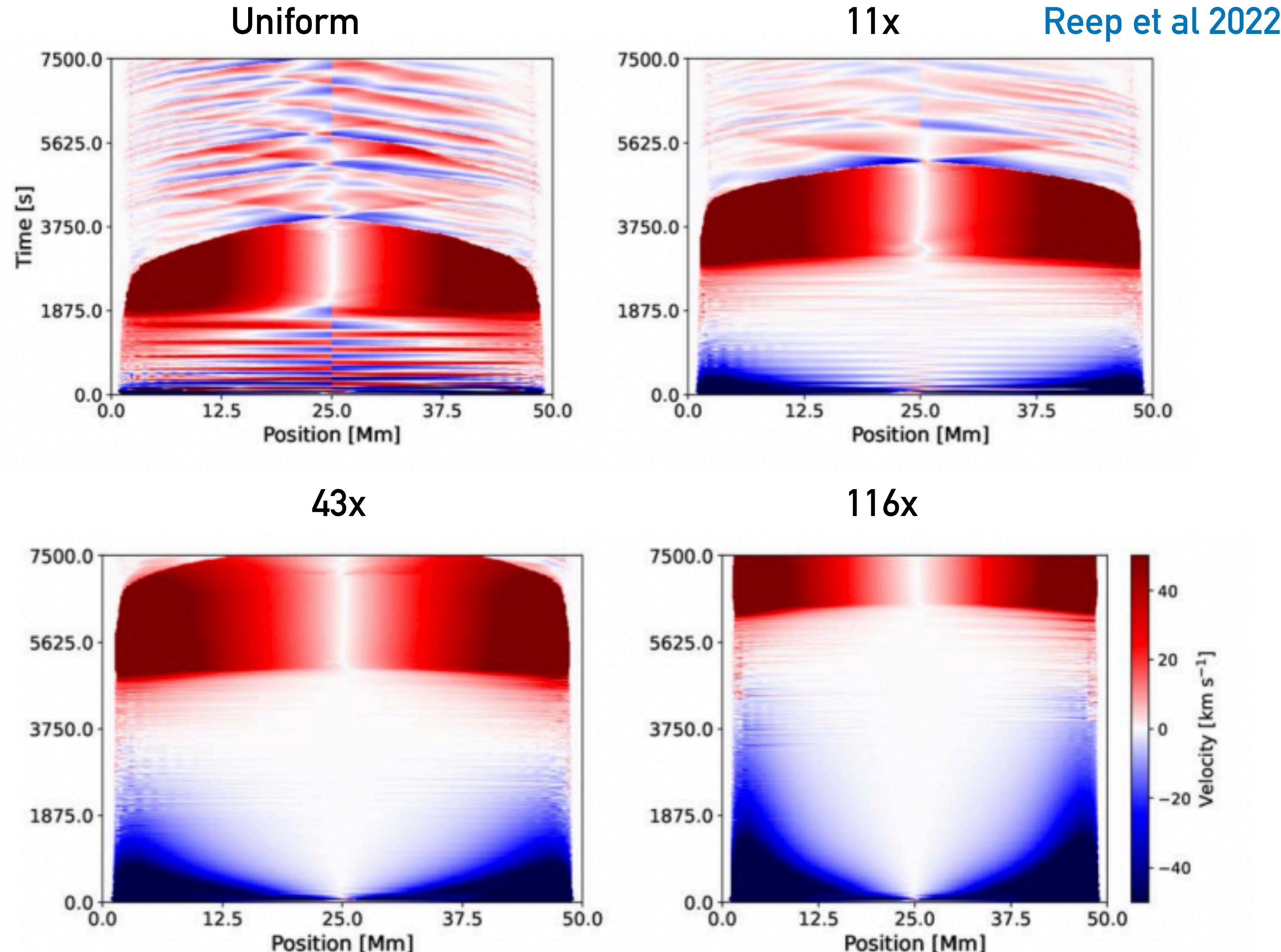
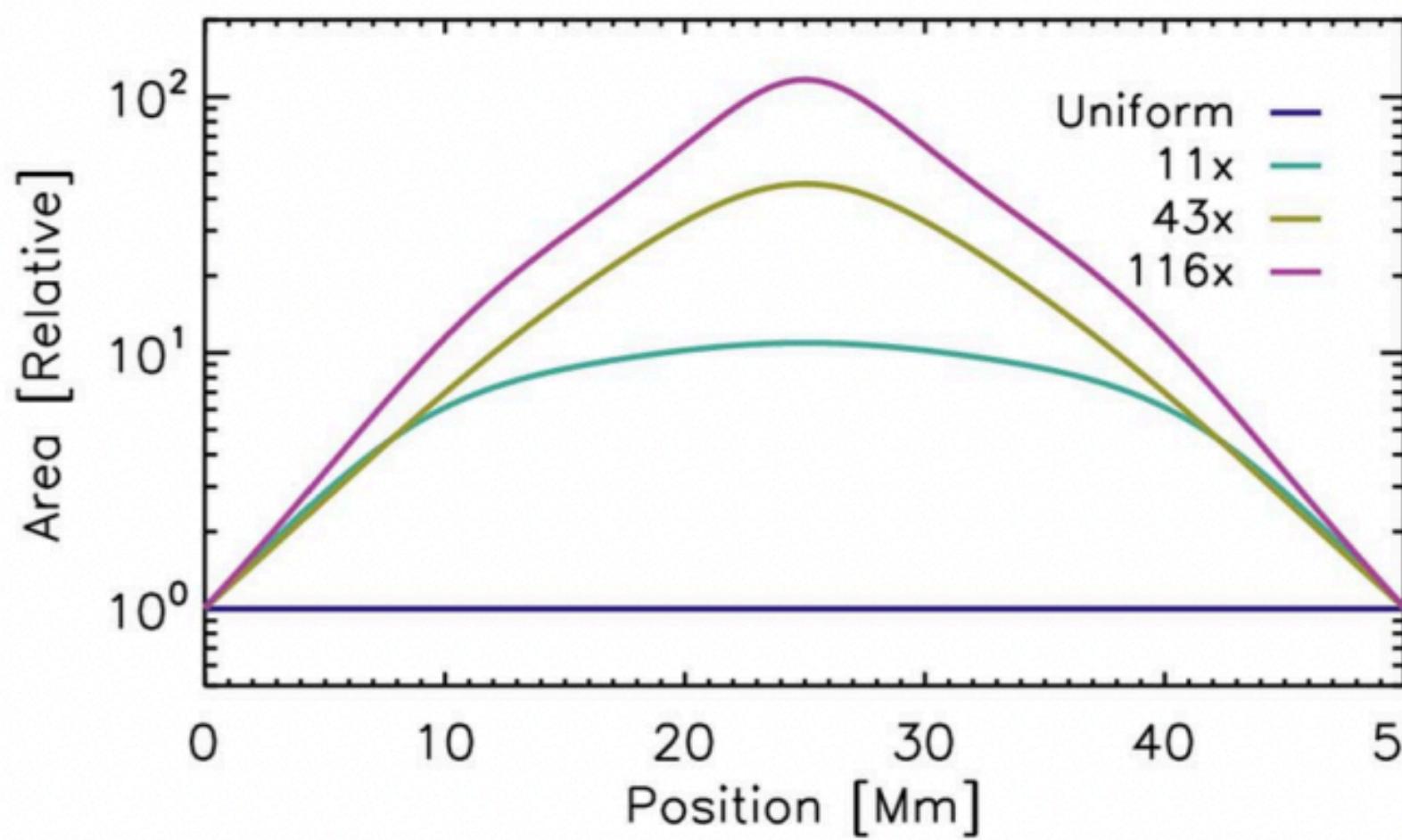
- ▶ Chromospheric evaporation (upflows) persist for 100s, **but** loop models predict the cessation of flows soon after we stop injecting energy.
- ▶ Multi-threaded modelling has proved successful – ~100 loops per IRIS pixel ( $<0.01''$ ), variable injection onto each loop for up to tens of seconds (or longer).
- ▶ However, total duration into each pixel still needs to be several minutes.



Reep et al 2016, 2018, 2019.

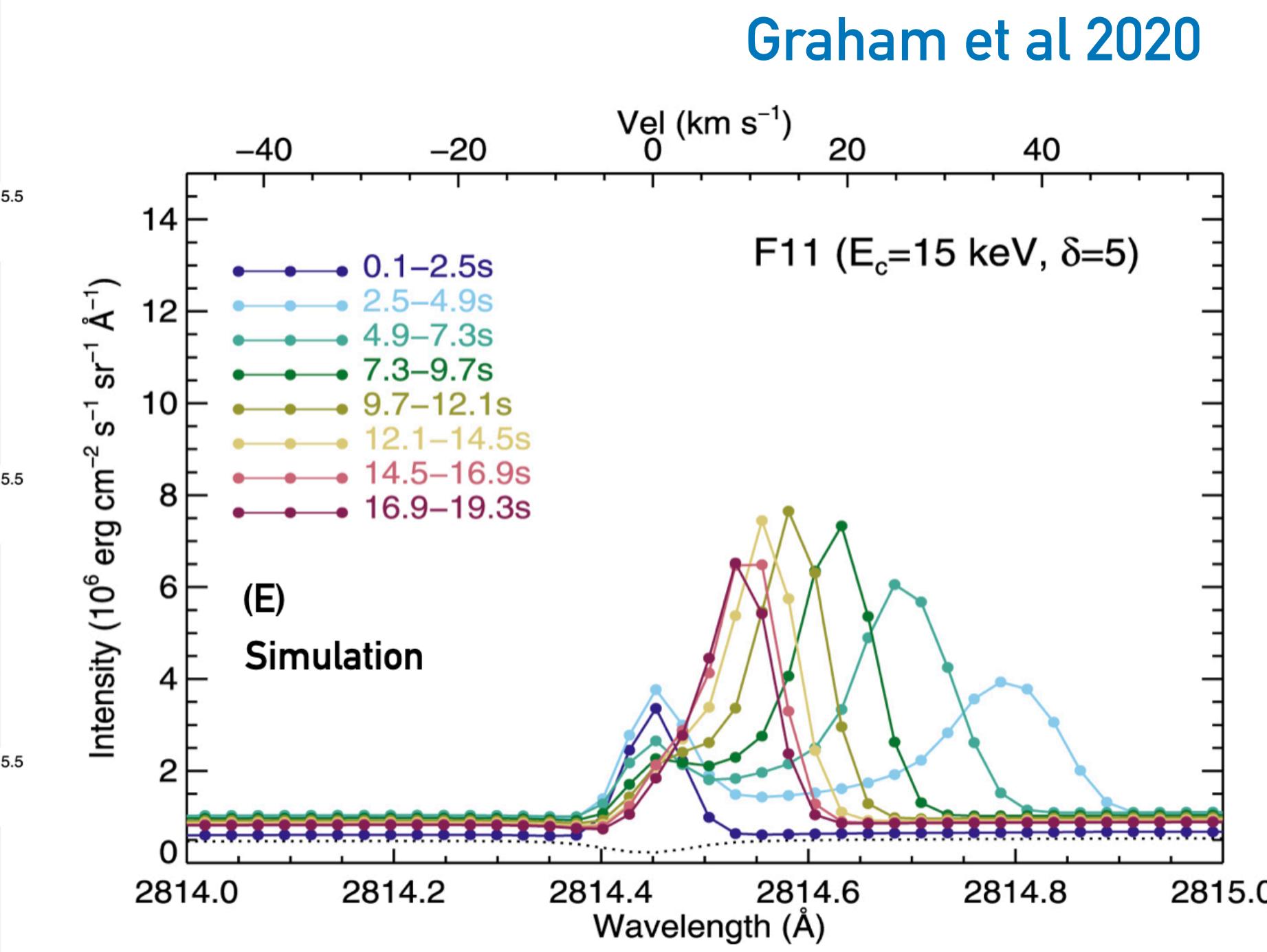
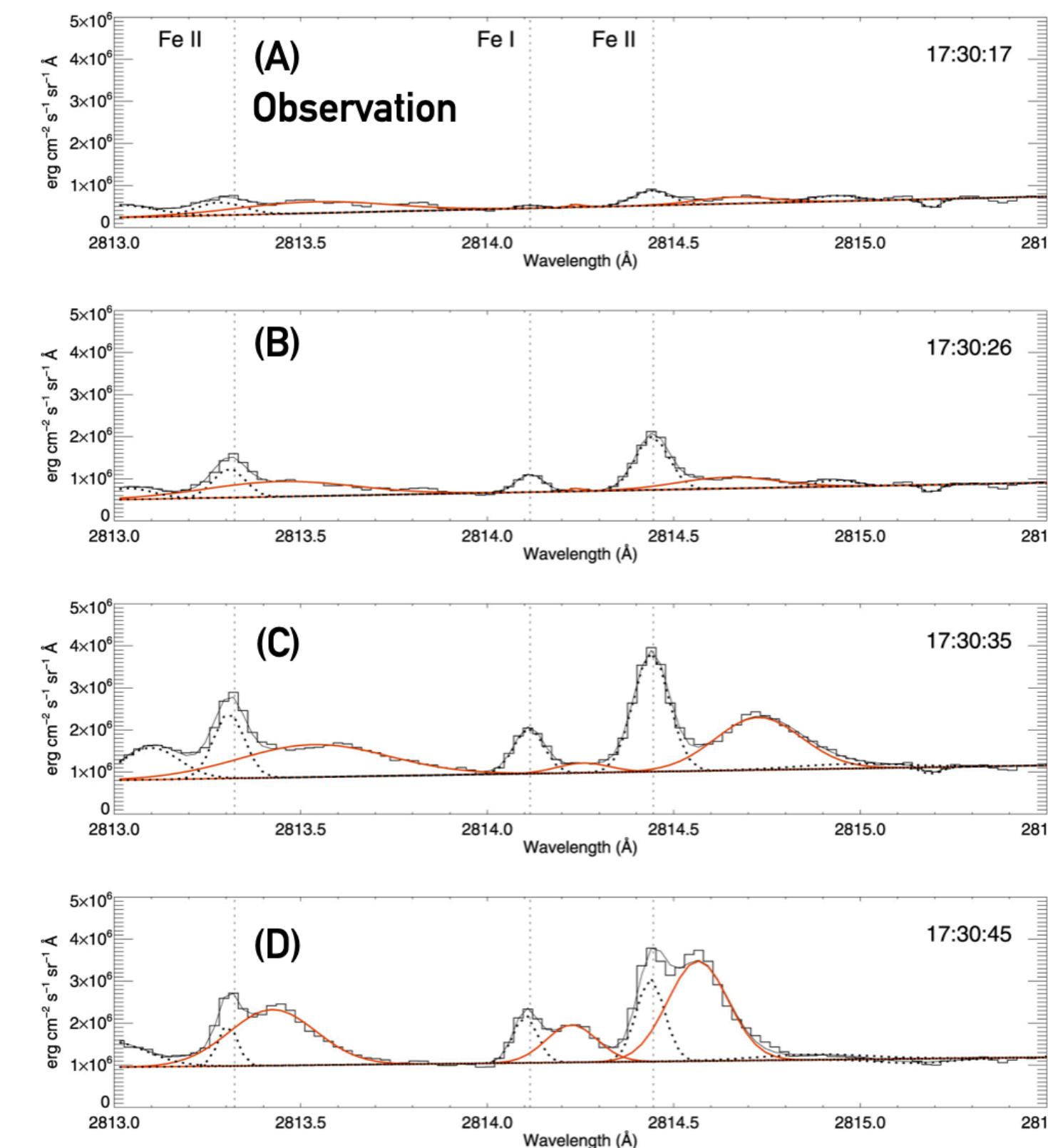
# MASS FLOWS IN SOLAR FLARES: LONG DURATION UPFLOWS

- ▶ Chromospheric evaporation (upflows) persist for 100s, **but** loop models predict the cessation of flows soon after we stop injecting energy.
- ▶ One other potential solution is to include area expansion along the loop.



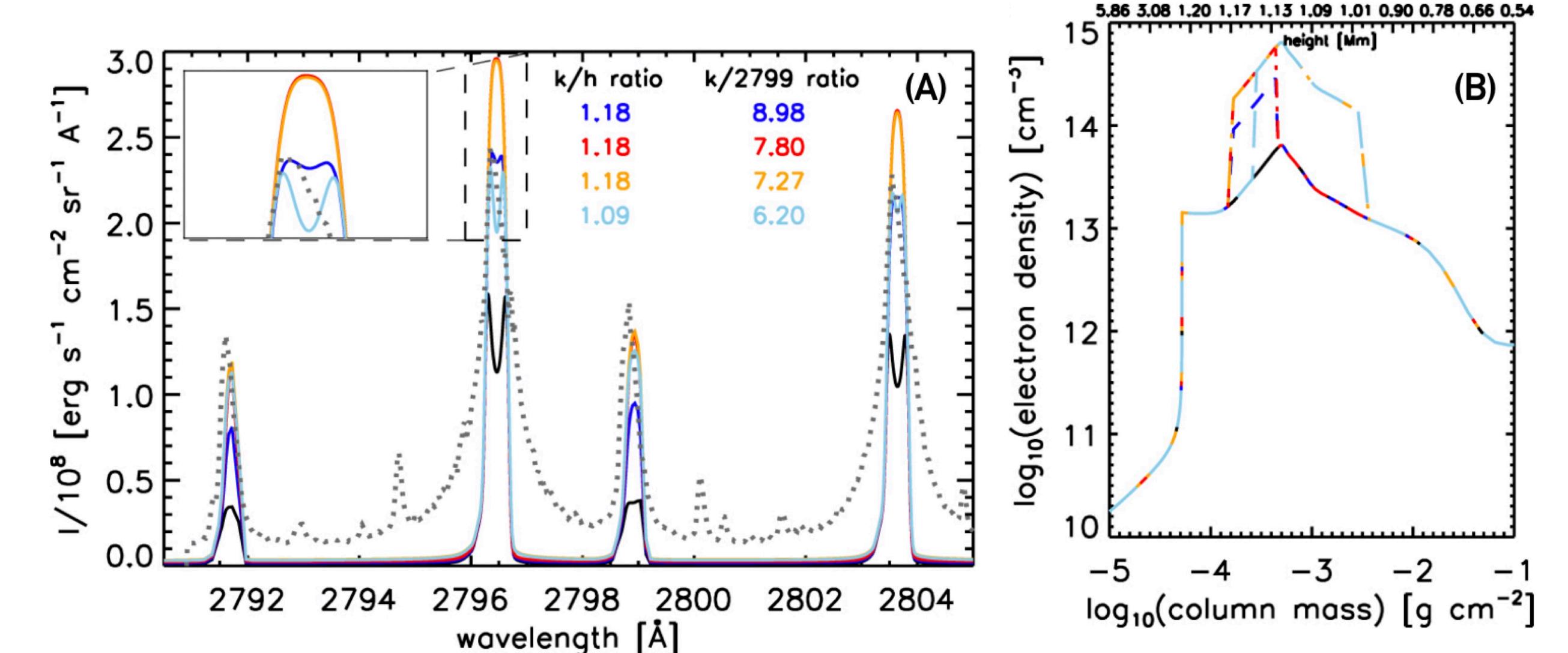
# MASS FLOWS IN SOLAR FLARES: RED WING ASYMMETRIES

- ▶ Curiously, however, chromospheric condensations can be reasonably well produced by models and do not seem to require repeated energy injection into the loop.
- ▶ A RADYN model of an IRIS flare reproduced the red-wing component behavior.
- ▶ Only a slightly shorter timescale, and it is likely their condensation was too dense.

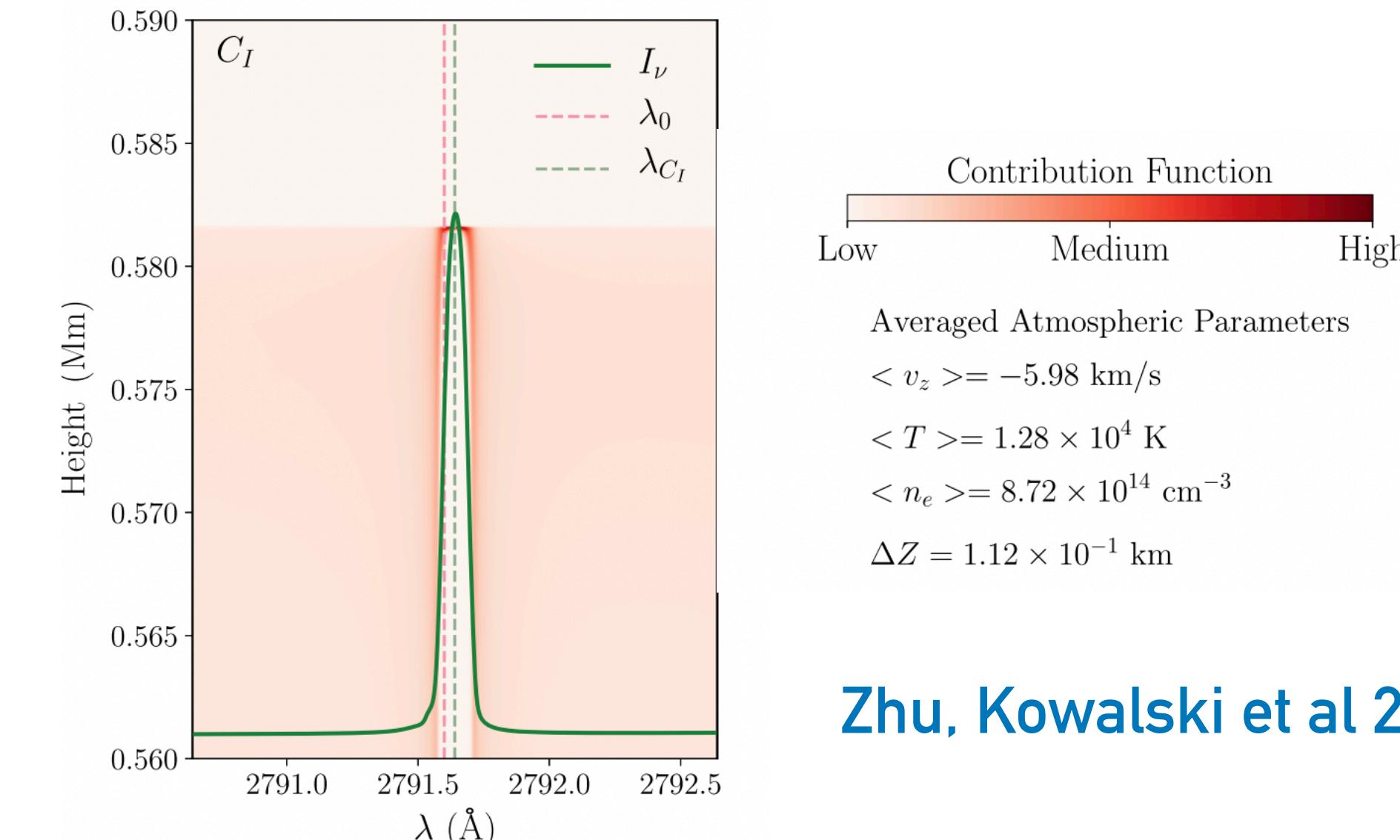


# SINGLE PEAKED Mg II PROFILES — INDICATES HIGH CHROMOSPHERIC DENSITY

- ▶ Mg II flare profiles are (almost) always single peaked, but most models have some level of central reversal.
- ▶ Parametric study by Rubio da Costa & Kleint (2017) indicated that a **very large electron density was required in the upper chromosphere** to achieve this:  
 $n_e > \sim 5 \times 10^{14} \text{ cm}^{-3}$
- ▶ Zhu et al 2019, using a high injected energy flux, produced this 'naturally' in a flare simulation, which **required a very large injected energy flux**  
 $F = 5 \times 10^{11} \text{ erg s}^{-1} \text{ cm}^{-2}$
- ▶ Mg II core formed over a vanishingly narrow region in the upper chromosphere, only a few tens of meters thick!



