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# A REVIEW OF THE MOLECULAR DIFFUSIVITIES OF H<sub>2</sub>O, CO<sub>2</sub>, CH<sub>4</sub>, CO, O<sub>3</sub>, SO<sub>2</sub>, NH<sub>3</sub>, N<sub>2</sub>O, NO, AND NO<sub>2</sub> IN AIR, O<sub>2</sub> AND N<sub>2</sub> NEAR STP

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Abstract— Accurate knowledge of the coefficients of molecular diffusivity (D) of trace gases has important application to global change, atmospheric chemistry, combustion science, studies of indoor air pollution and atmosphere—biosphere interactions. This study reviews and re-analyzes the historical data and some modeling results for D for  $H_2O$ ,  $CO_2$ ,  $CH_4$ , CO,  $O_3$ ,  $SO_2$ ,  $NH_3$ ,  $N_2O$ , NO, and  $NO_2$  in air,  $N_2$  and  $O_2$ . Because CO and  $N_2$  are isosteric molecules the self-diffusion coefficient for  $N_2$  and D for  $N_2$  in  $O_2$  are also evaluated as an independent check on the determination of D for CO. Also included is the coefficient of self-diffusion of  $O_2$  ( $D_{O_2,O_2}$ ) because an earlier model parameterized D for  $O_3$  in  $O_2$  in terms of  $D_{O_2,O_2}$ . Because D is a function of temperature and pressure, all modeling results and data are corrected to 1 atm pressure and then used with a one- and two-parameter regression model to determine optimal values for the temperature exponent and the value of D at  $0^{\circ}C$ . For most experimental data the loess regression technique is used to identify and eliminate outliers. To simplify comparisons and usage all results and recommendations are standardized to one regression curve. For some gases like  $H_2O$  and  $CO_2$  present results are similar to some (but not all) values cited in current sources. However, the coefficient of diffusivity in air for some very important gases, such as  $O_3$ , NO and  $NO_2$ , have never been measured and therefore, remain relatively uncertain. Based on present recommendations it is suggested that the ratio of the diffusivities of water vapor and ozone should be 1.51 rather than the often cited 1.65.  $\bigcirc$  1998 Elsevier Science Ltd. All rights reserved.

Key word index: Gaseous coefficients of diffusivity, gaseous binary diffusion.

#### INTRODUCTION

Molecular diffusion occurs as different species of a mixture (of gases or liquids, etc.) move under the influence of concentration gradients and it plays a key role in a variety of atmospheric and biospheric sciences. For example, because gases must diffuse through the boundary layer and stomata of plant leaves diffusion is thought to be the dominant exchange process during plant transpiration, photosynthesis, and uptake of ozone and other pollutants. Diffusion also affects the rate of movement of gases in soils and is central to much of atmospheric chemistry. In general the rate of diffusion (or flux) is determined by the product of a concentration gradient (a property of an individual species) and a coefficient of diffusivity (a property of the mixture). In turn, the coefficient of diffusivity, D(T, p), is known to be a function of temperature (T) and pressure (p), (e.g. Reid et al., 1987). The purpose of this study is twofold. The first is to review the experimental data and some of the models used for predicting the coefficients of diffusivity for 10 specific gases (H2O, CO2, CH4, CO, O3, SO2, NH3, N2O, NO and NO2) in air, N2 and O2 at temperatures and pressures typically encountered near the

earth's surface (i.e. temperatures approximately between  $-50^{\circ}$ C and  $100^{\circ}$ C and pressure near 1 atm or less). The second purpose is to derive and recommend estimates for D(T, p) for use in applications at conditions near STP (standard temperature and pressure). (For convenience D(T, p) at STP will be denoted by D(0, 1), because STP is defined as  $T = 0^{\circ}$ C and p = 1 atmosphere or 101.325 kPa.)

Although coefficients of diffusivity have been experimentally determined for over a century (e.g. Loschmidt, 1870a, b; Winklemann, 1884a, b) and a comprehensive review of D(T, p) for many gases is available (i.e. Marrero and Mason, 1972), there is still considerable variation in the values quoted by different researchers and references. The coefficient of diffusivity of water vapor in air  $(D_w)$  can serve as an example for this, although the coefficient of diffusivity of carbon dioxide in air would also serve nearly as well. Marrero and Mason (1972) suggest that the most precise determination for  $D_{\rm w}(0,1)$  is 0.209 cm<sup>2</sup> s<sup>-1</sup>. Monteith and Unsworth (1990) suggest that  $D_{\rm w}(0, 1)$ is 0.212 cm<sup>2</sup> s<sup>-1</sup>. Laisk et al. (1989), using the model of Chen and Othmer (1962), estimated  $D_{\rm w}(0,1)$  to be 0.195 cm<sup>2</sup> s<sup>-1</sup>. Extrapolating the value given by Hicks et al. (1985) to STP yields approximately 1112 W. J. MASSMAN

 $D_{\rm w}(0,1)=0.194~{\rm cm^2\,s^{-1}}$ . Tabulations by Andrussow (1969) suggest that  $0.161\leqslant D_{\rm w}(0,1)\leqslant 0.219~{\rm cm^2\,s^{-1}}$  and finally List (1949), in The Smithsonian Meteorological Tables, quotes a value of  $0.226~{\rm cm^2\,s^{-1}}$  from Schirmer (1938) for  $D_{\rm w}(0,1)$ . In other words there is an approximate  $\pm 15\%$  uncertainty in  $D_{\rm w}(0,1)$  that results solely from which source is being quoted. Consequently, this study is motivated by a desire to reduce this uncertainty in D(T,1) for the gases examined herein. As will be shown later, there is an appropriate statistical technique and enough historical data for many gases to estimate D(T,1) more accurately than +10%.

A further source of uncertainty in D(T, p) results from a misunderstanding of Graham's law. This occurs when estimating the diffusivity of a particular gas  $(D_u)$  in air (or other gas) by scaling a known diffusivity of another gas  $(D_k)$  by the square root of the ratio of molecular masses; i.e.  $D_u = D_k \sqrt{M_k/M_u}$  where  $M_k$  and  $M_u$  are the appropriately corresponding molecular masses. Graham's law applies to the ratio of the fluxes of the interdiffusing gases, not to their respective diffusivities in another gas (e.g. Mason and Evans (1969) or Mason (1971)).

The present study is not the first attempt to estimate coefficients of diffusivity by assembling and analyzing the historical data set. Marrero and Mason (1972) (henceforth denoted as MM72) provide a comprehensive review and critical evaluation of binary diffusivity coefficients. Yet, in spite of the detailed analysis provided by MM72, there remains a definite need to update the state of knowledge of diffusivity coefficients. Consequently, the present study differs from MM72 in several ways. First, the present analysis includes (a) results from more recent theoretical studies (viz., Bzowski et al., 1990a, b) and (b) several experimental determinations of D(T, p) subsequent to MM72 (e.g. Cowie and Watts, 1971; Barr and Watts, 1972; Prichard and Currie, 1982). Second, the present study emphasizes D(T, p) for  $O_3$ ,  $SO_2$ ,  $NH_3$ ,  $N_2O$ , NO, NO<sub>2</sub> in air,  $N_2$  and  $O_2$  more than do MM72. Third, the present study focuses only on temperatures between about  $-20^{\circ}$  and  $100^{\circ}$ C; whereas, MM72 included diffusivity data at much higher temperatures. The temperature range is restricted because the results of MM72's analysis clearly indicate a possible bias and/or considerably more variability in the measured values of  $D_{\rm w}(T,p)$  above 100°C than below 100°C. Therefore, to avoid this potential bias this study limits the diffusivity data to temperatures below 100°C. (Fortunately, this restriction has virtually no effect on any other data sets associated with the other gases examined herein.) Fourth, to facilitate comparisons between gases, the present analysis uses the same exponent to describe the temperature dependency. Fifth and lastly, the present study differs from MM72 in the method used to eliminate outliers from the data set. This study uses a local regression model technique known as loess (Cleveland et al., 1992; Cleveland, 1993) to separate outliers from the more congruous data; whereas, MM72 used their knowledge of theory and experimental procedures to evaluate the data quality.

More recent than MM72, Yaws (1995) also provides an analysis and compilation of binary diffusion coefficients in air. Like MM72, Yaws (1995) also uses a broad temperature range (200–1000 K). However, unlike MM72 and the present study, Yaws (1995) focuses on organic compounds. Consequently, the only overlap between the present study and Yaws (1995) are  $\rm CO_2$ ,  $\rm CH_4$  and  $\rm CO$ .

The next section of this study describes the data sources, the loess technique and general methods of analysis. In a following section results are tabulated and compared and where applicable, comparisons between Yaws (1995), MM72's results and loess results are also made.

