

$$\begin{aligned}
 u &= -\frac{g}{f_0} \frac{\partial \eta}{\partial y} - \frac{g}{f_0^2} \frac{\partial^2 \eta}{\partial x \partial t} + \frac{\beta_0 g}{f_0^2} y \frac{\partial \eta}{\partial y} \\
 v &= +\frac{g}{f_0} \frac{\partial \eta}{\partial x} - \frac{g}{f_0^2} \frac{\partial^2 \eta}{\partial y \partial t} - \frac{\beta_0 g}{f_0^2} y \frac{\partial \eta}{\partial x}
 \end{aligned}$$

Substitution into the continuity equation: $\frac{\partial \eta}{\partial t} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) = 0.$

$$\frac{\partial \eta}{\partial t} - R^2 \frac{\partial}{\partial t} \nabla^2 \eta - \beta_0 R^2 \frac{\partial \eta}{\partial x} = 0 \quad \boxed{R = \sqrt{gH}/f_0}$$

Apply a wave solution $\eta = Ae^{i(kx+ly-\omega t)}$, and the dispersion relation is:

$$\omega = -\beta_0 R^2 \frac{k}{1 + R^2(k^2 + l^2)}$$

$$\omega = -\beta_0 R^2 \frac{k}{1 + R^2(k^2 + l^2)}$$

If $\beta_0 = 0$, $\omega = 0$, geostrophic flow

If $R^2 K^2 \ll 1$, $L \gg R$, long wave:

$$\omega \sim \frac{\beta_0 R^2}{L} \ll \beta_0 L \ll f_0$$

$$\boxed{\frac{\beta_0 L}{f_0} \ll 1}$$

If $R^2 K^2 \geq 1$, $L \leq R$, short wave:

$$\omega \sim \beta_0 L \ll f_0$$

Planetary Rossby waves are subinertial (low frequency) waves

The zonal wave speed:

$$c = \frac{\omega}{k} = \frac{-\beta_0 R^2}{1 + R^2(k^2 + l^2)} < 0$$

Planetary Rossby waves always propagate westward in the zonal direction

For very long waves ($R^2 K^2 \ll 1$):

$$c_x = -\beta_0 R^2 \quad \text{maximum zonal wave speed}$$

The dispersion relation can be reorganized as:

$$\left(k + \frac{\beta_0}{2\omega}\right)^2 + l^2 = \left(\frac{\beta_0^2}{4\omega^2} - \frac{1}{R^2}\right)$$

$$|\omega|_{\max} = \frac{\beta_0 R}{2} \quad \text{maximum frequency}$$

The dispersion relation diagram

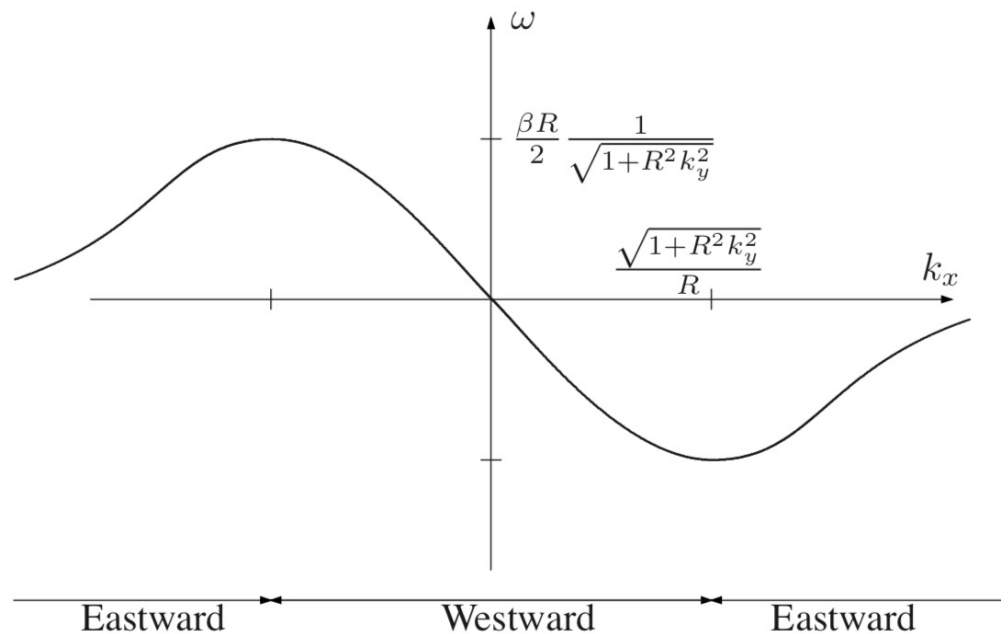
$$\omega = -\beta_0 R^2 \frac{k}{1 + R^2(k^2 + l^2)}$$

