

## PROBLEM 1: MECHANICS

### A. Determining the radius of curvature of a planar curve by means of Mechanics

**a. 1.0 point**

$$v_x = \frac{dx}{dt} = \dot{x}$$

$$v_y = \frac{dy}{dt} = \frac{dy}{dx} \frac{dx}{dt} = f'(x) \cdot \dot{x}$$

**b. 1.0 point**

$$a_x = \frac{dv_x}{dt} = \dot{v}_x = \ddot{x}$$

$$a_y = \frac{dv_y}{dt} = \frac{d}{dt}[f'(x) \cdot \dot{x}] = f''(x) \cdot \dot{x}^2 + f'(x) \cdot \ddot{x}$$

$$\vec{a}_t = \frac{\vec{a} \cdot \vec{v}}{v} \cdot \frac{\vec{v}}{v} = \frac{a_x v_x + a_y v_y}{\sqrt{v_x^2 + v_y^2}} \cdot \frac{v_x \vec{i} + v_y \vec{j}}{\sqrt{v_x^2 + v_y^2}} = \frac{v_x (a_x v_x + a_y v_y) \vec{i} + v_y (a_x v_x + a_y v_y) \vec{j}}{v_x^2 + v_y^2}$$

**c. 1.0 point**

$$\begin{aligned} \vec{a}_n &= \vec{a} - \vec{a}_t = \\ &= a_x \vec{i} + a_y \vec{j} - \frac{v_x (a_x v_x + a_y v_y) \vec{i} + v_y (a_x v_x + a_y v_y) \vec{j}}{v_x^2 + v_y^2} = \\ &= \frac{[a_x (v_x^2 + v_y^2) - v_x (a_x v_x + a_y v_y)] \vec{i} + [a_y (v_x^2 + v_y^2) - v_y (a_x v_x + a_y v_y)] \vec{j}}{v_x^2 + v_y^2} = \\ &= \frac{(a_x v_y^2 - a_y v_x v_y) \vec{i} + (a_y v_x^2 - a_x v_x v_y) \vec{j}}{v_x^2 + v_y^2} = \\ &= \frac{(a_x v_y - a_y v_x) (v_y \vec{i} - v_x \vec{j})}{v_x^2 + v_y^2} \\ |\vec{a}_n| &= \frac{|a_x v_y - a_y v_x|}{\sqrt{v_x^2 + v_y^2}} \end{aligned}$$

d. 0.5 point

$$\begin{aligned}
 R &= \frac{v^2}{a_n} = \\
 &= \frac{\sqrt{v_x^2 + v_y^2}^3}{|a_x v_y - a_y v_x|} = \\
 &= \frac{[\dot{x}^2 + (f'(x) \cdot \dot{x})^2]^{\frac{3}{2}}}{|\ddot{x} \cdot f'(x) \cdot \dot{x} - (f''(x) \cdot \dot{x}^2 + f'(x) \cdot \ddot{x}) \cdot \dot{x}|} = \\
 &= \frac{|\dot{x}|^3 [1 + f'^2(x)]^{\frac{3}{2}}}{|\dot{x}^3 \cdot f''(x)|} = \\
 &= \frac{[1 + f'^2(x)]^{\frac{3}{2}}}{|f''(x)|}
 \end{aligned}$$

e. 0.5 point

$$\left. \begin{aligned} f'(x_0) &= 2Ax_0 \\ f''(x_0) &= 2A \end{aligned} \right\} \Rightarrow R = \frac{(1 + 4A^2 x_0^2)^{\frac{3}{2}}}{2A}$$

f. 0.5 point

$$f(x_0) = -1 \Rightarrow \sin 2x_0 = -1 \Rightarrow x_0 = \frac{3}{4}\pi$$

$$\left. \begin{aligned} f'(x_0) &= 2 \cos \left( 2 \cdot \frac{3}{4} \pi \right) = 0 \\ f''(x_0) &= -4 \sin \left( 2 \cdot \frac{3}{4} \pi \right) = 4 \end{aligned} \right\} \Rightarrow R = \frac{1}{4} = 0.25 \text{ m}$$

