

$$(1) \quad H = \frac{1}{2} m L^2 \omega^2 + \frac{1}{2} m g L \theta^2 + \text{const} \quad \begin{matrix} \nearrow \\ \text{who} \\ \text{cares.} \end{matrix}$$

The forces on the pendulum are

$$\ddot{\theta} = -\frac{g}{L} \theta \quad (q = \theta, \quad p = mL^2 \omega)$$

The Euler method ~~can~~ solves this
via

$$\begin{pmatrix} \theta_{n+1} \\ \omega_{n+1} \end{pmatrix} = \begin{pmatrix} \theta_n \\ \omega_n \end{pmatrix} + \Delta t \begin{pmatrix} \omega_n \\ -\frac{g}{L} \theta_n \end{pmatrix}$$

$$\begin{aligned} \frac{1}{2} m L^2 \omega_{n+1}^2 &= \frac{1}{2} m L^2 \left(\omega_n - \frac{\Delta t g}{L} \theta_n \right)^2 \\ &= \frac{1}{2} m L^2 \omega_n^2 - \Delta t m g L \theta_n \omega_n + \frac{\Delta t^2}{2} m g^2 \theta_n^2 \end{aligned}$$

$$\frac{1}{2} m g L \theta_{n+1}^2 = \frac{1}{2} m g L \left(\theta_n - \Delta t \omega_n \right)^2$$

$$= \frac{1}{2} m g L \theta_n^2 + \Delta t m g L \omega_n + \frac{1}{2} \Delta t^2 m g L \omega_n^2$$

$$\text{So } E_{n+1} = E_n + \frac{\Delta t^2}{2} \frac{g}{L} \left(m L^2 \omega_n^2 + m g L \theta_n^2 \right)$$

$$= E_n \underbrace{\left(1 + \frac{\Delta t^2 g}{L} \right)}_{\text{positive \#}}$$

Note this implies

$$E_n = E_0 \cdot \left(1 + \frac{\Delta t g}{L} \right)^n,$$

monotonic exp. growth.

$$(2) (a) \quad \frac{d^2 x}{dt^2} = \frac{d^2 x}{d \left[\tau / \omega \right]^2} = \omega^2 \frac{d^2 x}{d \tau^2}.$$

$$\text{Then } \frac{d^2 x}{d \tau^2} = -\omega^2 x \quad \text{implies}$$

$$\frac{dx}{d\tau} = -x,$$

Init. velocity is $\frac{dx}{dt} = \omega \frac{dx}{d\tau}$,

so in our dimensionless units,
the init vel is v_0/ω .

2(b) The general trends are that
 g_2 increases faster than
 g_1 , roughly linearly on
log-log axes. This means
the error is a power
law in Δt . If we
had plotted Δt^{-2} &
 Δt^{-4} , we could see!

the agreement.

(3) a with $\eta = x - ct$,

$$\frac{\partial u}{\partial t} = \frac{\partial u}{\partial \eta} (-c)$$

$$\frac{\partial u}{\partial x} = \frac{\partial u}{\partial \eta}$$

so

$$-c u' + b u u' + u''' = 0.$$

(b) If $\eta = z/\sqrt{c}$, $\frac{\partial}{\partial \eta} = \sqrt{c} \frac{\partial}{\partial z}$, so,

If $v = cu$, then

$$-c \frac{\partial u}{\partial \eta} = -c^{5/2} \frac{\partial v}{\partial z}$$

$$b u \frac{\partial u}{\partial \eta} = b c^{5/2} v \frac{\partial v}{\partial z}$$

$$\frac{\partial^3 u}{\partial y^3} = c^{5/2} \frac{\partial^3 v}{\partial z^3},$$

Divide through by $c^{5/2}$ to
find

$$-v' + 6vv' + v''' = 0.$$

(C) There is good agreement.

(N) ^{repeated} Peaks are observed. Their
location depends on the
integration timestep.

Because the solutions are
time invariant,

$$w(z) = \frac{1}{2} \operatorname{sech}^2 \left(\frac{1}{2}z + a \right)$$

is a solution for all a .

Solitons will be observed
far from each other,

where the numerical
methods can't distinguish
between 0 & $\frac{1}{2}a$ a small
number.

$w(z) + w(z+a)$ is close
to a solution as long
as a is "large enough".

large enough dep. on
the alg.