#### DATA SOURCES, MODELS AND METHODS OF ANALYSIS

Near STP the binary diffusivity as a function of temperature and pressure can be expressed as

$$D(T, p) = D(0, 1)(p_0/p)(T/T_0)^{\alpha}$$
 (1)

where  $T_0 = 273.15 \text{ K}$  (= 0°C), T is temperature expressed in K, and the exponent  $\alpha$  is limited by theory to a value between 1.5 and 2.0 (e.g. Dunlop and Bignell, 1997). This formulation is valid providing that the pressure (p) is not too near the critical pressure for the gas mixture. (All critical pressures for gases examined herein exceed 30 atmosphere's pressure.) If the critical pressure is exceeded then the simple inverse relationship between D and p is not valid (Reid et al., 1987). Aside from this exception, equation (1), in general, is an extremely good approximation to the more exact formulation based on the kinetic theory of gases and the Chapman-Enskog theory (e.g. Reid et al., 1987). For this reason, equation (1) is also the basis for this study. However, this study focuses on the variation of D(T, p) with temperature, not with pressure. Consequently, all data used herein for analysis, but not originally obtained at  $p = p_0$ , are adjusted to  $p_0$  and all modeling results will assume  $p = p_0$ . Furthermore, all data used in this study are limited to pressures less than 25 atm which is well below the critical pressure for any gas discussed herein. In essence, equation (1), is used in this study as a regression model for D(T, 1) for the experimental data sets. Furthermore, it is used in both a twoparameter mode (when both D(0, 1) and  $\alpha$  are "fitted" parameters) and a one-parameter mode (when only D(0, 1) is "fit" to the data and  $\alpha$  is fixed at a value of 1.81). (The justification for this particular choice of  $\alpha$  is discussed below.) In addition to fitting the experimentally derived data, equation (1) is also used in both one- and two-parameter modes for some of the model-predicted temperature curves. Comparing the one- and two-parameter regressions is one method of assessing the impact that constraining  $\alpha$  to 1.81 has on the quality of the model fit. For now, suffice it to say that constraining  $\alpha$  to 1.81 causes no serious loss of information nor compromises the results in any significant way.

Using a fixed value of  $\alpha$  has the advantage of simplifying the comparisons between different gases and different models. Furthermore, the value of 1.81 represents a reasonable compromise for both observational and modeling purposes. For example,  $\alpha = 1.81$  is the basis of the model of Chen and Othmer (1962) and  $\alpha = 1.81$  (  $\pm 0.01$ ) results from two independent previous analyses for  $D_{\rm w}(T,1)$  (Schrimer (1938) as cited by List (1949) and Vasilevskaya (1961)) and one previous study of D(T, 1) for NH<sub>3</sub> in air (Wintergerst, 1930). Nevertheless, α could have been fixed at other values as indicated by other models. For example, Fuller et al. (1966) developed a model that indicates  $\alpha = 1.75$  is valid for a large number of binary diffusivity coefficients. The present study focuses on the model of Chen and Othmer (1962) and the concomitant value of 1.81 for  $\alpha$ , rather than the model of Fuller et al. (1966) because the Chen and Othmer (1962) model was developed in terms of critical temperatures and critical volumes, which are physical quantities, rather than in terms of elemental diffusion volumes which are determined from measured diffusion coefficients using least squares regression techniques. Consequently, for the present purposes the model of Chen and Othmer (1962) has greater predictive capabilities than does the model of Fuller et al. (1966).

### Data sources

As much as possible only original sources were used for reconstructing the historical data set or for obtaining model parameters for use with model estimations of the diffusion coefficients. Furthermore, all data sources were checked for multiple references to the same diffusivity values to ensure that no data from any single experiment were duplicated during the analysis. However, experimental determinations of D(T, 1) were averaged if, for any given experiment, D(T, 1) was evaluated at several temperatures that were within  $\pm 0.5$ °C of the mean temperature of the measurements. For convenience, all diffusivity values used throughout this study are expressed in cm<sup>2</sup> s<sup>-1</sup>.

The major source for references to data is the bibliographic list of gaseous diffusion coefficients for binary mixtures given by MM72 and Gordon (1977) with secondary sources supplied by Bzowski *et al.* (1990a), Reid *et al.* (1987), Andrussow (1969), Liley (1962) and the author's continuing search of the literature.

#### Absolute uncertainty

Absolute uncertainty plays an important role in this study when evaluating the quality of data and the goodness of the model regression. After any regression has been performed, the maximum percentage difference between the values predicted with the regression curve and the data values is then determined. Herein this maximum difference is termed the absolute uncertainty and it is used to assess the inherent variability of any given data set.

For this study it is possible a priori to estimate a lower bound on the absolute uncertainty. MM72 note that D(T,1) is a relatively weak function of composition; that is D(T,1) is weakly influenced by the mole fraction of the two interdiffusing gases. In most cases this effect is usually less than about 4% (MM72). Unlike MM72, the present study does not attempt to correct the data for composition dependence; rather it is assumed that the lower bound on the absolute uncertainty is  $\pm 4\%$ . Because in general any data set is likely to possess more than one source of variability, it is probably unreasonable to assume that this level of absolute uncertainty can be achieved in the present study.

The loess technique for analysis of experimentally determined data

Because Cleveland *et al.* (1992) and Cleveland (1993) provide a detailed account of the loess technique, it is sufficient for the present purposes to provide only a brief summary of how it was employed for evaluating D(0, 1). Loess, as implemented in this study, is a relatively simple iterative technique that provides a method for fitting regression functions to data and for detecting potential outliers in the data set that may need to be eliminated because they can skew the parameter values and bias the results. Except to weight the outliers with a value of 0, this study does not include a local weighting function for fitting the data.

As implemented in this study the loess technique consists of two iterated steps. First, all experimentally determined diffusivity values (N = total number ofvalues) are adjusted for  $p = p_0$  and then fit to equation (1) with a least squares minimization technique to determine the optimal values of D(0, 1) and  $\alpha$ . Second the residuals  $(e_i: i = 1, N)$  are calculated and the median value  $(e_{med})$  of the  $e_i$ 's is determined. Then all diffusivity values with a residual (e) at least three times the median value ( $e \ge 3e_{\text{med}}$ ) are removed from the data set and the process is repeated until one of two conditions is met: either (1) there are no remaining data points with residuals greater than 3 times the median value or (2) the absolute uncertainty of the data relative to the regression curve is not changed significantly by the removal of any more data points. (Note that  $e_{\text{med}}$  can change with each iteration.) Once the most consistent (congruous) data set has been identified (i.e. all probable outliers have been removed), then a one-parameter fit is performed for D(0, 1) (with  $\alpha = 1.81$ ). Finally, an estimate of the absolute uncertainty is determined for the one-parameter regression. Both D(0, 1) and the absolute uncertainty for the one-parameter regressions are tabulated in this study. As employed herein the absolute uncertainty does not include the outliers, but pertains only to the congruous data set.

As will be shown later this method works well for diffusivities of H2O and CO2 in air and N2 because these are relatively large data sets with many easily identified outliers. Nevertheless, results at each iteration are examined for all gases before the next iteration is performed. In general the loess technique is quite amenable to computer implementation and therefore is relatively objective, but it should not be used as a substitute for careful examination of any data set. In some cases the diffusivity data sets inherently possess large variability or they consist of only a few experimental determinations. In these and other cases it becomes difficult to determine which data points are useful for further analysis and which are outliers. Consequently, the analysis for each diffusivity data set includes a brief discussion of its unique features and whatever adaptations were necessary for implementing the loess technique.

#### Models

The models of Bzowski *et al.* (1990a, b), Chen and Othmer (1962) and the Chapman–Enskog theory with the Lennard–Jones 6–12 potential (e.g. Reid *et al.*, 1987) are included and compared with the data analyses and with each other in this study. (Henceforth, Bzowski *et al.* (1990a, b) will be denoted as Bea90, Chen and Othmer (1962) as CO62 and the Chapman–Enskog/Lennard–Jones model as LJ.)

Bea90 provides calculated values of D(T, 1) for  $CO_2$ ,  $CH_4$ , CO,  $N_2O$  and NO in  $N_2$  and  $O_2$  which are used for one- and two-parameter fits with equation (1). However, the fitting procedures are done only once and the loess procedure is not required, because the model is (by necessity) already quite similar to equation (1).