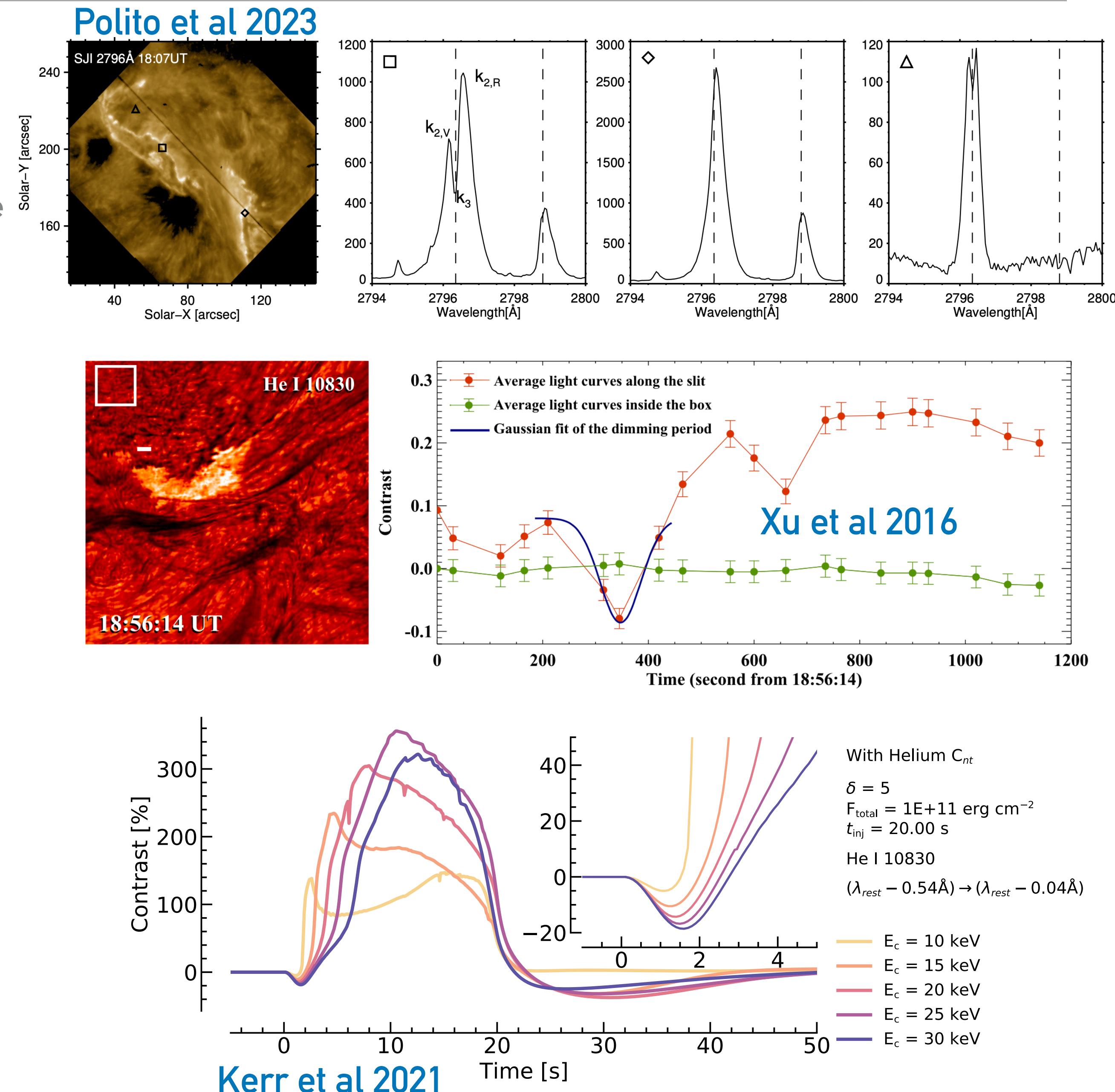
Rubio da Costa &amp; Kleint 2017



Zhu, Kowalski et al 2019

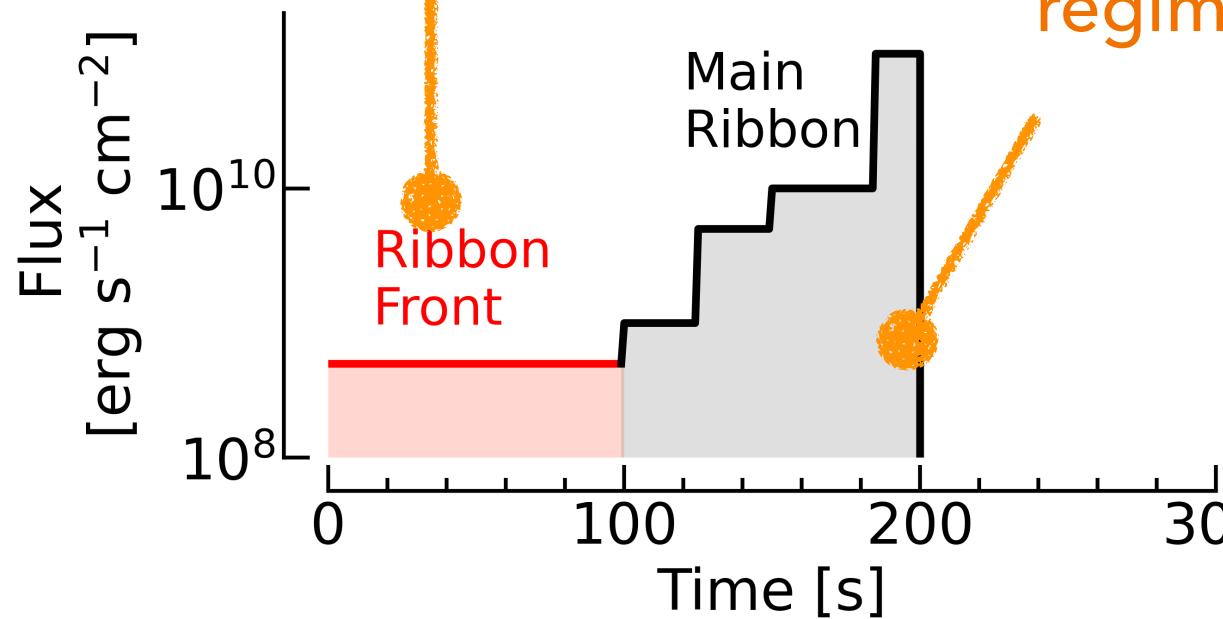
# RIBBON FRONTS

- ▶ At the leading edge of flare ribbons we have seen:
- He I 10830 dimmed for up to several hundred seconds before brightening (Xu et al 2016);
- Mg II exhibited deep central reversals, were slightly blueshifted, and extremely broad (Xu et al 2016, Panos et al 2018, 2021ab).
- That NUV/FUV intensities took several tens of seconds and up to 2 minutes to reach peak after activation (Naus et al 2022).
- ▶ RADYN modelling (Kerr et al 2021, Polito et al 2022, Kerr et al 2023 *in prep*) has gone some way to explaining these features: **ribbon fronts require a very weak flux of non-thermal electrons compared to the main ribbons that produce the brighter emission. Could suggest up to 2-3 minutes of 'pre-heating' by a weak flux of electrons in some locations.**
- ▶ Ribbon fronts have  $T < 20\text{kK}$ ,  $n_e < 5 \times 10^{12} \text{ cm}^{-3}$ , and are undergoing gentle evaporation.



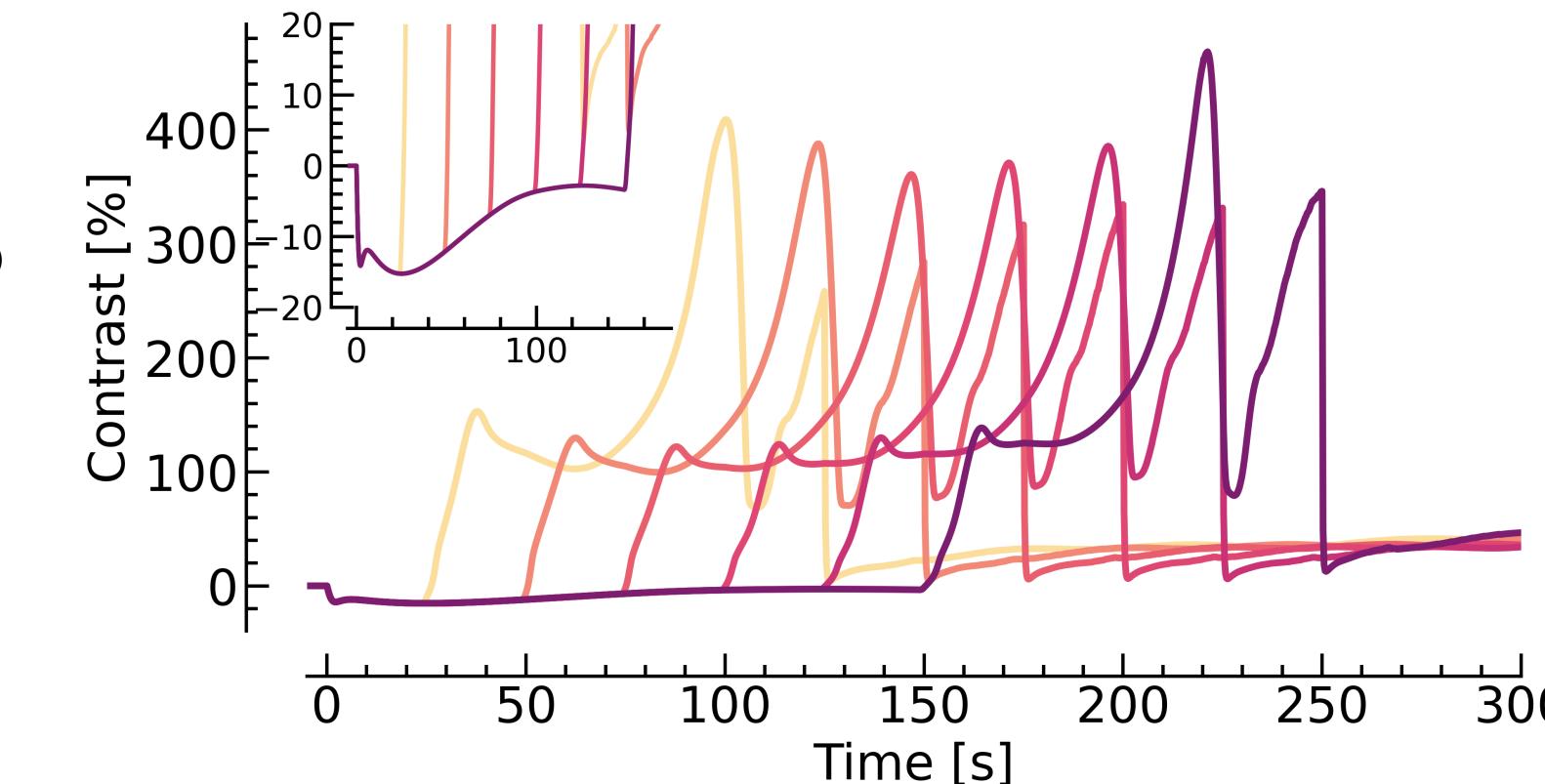
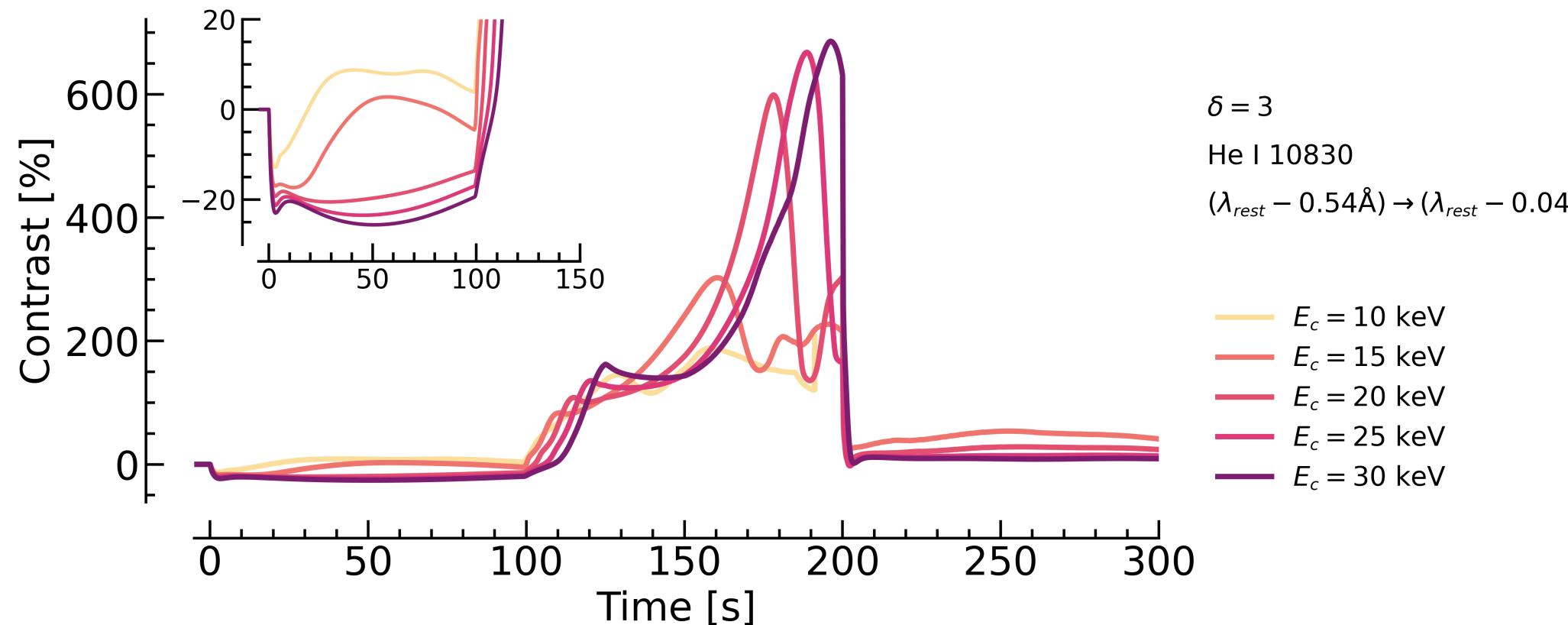
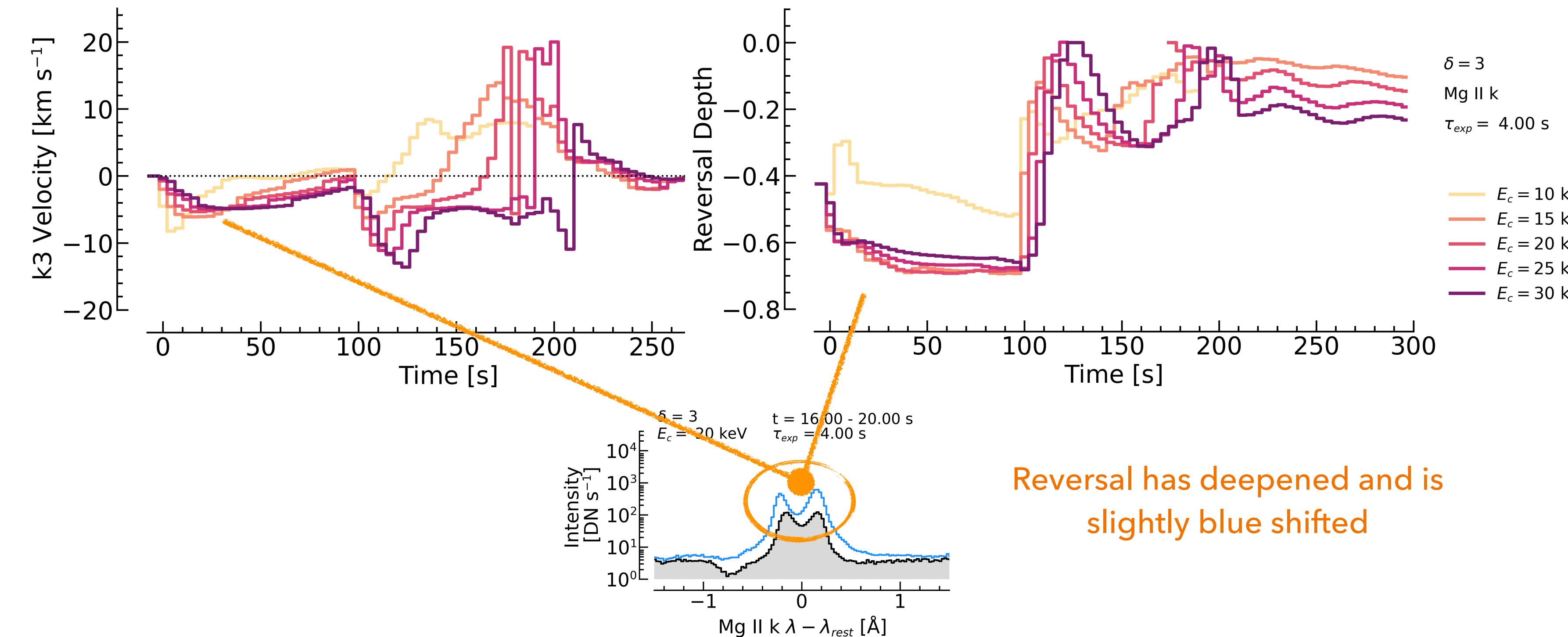
# RIBBON FRONTS

Time based  
on ribbon  
front  
lifetime



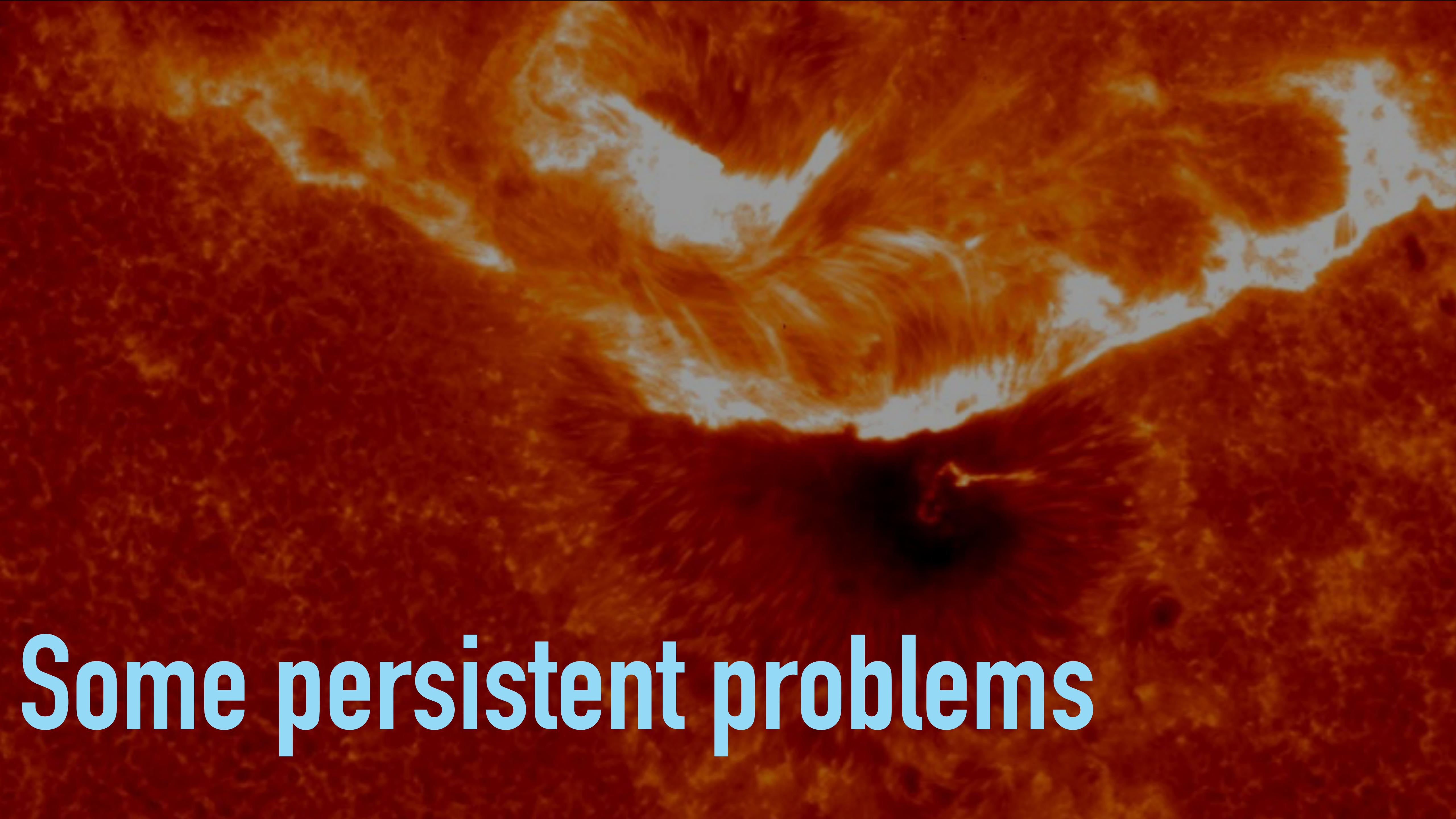
Ad-hoc to  
experiment  
with  
different  
regimes

Kerr et al 2023 (in prep)



Reversal has deepened and is slightly blue shifted

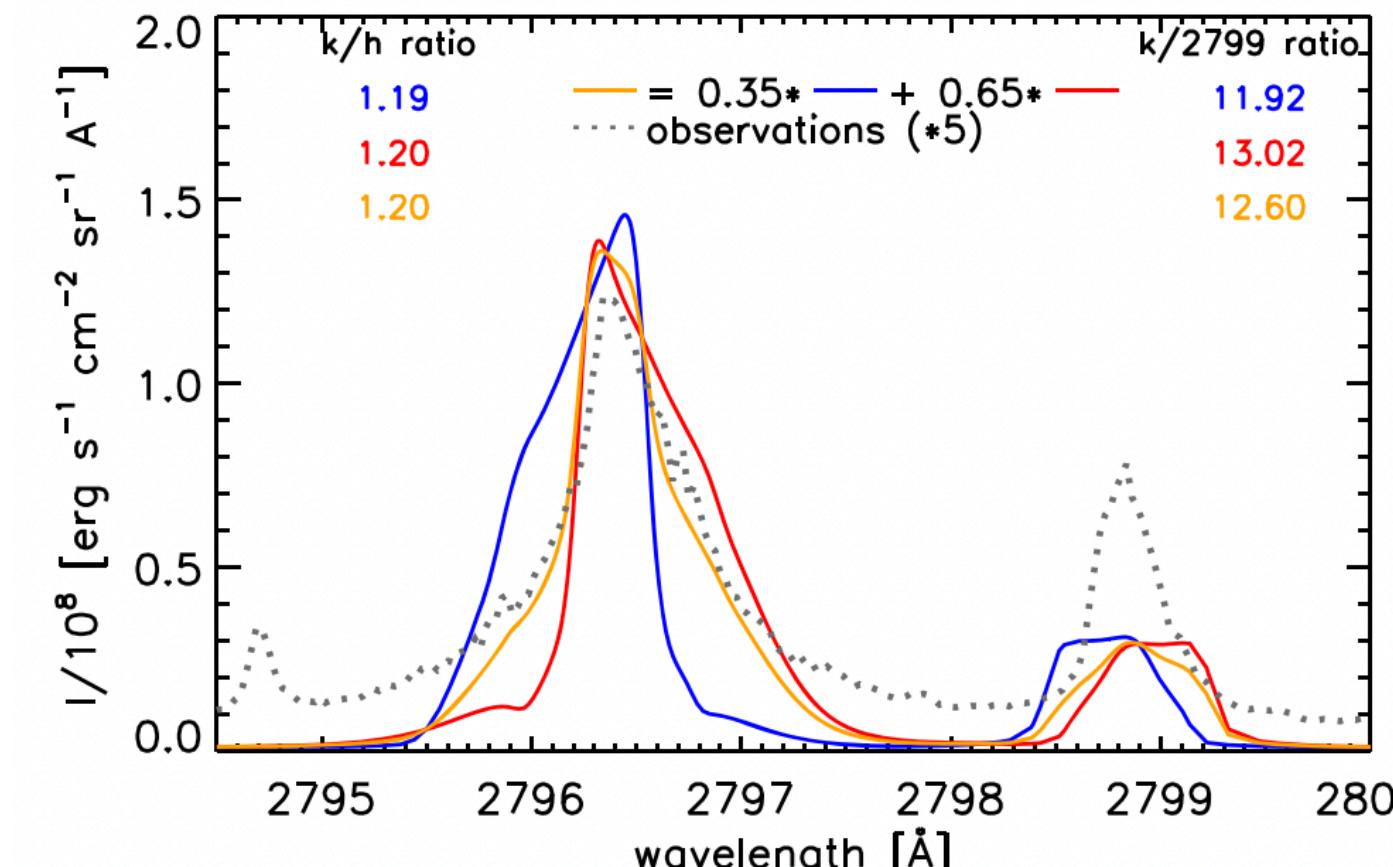
He I 10830 $\text{\AA}$  remains dimmer for as long as we inject a weak flux of electrons

The background of the slide features a dramatic sunset or sunrise over a body of water. The sky is filled with large, billowing clouds in shades of orange, red, and white. The horizon line is visible, showing the calm water meeting the sky. The overall atmosphere is warm and serene.