$$k = 0, \quad \omega = 0$$

$$k > 0, \quad \omega < 0; \quad k < 0, \quad \omega > 0$$

$$k \rightarrow \infty, \quad \omega \rightarrow 0$$

$$\frac{d\omega}{dk} = 0, \quad k = \pm \frac{\sqrt{1+R^2 l^2}}{R}, \quad |\omega| = \frac{\beta R}{2} \frac{1}{\sqrt{1+R^2 l^2}}$$



for long Rossby waves, energy propagates westward

for short Rossby waves, energy propagates eastward
(opposite to wave propagation)

Topographic Rossby Waves

Assumptions: sloping topography

$$H = H_0 + \alpha_0 y$$

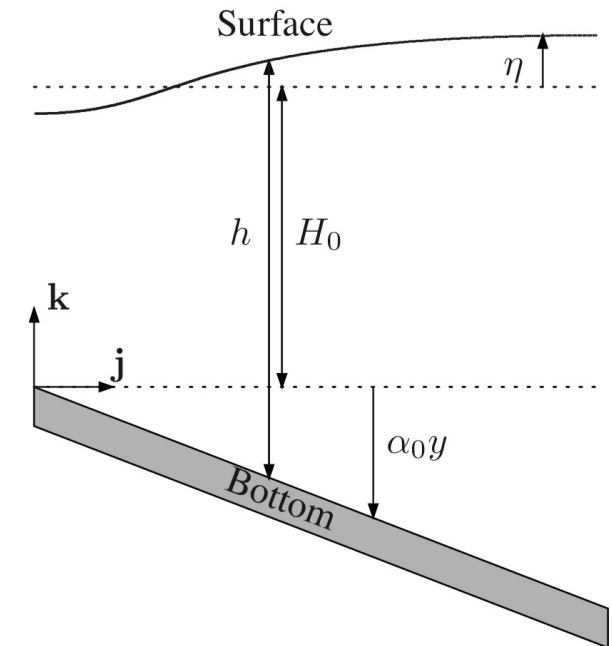
$$\alpha = \frac{\alpha_0 L}{H_0} \ll 1$$

$$h(x, y, t) = H_0 + \alpha_0 y + \eta(x, y, t)$$

The continuity equation:

$$\frac{\partial \eta}{\partial t} + h \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + u \frac{\partial h}{\partial x} + v \frac{\partial h}{\partial y} = 0$$

$$\begin{aligned} \frac{\partial \eta}{\partial t} + \left(u \frac{\partial \eta}{\partial x} + v \frac{\partial \eta}{\partial y} \right) + (H_0 + \alpha_0 y) \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \\ + \eta \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \alpha_0 v = 0. \end{aligned}$$



Small-amplitude waves:

$$\Delta H \ll H$$

The governing equations: $\frac{\partial \eta}{\partial t} + H_0 \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \alpha_0 v = 0$

$$\begin{aligned} \frac{\partial u}{\partial t} - f v &= -g \frac{\partial \eta}{\partial x} \\ \frac{\partial v}{\partial t} + f u &= -g \frac{\partial \eta}{\partial y} \end{aligned}$$

To 1st order approximation – **geostrophic balance**:

$$-f v = -g \frac{\partial \eta}{\partial x}$$

$$f u = -g \frac{\partial \eta}{\partial y}$$

Substitute the solutions into the small terms of the governing equations:

$$\begin{aligned} u &= -\frac{g}{f} \frac{\partial \eta}{\partial y} - \frac{g}{f^2} \frac{\partial^2 \eta}{\partial x \partial t} \\ v &= +\frac{g}{f} \frac{\partial \eta}{\partial x} - \frac{g}{f^2} \frac{\partial^2 \eta}{\partial y \partial t} \end{aligned}$$

ageostrophic flow

Substitution in continuity equation: $\frac{\partial \eta}{\partial t} + H_0 \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \alpha_0 v = 0$

$$\frac{\partial \eta}{\partial t} - R^2 \frac{\partial}{\partial t} \nabla^2 \eta + \frac{\alpha_0 g}{f} \frac{\partial \eta}{\partial x} = 0.$$