The model of CO62 is as follows:

$$D(T, 1) =$$

$$\frac{2.616\sqrt{1/M_i + 1/M_j}(T/T_0)^{1.81}}{(Tc_iTc_j/10,000)^{0.1405}[(Vc^i/100)^{0.4} + (Vc_j/100)^{0.4}]^2} (2)$$

where, henceforth the subscripts "i" and "j" refer to the specific gases under consideration,  $M_i$  and  $M_j$  are their respective molecular masses (gm mol<sup>-1</sup>), T $c_i$  and T $c_j$  are their respective critical temperatures (K), and V $c_i$  and V $c_j$  are their respective critical volumes (cm<sup>3</sup> mol<sup>-1</sup>). The CO62 model is basically an empirical approximation to the more rigorous Chapman–Enskog theory, given next.

$$D(T, 1) = \frac{8.380\sqrt{1/M_i + 1/M_j}(T/T_0)^{3/2}}{\sigma_{ij}^2 \Omega_{ij}}$$
(3)

where  $\sigma_{ij}$  is a typical length scale characterizing the molecular size or range-of-force between molecules of the two gases during the diffusive process (Å) and  $\Omega_{ij}$  is the (dimensionless) reduced collision integral for diffusion which, in general, is a weak function of temperature. For binary diffusion,  $\sigma_{ij}$  is usually evaluated as the arithmetic average of the appropriate

 $\sigma$ -values for each gas (i.e.  $\sigma_{ij} = (\sigma_i + \sigma_j)/2$ ) (e.g. Reid et~al., 1987) and  $\Omega_{ij}$  is determined from a model of the potential function describing the potential field between molecules of the two gases during the diffusive process (e.g. Reid et~al., 1987; MM72). For convenience, this study assumes the Lennard–Jones 6–12 potential (e.g. Reid et~al., 1987; MM72) which yields the following form for the temperature dependence for  $\Omega_{ij}$ 

$$\Omega_{ii} = \Omega_{ii}(kT/\varepsilon_{ii}) \tag{4}$$

where k is Boltzmann's constant and  $\varepsilon_{ij}$  is the maximum energy of attraction between the two molecules.

In general  $\varepsilon_{ij}$  is approximated by the geometric mean of the individual  $\varepsilon$ 's, i.e.  $\varepsilon_{ij} = \sqrt{\varepsilon_i \varepsilon_j}$ , and the  $\varepsilon$ 's, like the  $\sigma$ 's above and the critical constants Tc and Vc earlier, have often been previously determined for many gas species and are often available from tabulations (e.g. Reid et al., 1987; Svehla, 1962; Yaws, 1972; Ambrose, 1992; Kee et al., 1986; Dandy, 1997). For this study Table 1 summarizes all values of M, Tc, Vc,  $\sigma$  and  $\varepsilon/k$  used with the CO62 and LJ models. Except for  $\sigma$  and  $\varepsilon/k$  for  $O_3$  and Tc and Vc for  $NO_2$  all values presented in Table 1 are well documented. The discussion of these four exceptions is deferred until later.

For this study  $\Omega_{ij}(kT/\varepsilon_{ij})$  is approximated by the (highly accurate) analytical expression of Neufeld *et al.* (1972). For comparison with other models and data results the LJ model is treated the same way as Bea90. That is, the LJ model is used to predict D(T, 1) at several different temperatures and then these model-generated diffusivities are used as input data for one-parameter and two-parameter regressions with equation (1).

#### Blanc's law

For all model calculations and some of the databased results it is necessary to estimate the coefficients of diffusivity for diffusion in air from the diffusivities appropriate to  $N_2$  and  $O_2$ . This study uses Blanc's law for combining coefficients of binary diffusion to estimate D(T, p) of a gas diffusing into a mixture of gases. Blanc's law (e.g. Reid *et al.*, 1987; MM72) applied to the conditions of the present study yields:

$$D_{i, \text{air}}(T, p) = \left[\chi_{N_2}/D_{i, N_2}(T, p) + \chi_{O_2}/D_{i, O_2}(T, p)\right]^{-1} (5)$$

where  $D_{i, \text{air}}$  is the diffusivity of gas "i" in air,  $D_{i, \text{N}_2}$  and  $D_{i, \text{O}_2}$  are the respective diffusivities of the same gas in N<sub>2</sub> and O<sub>2</sub>,  $\chi_{\text{N}_2}$  (=0.79) and  $\chi_{\text{O}_2}$  (=0.21) are the approximate respective mole fractions of N<sub>2</sub> and O<sub>2</sub> in the atmosphere.

Blanc's law was used in a few cases to estimate the diffusivity of water vapor and  $CO_2$  in air from experimental determinations of D(T, 1) in  $N_2$  and  $O_2$ , but only if the temperatures of the respective D(T, 1) values were within about  $0.5^{\circ}$ C. Furthermore, any diffusivity data appropriate to  $N_2$  and  $O_2$  combined using Blanc's law were not used again in any later

Vc $\varepsilon/k$ σ (Å)  $(gm \, mol^{-1})$  $(cm^3 mol^{-1})$ Gas (K) (K) 3.798  $N_2$ 28.01 126.2 89.49 71.4  $O_2$ 32.00 154.6 73.39 3.467 106.7 H<sub>2</sub>O 18.02 647.1 55.96 2.641 809.1  $CO_2$ 44.01 94.04 3.941 304 1 1952  $CH_4$ 16.04 190.5 99.01 3.758 148.6 CO 132.9 93.06 3.690 91.7 28.01  $SO_2$ 64.06 430.8 122.20 4.026a 363.0a  $3.875^{b}$ 208.4b O<sub>3</sub> 48.00 261.1 88.89 NH<sub>3</sub> 405.5 72.50 2.900 17.03 558 3 97.37  $N_2O$ 44.01 309.6 3.828 232.4 NO 30.01 180.0 57.71 3.492 116.7 NO<sub>2</sub> 431.0b 83.64<sup>b</sup> 3.765° 46.01 210.0°

Table 1. Parameter values used for CO62 and LJ models

Note. Critical temperature (Tc) and critical volume (Vc) used with CO62. Molecular size ( $\sigma$ ) and energy of attraction ( $\varepsilon/k$ ) used with LJ model. Molecular mass (M) used with both models. Unless noted otherwise M and Tc are taken from Ambrose (1992) and Vc is computed from critical density,  $\rho_c$ , which is also taken from Ambrose (1992); i.e.  $Vc = M/\rho_c$ . Unless noted otherwise  $\sigma$  and  $\varepsilon/k$  are taken from Reid et al. (1987).

Table 2. Coefficients of diffusivity (cm<sup>2</sup> s<sup>-1</sup>) of selected gases in air

Gas	This study	CO62ª	Bea90a	LJa	Other
H <sub>2</sub> O	0.2178	0.1934	NA	0.1837	$0.2090(T/T_0)^{2.072^{b}}$
$CO_2$	0.1381	0.1404	0.1391	0.1314	$0.2020 \exp(-\frac{0.3738}{T/T_0})(T/T_0)^{1.590^{b}}$
$CH_4$	0.1952	0.1909	0.1951	0.1862	$0.1859(T/T_0)^{1.747^{b}}$
CO	$0.1807^{a}$	0.1755	0.1794	0.1716	$0.1838(T/T_0)^{1.730^{b}}$
$SO_2$	0.1089°	0.1122	NA	0.1087	NI
$O_3$	NM	0.1444	NA	0.1301	NI
$NH_3$	$0.1978^{d}$	0.1908	NA	0.1866	NI
$N_2O$	0.1436	0.1382	0.1401	0.1318	NI
NO	NM	0.1988	0.1802	0.1729	NI
$NO_2$	IM	0.1361	NA	0.1349	0.1222 <sup>e</sup>

*Note*: Unless otherwise noted only D(0, 1) is given and  $D(T, 1) = D(0, 1)(T/T_0)^{1.81}$  is assumed. NA = not available. NM = never measured. IM = impossible to measure. NI = not included herein.

analysis of D(T, 1) in  $N_2$  and  $O_2$ . This is done to avoid introducing any unnecessary biases into the separate analyses for air,  $N_2$  and  $O_2$ .

## RESULTS

All results are summarized in Tables 2–7. Tables 2–4 list the results of the loess and regression analysis for the diffusivities of the 10 selected gases in air (Table 2),  $N_2$  (Table 3) and  $O_2$  (Table 4). Also listed in these three tables are the respective diffusivi-

ties predicted by the Bea90 model, the CO62 model, the LJ model, the regression results of MM72 and a few miscellaneous data/modeling results. Tables 5–7 present the absolute uncertainties associated with the data and loess regression, as well as the percent difference between the data-based estimates of diffusivity and the (Bea90, CO62 and LJ) model predictions for air (Table 5),  $N_2$  (Table 6) and  $O_2$  (Table 7). The rest of this section includes a brief description and explanation of the entries in these tables and brief discussions of the data sets and figures.

<sup>&</sup>lt;sup>a</sup> From Monchick and Mason (1961).

<sup>&</sup>lt;sup>b</sup> See discussion in text.

<sup>&</sup>lt;sup>c</sup> From Svehla and Brokaw (1966).

<sup>&</sup>lt;sup>a</sup> From Blanc's law and D(0, 1) for gas in  $N_2$  and in  $O_2$ .

<sup>&</sup>lt;sup>b</sup> Marrero and Mason (1972).

<sup>&</sup>lt;sup>c</sup> Based on combining two data points: one for air and one for N<sub>2</sub>.

<sup>&</sup>lt;sup>d</sup> Wintergerst (1930).