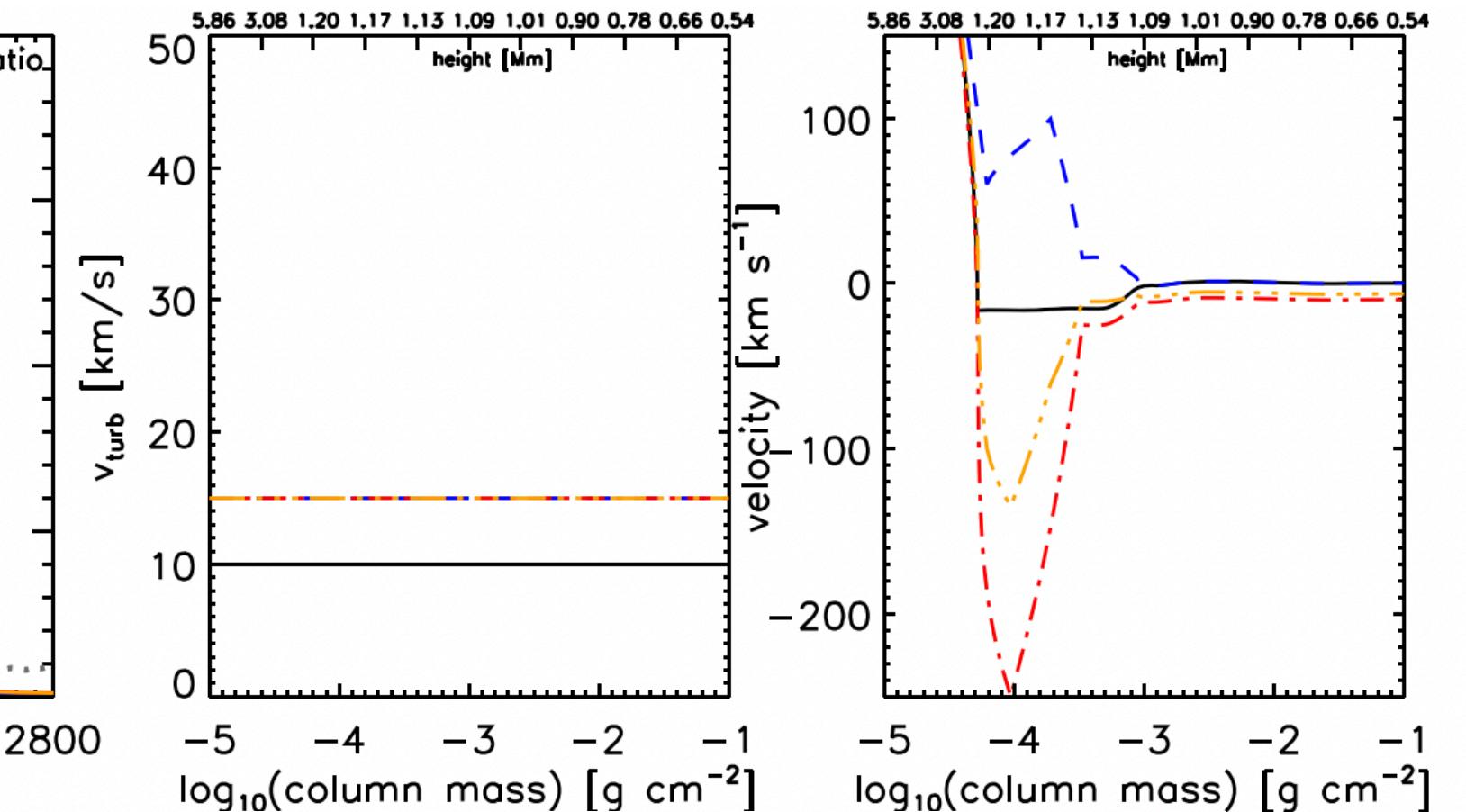
# Some persistent problems

# LINE WIDTHS! - CHROMOSPHERIC LINES (Mg II)

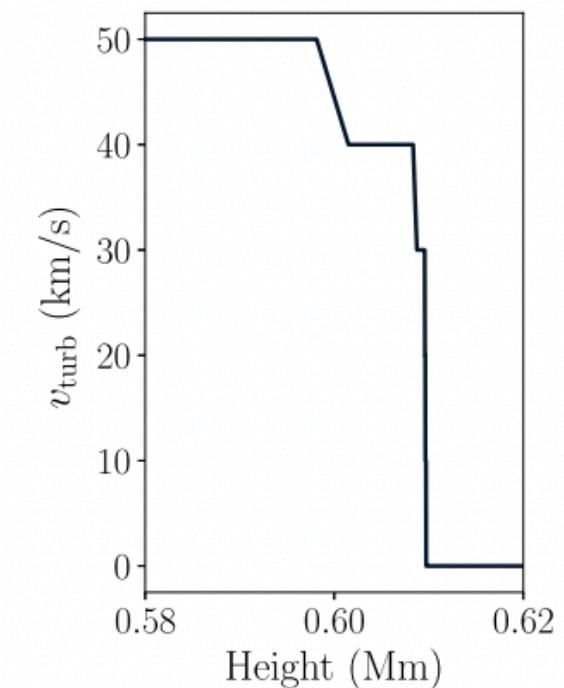
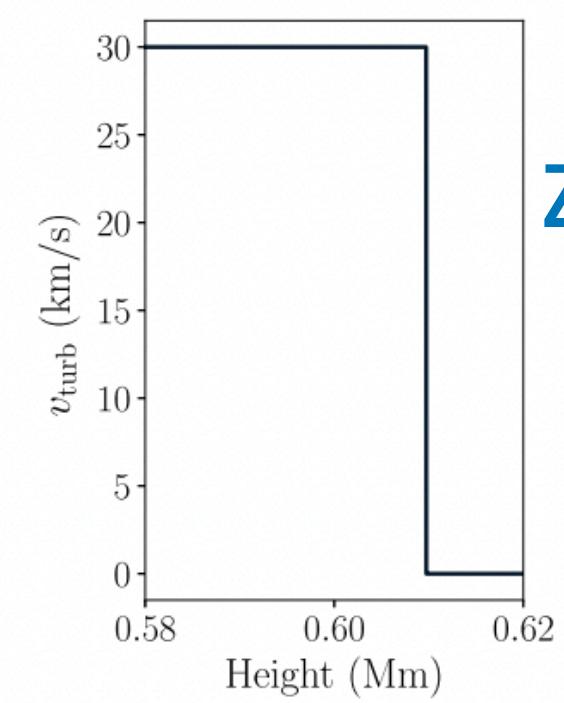
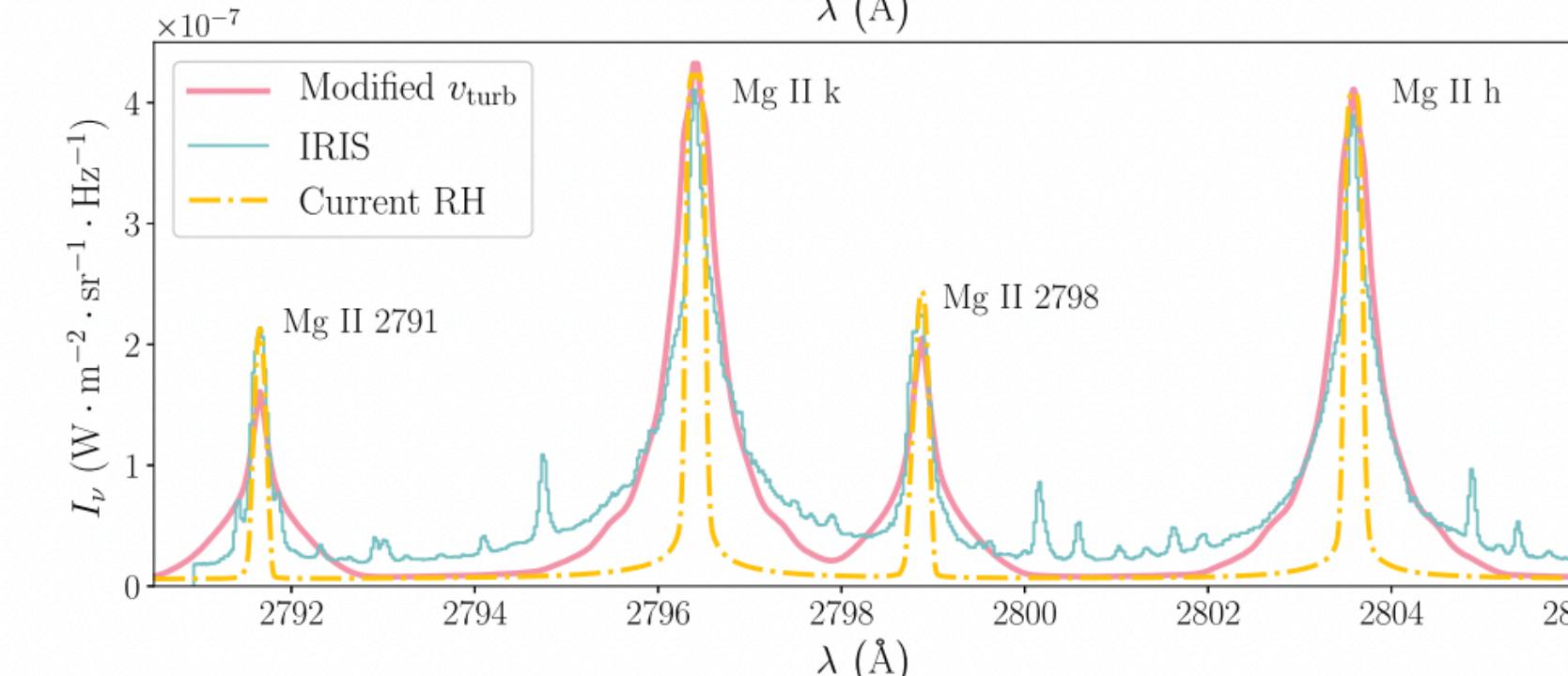
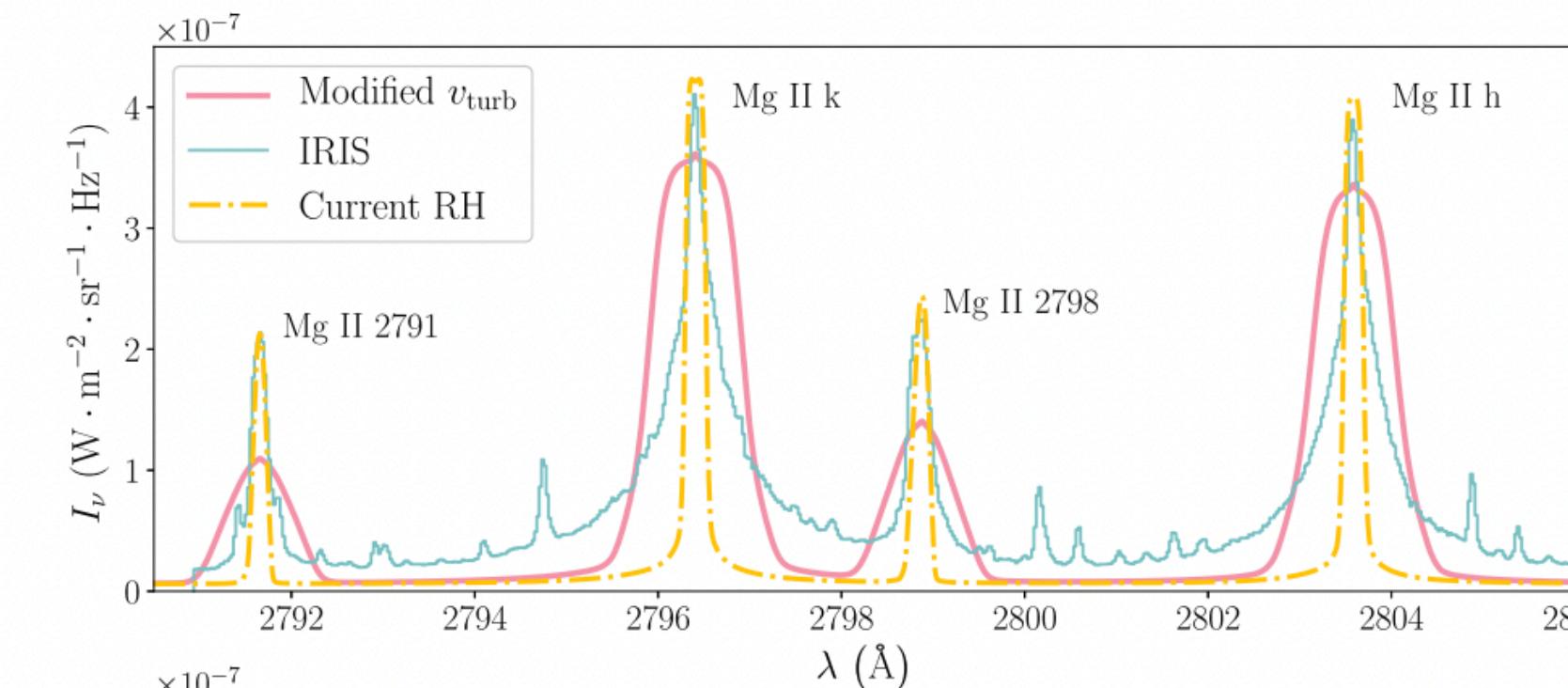
- ▶ Chromospheric and transition region lines are much broader than observed.
- ▶ Suggestions include:
  - Extreme bi-directional flows (any evidence?)
  - Large amounts of micro-turbulent broadening (up to 30-50 km/s). However, observations suggest this is unlikely.



Rubio da Costa & Kleint 2017



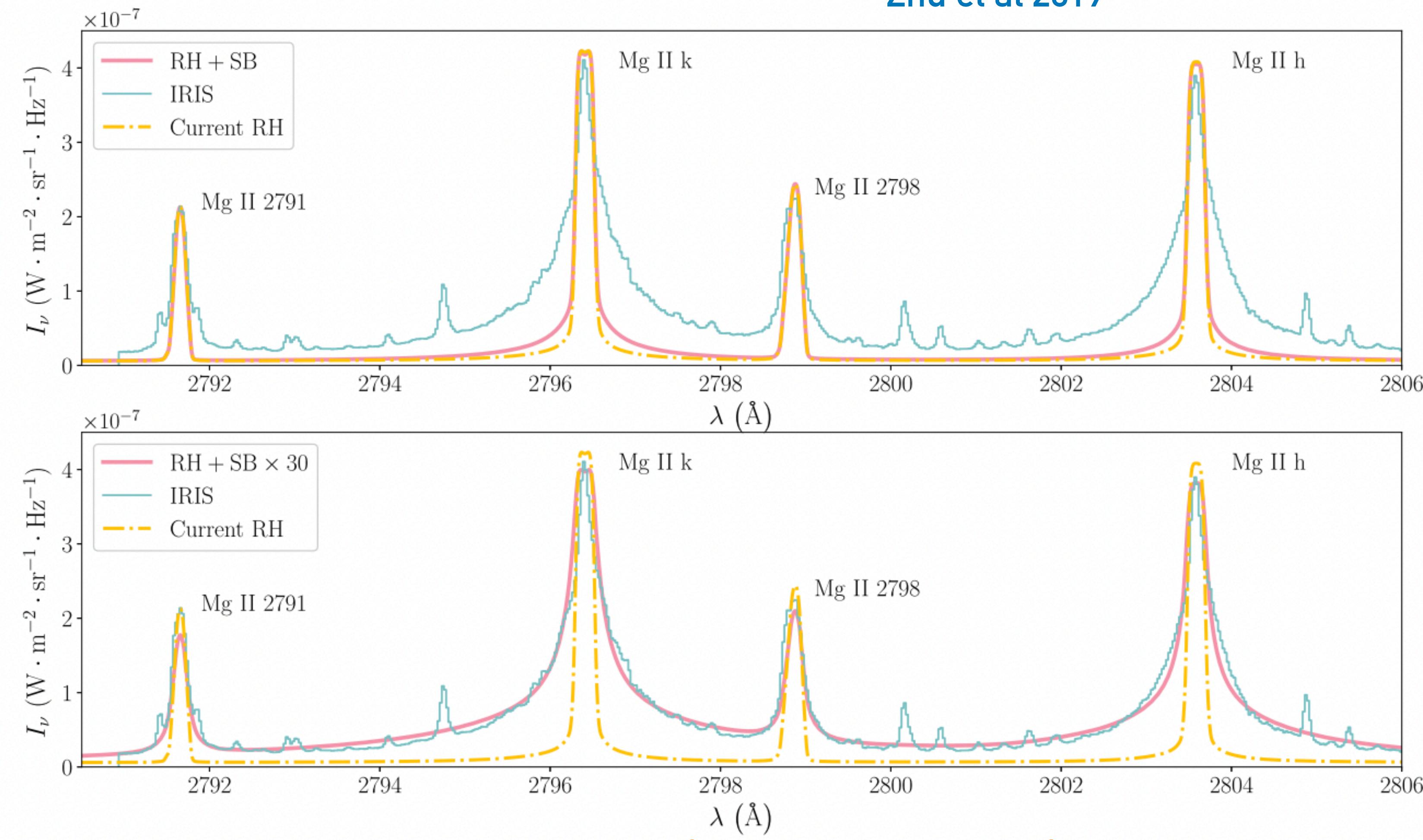
Zhu et al 2019



# LINE WIDTHS! - CHROMOSPHERIC LINES (Mg II)

Zhu et al 2019

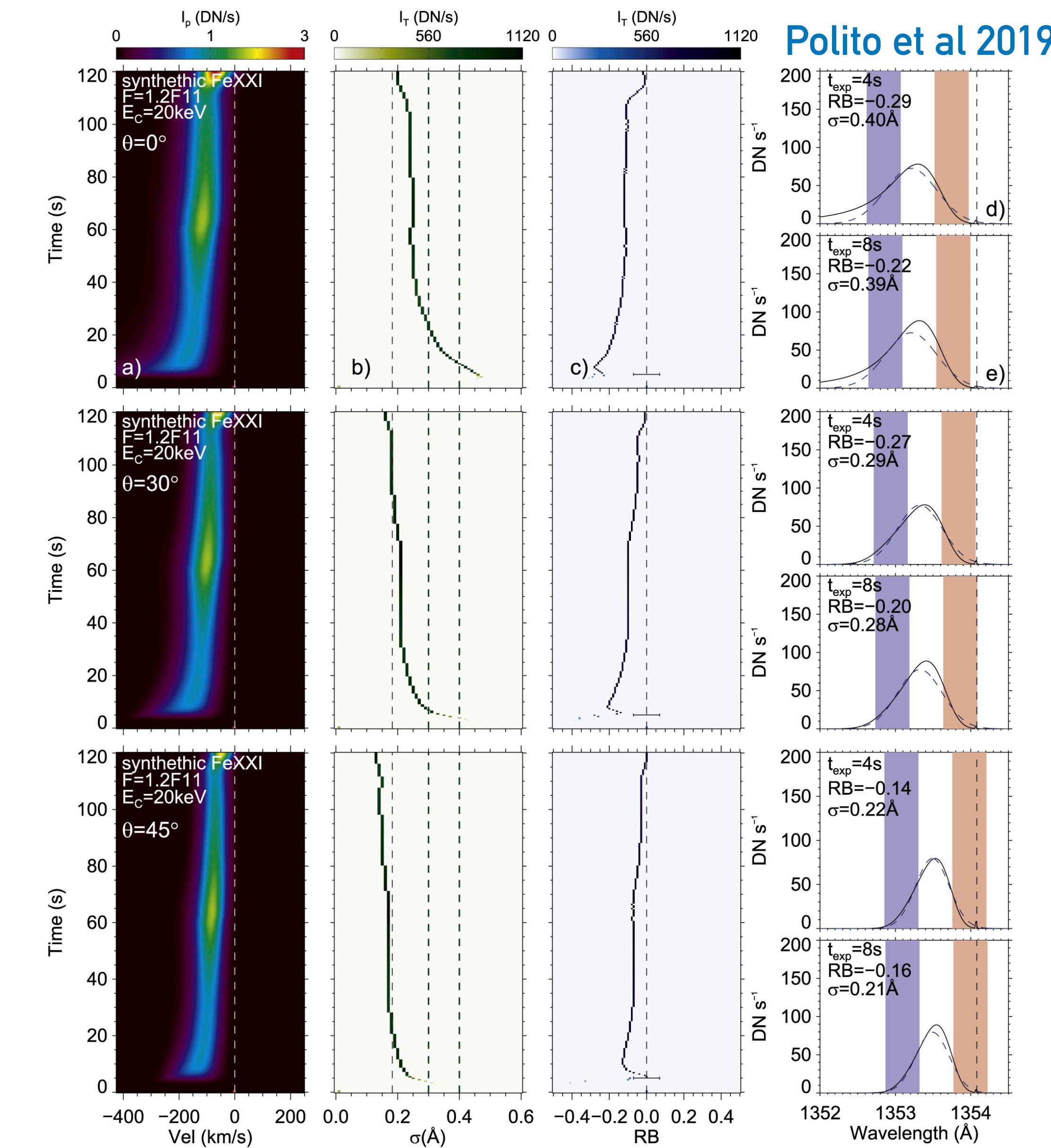
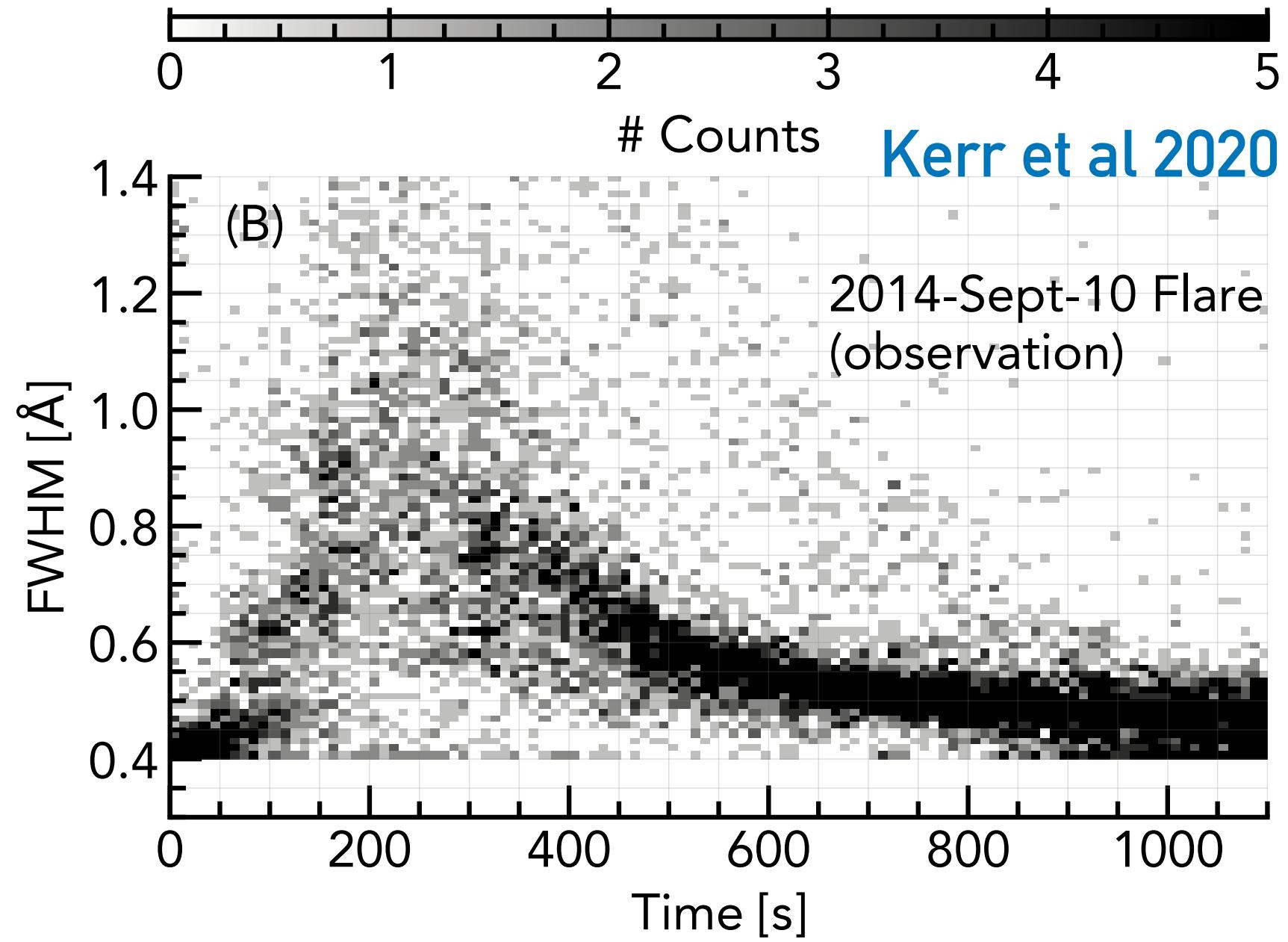
- ▶ Improved the treatment of Stark broadening for Mg II, using the STARK-B database broadens the lines but still not enough – **another factor of 30x was needed.**
- ▶ One possibility is that we are not heating the lower atmosphere sufficiently.