Apply a wave solution $\eta = Ae^{i(kx+ly-\omega t)}$:

$$\omega = \frac{\alpha_0 g}{f} \frac{k}{1 + R^2(k^2 + l^2)}$$

$$\omega = \frac{\alpha_0 g}{f} \frac{k}{1 + R^2(k^2 + l^2)}$$

If $\alpha_0 = 0$, $\omega = 0$, geostrophic flow

If $R^2 K^2 \ll 1$, $L \gg R$, long wave:

$$\alpha = \frac{\alpha_0 L}{H_0} \ll 1$$

$$\omega = \frac{\alpha_0 g k}{f} \sim \frac{\alpha_0 L g}{f L^2} \ll \frac{\alpha_0 L g}{f R^2} \left(\sim \frac{\alpha_0 L g f^2}{f g H} \right) \ll f$$

If $R^2 K^2 \geq 1$, $L \leq R$, short wave:

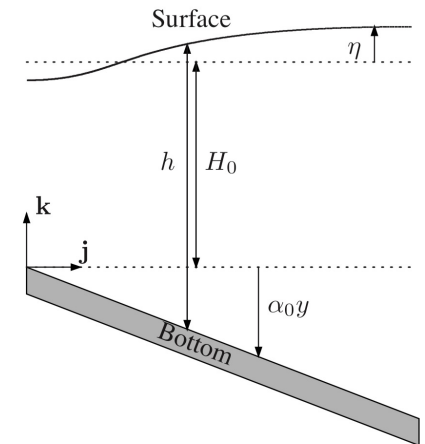
$$\omega = \frac{\alpha_0 g}{f} \frac{k}{R^2 K^2} \sim \frac{\alpha_0 g}{f} \frac{L}{R^2} \sim \frac{\alpha_0 L g f^2}{f g H} \ll f$$

Topographic Rossby waves are also subinertial (low frequency) waves

The zonal wave speed:

$$c = \frac{\omega}{k} = \frac{\alpha_0 g}{f} \frac{1}{1 + R^2(k^2 + l^2)} > 0$$

topographic Rossby waves propagate with shallow water on the right (left) in the northern (southern) hemisphere



For very long waves ($R^2 K^2 \ll 1$):

$$c_x = \frac{\alpha_0 g}{f} \quad \text{maximum zonal wave speed}$$

The dispersion relation can be reorganized as:

$$\left(k - \frac{\alpha_0 g}{2f\omega R^2}\right)^2 + l^2 = \left(\frac{\alpha_0^2 g^2}{4f^2 R^4 \omega^2} - \frac{1}{R^2}\right)$$

$$|\omega|_{\max} = \frac{|\alpha_0|g}{2|f|R}$$

Analogy between planetary and topographic Rossby waves

“North-Shallow” analogy (northern hemisphere)

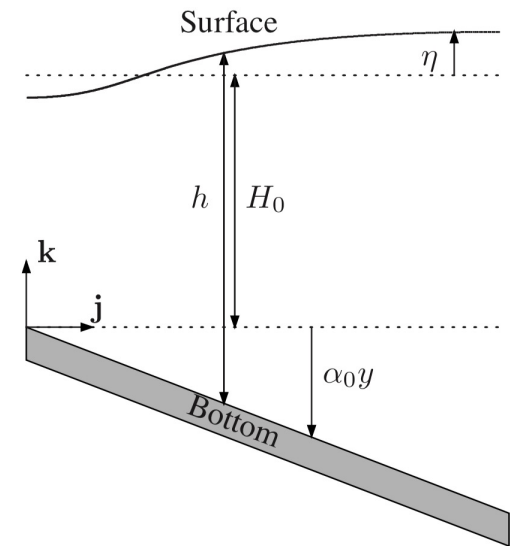
The potential vorticity:

planetary wave $q = \frac{f_0 + \beta_0 y + \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}}{H_0}$

toward the north, $y \uparrow$, q increases

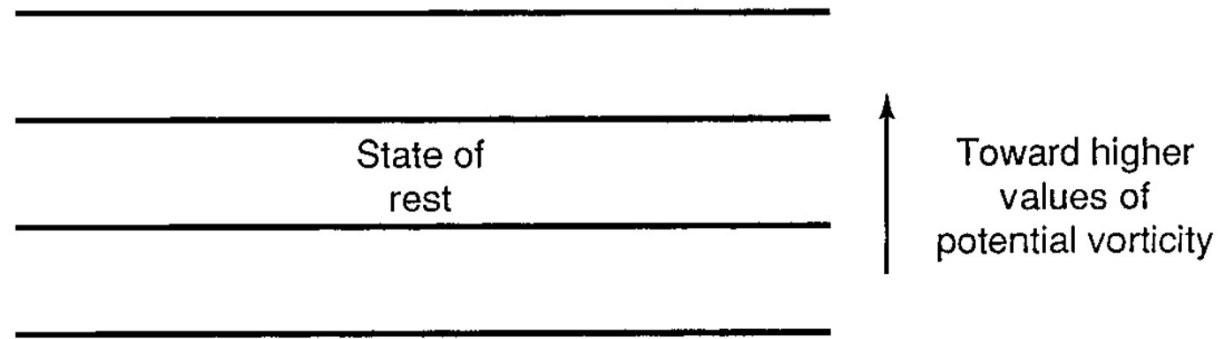
topographic wave $q = \frac{f_0 + \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}}{H_0 + \alpha_0 y + \eta}$