$$T = 2\pi \sqrt{\frac{R}{g}} \approx 1 \text{ s}$$

**B. Springs on a circle**

**g. 0.5 point**

According to the cosine rule:

$$l_n^2 = R^2 + r^2 - 2Rr \cos\left(\frac{2n\pi}{N+1} - \alpha\right)$$

**h. 0.25 point + 0.25 point**

$$\begin{aligned} E_{\text{kin}} &= \frac{m(\dot{r}^2 + r^2 \dot{\alpha}^2)}{2} \\ E_{\text{pot}} &= \sum_{n=0}^N \frac{kl_n^2}{2} = \\ &= \frac{k}{2} \sum_{n=0}^N \left[ R^2 + r^2 - 2Rr \cos\left(\frac{2n\pi}{N+1} - \alpha\right) \right] \end{aligned}$$

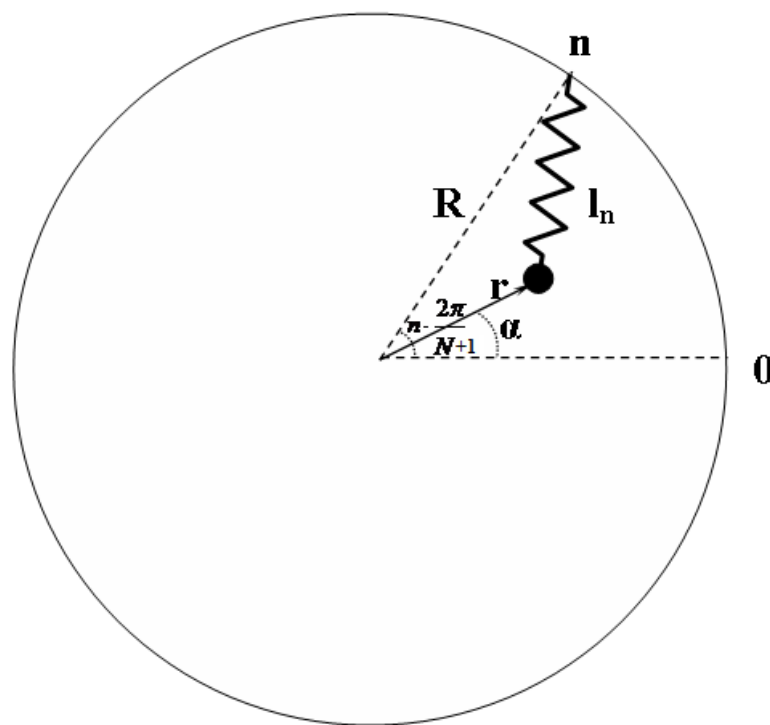
**i. 0.25 point + 0.25 point**

One can use the sum of a geometric progression having complex terms.

$$\begin{aligned} \sum_{n=0}^N \left[ \cos\left(\frac{2n\pi}{N+1}\right) + i \sin\left(\frac{2n\pi}{N+1}\right) \right] &= \sum_{n=0}^N e^{i \frac{2n\pi}{N+1}} = \\ &= \frac{\left( e^{i \frac{2\pi}{N+1}} \right)^{N+1} - 1}{e^{i \frac{2\pi}{N+1}} - 1} = \\ &= \frac{\cos 2\pi + i \sin 2\pi - 1}{\cos\left(\frac{2\pi}{N+1}\right) + i \sin\left(\frac{2\pi}{N+1}\right) - 1} = \\ &= 0 \end{aligned}$$

So

$$\begin{aligned} \sum_{n=0}^N \cos\left(\frac{2n\pi}{N+1}\right) &= \sum_{n=0}^N \sin\left(\frac{2n\pi}{N+1}\right) = 0 \\ E_{\text{pot}} &= \frac{k}{2} \left[ \sum_{n=0}^N (R^2 + r^2) - 2Rr \sum_{n=0}^N \cos\left(\frac{2n\pi}{N+1} - \alpha\right) \right] = \frac{N+1}{2} k (R^2 + r^2) \end{aligned}$$