\*\* See also Kowalski et al 2022 for improved treatment of H Stark broadening in RADYN, achieving much better model-data consistency

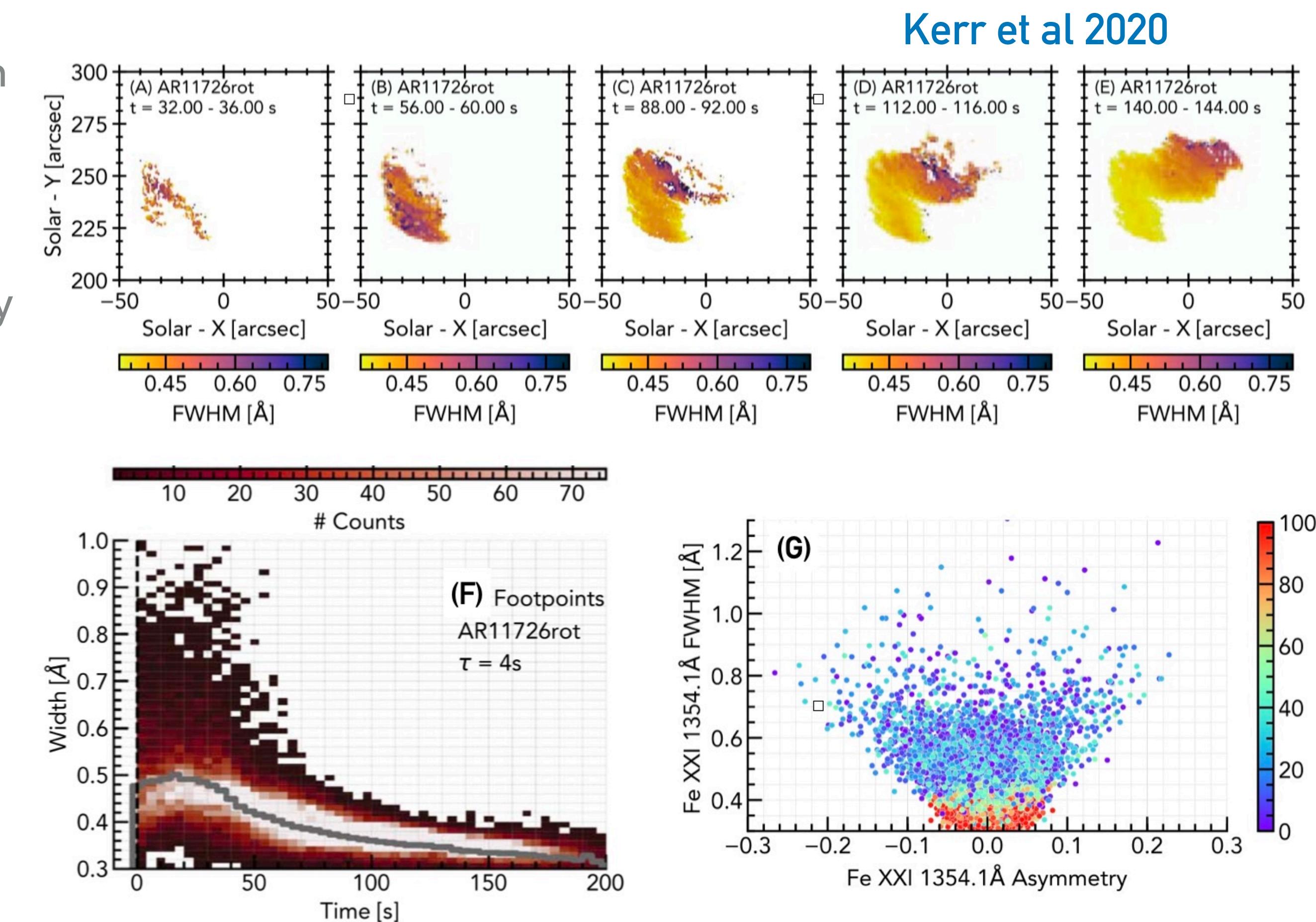
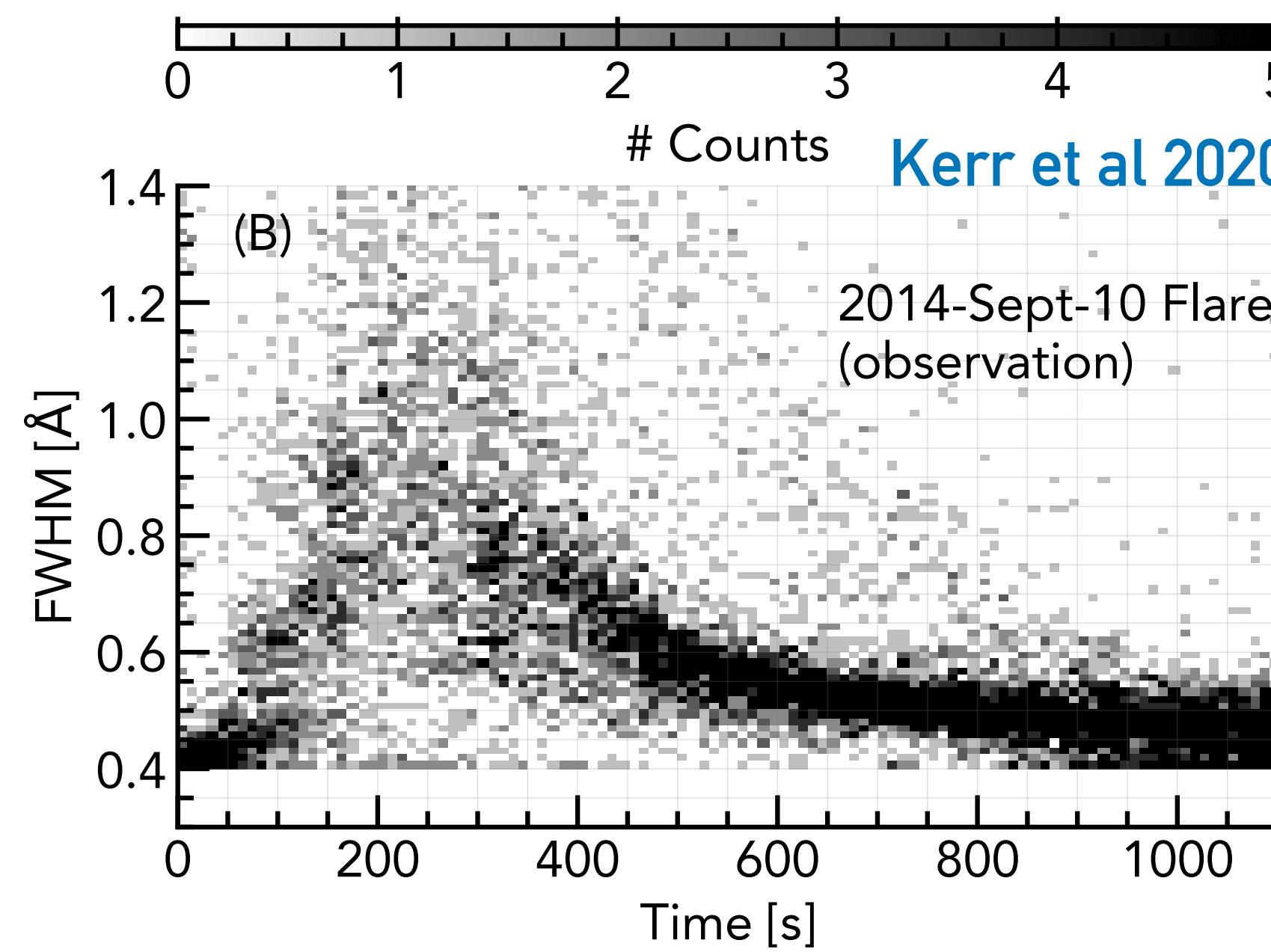
# LINE WIDTHS! - FE XXI

- ▶ Line widths of coronal species, e.g Fe XXI aren't much better.
- ▶ Modelling suggests that superposition of loops is not the answer to the Fe XXI line broadening observed by IRIS.



# LINE WIDTHS! - FE XXI

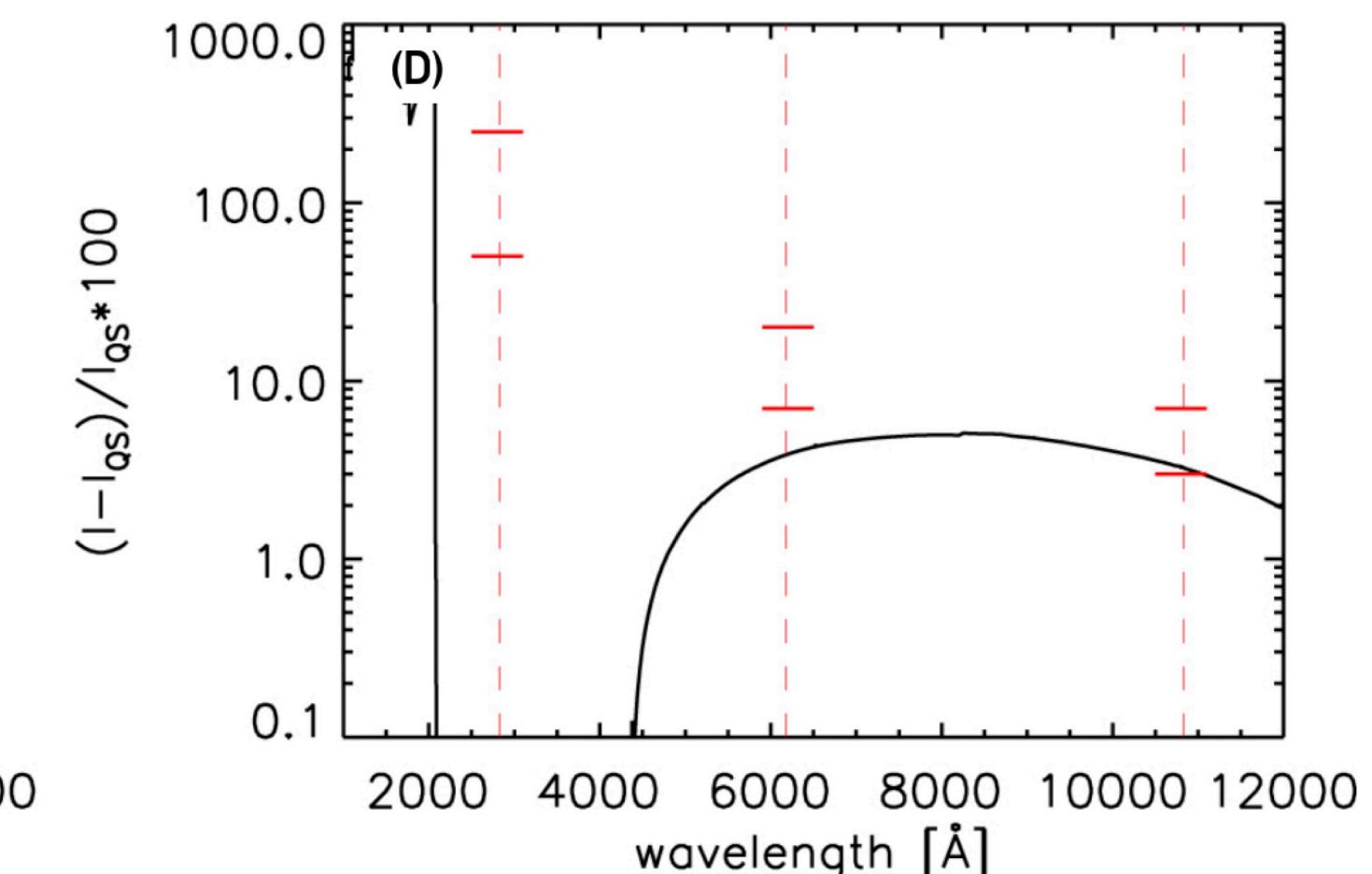
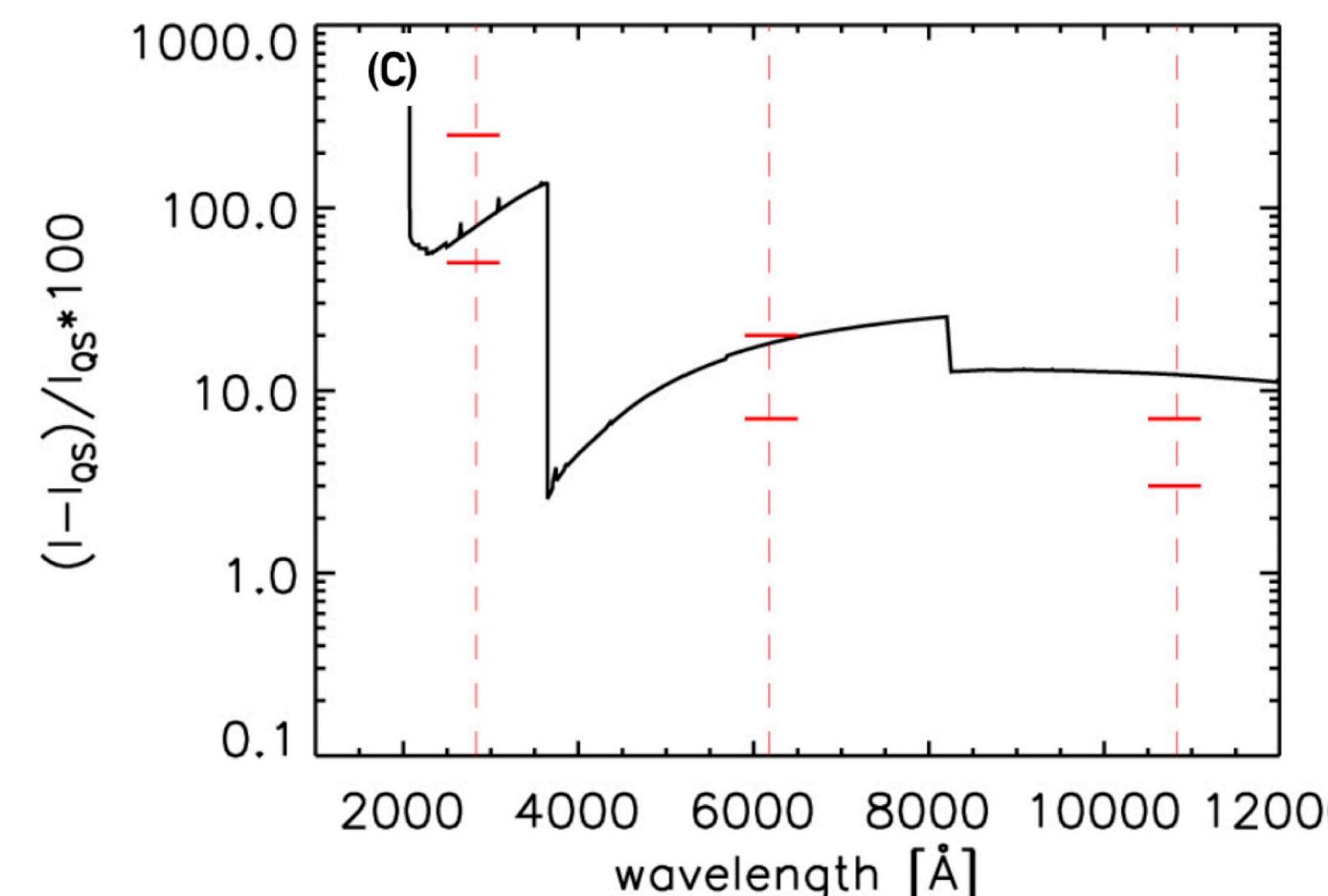
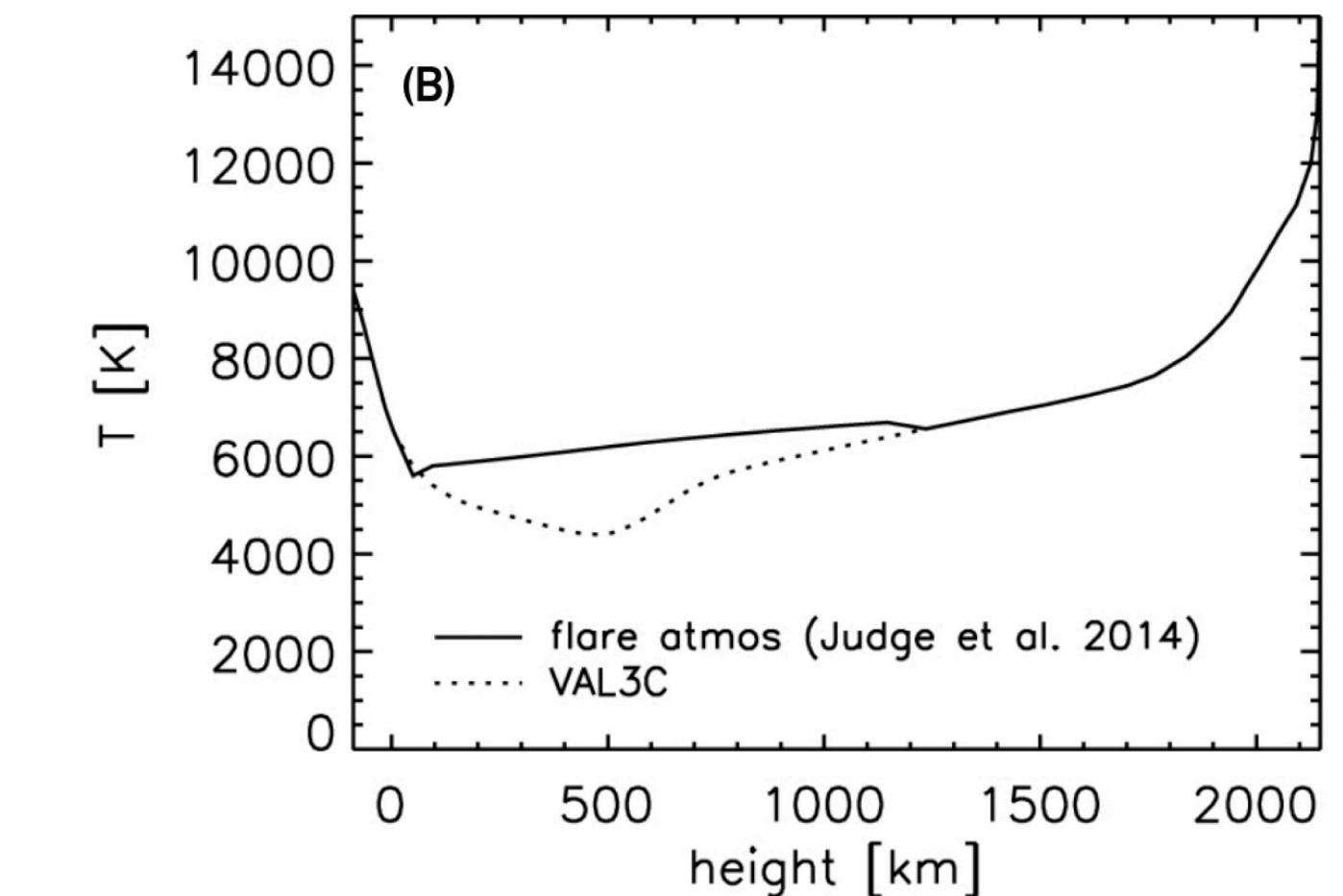
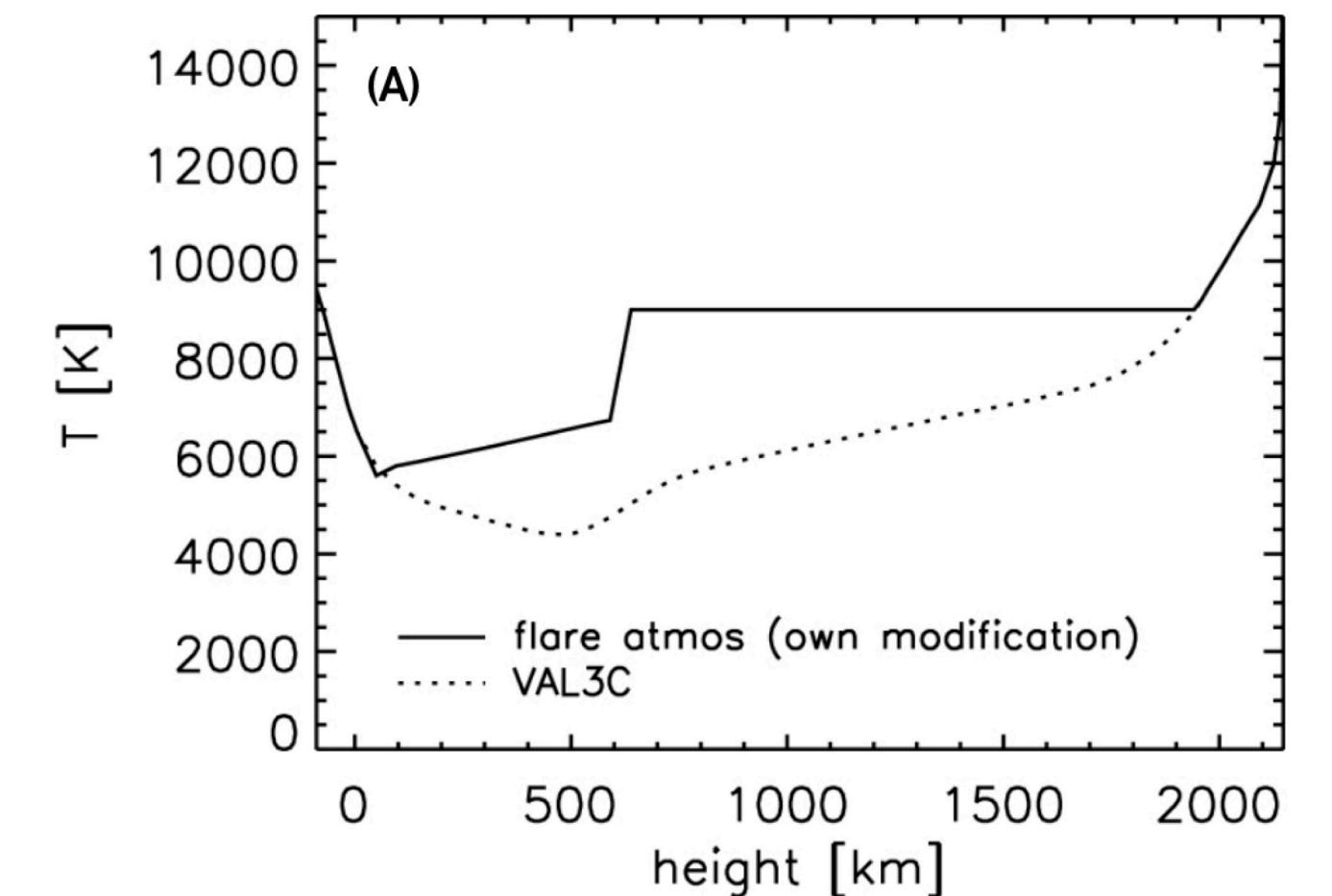
- ▶ Line widths of coronal species, e.g Fe XXI aren't much better.
- ▶ Modelling suggests that superposition of loops is not the answer to the Fe XXI line broadening observed by IRIS.