toward the shallow water, $y \downarrow$, q increases

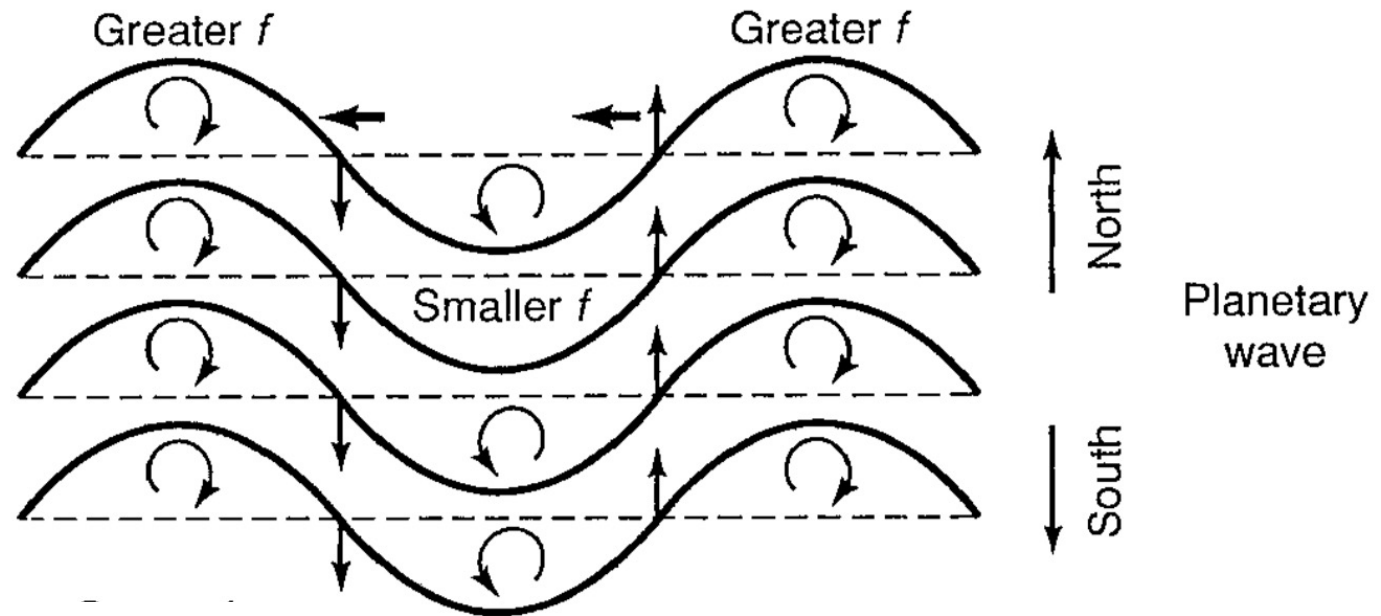


Mechanisms for Rossby Wave propagation

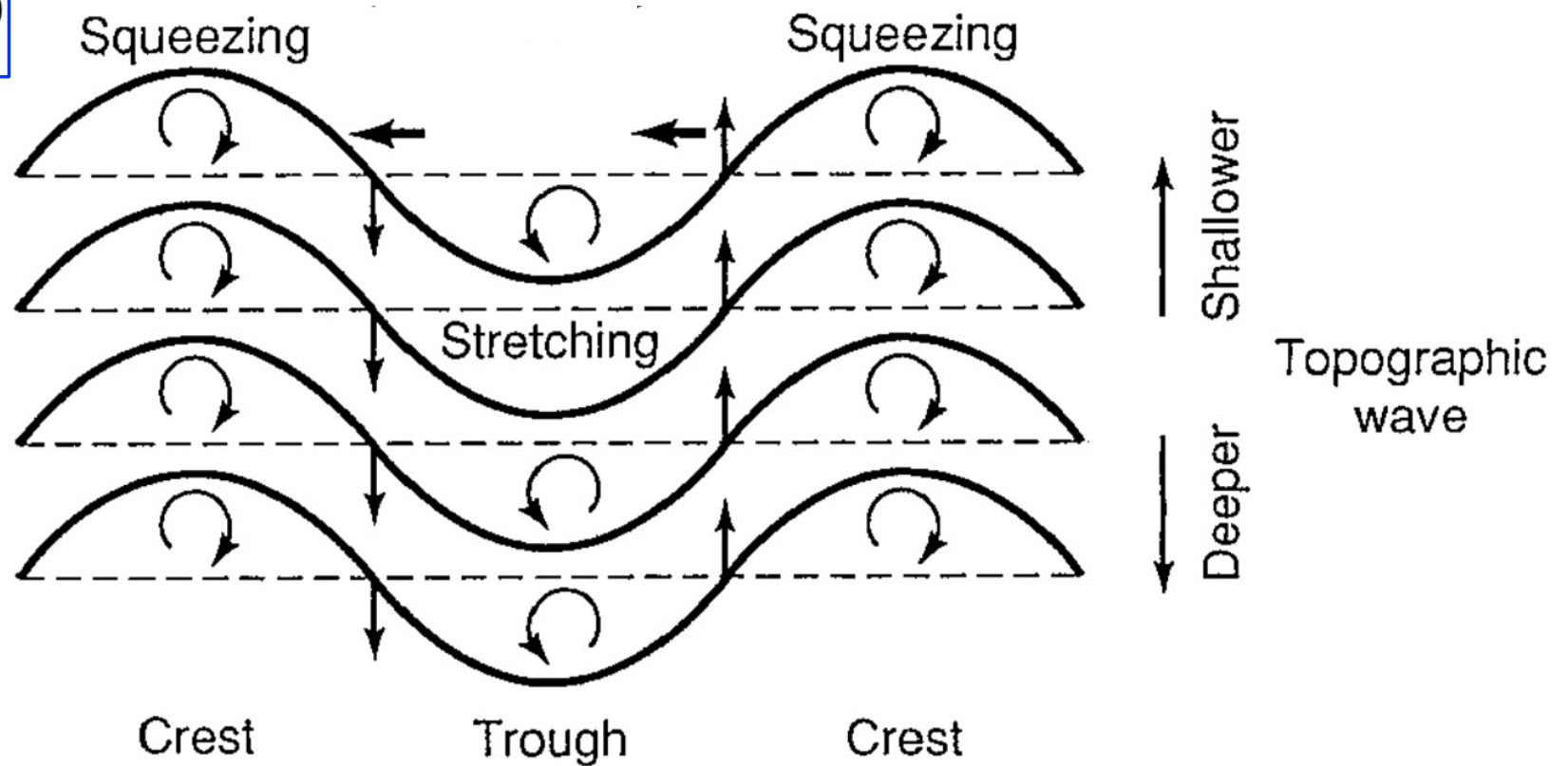
$$\frac{d}{dt} \left(\frac{f + \zeta}{h} \right) = 0$$