 $<sup>^{\</sup>rm c}$  Blanc's law and results from Fan and Mason (1962) for NO<sub>2</sub> and O<sub>2</sub> and (indirectly) from Sherwood and Chambers (1937) for NO<sub>2</sub> and N<sub>2</sub>.

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Table 3. Coefficients of diffusivity (cm2 s-1) of selected gases in N2

Gas	This study	CO62	Bea90	LJ	Other
H <sub>2</sub> O	0.2190	0.1922	NA	0.1841	$0.2100(T/T_0)^{2.072^a}$
$CO_2$	0.1429	0.1401	0.1392	0.1326	$0.21072 \exp(-\frac{0.1458}{T/T_0})(T/T_0)^{1.570^{a}}$ $0.1836(T/T_0)^{1.750^{a}}$ $30.433(e1/e2)(T/T_0)^{1.576^{a,b}}$
CH <sub>4</sub>	0.1892	0.1899	0.1943	0.1858	$0.1836(T/T_0)^{1.750^a}$
CO	0.1804	0.1749	0.1789	0.1713	$30.433(e1/e2)(T/T_0)^{1.576^{a,b}}$
$SO_2$	0.1089°	0.1122	NA	0.1094	NI
$O_3$	NM	0.1441	NA	0.1306	NI
$NH_3$	0.1978	0.1896	NA	0.1869	NI
$N_2O$	NM	0.1379	0.1402	0.1322	NI
NO	0.1809	0.1979	0.1798	0.1727	NI
$NO_2$	IM	0.1359	NA	0.1352	$0.1210^{d}$
-					0.1460°

Note: Unless otherwise noted only D(0, 1) is given and  $D(T, 1) = D(0, 1)(T/T_0)^{1.81}$  is assumed. NA = not available. NM = never measured. IM = impossible to measure. NI = not included herein.

Table 4. Coefficients of diffusivity (cm<sup>2</sup> s<sup>-1</sup>) of selected gases in O<sub>2</sub>

Gas	This study	CO62	Bea90	LJ	Other
H <sub>2</sub> O	0.2457a	0.1984	NA	0.1821	$0.2110(T/T_0)^{2.072^{b}}$
$CO_2$	0.1402	0.1414	0.1386	0.1300	$0.17383 \exp(-\frac{0.2244}{T/T_0})(T/T_0)^{1.661^{b}}$
$\mathrm{CH_4}$	0.1942	0.1946	0.1978	0.1877	$0.2265 \exp(-\frac{0.1618}{T/T_0})(T/T_0)^{1.695^{b}}$
CO	0.1817	0.1777	0.1813	0.1729	$0.1793(T/T_0)^{1.724^{b}}$
$SO_2$	NM	0.1121	NA	0.1059	NI
$O_3$	NM	0.1454	NA	0.1285	$\frac{5}{6}D_{O_2,O_2}(T,1)^{d}$
					$\frac{0.2072 (T/T_0)^{3/2}}{1 + 0.622 (T/T_0)}$ e
$NH_3$	0.1992	0.1951	NA	0.1857	NI
$N_2O$	NM	0.1391	0.1396	0.1301	NI
NO	NM	0.2024	0.1815	0.1739	0.1749°
$NO_2$	IM	0.1369	NA	0.1334	0.1271°

*Note*: Unless otherwise noted only D(0, 1) is given and  $D(T, 1) = D(0, 1)(T/T_0)^{1.81}$  is assumed. NA = not available. NM = never measured. IM = impossible to measure. NI = not included herein.

Table 5. Absolute error associated with D(T, 1) from data analysis and difference between model predictions of D(T, 1) and data-derived estimates

Gas pair	Absolute error (N)	CO62	Bea90	LJ
H <sub>2</sub> O/air CO <sub>2</sub> /air CH <sub>4</sub> /air CO/air SO <sub>2</sub> /air O <sub>3</sub> /air NH <sub>3</sub> /air N <sub>2</sub> O/air NO/air NO <sub>2</sub> /air	$ \pm 7\% (N = 41)  \pm 5\% (N = 56)  \pm 5\% (N = 11)  \pm 5\% (N = 0)^{a}  I(N = 2)  I(N = 0)  \pm 9\% (N = I)  I(N = 4)^{b}  I(N = 0)  I(N = 0)$	- 11.2% + 1.7% - 2.2% - 2.9% + 3.0% I - 3.5% - 3.8% I	NA + 0.7% - 0.1% - 0.7% NA NA NA - 2.4% I	- 15.6% - 4.9% - 4.6% - 5.0% - 0.2% I - 5.7% - 8.2% I

*Note*: N = number of data points. NA = not applicable. I = indeterminate.

<sup>&</sup>lt;sup>a</sup> Marrero and Mason (1972). <sup>b</sup> e1 =  $\exp(\frac{0.135}{T/T_0} - \frac{0.051247}{(T/T_0)^2})$ , e2 =  $(-13.2616 + \ln(T/T_0))^2$ <sup>c</sup> Based on combining two data points: one for air and one for N<sub>2</sub>.

<sup>&</sup>lt;sup>d</sup> Adapted from Sherwood and Chambers (1937).

eAdapted from Sviridenko et al. (1973). See Text.

<sup>&</sup>lt;sup>a</sup> From a single determination given by Walker and Westenberg (1960).

<sup>&</sup>lt;sup>b</sup> Marrero and Mason (1972).

<sup>&</sup>lt;sup>c</sup> From Fan and Mason (1962).

<sup>&</sup>lt;sup>d</sup> From Hirschfelder et al. (1953).

e From Campbell (1965).

<sup>&</sup>lt;sup>a</sup> From Blanc's law. Absolute uncertainty assumed to be the same as that determined for diffusivity in  $N_2$  and  $O_2$ .

<sup>&</sup>lt;sup>b</sup> Data from only one experiment showed very little variability; therefore, absolute uncertainty could not be evaluated.

LJ Gas pair Absolute error (N) CO62 Bea90 -2.2%-15.9% $H_2O/N_2$  $\pm 6\% (N = 10)$ NA - 2.6% +6% (N = 33) $CO_2/N_2$ -2.0%-7.2% $\pm 5\% (N = 9)$  $CH_4/N_2$ +2.7%+0.4%-1.8% $\pm 5\% (N = 12)$ -0.8% $CO/N_2$ -30%-50% $SO_2/N_2$ I(N = 2)+4.4%NA +0.5%I(N=0)NA  $O_3/N_2$ I I +5% (N=5)- 4.1% - 5.5% NH<sub>3</sub>/N<sub>2</sub> NA  $N_2O/N_2$ I(N=0)I I I  $NO/N_2$  $\hat{I(N} = 2)$ + 9.4% - 0.6% -4.5% $NO_2/N_2$ I(N = 0)Ι NA Ι

Table 6. Absolute error associated with D(T, 1) from data analysis and difference between model predictions of D(T, 1) and data-derived estimates

Note: N = number of data points. NA = not applicable. I = indeterminate.

Table 7. Absolute error associated with D(T, 1) from data analysis and difference between model predictions of D(T, 1) and data-derived estimates

Gas pair	Absolute error (N)	CO62	Bea90	LJ
H <sub>2</sub> O/O <sub>2</sub> CO <sub>2</sub> /O <sub>2</sub> CH <sub>4</sub> /O <sub>2</sub> CO/O <sub>2</sub> SO <sub>2</sub> /O <sub>2</sub> O <sub>3</sub> /O <sub>2</sub> NH <sub>3</sub> /O <sub>2</sub> N <sub>2</sub> O/O <sub>2</sub> NO <sub>2</sub> /O <sub>2</sub>	$I(N = 1)$ $\pm 5\% (N = 12)$ $I(N = 1)$ $\pm 5\% (N = 6)$ $I(N = 0)$ $I(N = 0)$ $I(N = 2)$ $I(N = 0)$ $I(N = 0)$ $I(N = 0)$	- 19.3% + 0.9% + 0.2% - 2.2% I I - 2.1% I	NA - 1.1% + 1.9% - 0.2% NA NA NA I I I	- 25.9% - 7.3% - 3.3% - 4.8% I I - 6.8% I

*Note*: N = number of data points. NA = not applicable. I = indeterminate.