# DEEP HEATING: WHITE LIGHT FLARES

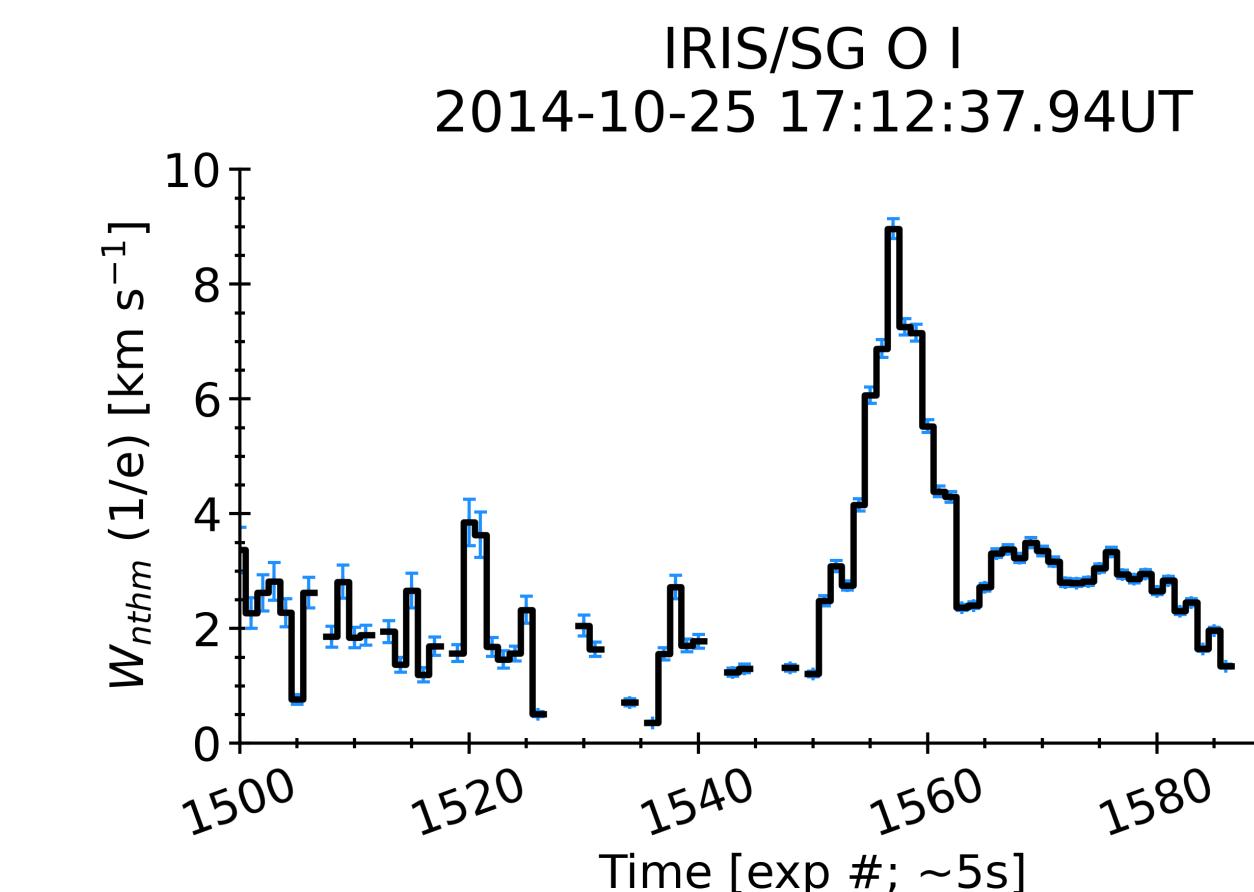
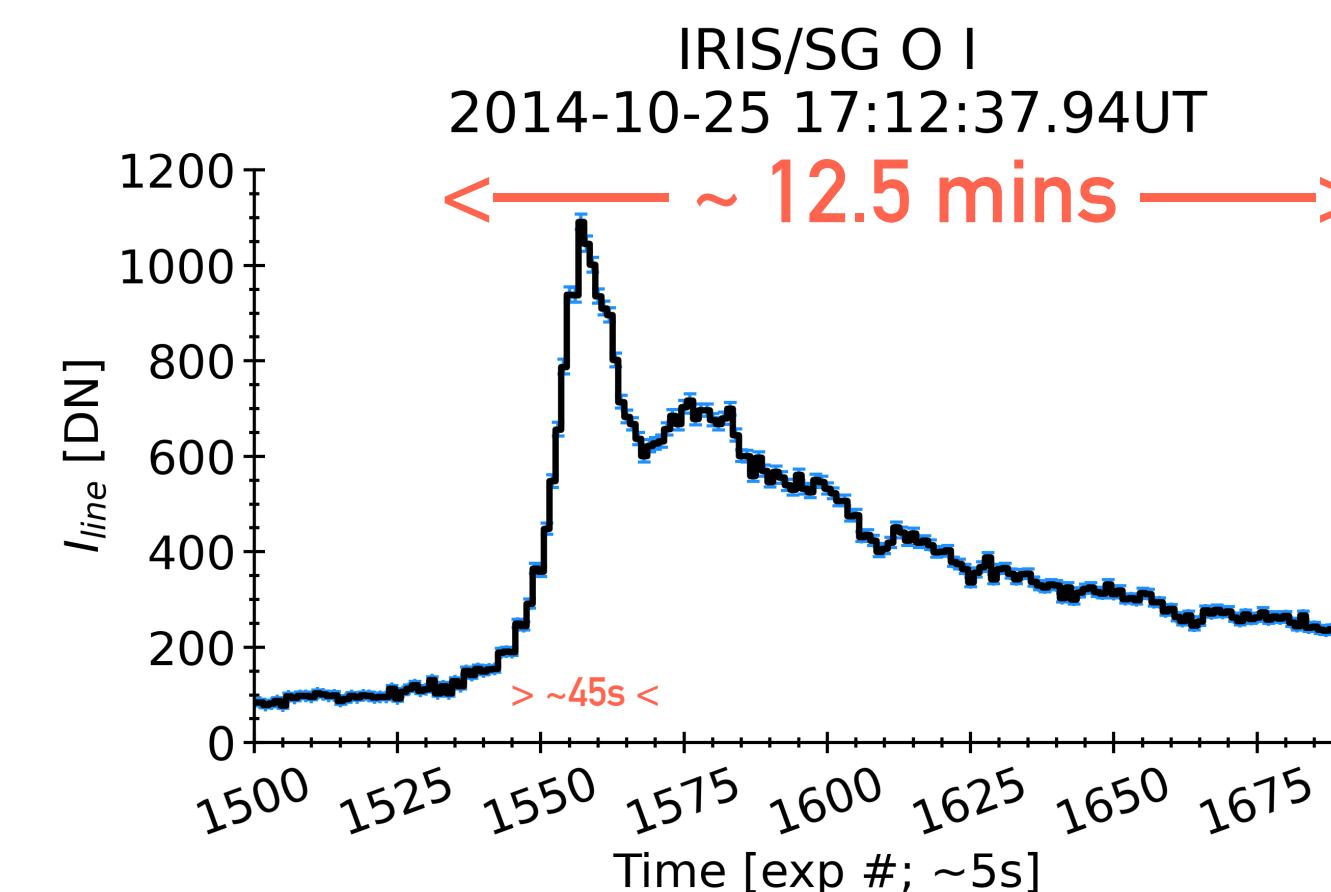
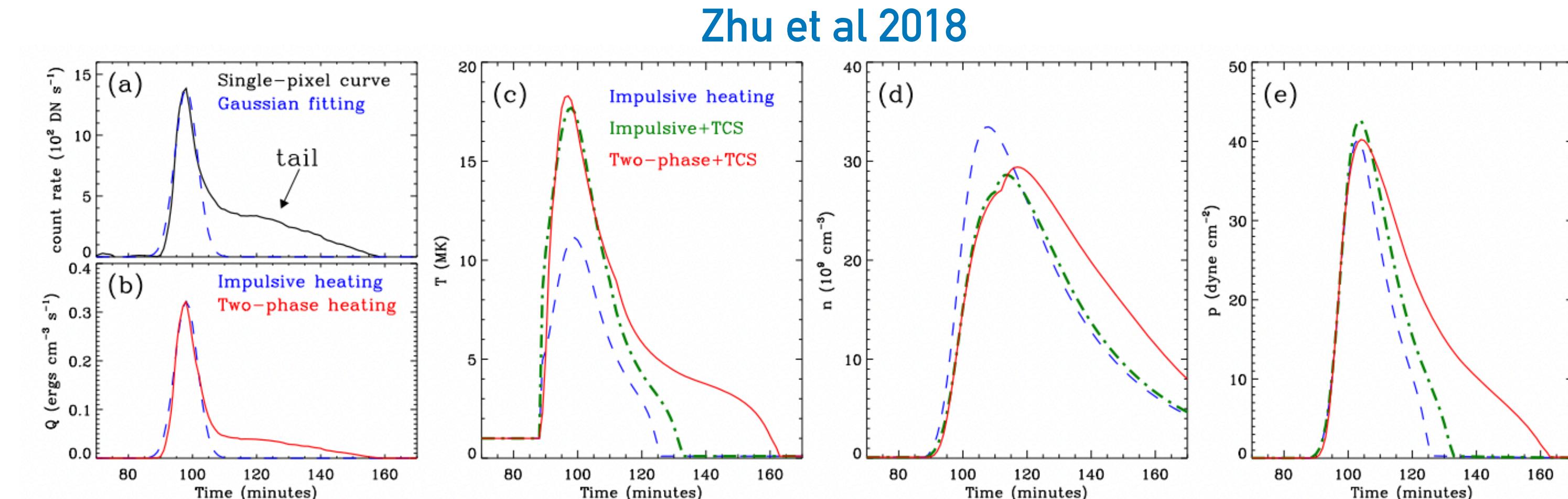
- ▶ Evidence from continuum observations of the optical and NUV that there could be emission coming from deep layers.
- ▶ Emission mechanisms behind white light flares are ambiguous, with some observations consistent with upper photospheric heating and others with mid-upper chromosphere (models predict primarily upper chromospheric).
- ▶ Fe II - to - NUV continuum also suggests heating at high column mass.
- ▶ How do we heat deeper than models currently predict? – missing energy transport ingredients.

Kleint et al 2016



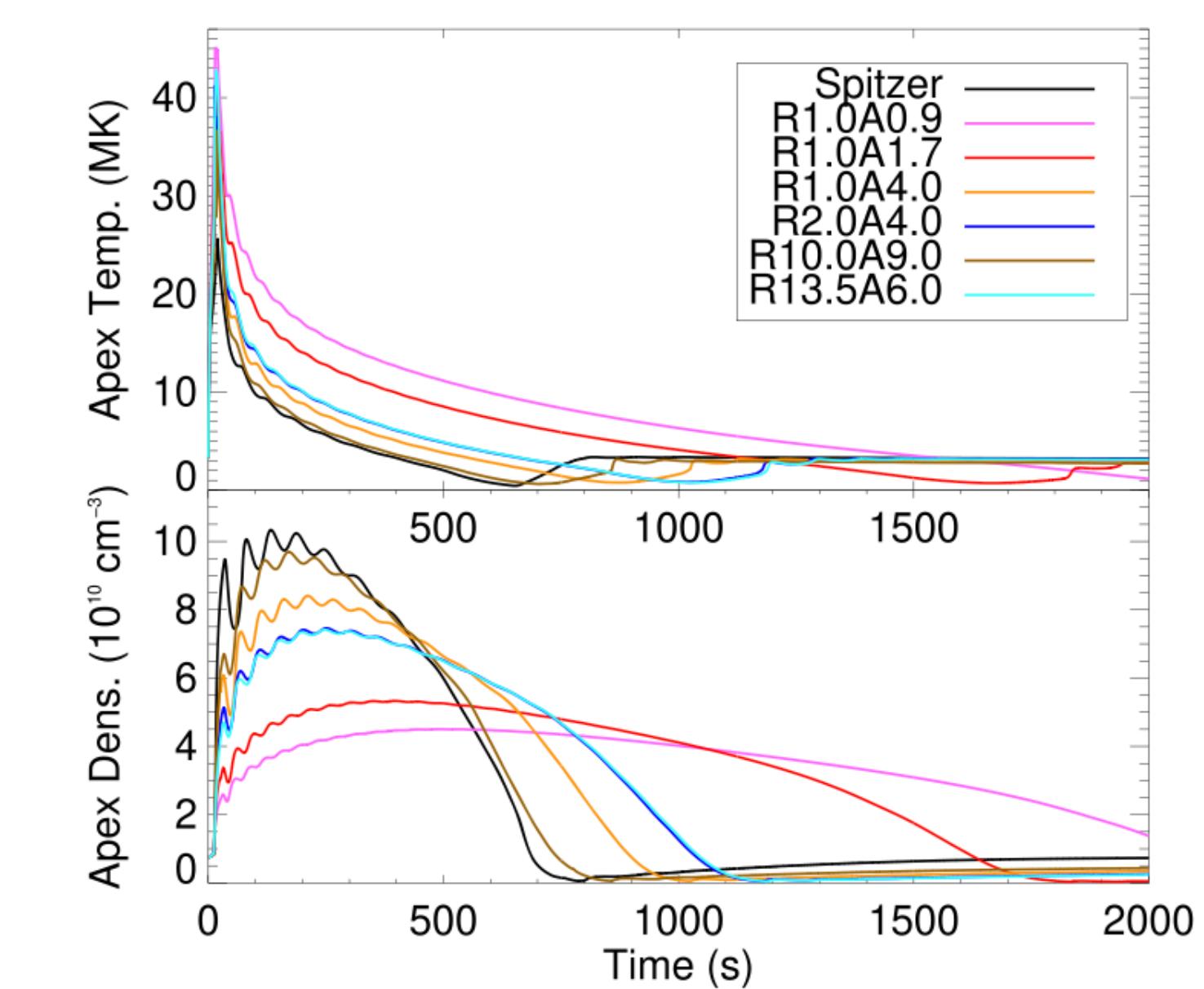
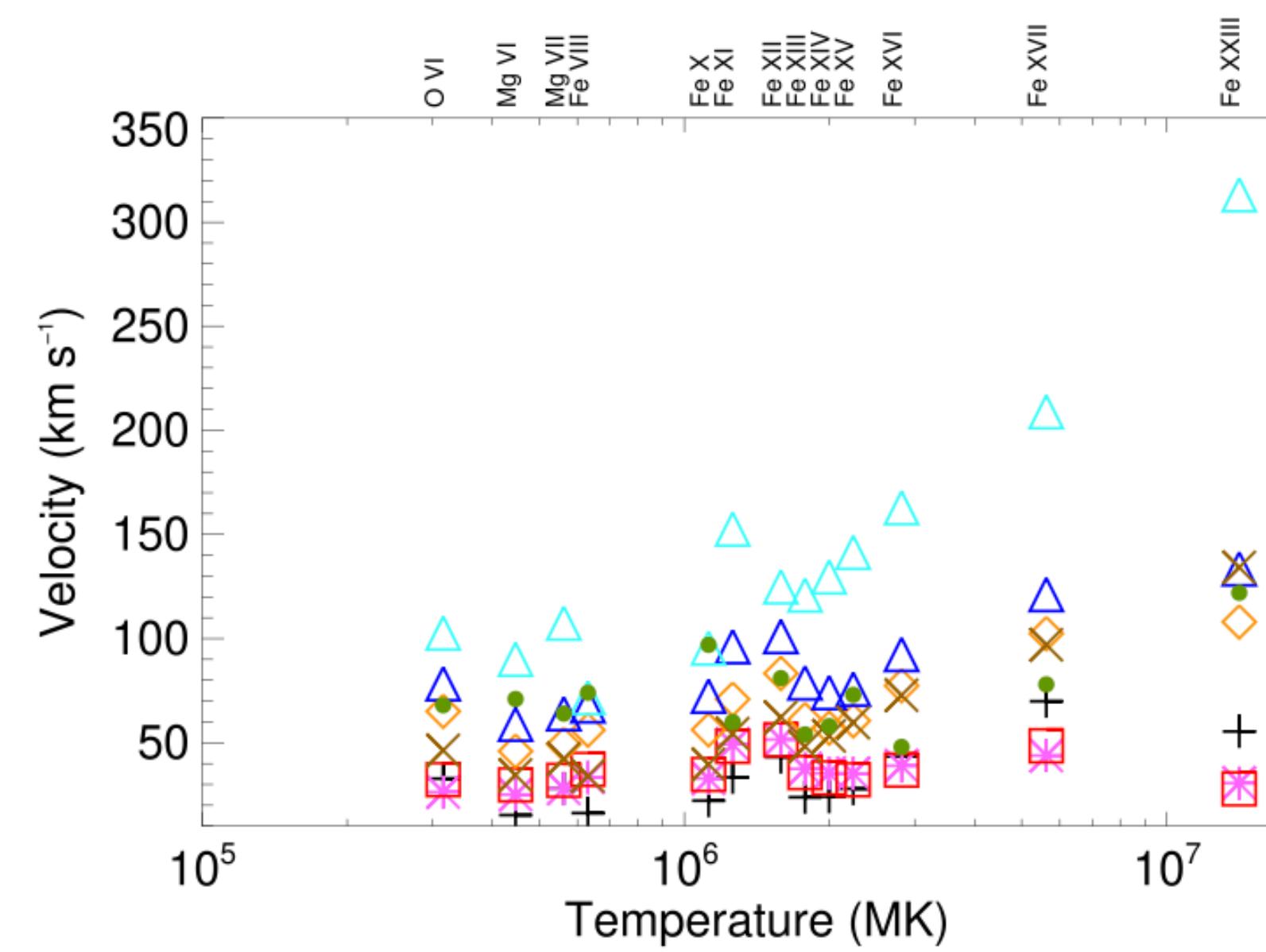
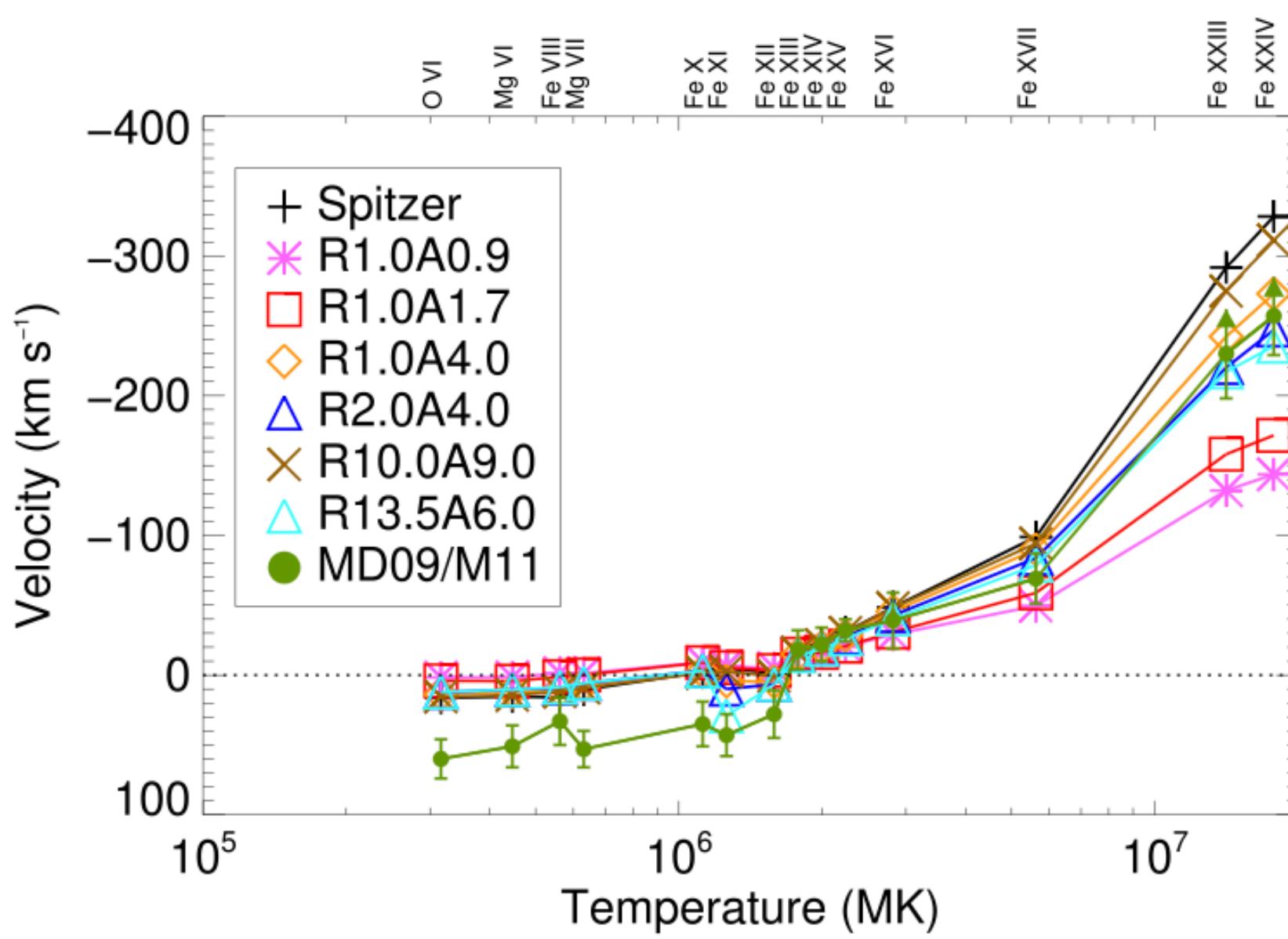
# DURATION OF THE FLARE GRADUAL PHASE

- Both the global and local gradual phase predicted by loop models are far too short.
- Observations of the decay in individual flare footpoints in the chromosphere suggest that after the initial rapid decay (~ one minute?) there is an extended period of up to several minutes (sometimes longer!).
- Models show that as soon as energy injection ceases the atmosphere cools catastrophically and that line intensities plummet.
- Two schools of thought: (1) continued energy injection through the gradual phase (with magnitudes rivaling the impulsive phase) or (2) thermal conduction is suppressed.