North on the right

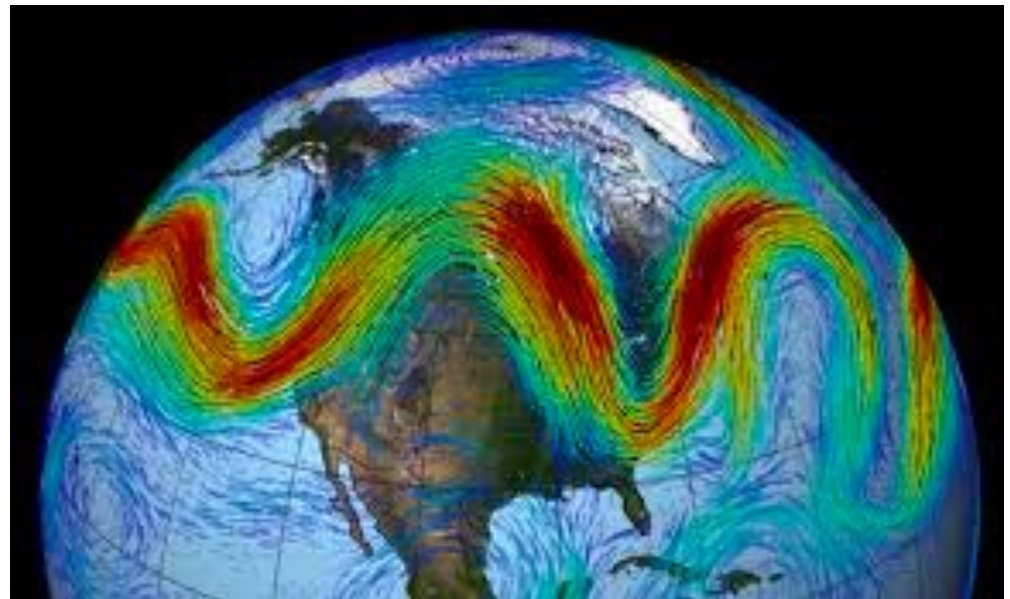
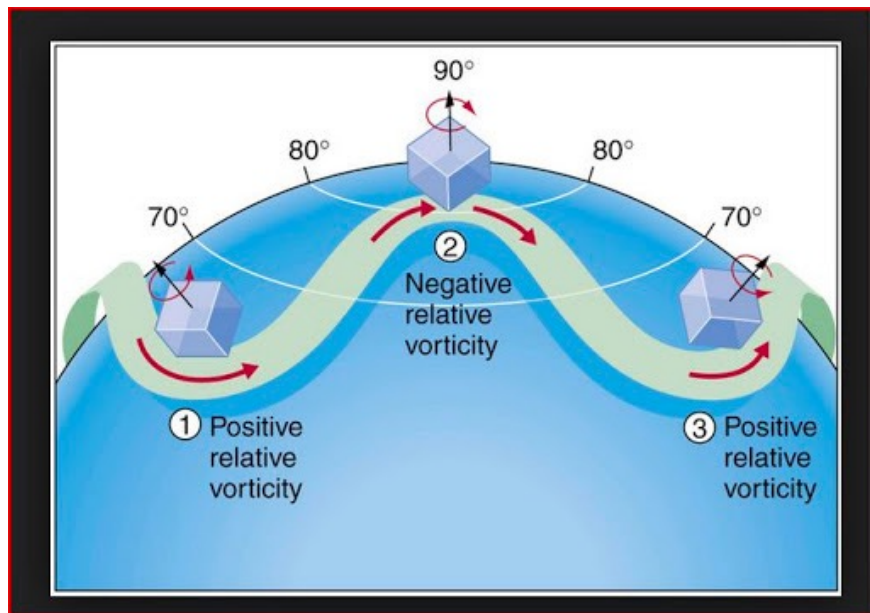