### $H_2O$

Figure 1 summarizes the results of this study for the diffusivity of water vapor in air. The final recommended value for  $D_{\rm w}(T,1)$  shown in this figure is  $0.2178(T/T_0)^{1.81}$  with an absolute uncertainty of  $\pm$  7% and a regression  $R^2 = 0.972$ . ( $R^2 = 1 - [Sum]$ of squared regression residuals]/[Variance of data relative to the mean of the data].) Before applying the loess technique to  $D_{\rm w}(T,1)$ , the two-parameter regression yielded  $D(T, 1) = 0.2043(T/T_0)^{2.2172}$  with an  $R^2 = 0.75$  and an absolute uncertainty of about  $\pm$  25%. After the first application of the loess technique the two-parameter fit yielded D(T, 1) = $0.2191(T/T_0)^{1.8464}$  with an  $R^2 = 0.91$  and an absolute uncertainty of about  $\pm 12\%$ . The second application of the loess technique yielded some improvements over the first and the third and final application resulted in  $D(T, 1) = 0.2188(T/T_0)^{1.7768}$ with an  $R^2 = 0.972$ . The subsequent one-parameter regression yielded  $D(T, 1) = 0.2178(T/T_0)^{1.81}$  with an  $R^2 = 0.972$  and an absolute uncertainty of about  $\pm$  7%. The standard error of the estimate for  $D_{\rm w}(0, 1)$ with the one-parameter regression is  $\pm 0.001$ .

In all 17 data points of a total of 58 determinations from 28 experiments were eliminated on the basis of the loess technique. But, the data identified as outliers

were not randomly distributed throughout the full data set. Rather the outliers encompassed 7 of the 27 different experiments for  $D_{\rm w}(T,1)$ . The fact that these data points are associated with specific experiments suggests (at least to the author) that there were problems with the experimental techniques or their implementation as employed by the researchers. All data from Summerhays (1930), Houghton (1933), Ackermann (1934), Klibanova *et al.* (1942), Bose and Chakraborty (1955–56), Richardson (1959) and Petit (1965) were eliminated for the final regression as indicated by the loess technique.

Figure 2 summarizes the results of this study for the diffusivity of water vapor in N<sub>2</sub>. For comparison, Fig. 2 also includes some of the measured diffusivities for water vapor in O2 (which were not used in the analysis). Although these diffusivities have been measured far fewer times than  $D_{w}(T, p)$  (Fig. 1), the results of the loess analysis and the subsequent one- and two-parameter regressions are very similar to  $D_{\rm w}(T,1)$ . The one independent determination of water vapor in O2 (Walker and Westenberg, 1960), when used to estimate D(0, 1) for water vapor in  $O_2$ (equation (1) with  $\alpha = 1.81$ ),  $D(0, 1) = 0.2457 \text{ cm}^2 \text{ s}^{-1}$ . This value is sufficiently different from theoretical predictions (Table 4) that it is

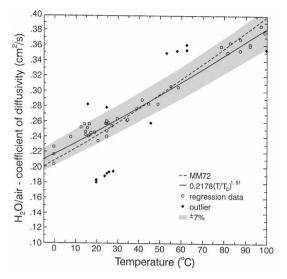


Fig. 1. Observed and modeled diffusivities for water vapor in air. Data sources are: Guglielmo (1881, 1882) as tabulated in Andrussow (1969); Winkelmann (1884a, b, 1888, 1889) including corrections as discussed in Trautz and Müller (1935a) and Mason and Monchick (1962); Houdaille (1897), Brown and Escombe (1900); Mache (1910); LeBlanc and Wupperman (1916); Summerhays (1930); Houghton (1933); Ackermann (1934), data not shown because values exceed 0.4 cm<sup>2</sup> s<sup>-1</sup>; Gilliland (1934); Schirmer (1938) as discussed in List (1949); Klibanova et al. (1942); Brookfield et al. (1947); Schwertz and Brow (1951) with Blanc's law; Kimpton and Wall (1952); Rossié (1953); Lee and Wilke (1954); Bose and Chakraborty (1955-56) including some estimates from Blanc's law; Narsimhan (1955–1956); Nelson (1956); Richardson (1959); Petit (1965); Andrussow (1969) quoting an unidentified source no. 156; Wilhelm and Battino (1972).

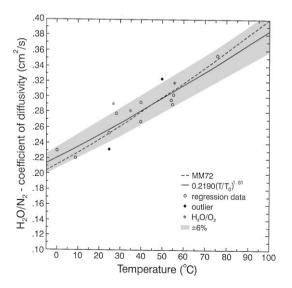


Fig. 2. Observed and modeled diffusivities for water vapor in  $N_2$ . Data are from Hippenmeyer (1949), Crider (1956), Nelson (1956), and O'Connell *et al.* (1969). Included but not used in the analysis are some observed diffusivities of water vapor in  $O_2$  for which the data sources are Schwertz and Brow (1951) and Walker and Westenberg (1960). Not shown here are the  $O_2$  data from Bose and Chakraborty (1955–1956) because they were used with Blanc's law to estimate D(T,1) for water vapor in air and subsequently eliminated by the loess analysis.

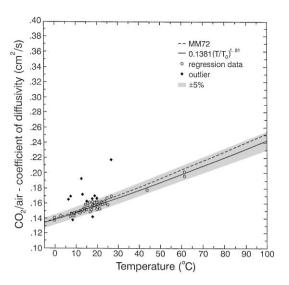


Fig. 3. Observed and modeled diffusivities for CO<sub>2</sub> in air. Data sources are: Loschmidt (1870a, b); von Obermayer (1880, 1882a, b, 1883, 1887), including corrections to earlier results discussed by von Obermayer (1887); Waitz (1882a, b); Brown and Escombe (1900); Buckingham (1904); Coward and Georgeson (1937); Klibanova *et al.* (1942); Waldmann (1944, 1947) with Blanc's law; Andrew (1955); Gwertsiteli *et al.* (1961); Suemin and Ivakin (1961); Holsen and Strunk (1964); Barr and Watts (1972); Pritchard and Currie (1982); Robjohns and Dunlop (1982) with Blanc's law.

probably an outlier. This notion is further supported by the comparisons in Fig. 2 and the results of MM72 who suggest that the diffusivity of water vapor in  $N_2$  and  $O_2$  are nearly identical to  $D_w(T, 1)$ . Consequently, present results suggest that the diffusivity of water vapor in  $O_2$  is relatively uncertain and further experimental determinations are needed.

Also shown in Figs 1 and 2 are the respective regression models of MM72. As can be seen, the results from MM72 and the present study are in good agreement. In fact, it may be said that the two models agree to within the absolute uncertainty of the (congruous) data. The present results for  $D_{\rm w}(T,1)$  are also in good agreement with Vasilevskaya (1973), who found, for temperatures exceeding  $100^{\circ}{\rm C}$ , that  $D_{\rm w}(T,1)=0.2106(T/T_0)^{1.82}$ .

 $CO_2$ 

Figure 3 shows the results of the analysis for  $CO_2$  and air. Unlike the  $D_w(T, 1)$  data, the  $CO_2$  shows very little inherent variability, so that the loess iterations were terminated after two applications and from a total of 67 data points only 11 were excluded from the final regression. But, as with  $D_w(T, 1)$  the excluded data points were not randomly distributed; rather they basically comprised the data of Brown and Escombe (1900), Buckingham (1904), Andrew (1955) and Kosov (1957). The final two-parameter regression yielded  $D(T, 1) = 0.1369(T/T_0)^{1.9206}$  with an  $R^2 = 0.958$  and the one-parameter regression yielded

D(0, 1) = 0.1381 ( $\pm 0.0005$ ) with an  $R^2 = 0.952$ . (Note that the standard error of the estimate is enclosed in parentheses in the last sentence.) Figure 3 also indicates good agreement between the present analysis and the results of MM72. Although not shown, the model of Yaws (1995) is in agreement with the present result to within the absolute uncertainty of the loess analysis.

The analysis of the data for CO<sub>2</sub> and N<sub>2</sub> did not follow the pattern outlined above for the three previous gas combinations. The first regression produced only one data point that might be considered an outlier by the loess technique. In other words the data set which is comprised of 46 data points from 23 experiments, was extremely consistent. The two-parameter regression gave the following result: D(T, 1) = $0.1347(T/T_0)^{2.2240}$  with an  $R^2 = 0.912$ . However, the exponent ( $\alpha$ ) is somewhat too high since kinetic theory and virtually all model and data analyses suggest that  $1.5 \le \alpha \le 2$ . This suggests that the regression may be slightly skewed. A careful examination of the data did suggest that at relatively low temperatures the data of Schäfer et al. (1951) and Ivakin and Suetin (1964a) slightly underestimated the actual value for D(T, 1)while at relatively high temperatures the data of Nagata and Hasegawa (1970) overestimated the actual value of D(T, 1). This combination of biases is not random and tends to bias the results toward a higher

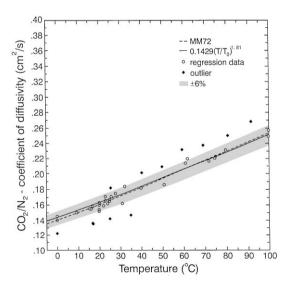


Fig. 4. Observed and modeled diffusivities for CO<sub>2</sub> in N<sub>2</sub>. Data sources are: Boardman and Wild (1937); Wicke and Kallenbach (1941); Boyd et al. (1951); Schäfer et al. (1951); Westenberg and Walker (1957); Walker and Westenberg (1958); Walker (1958); Walker et al. (1960); Bohemen and Purnell (1961); Giddings and Seager (1962); Ivakin and Suetin (1964b); Weissmann (1964); Vyshenskaya and Kosov (1965); Kestin et al. (1966); Pakurar and Ferron (1966); Coates and Mian (1967); Mian (1967); Ellis and Holsen (1969); Schneider and Schäfer (1969); Lannus (1970); Nagata and Hasegawa (1970).

value of  $\alpha$  than would otherwise be expected. After removing these (12) potentially biased data points, the regression was performed again and one data point from Giddings and Seager (1967) was identified as an outlier. The final regression result is shown in Fig. 4. The two-parameter regression yielded  $D(T, 1) = 0.1405(T/T_0)^{1.9212}$  with an  $R^2 = 0.972$  and the one-parameter regression produced  $D(0, 1) = 0.1429(\pm 0.0008)$  with an  $R^2 = 0.969$ . This one-parameter regression curve is virtually identical with the regression curve of MM72 (Fig. 4).

