Solution Problem 2

Detection of Alpha Particles

a. From the given range-energy relation and the data supplied we get

$$E = \left(\frac{R\alpha}{0.318}\right)^{\frac{2}{3}} MeV = \left(\frac{5.50}{0.318}\right)^{\frac{2}{3}} = 6.69 MeV$$
(0.5 point)

since $W_{ion-pair} = 35 \text{ eV}$, then

$$N_{ion-\rho air} = \frac{6.69 \times 10^6}{35} = 1.9 \times 10^5$$
(0.5 point)

Size of voltage pulse:

$$\Delta V = \frac{\Delta Q}{C} = \frac{N_{air-pair}e}{C}$$
with $C = 45 \, pF = 4.5 \times 10^{-11}$
(0.5 point)

Hence

$$\Delta V = \frac{1.9 \times 10^5 \times 1.6 \times 10^{-19}}{4.5 \times 10^{-11}} V = 0.68 mV$$
(0.5 point)

b. Electrons from the ions-pairs produced by α particles from a radioactive sources of activity A (=number of α particles emitted by the sources per second) which enter the detector with detection efficiency 0.1, will produce a collected current.

$$I = \frac{Q}{t} = 0.1 \times AN_{ion-\rho\alpha ir}e$$

= 0.1 \times A \times 1.9 \times 10^5 \times 1.6 \times 10^{-19} A (1.0 point)

With $I_{min}=10^{-12}$ A, the

$$A_{\min} = \frac{10^{-12} \operatorname{dis} s^{-1}}{1.6 \times 1.9 \times 10^{-15}} = 330 \operatorname{dis} s^{-1}$$
(1.0 point)

Since 1 Ci = $3.7 \times 10^{10} \text{ dis s}^{-1} \text{ then}$

$$A_{\min} = \frac{330}{3.7 \times 10^{10}} Ci = 8.92 \times 10^{-9} Ci$$
(1.0 point)

c. With time constant

$$\tau = RC \left(\text{with C} = 45 \times 10^{-12} \text{F} \right) = 10^{-3} s$$

$$R = \left(\frac{1000}{45} \right) M \Omega = 22.22 M \Omega$$
(0.5 point)

For the voltage signal with height Δ V = 0.68 mV generated at the anode of the ionization chamber by 6.69 MeV α particles in problem (a), to achieve q 0.25 V= 250mV voltage signal, the necessary gain of the voltage pulse amplifier should be

$$G = \frac{250}{0.68} = 368 \tag{0.5 point}$$

d. By symmetry, the electric field is directed radially and depends only on distance the axis and can be deducted by using Gauss' theorem.

If we construct a Gaussian surface which is a cylinder of radius τ and length l, the charge contained within it is σl .

The surface integral

$$\int E.dS = 2\pi r l E$$

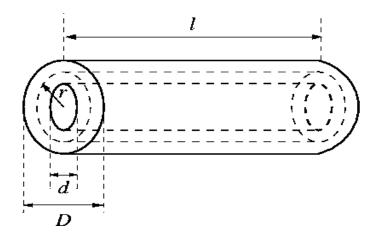


Figure 1 : The Gaussian surface used to calculate the electric field E. (1.0 point)

Since the field E is everywhere constant and normal to the curved surface. By Gauss's theorem :

$$2\pi r l E = \frac{\lambda l}{\varepsilon_0}$$

so

$$E(r) = \frac{\lambda}{2\pi\varepsilon_0 r}$$

Since E is radial and varies only with τ , then $E=-\frac{dV}{dr}$ and the potential V can be found by integrating E(τ) with respect to τ , if we call the potential of inner wire V₀, we have

$$V(r) - V_0 = -\frac{\lambda}{2\pi\varepsilon_0} \int_{\frac{d}{2}}^{\tau} \frac{dr}{r}$$

Thus

$$V(r) - V_0 = -\frac{\lambda}{2\pi\varepsilon_0} \ln\left(\frac{2r}{d}\right)$$
(1.0 point)

We can use this expression to evaluate the voltage between the capacitor's conductors by setting $r = \frac{D}{2}$, giving a potential difference of

$$V = \frac{\lambda}{2\pi\varepsilon_0} \ln\left(\frac{D}{d}\right)$$

since the charge Q in the capacitor is σl , and the capacitance C is defined by Q=CV, the capacitance per unit length is

$$\frac{2\pi\varepsilon L_0}{\ln\frac{D}{d}}$$
(1.0 point)

The maximum electric field occurs where r minimum, i.e. at $r=\frac{d}{2}$ if we set the field at $r=\frac{d}{2}$ equal to the breakdown field E_b , our expression for E® shows that the charges per unit length σ in the capacitor must be $E_b\pi_0 d$. Substituting for the potential difference V across the capacitor gives

$$V = \frac{1}{2} E_b d \ln \left(\frac{D}{d} \right)$$

Taking E $_b$ = 3 x 10 6 V , d= 1mm, and D= 1 cm , gives V = 3.453.45 kV. (1.0 point).