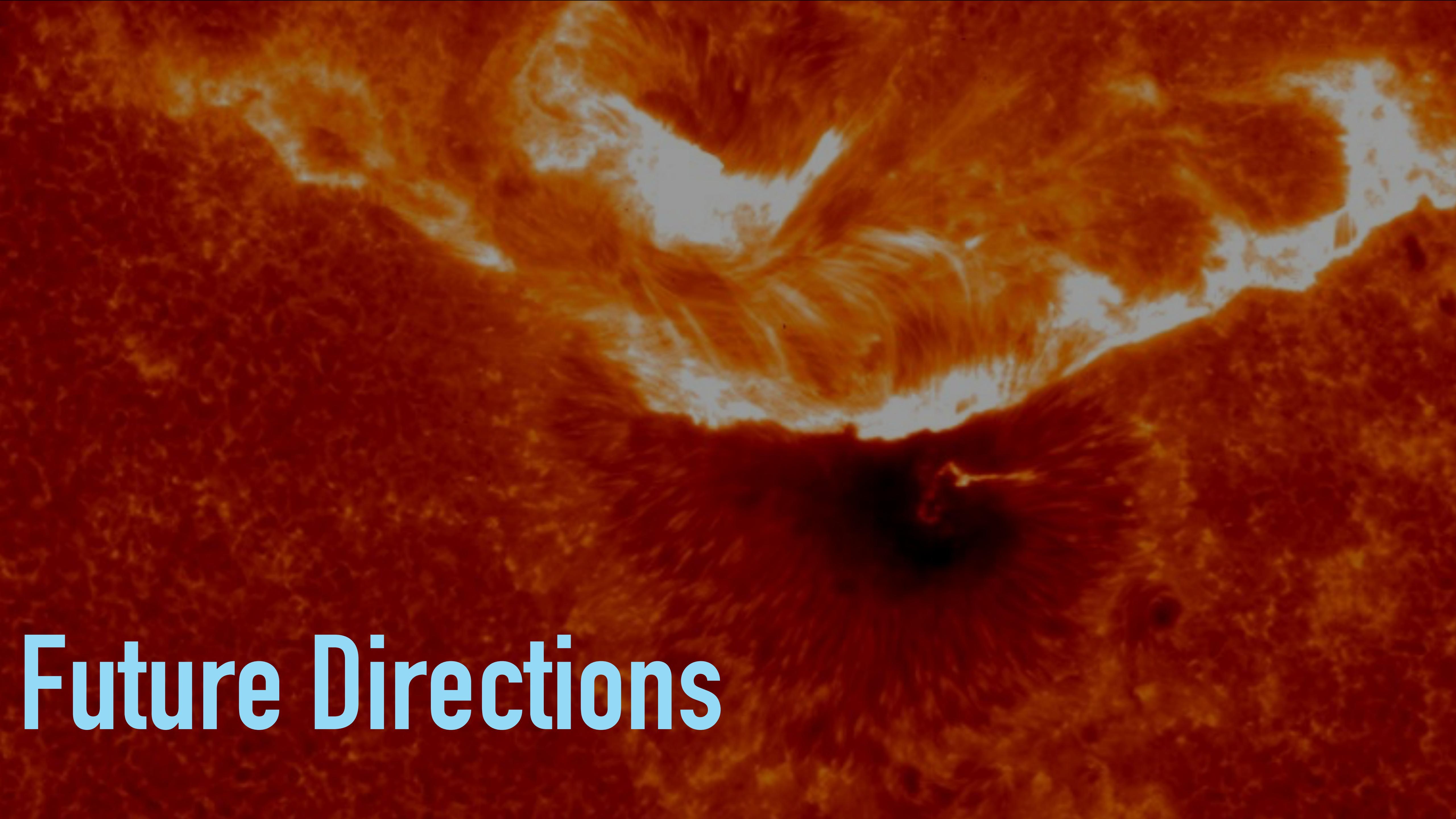


# DURATION OF THE FLARE GRADUAL PHASE

- ▶ Allred et al 2022 included suppression of thermal conduction when modelling the flare from Milligan and Dennis (2009), using the theory of Emslie & Bian 2018.
- ▶ Cooling times were extended dramatically, magnitude of flows were reduced, AND the non-thermal widths were more consistent with observations (calculated using the turbulent mean free path that causes the conduction suppression).



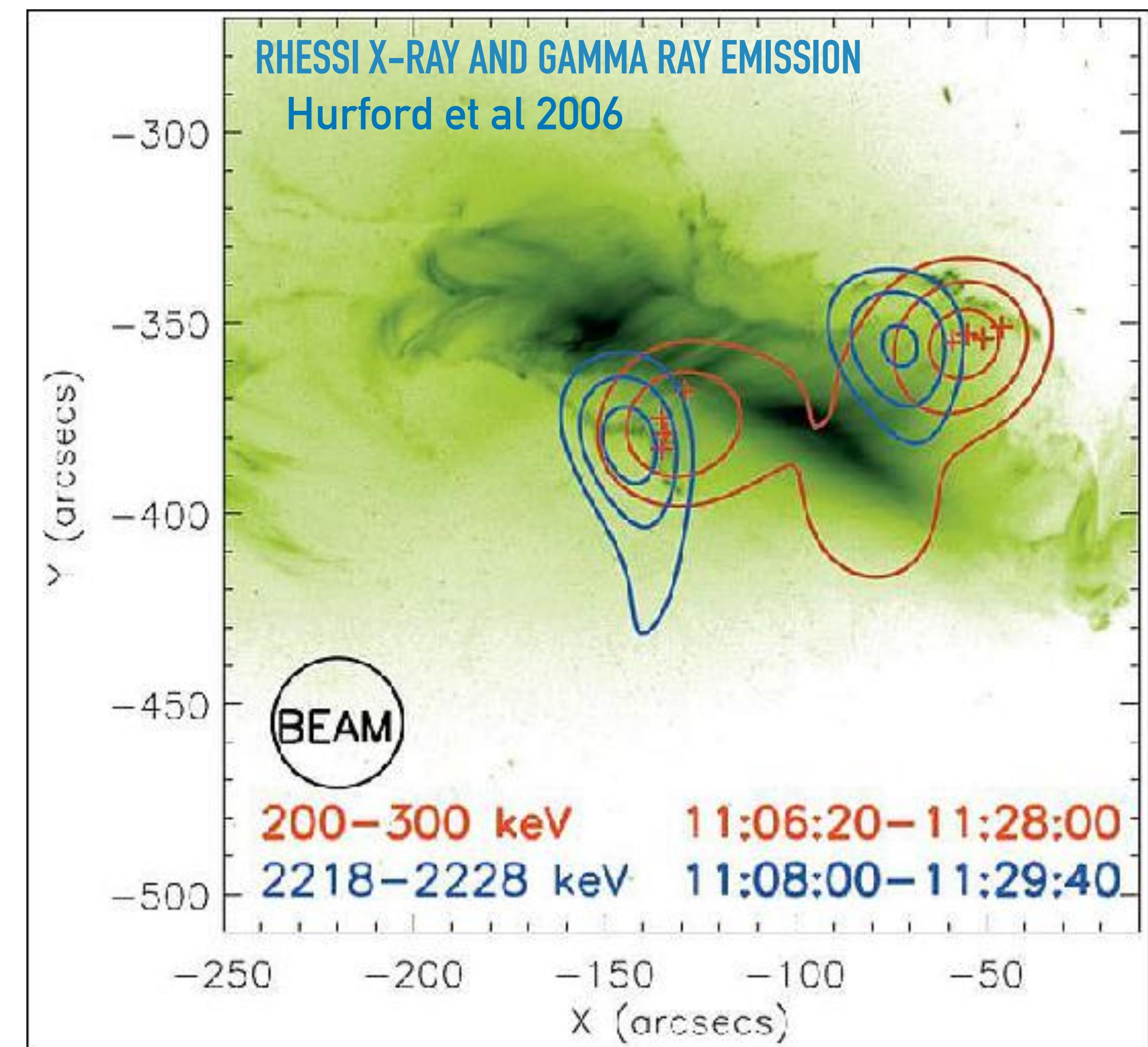
Allred et al 2022



# Future Directions

# PROTON BEAMS IN SOLAR FLARES

- ▶ Very likely that protons are accelerated in flares too.
- ▶ They may even carry a substantial fraction (potentially equal to that of electrons!) of the energy released in a flare and can penetrate deeper.
- ▶ However they are generally ignored in flare models (partly due to very poor constraints of their energy distribution, particularly at low energy).
- ▶ There is a danger that we are missing up to half of the flare energy transported to the lower atmosphere!
- ▶ RADYN+FP now has the ability to model these and we have started those investigations... see Kerr et al 2023 ApJ for the first of those.



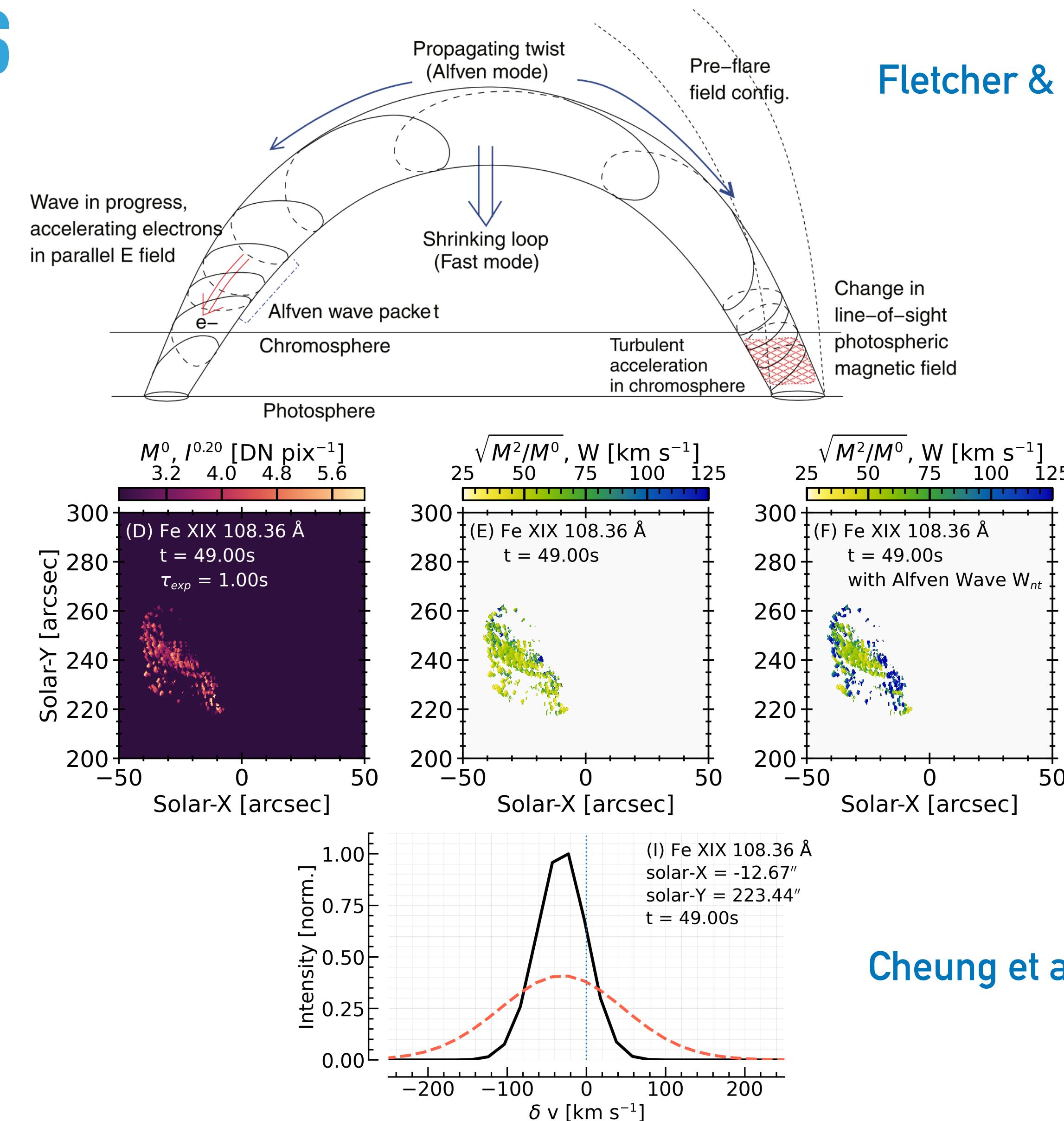
# ALFVEN WAVES IN SOLAR FLARES

Fletcher & Hudson, 2008

- ▶ Flares are fundamentally a large scale magnetic disturbance, so it is very reasonable to expect that MHD waves are produced, and logical to study their role in energy transport.

- ▶ There has been interest in energy transport mechanisms 'beyond the standard model' for some time, but this was recently renewed by Fletcher & Hudson (2008).

- ▶ Dissipation of Alfvén waves has been suggested as either an alternative, or complementary mechanism to deliver energy from the coronal release site to the chromosphere: Reep et al 2016, 2018; Kerr et al 2016.

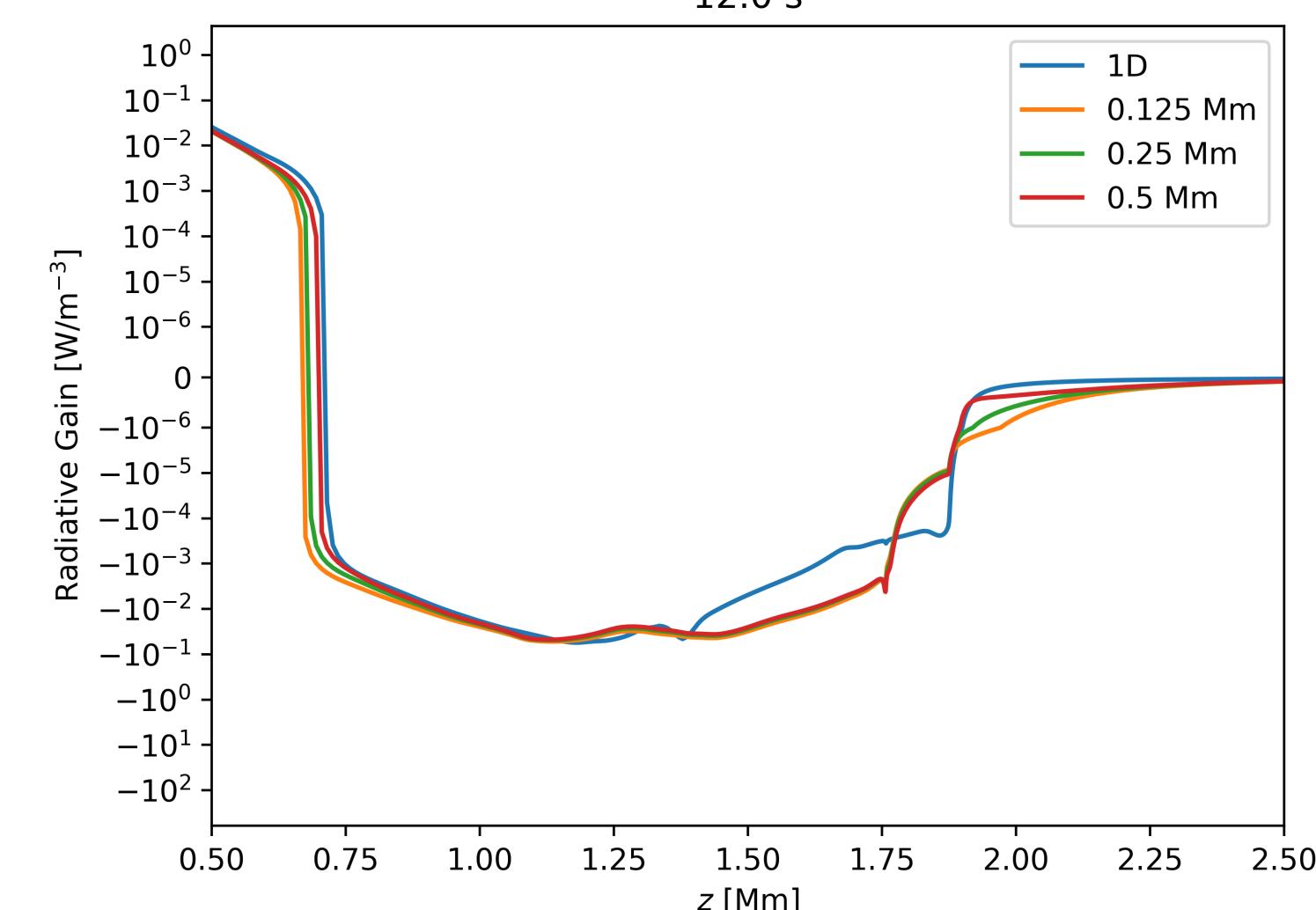
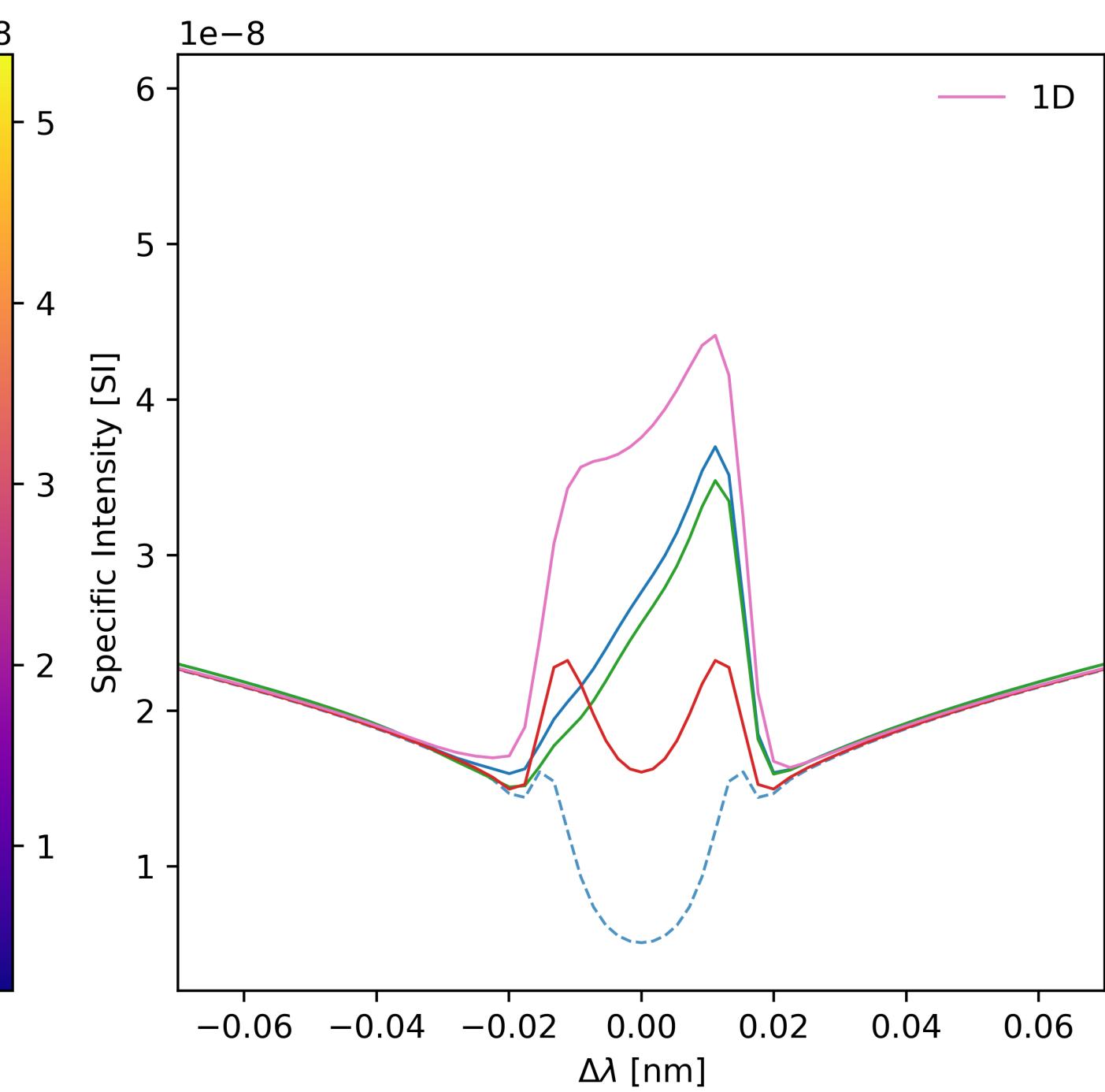
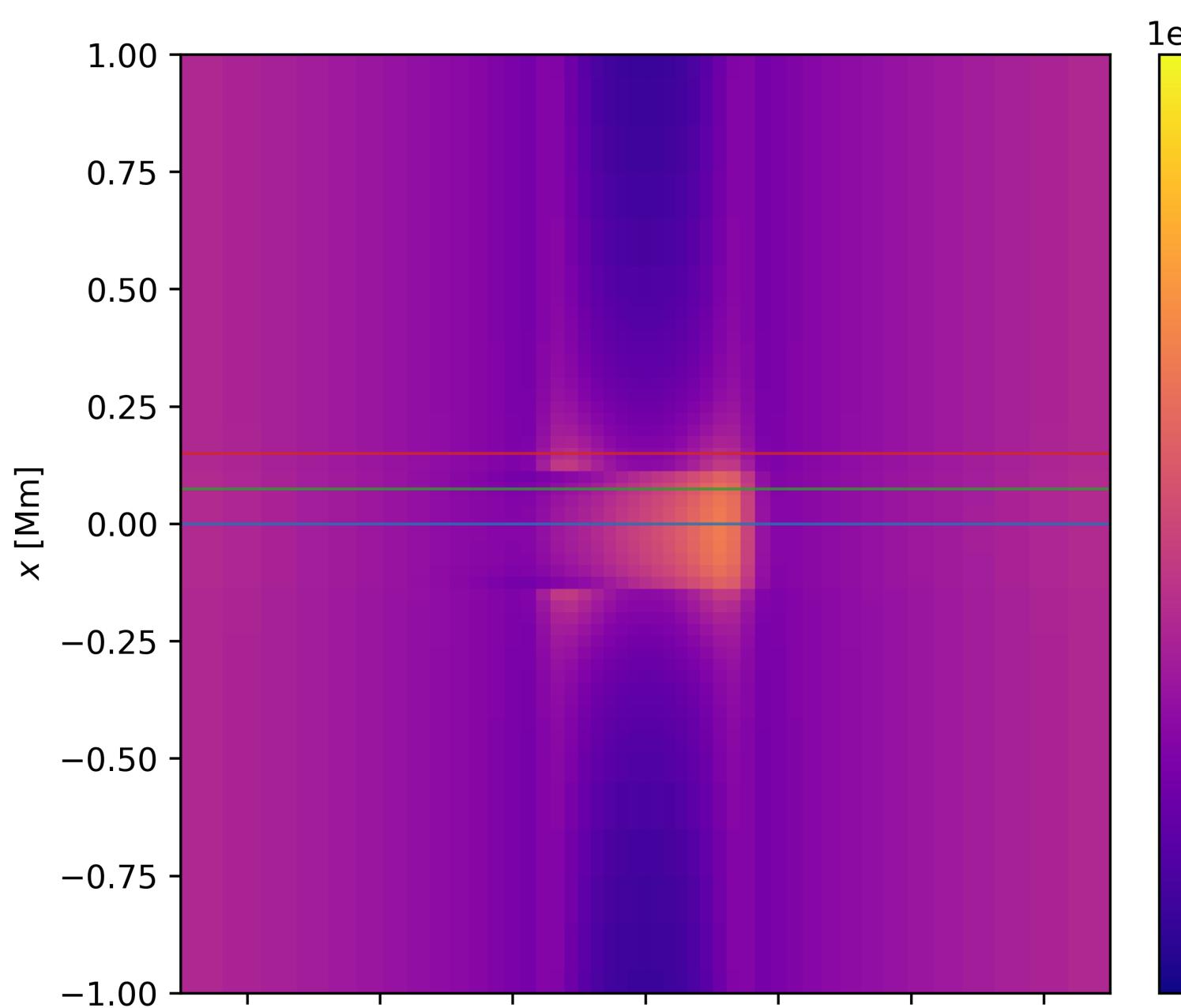


Cheung et al 2022

# MULTI-D RT

- ▶ Implant a kernel of RADYN flare atmosphere of width  $k$  in a 2D periodic volume and solve for time-dependent NLTE populations by “replaying” the thermodynamic evolution of the RADYN model.