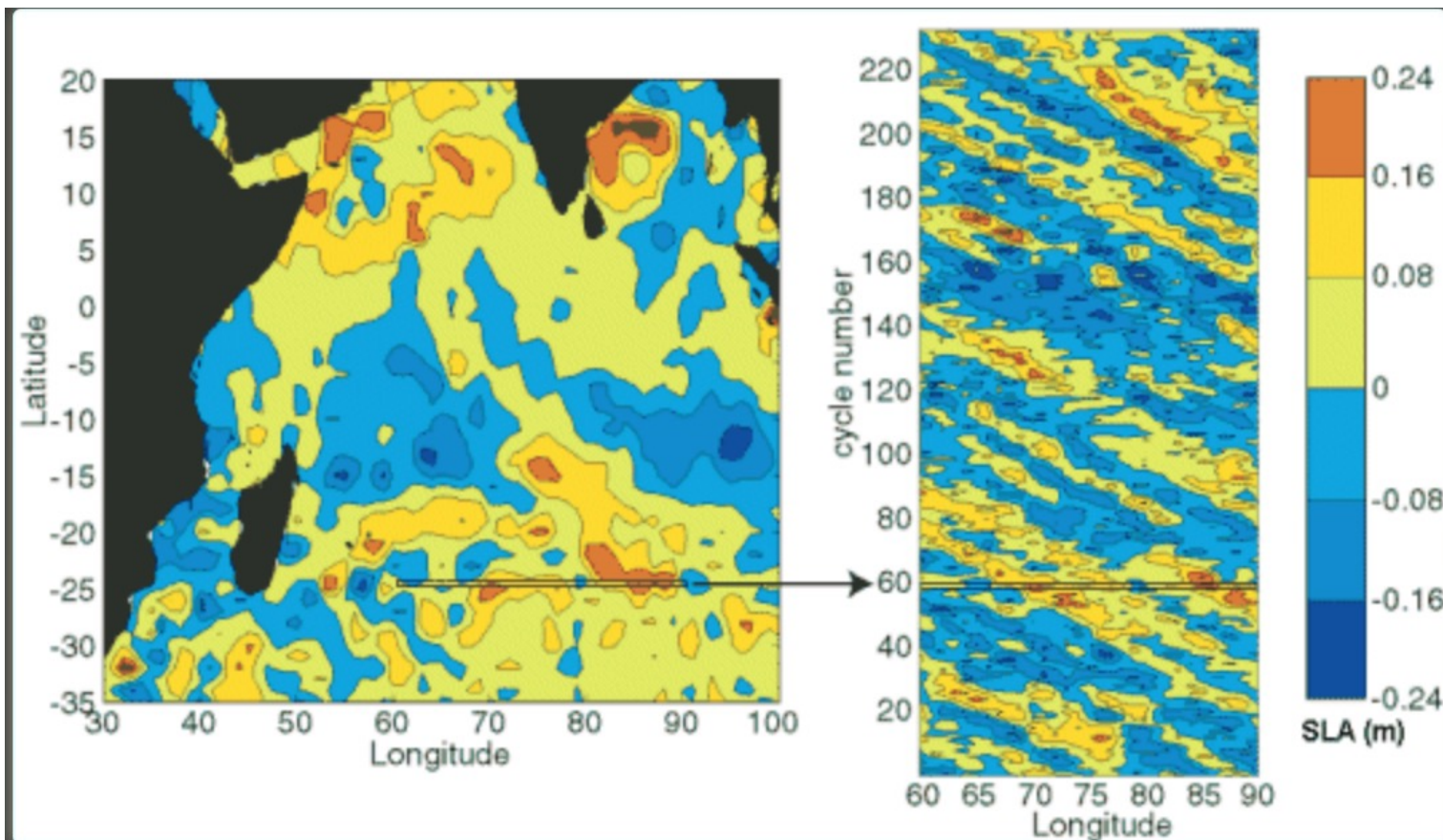
$$\frac{d}{dt} \left(\frac{f + \zeta}{h} \right) = 0$$