The results and comparisons for  $CO_2$  and  $O_2$  are presented in Tables 4 and 7; they are not presented graphically. The data, taken from Loschmidt (1870b), Wrestchko (1870), von Obermayer (1882a), Kosov and Zhalgasov (1959), Walker and Westenberg (1960) and Suetin and Ivakin (1961), were all within  $\pm 5\%$  of  $0.1402(T/T_0)^{1.81}$  and displayed no obvious outliers. The agreement between the present one-parameter regression curve and the curve given by MM72 is even closer than shown in Fig. 4 for  $CO_2$  and  $N_2$ .

#### $CH_{\Lambda}$

Figure 5 shows the results of the present analysis for  $CH_4$  in air and compares them with MM72's curve. No loess iteration was performed on the data because the absolute uncertainty was already relatively small ( $\pm 5\%$ ). The present results differ nearly uniformly from the results of MM72 by about 5.5%. However MM72's result is based on Blanc's law, not on a regression determination. Yaws' (1995) result (not shown) differs from the present result by approximately -3.3%. The results and comparisons for  $CH_4$  and  $N_2$  and  $O_2$  are presented in Tables 3, 4, 6 and 7; they are not presented graphically. The data for

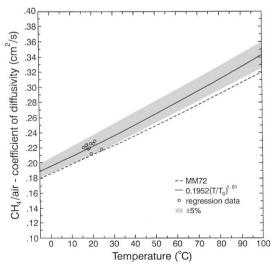


Fig. 5. Observed and modeled diffusivities for CH<sub>4</sub> in air. Data are from Coward and Georgeson (1937), Gwertsiteli et al. (1961) and Cowie and Watts (1971).

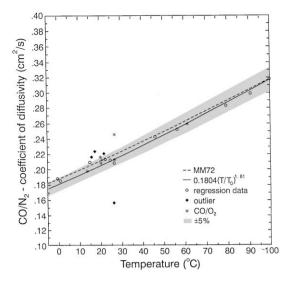


Fig. 6. Observed and modeled diffusivities for CO in N<sub>2</sub>. Data sources are: Trautz and Müller (1935b); Boardman and Wild (1937); Wicke and Hugo (1961); Weismann and Mason (1962); Amdur and Schuler (1963); Ivakin and Suetin (1964a); Vugts *et al.* (1970).

 ${\rm CH_4/N_2}$  were taken from Berry and Koeller (1960), Mueller and Cahill (1964), Arai et al. (1967), Jacobs et al. (1970), Engel and Knapp (1973) and Wakeham and Slater (1973). Only one data point from Wakeham and Slater (1973) was identified as an outlier. For  ${\rm CH_4/O_2}$  the author is aware of only one observation and it is given by Walker and Westenberg (1960). The comparisons between MM72 and the present analysis is much better for  ${\rm CH_4/N_2}$  and  ${\rm CH_4/O_2}$  than for  ${\rm CH_4/air}$ .

#### CO

The diffusivity coefficient for CO in air has never been measured. Figure 6 shows the results of the present analysis and MM72's result for CO/N<sub>2</sub>. Of the 16 data points 4 were removed with two iterations of the loess technique. These outliers came from the data of Trautz and Müller (1935b), Boardman and Wild (1937) and Wicke and Hugo (1961). The results for CO/O<sub>2</sub> are not shown graphically, but are included in Tables 4 and 7. Data for CO/O<sub>2</sub> are taken from Loschmidt (1870b), von Obermayer (1883), Trautz and Müller (1935b), Walker and Westenberg (1960) and Weissman and Mason (1962) and the analysis suggested that the datum of Trautz and Müller (1935b) is an outlier. The agreement between MM72 and the present result for CO/O2 is similar to that shown in Fig. 6 for CO/N<sub>2</sub>. The results for CO/N<sub>2</sub> and CO/O2 were combined using Blanc's law to provide an estimate of D(0, 1) for CO/air. Although not shown a comparison of the present determination of D(T, 1) for CO/air with MM72 and Yaws (1995) indicates that the present analysis and MM72 are nearly identical, whereas for the temperature range con-

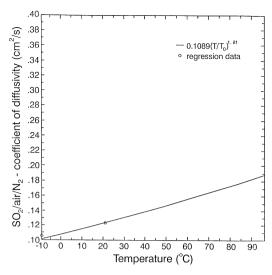


Fig. 7. Observed and modeled diffusivities for  $SO_2$ . The diffusivity datum for  $SO_2$  in air is from Andrew (1955) ( $T = 20^{\circ}$ C) and the diffusivity of  $SO_2$  in  $N_2$  is from Schäfer (1959) ( $T = -10^{\circ}$ C). These diffusivity values are used in a combined data set for the model regression.

sidered herein Yaws (1995) is approximately 10% less than the present results.

## $SO_2$

Because the diffusivity of  $SO_2$  has only been measured twice (once for  $SO_2$ /air and once for  $SO_2/N_2$ ), these two experimental determinations have been combined for the purposes of analysis (Fig. 7). Consequently, the resulting value of D(0,1) is used for both  $SO_2$ /air and  $SO_2/N_2$ . This combined approach yielded results that more closely approximated the modeling predictions than did treating the determinations of diffusivity separately. Finally, D(0,1) for  $SO_2/O_2$  is assigned the same value in accordance with Blanc's law. Nevertheless, the paucity of diffusivity data for  $SO_2$  suggests that more experimental determinations need to be made for this gas.

 $O_3$ 

Apparently, the diffusivity of ozone in any gas has never been measured. However, the critical constants have been measured (Ambrose (1992), who ultimately cites Mathews (1972)). Consequently, CO62 can easily be used for estimating the diffusivity of  $O_3$  in air (Table 2),  $O_2$  (Table 3) and  $O_2$  (Table 4). For the LJ model  $\sigma$  and  $\varepsilon/k$  were estimated from the following relations discussed by Reid *et al.* (1987):  $\sigma = 1.18V_b^{1/3}$  and  $\varepsilon/k = 1.15T_b$ ; where  $V_b$  is the liquid molar volume (cm<sup>3</sup> mol<sup>-1</sup>) and  $V_b$  is the boiling point at 1 atm pressure (K). For ozone Reid *et al.* (1987) give  $T_b = 181.2$  K and  $V_b = (48.0 \text{ gm mol}^{-1})/(1.356 \text{ gm cm}^{-3}) = 35.4 \text{ cm}^3 \text{ mol}^{-1}$ . The resulting estimates for  $\sigma$  and  $\varepsilon/k$  are given in Table 1 and the corresponding CO62 and LJ model estimates for D(0, 1) are given in Tables 2–4.

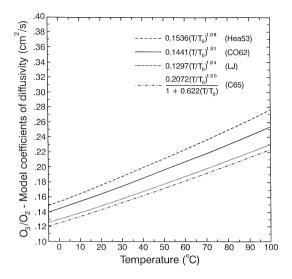


Fig. 8. Model predictions of the coefficient of diffusivity for  $O_3$  in  $O_2$ . The LJ and CO62 results are from this study with the LJ model being the two parameter fit. Hea53 is adapted from the model of Hirschfelder *et al.* (1953) and C65 is the model of Campbell (1965). See the text for a discussion of the Hea53 and C65 models.

In addition to the LJ and CO62 models of  $D_{O_3,O_2}(T,1)$  discussed in the present study, there are two other models that are included in Table 4. These are the models of Herschfelder et al. (1953) (Hea53) and Campbell (1965) (C65). The assumptions made Hea53 lead to the following relation:  $D_{O_3,O_2}(T,1) = 5D_{O_2,O_2}(T,1)/6$ . For use with Hea53  $D_{\Omega_2,\Omega_2}(T,1)$  was estimated from the data of Winn (1950) and Winter (1951) using the loess technique. The results of the two parameter fit is  $0.1843(T/T_0)^{1.88}$  with N = 5 and  $R^2 = 0.956$ . The one parameter fit yielded  $D(0, 1) = 0.1862 \text{ cm s}^{-1} \text{ with } N = 5, R^2 = 0.955 \text{ and}$ an absolute uncertainty of  $\pm 7\%$ . The LJ and CO62 models of  $D_{\Omega_3,\Omega_2}(T,1)$  are compared in Fig. 8 with the models of Hea53 and C65. Hea53 predict the highest values for  $D_{O_3,O_2}(T,1)$  and C65 predicts the lowest values. The variability between the four models is about  $\pm 10\%$  about the mean of the two extremum models and because the CO62 model is very close to this mean, this  $\pm 10\%$  variability is likely to represent the uncertainty of the CO62 model predictions for the diffusivity of ozone.