Osborne & Fletcher 2022



**Ca II 8542 Å**  
 **$k = 250 \text{ km}$**   
 **$t = 12 \text{ s}$**   
**F9 flat deposition for 10s**  
 **$\delta = 5, E_c = 20 \text{ keV}$**

# TOWARDS A COMPREHENSIVE SOLAR ERUPTIVE EVENT MODEL FRAMEWORK

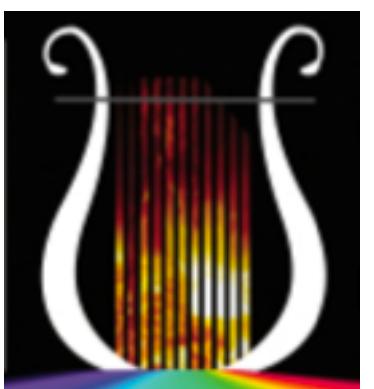
Upcoming game-changing observations demand new ways of thinking about flare modelling.



DKIST ~ now!



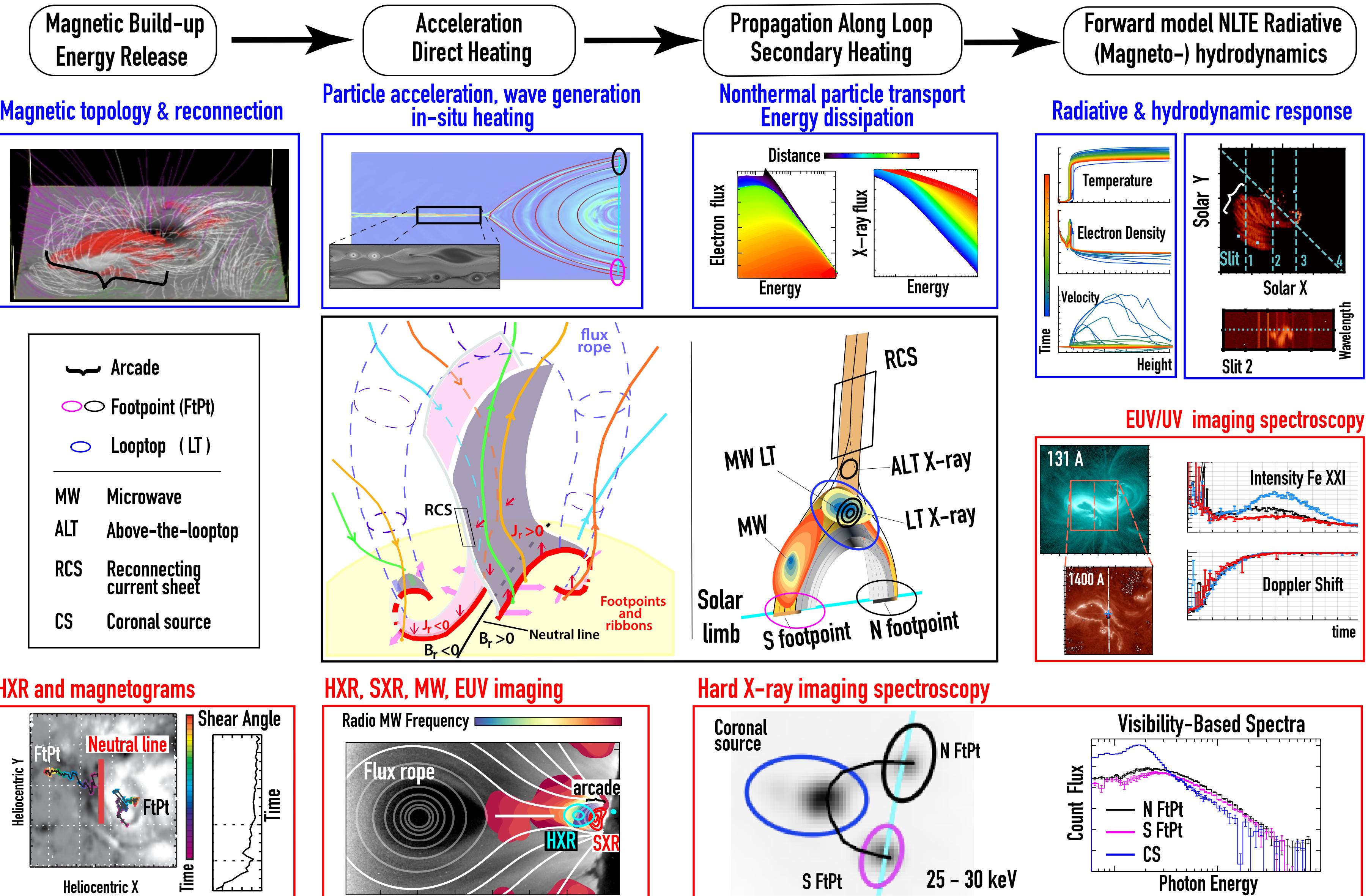
EUVST ~ 2027

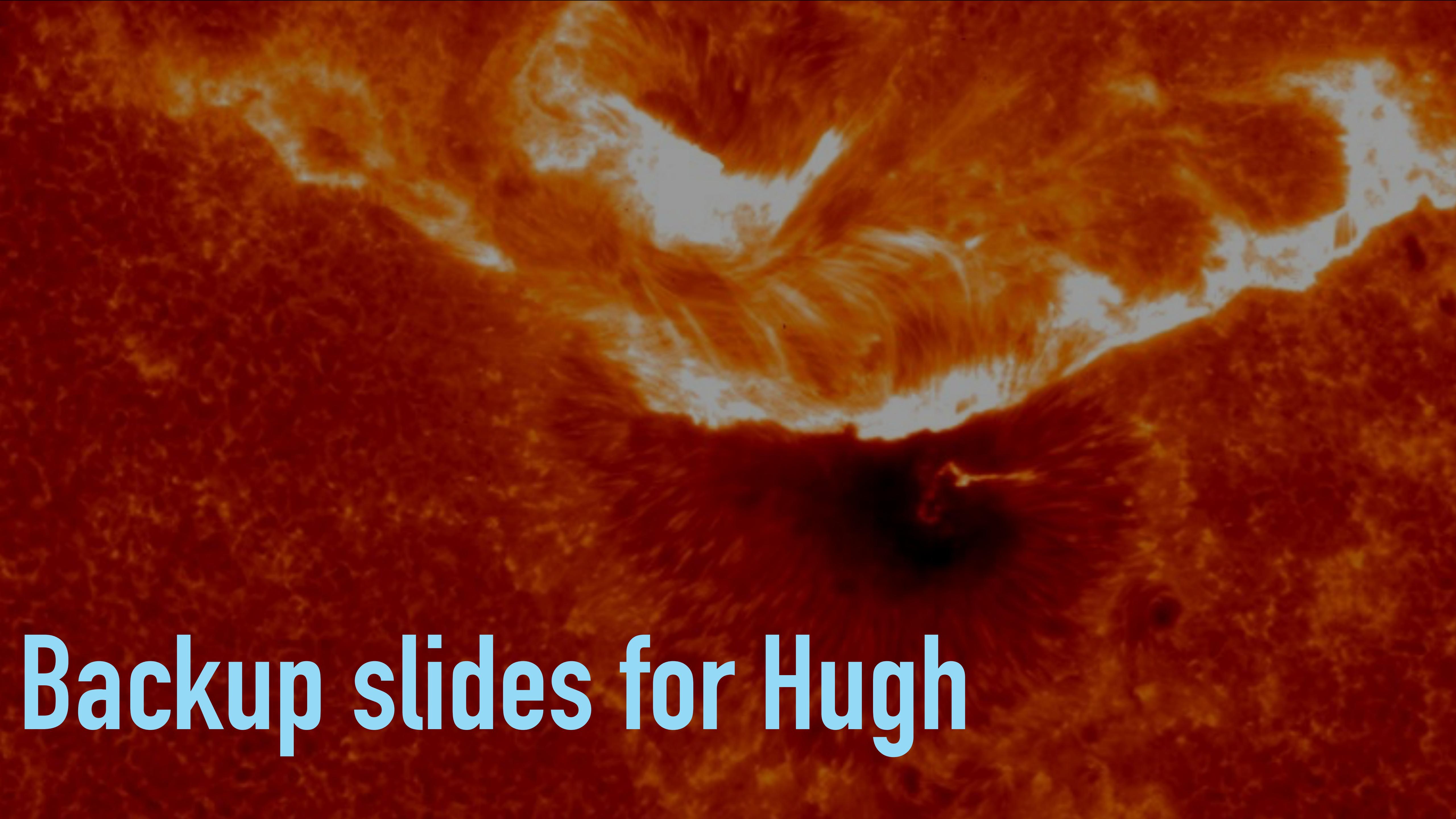


MUSE ~ 2027

Allred et al  
NASEM 2023  
Decadal  
Survey white  
paper

(See also  
Ruan et al  
2020 for a  
hybrid model)

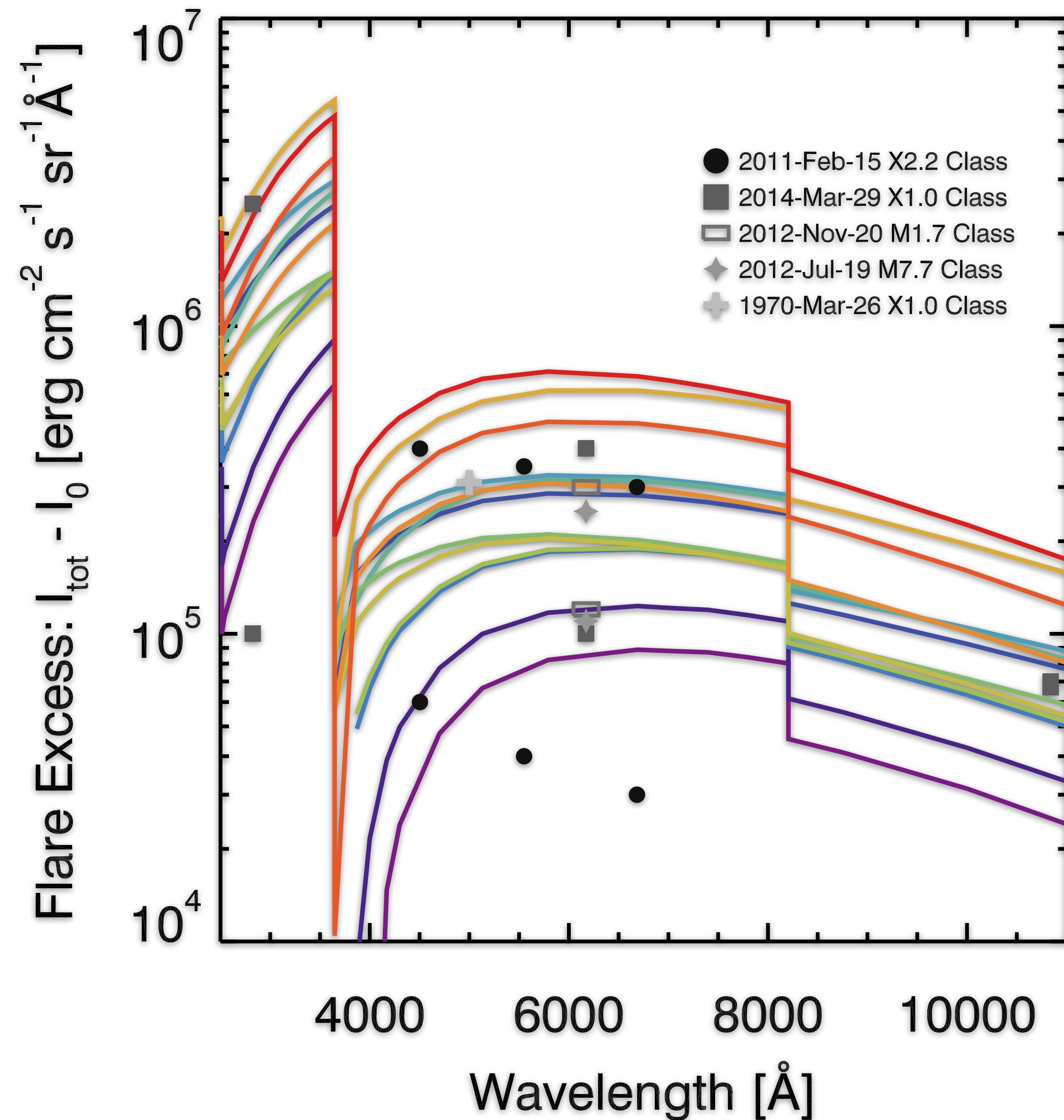




Backup slides for Hugh

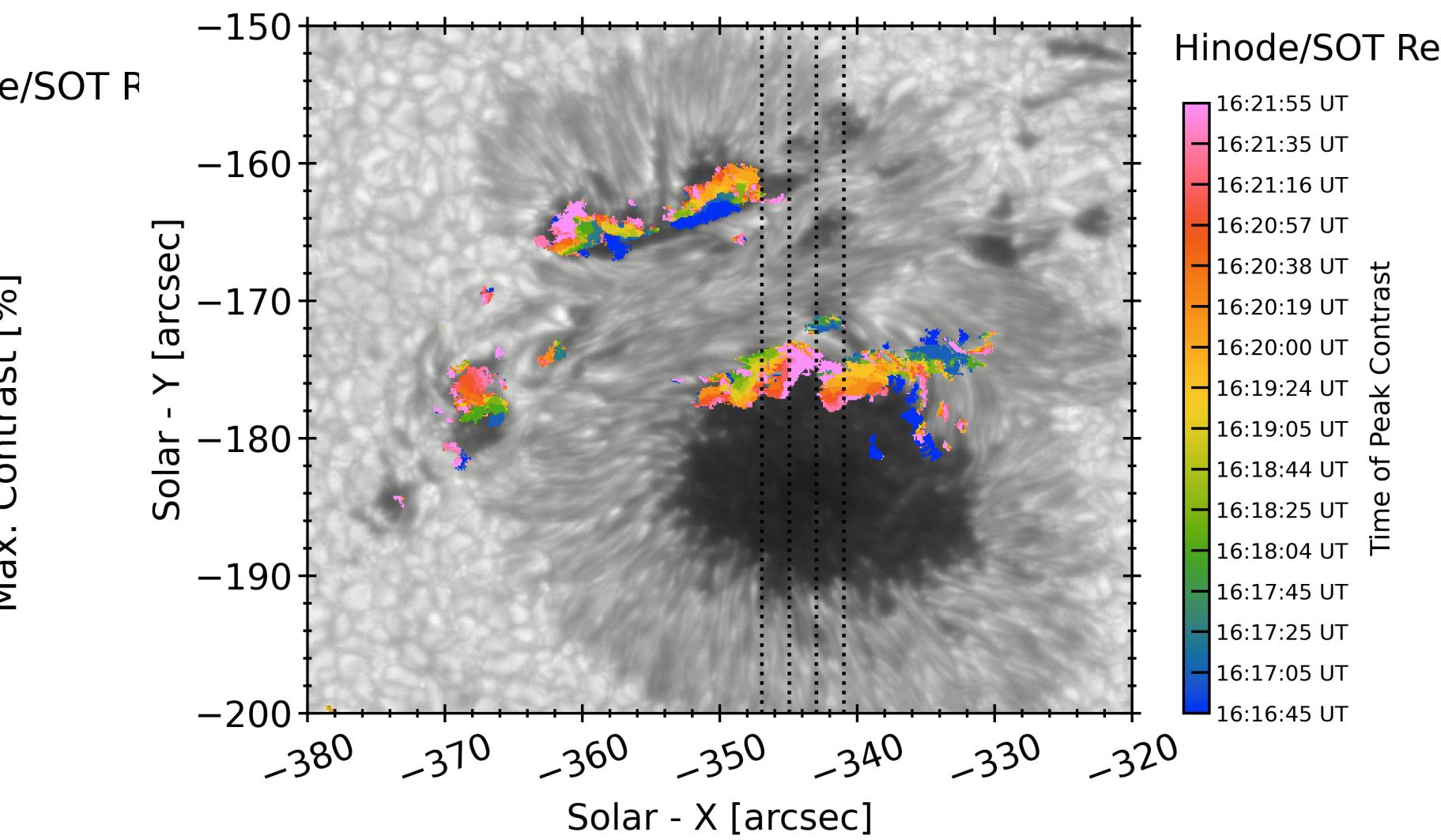
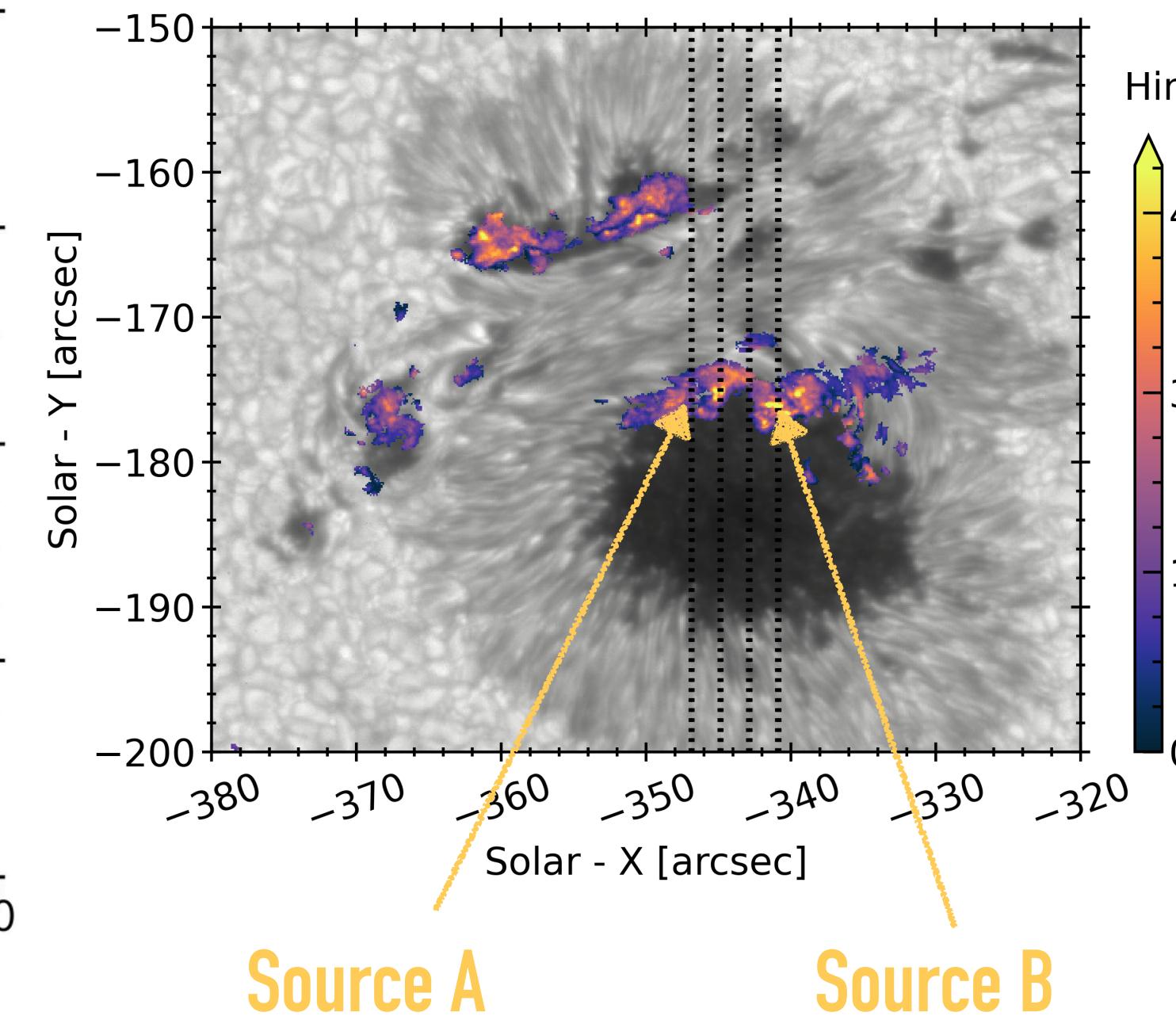
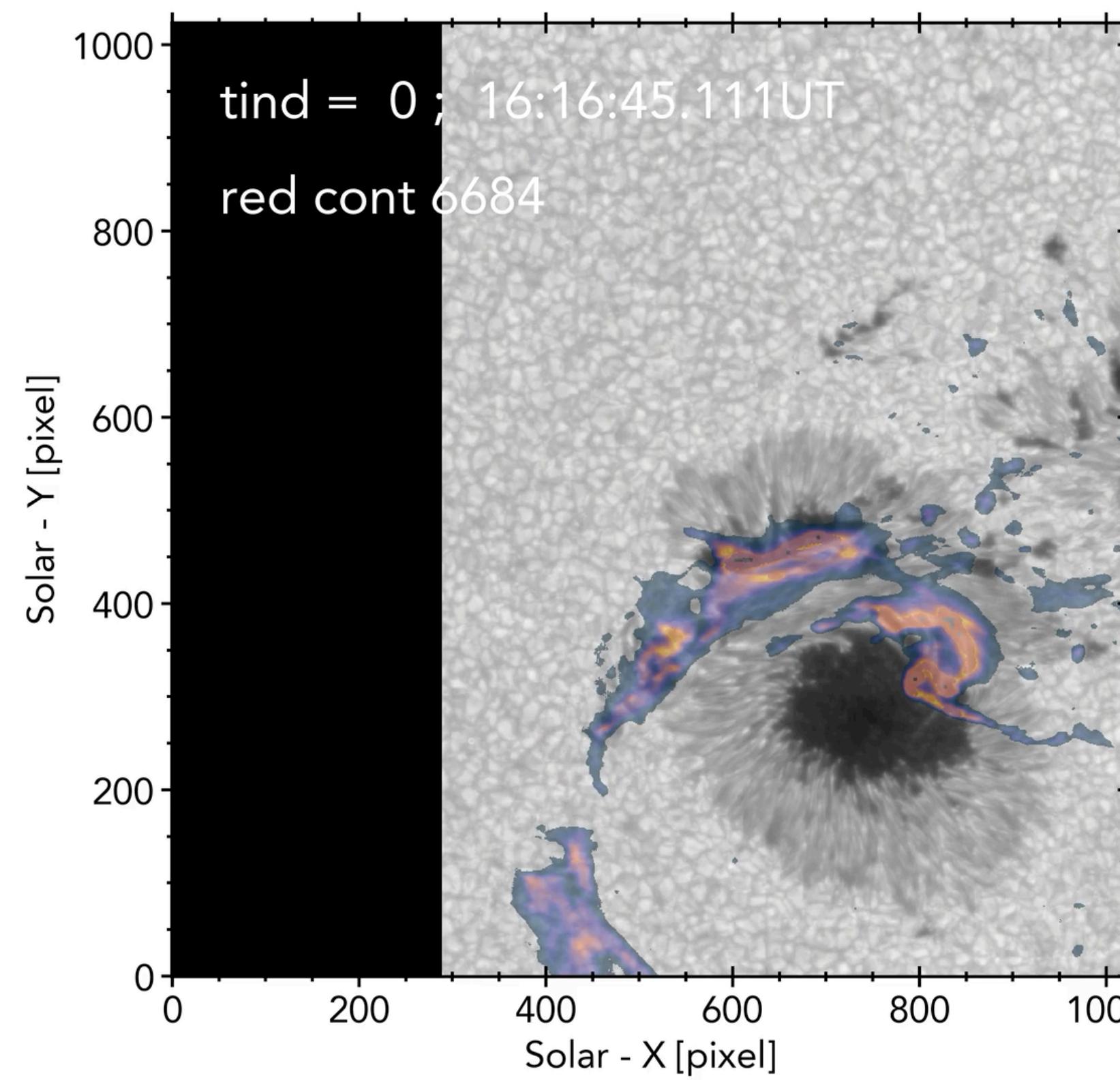
## WLF FORMATION IN RHD MODELS