**j. 0.5 point**

The force with which the  $n$ -th spring is acting on the object is

$$F_n = kl_n$$

According to the sine rule,

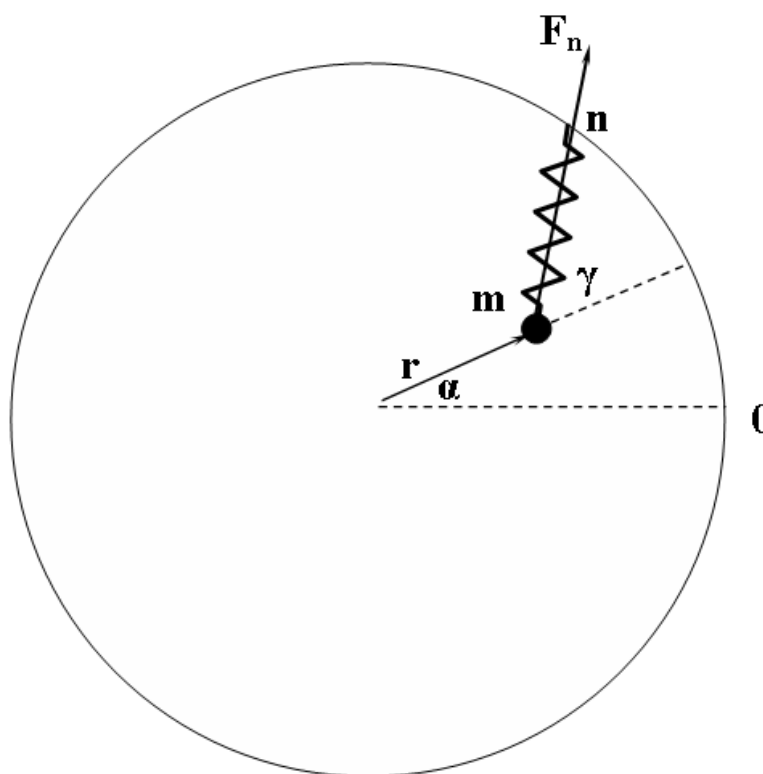
$$\frac{\sin(\pi - \gamma)}{R} = \frac{\sin\left(\frac{2n\pi}{N+1} - \alpha\right)}{l_n}$$

The torque of  $F_n$  equals

$$M_n = r \cdot F_n \sin \gamma =$$

$$= r \cdot kl_n \cdot \frac{R \sin\left(\frac{2n\pi}{N+1} - \alpha\right)}{l_n} =$$

$$= kRr \sin\left(\frac{2n\pi}{N+1} - \alpha\right)$$



The total torque equals

$$M = \sum_{n=0}^N M_n = \sum_{n=0}^N kRr \sin\left(\frac{2n\pi}{N+1} - \alpha\right) = 0$$

So the angular momentum is indeed constant.

$$L = mr^2 \dot{\alpha}$$

**k. 0.5 point**

$$E = E_{\text{kin}} + E_{\text{pot}} = \frac{1}{2} m \left( \dot{r}^2 + \frac{L^2}{m^2 r^2} \right) + \frac{1}{2} k(N+1)(R^2 + r^2) = \text{constant} \Rightarrow$$

$$\dot{E} = m \left( \ddot{r} - \frac{L^2}{m^2} r^{-3} \dot{r} \right) + k(N+1)r\dot{r} = 0 \Rightarrow$$

$$\ddot{r} - \mathcal{L} r^{-3} + \omega^2 r = 0$$

**l. 0.5 point**

$$\dot{r} = \frac{2z\dot{z}}{2\sqrt{z^2 + K}} = \frac{z\dot{z}}{r} \Rightarrow$$

$$\ddot{r} = \frac{(\dot{z}^2 + z\ddot{z})r - z\dot{z}\dot{r}}{r^2} = \frac{\dot{z}^2 + z\ddot{z}}{r} - \frac{z\dot{z}}{r^2} \cdot \frac{z\dot{z}}{r} \Rightarrow$$