## $NH_3$

The results of the analysis for NH<sub>3</sub>/air are shown in Fig. 9. However, no regression was done on these data for this study. Rather Fig. 9 combines Wintergerst's (1930) estimate of D(0, 1) and  $\alpha = 1.81$  with more recent determinations of D(T, 1). Wintergerst (1930) published only the regression analysis for his experimental determinations of D(T, 1). He did not tabulate any of the data on which the regression was based. Nevertheless, Wintergerst (1930) found that  $\alpha = 1.80$  from a

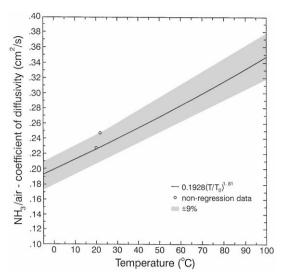


Fig. 9. Observed and modeled diffusivities for NH<sub>3</sub> in air. The model is the regression model of Wintergerst (1930). Otherwise the data are from Andrew (1955) and Ivakin and Suetin (1964a).

regression of his observations. Consequently, the value of 1.81 used in Fig. 9 is a very good approximation to his findings. The absolute uncertainty of  $\pm 9\%$  shown in Fig. 9 was determined from the more recent data of Andrew (1955) and Ivakin and Suetin (1964a) as compared with the curve presented in Fig. 9. Since no information is available from the Wintergerst (1930) concerning the variability of his original data, no data were eliminated as outliers.

Results for NH<sub>3</sub>/N<sub>2</sub> and NH<sub>3</sub>/O<sub>2</sub> are not presented graphically, but are summarized in Tables 3, 4, 6 and 7. The data for these analysis are taken from Trautz and Müller (1935b), Bunde (1955) [data also reported by Mason and Monchick (1962)], Ivakin and Suetin (1964a) and Weissman (1964). For NH<sub>3</sub>/O<sub>2</sub> the data of Trautz and Müller (1935b) exceed the data of Weissman (1964) and the LJ and CO62 models by 15 to 25%, therefore the data of Trautz and Müller (1935b) are assumed to be outliers. For NH<sub>3</sub>/N<sub>2</sub> one datum from Ivakin and Suetin (1964a) and one datum from Trautz and Müller (1935b) were eliminated by the loess technique as outliers. Although the other datum from Trautz and Müller (1935b) was not eliminated as an outlier by the loess technique, for the present study it is also eliminated as untrustworthy because all other data from Trautz and Müller (1935b) have either been identified as outliers by the loess technique or have significantly disagreed with other data or the model predictions (also see discussions of CO and NO).

## $N_2O$

The results of the present analysis for the diffusivity of  $N_2O$  in air are presented in Fig. 10. For this pair of gases there are a total of five experimental determinations of D(T, 1) from two experiments and these two

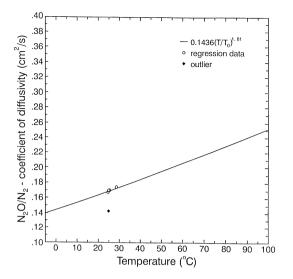


Fig. 10. Observed and modeled diffusivities for  $N_2O$  in air. Data are from Barr and Watts (1972) and Pritchard and Currie (1982).

experiments seem to give somewhat differing results. The single data point of Barr and Watts (1972) was eliminated as an outlier, but only because the other four points were in much better agreement with one another. However, these four "acceptable" data points are from the same experiment (Pritchard and Currie, 1982) and could conceivably be outliers relative to the data of Barr and Watts (1972). But, the data of Pritchard and Currie (1982) agree more closely with modeling results than does the datum of Barr and Watts (1972), so they are assumed to be more accurate. Nevertheless, given the paucity of data for the diffusivity of N<sub>2</sub>O in air and the apparent conflicts, more diffusion experiments for N<sub>2</sub>O/air need to be made. The same thing can be said for the diffusivity of N<sub>2</sub>O in N2 and O2 because no data are available for these gas pairs.

### NO

Neither the diffusivity of NO/air nor of NO/O<sub>2</sub> has been measured. However, there are four data for  $NO/N_2$ , which are shown in Fig. 11. But, the two sources of these data do not appear to be in reasonable agreement, as the data of Trautz and Müller (1935b) are systematically about 12% higher that the data of Weissman (1964). To further complicate matters, the data of Weissman (1964) are in very good agreement with the results of Bea90 (Tables 3 and 6), whereas the data of Trautz and Müller (1935b) are in equally good agreement with the model of CO62 (not shown). Considering that in general the Bea90 model performs slightly better than CO62 (Tables 5-7, excluding the NO comparisons) and that other datum of Trautz and Müller (1935b) have been previously identified as outliers by the present method of analysis (see discussions concerning CO and NH<sub>3</sub> above), it is assumed that the

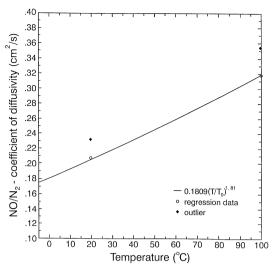


Fig. 11. Observed and modeled diffusivities for NO in N<sub>2</sub>. Data are from Trautz and Müller (1935b) and Weissman (1964).

data of Weissman (1964) is the more reliable of the two data sources. Figure 11 shows the results of the one-parameter regression using only the Weissman (1964) data.

Also included in Table 4 is a model based estimate of the diffusivity of NO/O<sub>2</sub> from Fan and Mason (1962). It is based on the LJ model but with slightly different values for  $\sigma$  and  $\varepsilon/k$ . It is included here for the sake of completeness and because very few results of any type are published for the diffusivity of NO.

## $NO_2$

The diffusivity of NO<sub>2</sub> has never been directly measured with any other gas. In fact D(T, 1) for  $NO_2$  may be impossible to measure directly because at high concentrations NO2 dimerizes to N2O4. Therefore, it may not be possible to obtain a pure sample of NO<sub>2</sub> for direct experimental determinations of D(T, 1). How-Sherwood and Chambers (1937) made indirect estimations of D(0, 1) and D(10, 1) for N<sub>2</sub>O<sub>4</sub> in N<sub>2</sub> from which it is possible to estimate D(0, 1) for NO<sub>2</sub> in N<sub>2</sub>. In their analysis Chambers and Sherwood (1937) assumed that the ratio  $D_{\text{NO}_2, \text{N}_2}(T, 1) = 1.43 D_{\text{N}_2\text{O}_4, \text{N}_2}(T, 1)$ , which from their tabulated values for  $D_{N_2O_4, N_2}(T, 1)$  $D_{\text{NO}_2, \text{N}_2}(0, 1) = 0.1210 \text{ cm}^2 \text{ s}^{-1}$ . Although Sviridenko et al. (1973) do not discuss the dimerization of NO<sub>2</sub> with  $N_2O_4$  with their data for  $D_{N_2O_4, N_2}(T, 1)$ , their results will also be used to estimate  $D_{NO_2,N_2}(T, 1)$ . Using a factor of 1.46 from CO62 for the ratio of the diffusivities of these two dimers, the results from Sviridenko et al. (1973) yield  $D_{NO_2,N_2}(T,1) =$  $0.1460(T/T_0)^2$ . These two estimates of  $D_{NO_2,N_2}(0,1)$  are included in Table 3 of this study.

The critical temperature and critical volume for  $NO_2$  used in this study (Table 1) for the model of

CO62 are derived from the critical constants for N<sub>2</sub>O<sub>4</sub> given by Ambrose (1992). It is assumed that the critical temperature is the same for both gases; but that the critical volume for NO2 is half the critical volume for N<sub>2</sub>O<sub>4</sub>. This last assumption is made because the molecular mass of NO2 is half that of N2O4. Also included in Table 4 is the model-derived estimate of D(0, 1) for NO<sub>2</sub> in O<sub>2</sub> from Fan and Mason (1962). Comparing the various estimates for NO<sub>2</sub> (Tables 2-4) indicate that the estimate derived from Sviridenko et al. (1973) is about 7% greater than CO62, that the CO62 and LJ models are in close agreement, that Sherwood and Chambers (1937) and Fan and Mason (1962) are in close agreement and that these last two pairs of estimates differ from one another by approximately 11%. As with O<sub>3</sub> the CO62 model appears to be a reasonable median value with an approximately ±10% uncertainty. However, present results also suggest that the diffusivity for NO2 is relatively uncertain and that experiments like those of Sherwood and Chambers (1937) and Sviridenko (1973) should be repeated.