- ▶ Radiation hydrodynamics (RHD) flare modelling suggests that in strong flares the Balmer jump is always present.
- ▶ Synthetic WLFs are the result, primarily, of optically thin hydrogen recombination radiation produced in the mid-upper flaring chromosphere.
- ▶ Plasma properties in the emitting volume are typically  $T \sim 8\text{-}15\text{ kK}$  and  $n_e \sim 10^{12\text{-}14} \text{ cm}^{-3}$
- ▶ Width (vertical extent) of the emitting region varies from a few dozen to several hundred km.
- ▶ How well do observations agree with this?



# WLF OBSERVATIONS: HINODE/SOT OPTICAL IMAGES

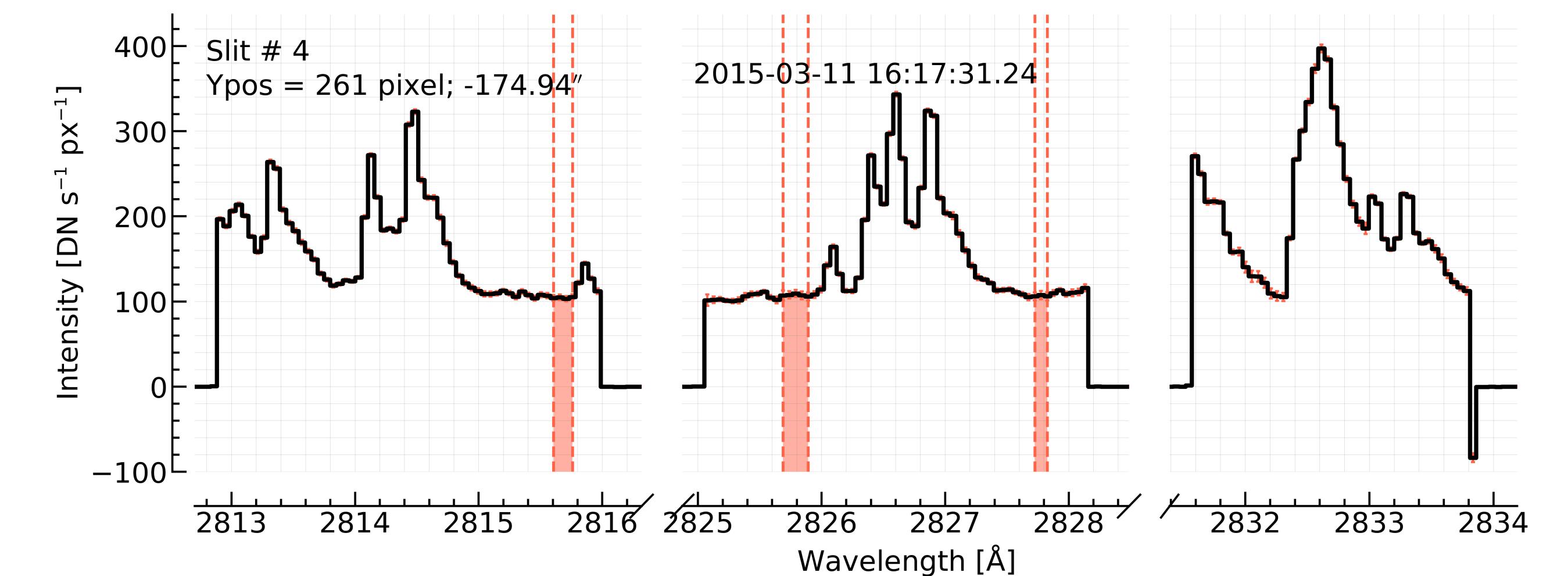
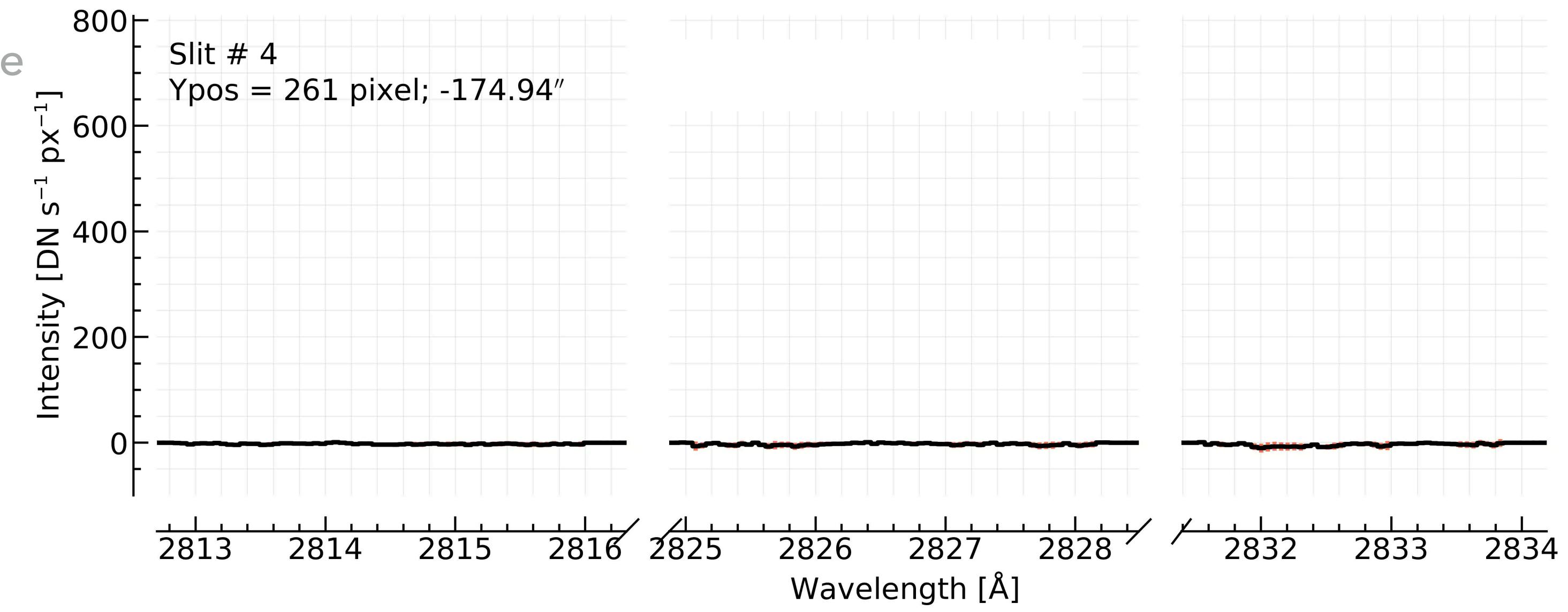
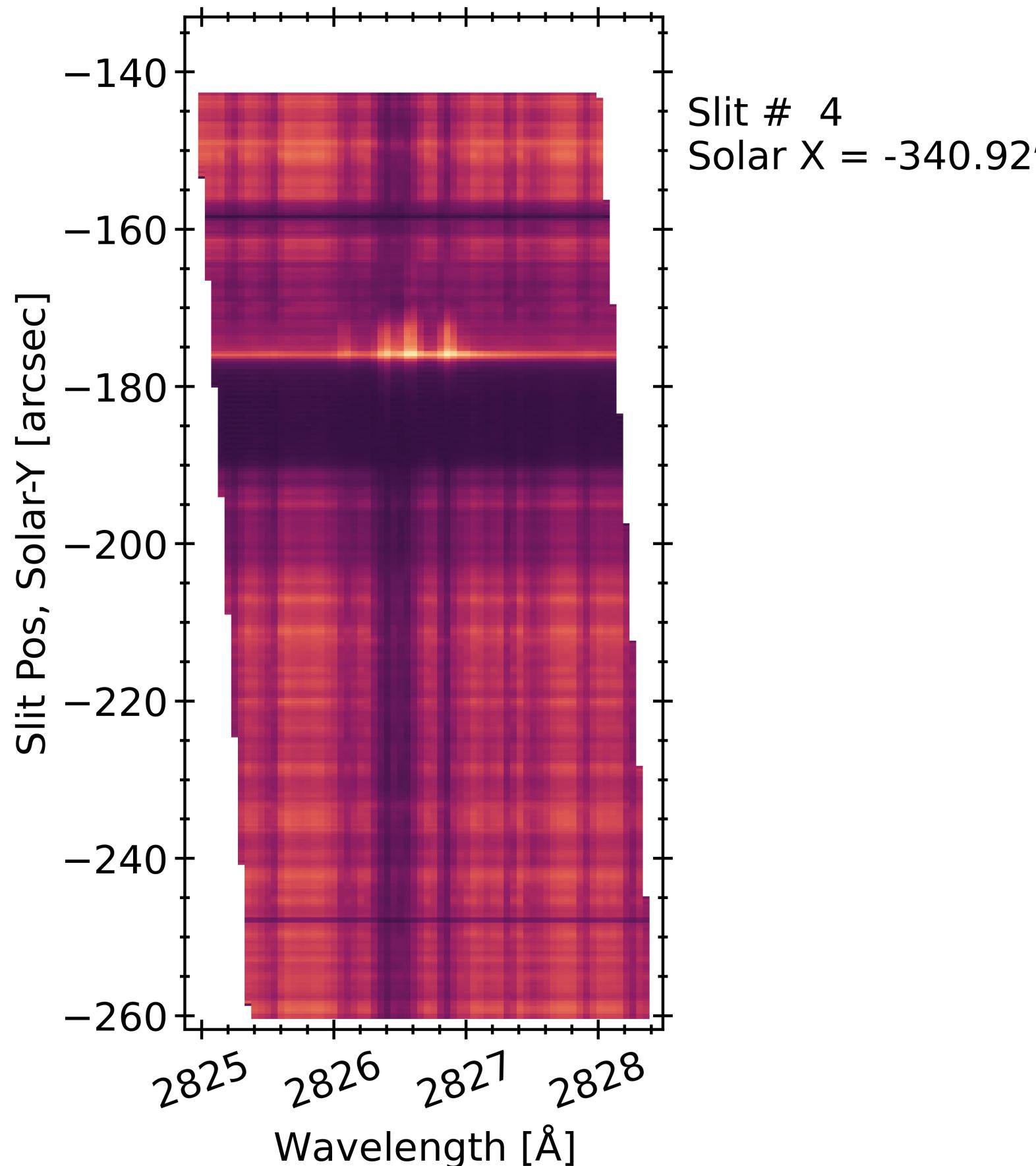
- ▶ Three passbands: **red** (6684Å), **green** (5550Å), **blue** (4504Å).
- ▶ Running difference of those images was used to identify the most likely WLF sources.
- ▶ Those sources were used as seeds in a region growing algorithm to identify weaker sources.



Black dashed lines are the locations of the IRIS slit.

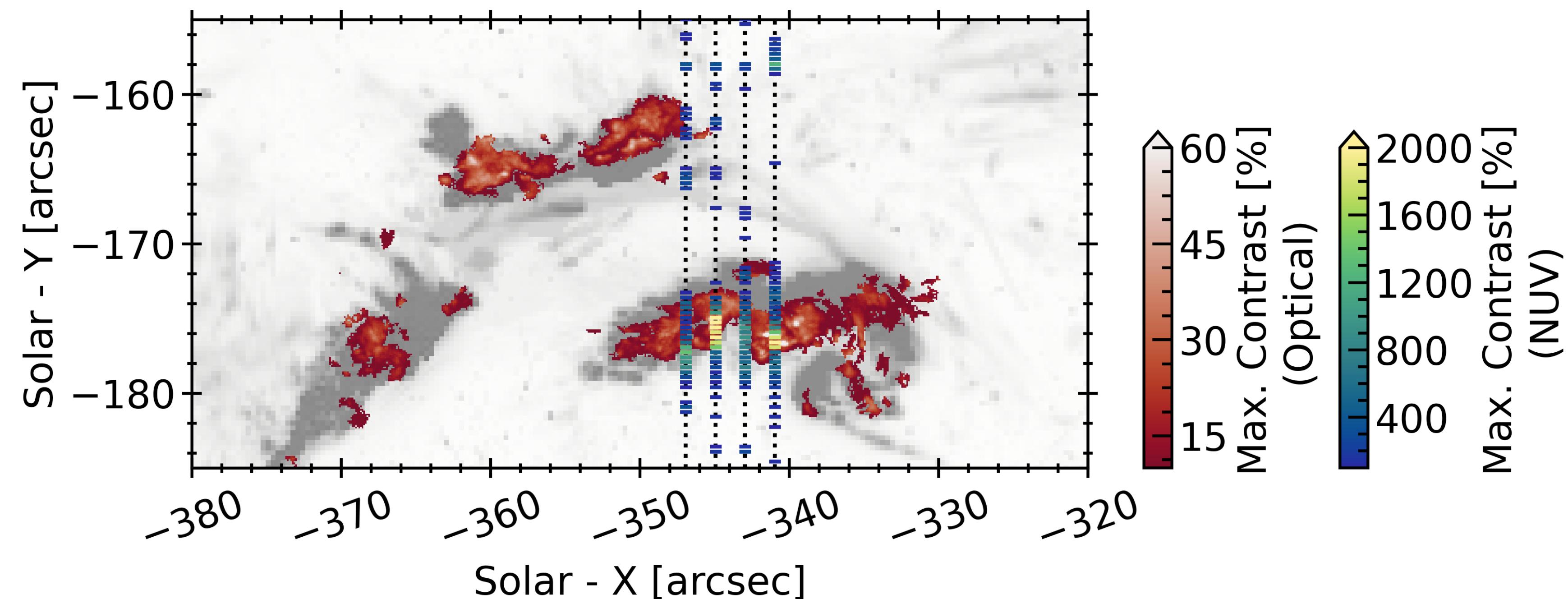
# WLF OBSERVATIONS: IRIS NEAR-UV SPECTRA

- IRIS NUV spectra, allows us to extract line free portions to measure the continuum enhancement near  $2826\text{\AA}$ .

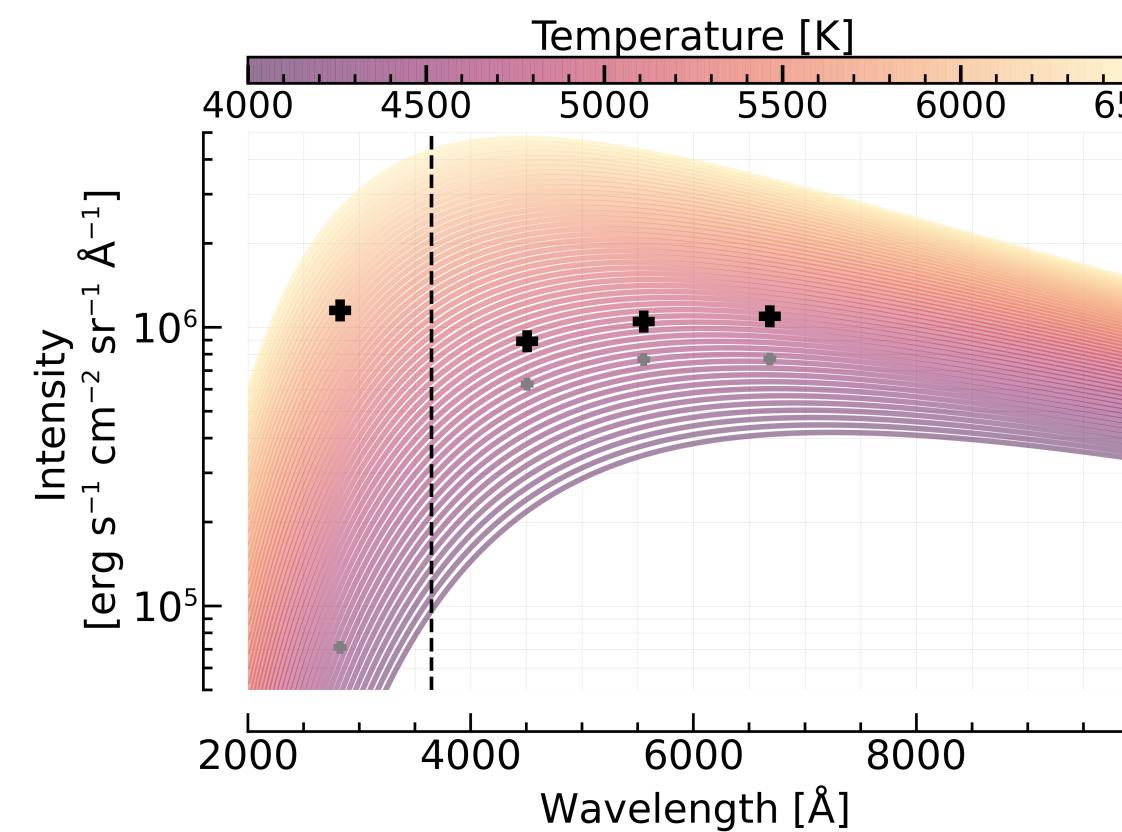
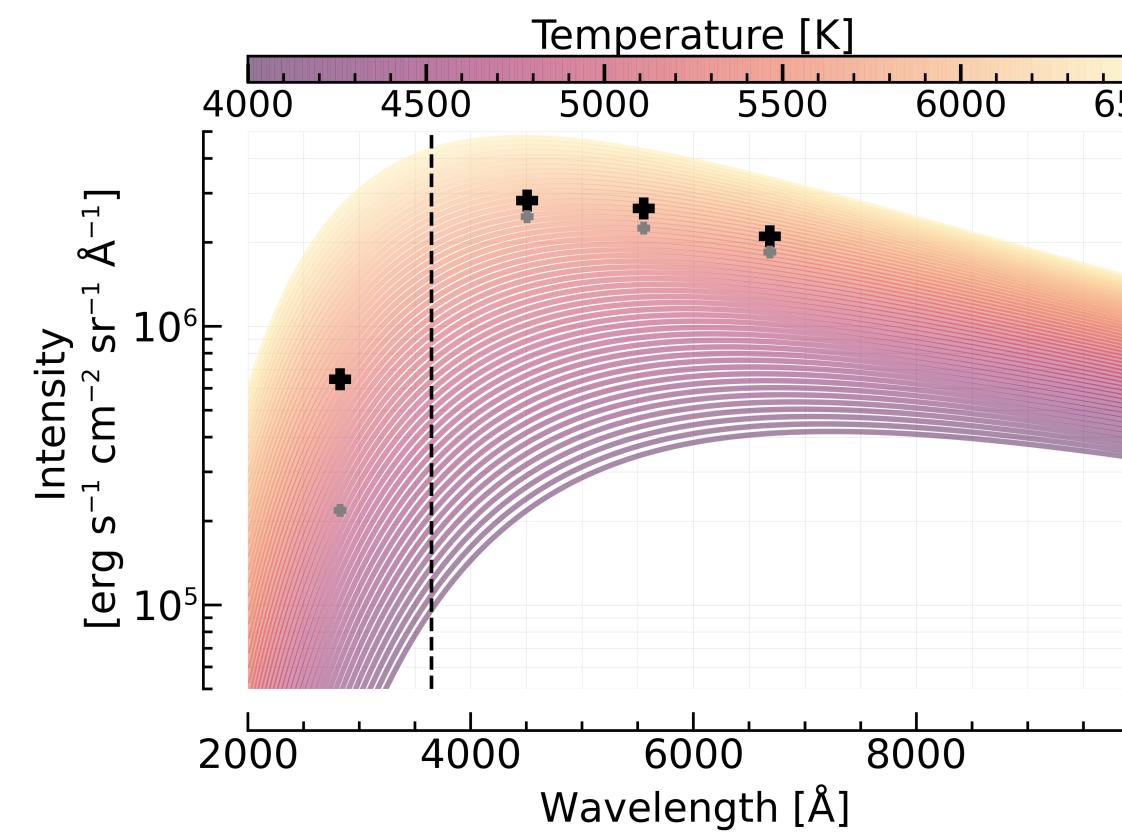


# WLF OBSERVATIONS: COMBINING NUV + OPTICAL

- ▶ Optical continuum enhancements and NUV continuum enhancements are largely co-spatial, but there are NUV sources with no optical counterpart.
- ▶ NUV contrast is very much larger than optical (it's hard to outshine the optical photosphere).



# WLF OBSERVATIONS: COMBINING NUV + OPTICAL



Source A

Source B

- ▶ The shape of the spectra varies in different sources.
- ▶ Assuming that NUV is recombination Balmer radiation means we can estimate the Paschen emission for different temperatures.
- ▶ There is a large excess of optical emission – Paschen recombination can't explain it all.

**Our models don't predict this – are we missing heating deep in the atmosphere?**