$$\frac{\dot{z}^2 + z\ddot{z}}{r} - \frac{z^2 \dot{z}^2}{r^3} - \frac{\mathcal{L}^2}{r^3} + \omega^2 r = 0$$

$$(\dot{z}^2 + z\ddot{z})(z^2 + K) - z^2 \dot{z}^2 - \mathcal{L}^2 + \omega^2 (z^2 + K)^2 = 0$$

$$(\ddot{z} + \omega^2 z)z^3 + K\ddot{z}z + K\dot{z}^2 + 2K\omega^2 z^2 + K^2\omega^2 - \mathcal{L}^2 = 0$$

**m. 0.5 point**

The first term of the equation cancels anyway.

$$-K\omega^2 A^2 \cos^2(\omega t + \phi_0) + K\omega^2 A^2 \sin^2(\omega t + \phi_0) + 2K\omega^2 A^2 \cos^2(\omega t + \phi_0) + K^2\omega^2 - \mathcal{L}^2 = 0$$

$$K\omega^2 A^2 + K^2\omega^2 - \mathcal{L}^2 = 0 \Rightarrow A^2 = \frac{\mathcal{L}^2}{K\omega^2} - K$$

**n. 0.5 point**

$$\begin{aligned} r(t) &= \sqrt{\frac{\mathcal{L}^2 - K^2\omega^2}{K\omega^2} \cos^2(\omega t + \phi_0) + K} = \\ &= \frac{1}{\omega} \sqrt{\frac{\mathcal{L}^2 \cos^2(\omega t + \phi_0) + K^2\omega^2 \sin^2(\omega t + \phi_0)}{K}} = \\ &= \frac{1}{\omega\sqrt{2K}} \sqrt{\mathcal{L}^2 [1 + \cos(2\omega t + 2\phi_0)] + K^2\omega^2 [1 - \cos(2\omega t + 2\phi_0)]} = \\ &= \frac{1}{\omega\sqrt{2K}} \sqrt{(\mathcal{L}^2 + K^2\omega^2) + (\mathcal{L}^2 - K^2\omega^2) \cos(2\omega t + 2\phi_0)} \\ T_r &= \frac{2\pi}{2\omega} \end{aligned}$$

But since the values of  $r$  repeat after every rotation by  $\pi$ , it means that the period of the motion is

$$T = 2\pi \sqrt{\frac{1}{N+1} \frac{m}{k}}$$

**o. 0.5 point**

$$L = 0 \Rightarrow \mathcal{L} = 0 \Rightarrow r(t) = \sqrt{K} |\sin(\omega t + \phi_0)|$$

$$L = 0 \Rightarrow \dot{\alpha} = 0 \Rightarrow \alpha = \text{constant}$$

According to the previous part, except for the moments when  $r = 0$  and  $\alpha$  is not defined, there will be two constant values for  $\alpha$ , corresponding to two opposing directions. So in this case the object will perform simple harmonic motion with amplitude  $\sqrt{K}$ .

**p. 0.5 point**

$$r = \text{constant} = \sqrt{K} \Rightarrow \mathcal{L}^2 - K^2 \omega^2 = 0 \Rightarrow \left| \frac{L}{m} \right| = K \sqrt{\frac{k}{m} (N+1)} \Rightarrow r = \sqrt{\frac{|L|}{\sqrt{(N+1)mk}}}$$

**q. 0.5 point**

We can think of this situation as having  $N + 1$  springs of positive constant  $k$  and another  $(N + 1)/d$  springs of “negative” stiffness  $-k$ . The potential energy  $E_{\text{pot}}$  derived in **i.** holds for both ensembles (with appropriate  $N$  and  $k$ ) and angular momentum is still conserved. Thus all the results from part **j.** onwards continue to hold if we replace  $(N + 1)$  by  $(N + 1) + [-(N + 1)/d]$ . Thus

$$\omega'^2 = \left[ N + 1 - \frac{N+1}{d} \right] \frac{k}{m} \Rightarrow \omega' = \omega \sqrt{1 - \frac{1}{d}}$$