

# Coherent and incoherent photoproduction of $J/\psi$ in Pb-Pb UPCs at 5.02 TeV

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In this note we update our predictions of Ref. [1] to Run 2 Pb-Pb UPCs at  $\sqrt{s_{NN}} = 5.02$  TeV. At  $y = 0$ , this corresponds to  $W_{\gamma p} = 125$  GeV and  $x = 6 \times 10^{-4}$ .

Using our predictions for the leading-twist gluon nuclear shadowing [2, 3], we can estimate the nuclear suppression factor at this value of  $x$  (see Eq. (11) of Ref. [1])

$$\begin{aligned} R(x = 6 \times 10^{-4}) &= \frac{\sigma_{\gamma A \rightarrow J/\psi Y}(W_{\gamma p})}{A(\sigma_{\gamma p \rightarrow J/\psi p}(W_{\gamma p}) + \sigma_{\gamma p \rightarrow J/\psi Y}(W_{\gamma p}))} \\ &= \int d^2b T_A(b) \left( 1 - \frac{\sigma_2}{\sigma_3} + \frac{\sigma_2}{\sigma_3} e^{-\frac{\sigma_3}{2} A T_A(b)} \right)^2 = 0.10 - 0.26, \end{aligned} \quad (1)$$

where  $T_A(b)$  is normalized to unity, i.e.  $\int d^2b T_A(b) = 1$ . These values are expectedly smaller than those for  $x = 10^{-3}$  in Ref. [1] because of a somewhat larger suppression due to nuclear shadowing at higher energies.

Combining this with the parametrization of the cross section of elastic  $J/\psi$  photoproduction on the proton  $\sigma_{\gamma p \rightarrow J/\psi p}(W_{\gamma p})$  and proton-dissociative  $J/\psi$  photoproduction on the proton  $\sigma_{\gamma p \rightarrow J/\psi Y}(W_{\gamma p})$ , see Ref. [4], one can readily make predictions for the elastic and nucleon-dissociation contributions to the  $d\sigma_{\gamma A \rightarrow J/\psi Y}(W_{\gamma p})/dt$  cross section as a function of  $|t|$ . The result is presented in Fig. 1 in next page.

Note that our predictions for the elastic contribution for  $|t| \geq 1.2$  GeV<sup>2</sup> is an extrapolation of the H1 parametrization into the unmeasured range of  $t$  [4].

One can see from the figure that our predictions for Run 2 are rather close (on the log scale) to our predictions for Run 1. This is a result of a partial compensation of two opposite effects: a decrease of  $R$  in Eq. (1) with an increase of energy is partially compensated by an increase of the proton  $\sigma_{\gamma p \rightarrow J/\psi p}(W_{\gamma p})$  and  $\sigma_{\gamma p \rightarrow J/\psi Y}(W_{\gamma p})$ .

## References

- [1] V. Guzey, M. Strikman and M. Zhalov, Phys. Rev. C **99** (2019) no.1, 015201 [arXiv:1808.00740 [hep-ph]].
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- [3] V. Guzey and M. Zhalov, JHEP **10** (2013), 207 [arXiv:1307.4526 [hep-ph]].
- [4] C. Alexa *et al.* [H1], Eur. Phys. J. C **73** (2013) no.6, 2466 [arXiv:1304.5162 [hep-ex]].

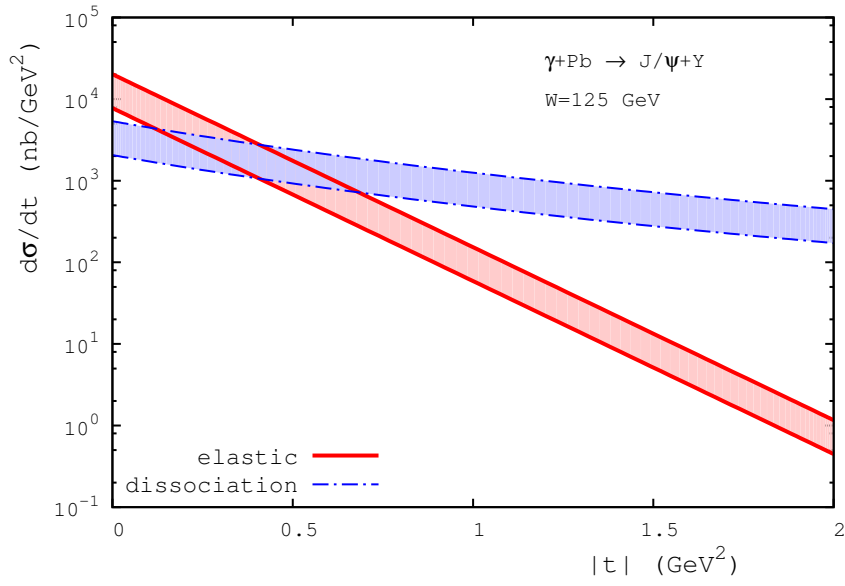


Figure 1: The elastic and nucleon-dissociation contributions to  $d\sigma_{\gamma A \rightarrow J/\psi Y}(W_{\gamma p})/dt$  cross section as a function of  $|t|$ . The shaded error bands quantify the uncertainty in the calculation of nuclear shadowing;  $W_{\gamma p} = 125 \text{ GeV}$  corresponds to Pb-Pb UPCs at  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$  and  $y = 0$ .