## Isosteric molecules — CO and $N_2$

Because CO and N<sub>2</sub> have identical molecular mass, molecular shape and size and electronic configuration they are isosteric molecules. According to kinetic theory therefore, CO and N<sub>2</sub> should have virtually identical coefficients of self-diffusion, D(T, 1). For example, in a companion study to Bea90, Boushehri et al. (1987) suggest that  $D_{N_2,N_2}(T, 1) = D_{CO,CO}(T, 1)$  is accurate to greater than 99.99%. Likewise the diffusion coefficients for N2 and CO interdiffusing with nonpolar gases like N<sub>2</sub> and O<sub>2</sub> can also be expected to be very similar. For example, the modeling results of Boushehri et al. (1987) and Bea90 suggest that  $D_{N_2,N_2}(T,1) =$  $D_{\text{CO},N_2}(T,1)$  is accurate to within  $\pm 0.5\%$  and that  $D_{N_2,O_2}(T, 1) = D_{CO,O_2}(T, 1)$  is accurate to within 0.1%. As an independent check of D(T, 1) for CO, this study uses the loess regression technique to evaluate the coefficient of self-diffusion for N2 in N2 and the coefficient of diffusivity of  $N_2$  in  $O_2$ .

The data for the coefficient of self-diffusion for  $N_2$  are taken from Winn (1948, 1951), Winter (1951), DeLuca (1954) and Vugts *et al.* (1970). No data were eliminated as outliers and only one regression was necessary with N=13. The results for the two parameter fit are  $D(T,1)=0.1803(T/T_0)^{1.7288}$  with an  $R^2=0.989$ . The results for the one parameter fit are  $D(0,1)=0.1783(\pm0.001)$  with and  $R^2=0.987$  and an absolute uncertainty of  $\pm5\%$ . The value of D(0,1) for CO in  $N_2$  found earlier is  $0.1804~\rm cm^2~s^{-1}$  (Table 4) or 1.2% greater than the present value for the self-diffusion coefficient of  $N_2$ . This confirms the expectations that these isosteric molecules should have very similar values of D(0,1) and provides greater confidence in the present methods and results.

The data for evaluating  $D_{N_2,O_2}(T, 1)$  are taken from von Obermayer (1880, 1882b), Lonnius (1909),

Waldmann (1944, 1947), Bohemen and Purnell (1961), Giddings and Seager (1962), Weissman and Mason (1962), Arnikar *et al.* (1967a, b), Saxena and Gupta (1970), Hellemans *et al.* (1973) and Shankland and Dunlop (1980). The data point of Giddings and Seager (1962) was identified as the only outlier. The final regression results are for N=21: D(T,1)=0.1847  $(T/T_0)^{1.6522}$  with an  $R^2=0.927$  for the two parameter fit and for the one parameter fit D(0,1)=0.1809 ( $\pm 0.002$ ) with an  $R^2=0.913$  and an absolute uncertainty of  $\pm 9\%$ . The value for  $D_{\text{CO},O_2}(0,1)$  discussed earlier is 0.4% greater than the present value of  $D_{N_2,O_2}(0,1)$ , thereby again confirming expectations.

## Synthesis and recommended values for D(0, 1)

The recommended values for D(0, 1) are listed in Table 8. The recommended values of D(0, 1) for water vapor in air, N2 and O2 are the same  $(=0.2178 \text{ cm}^2 \text{ s}^{-1})$  primarily because the difference between  $D_{\text{H}_2\text{O},\text{air}}(0, 1)$  and  $D_{\text{H}_2\text{O},\text{N}_2}(0, 1)$  is so small and because  $D_{H_2O,O_2}(0,1)$  is imprecisely known. Likewise the recommended value for  $D_{SO_2,O_2}$  is taken to be the same as D(0, 1) for  $SO_2$  in air and  $N_2$ . Otherwise the recommended values for D(0, 1) are taken from the corresponding tables (i.e. Tables 5-7). As much as possible the recommended values are derived from data. But in some cases the recommended values result from model calculations. These model-based recommendations were made only after the models were compared and evaluated relative to the data analysis. The model comparisons shown in Tables 5-7 suggest that the model of Bea90 tends to be slightly more accurate than the model of CO62, and both of these models tend to be more accurate than LJ. Therefore, the order of preference for model-based recommendations is Bea90 first, and then CO62 over LJ. However, recommending one model over another is not intended to suggest that one model is necessarily significantly better than another. With the exception of water vapor, no model results differ from the data analysis by more

Table 8. Recommended coefficients of diffusivity (cm  $^2\,s^{-1})$  of selected gases in air,  $N_2$  and  $O_2$ 

Gas	Air	$N_2$	$O_2$
H <sub>2</sub> O	0.2178	0.2178	0.2178
$CO_2$	0.1381	0.1429	0.1402
$CH_4$	0.1952	0.1892	0.1942
$CO^a$	$0.1807^{b}$	0.1804	0.1817
$SO_2$	0.1089	0.1089	0.1089
$O_3$	0.1444	0.1441	0.1454
NH <sub>3</sub>	0.1978	0.1978	0.1992
$N_2O$	0.1436	0.1402	0.1396
NO	0.1802	0.1809	0.1815
$NO_2$	0.1361	0.1359	0.1369
$N_2^a$	0.1788 <sup>b</sup>	0.1783	0.1809
$O_2$	0.1820 <sup>b</sup>	0.1809	0.1862

Note: Only D(0, 1) is given, D(T, p) = D(0, 1)  $(p_0/p)$   $(T/T_0)^{1.81}$  is assumed.

<sup>&</sup>lt;sup>a</sup> Isosteric molecules.

<sup>&</sup>lt;sup>b</sup> From Blanc's law and D(0, 1) for gas in  $N_2$  and in  $O_2$ .

than about  $\pm 10\%$  and usually the difference is much less. Therefore, in these cases the data and modeling results support one another because where data are available all modeling and data results are about equally uncertain.

The situation is unfortunately somewhat different for the cases with no data. Given the performances of the models for water vapor and  $NO/N_2$ , it seems possible that the model estimates of D(0, 1) for  $O_3$ , NO, and  $NO_2$  could be in error by as much as  $\pm 25\%$ , and indeed such large discrepancies between observed and modeled values for D(T, 1) have been previously noted for other gases (Reid *et al.*, 1987). Consequently, all recommendations for D(0, 1) for NO,  $NO_2$  and  $O_3$  must be viewed as quite uncertain and need to be verified experimentally if at all possible.

#### CONCLUSIONS

The loess regression technique is very effective for identifying outliers in the diffusivity data sets and worked especially well for the coefficients of diffusivity of H<sub>2</sub>O and CO<sub>2</sub> which had relatively large data sets. For smaller data sets, the loess method is not necessarily as useful because the inherent variability of the data is sometimes too great. Under these circumstances it is impossible to evaluate the data in a completely objective manner, because the disparities within the data set are often so great as to suggest that the data are contradictory. Nevertheless, for H<sub>2</sub>O, CO<sub>2</sub>, CH<sub>4</sub> and CO the final regression curve developed with the loess technique gave results that were extremely close to the regression curves developed by Marrero and Mason (1972). Because Marrero and Mason (1972) used an understanding of theory and experimental techniques to evaluate data quality, the close agreement between their results and the present results suggests that those specific experiments identified as outliers by the loess technique should be considered the least trustworthy of the experimental determinations of D(T, p).

The recommended values for the diffusivity are in many cases relatively precise and are valid to within an absolute uncertainty of  $\pm 5\%$  to  $\pm 9\%$ . Furthermore, when the model predictions and data analysis are in close agreement they tend to mutually support one another. However, in spite of over a century of diffusivity measurements for a variety of gases, the diffusivities of O3, NO and NO2 in air remain relatively uncertain. Given the importance of these gases to atmospheric chemistry and atmosphere-biosphere interaction studies, there is an obvious need to verify the recommended values presented in Table 8 of this study. For example, the recommended ratio of the diffusivity of water vapor in air to the diffusivity of ozone in air is 1.51 which is about 10% less that the often quoted value of 1.65.

Finally the present results demonstrate that the misuse of Graham's law to estimate D(T, p) can lead

to significant errors. For example, using  $D_{\rm w}(0,1)$  as the known diffusivity, the tabulations in Table 8 suggest that this approach can lead to errors of as much as 23%. Such inaccuracies should not be unexpected because the misapplication of Graham's law is in opposition to all theoretical results (e.g. equations (2) and (3), Reid *et al.*, 1987) which demonstrate that  $D(0,1) \propto \sqrt{1/M_i + 1/M_i}$ , not to  $\sqrt{M_i/M_i}$ .

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