Shallow-water on the right

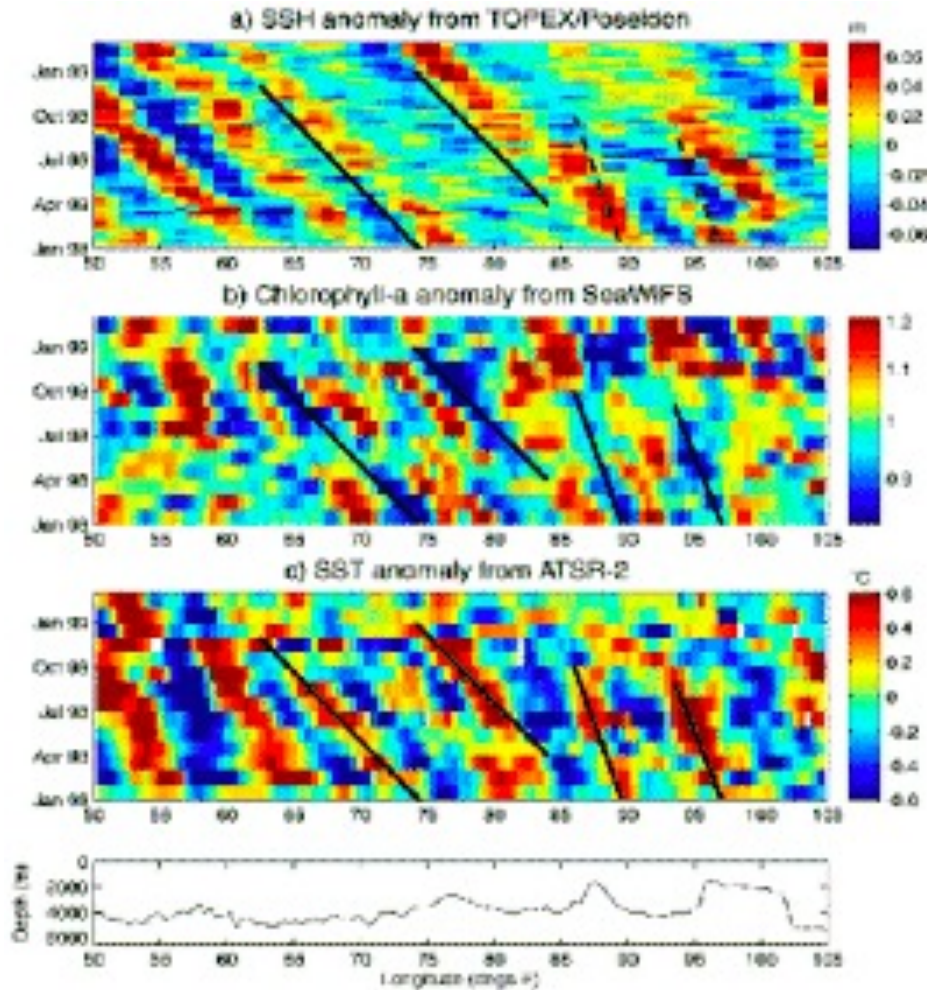
Rossby waves in the atmosphere





Map of sea level anomalies, on the left, and longitude-time diagram on the right (the time is indicated in Topex/Poseidon cycles). Diagrams like these bring out the variations of sea level over time along a particular parallel. The elevation at 90°E in cycle 20 (in yellow) can be found about 3 months (10 cycles) later at 85°E, 6 months later at 80°E, and so on. The westward motion of this elevation can be seen on the diagram as a sloped line. Another elevation follows the same path with a 3 month offset, creating a parallel trace. (Credits Southampton Oceanography Center).

<https://www.aviso.altimetry.fr/news/idm/2002/may-2002-planetary-waves-small-amplitudes-large-effects.html>

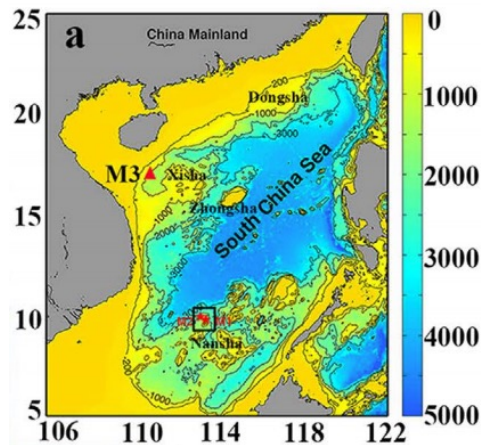


SSH

Chl-a

SST

Comparison of longitude-time diagrams from three different sensors (altimeter, water colour, surface temperature). (Credits Southampton Oceanography Center).



Shu et al. (2015)

Topographic Rossby waves in the South China Sea

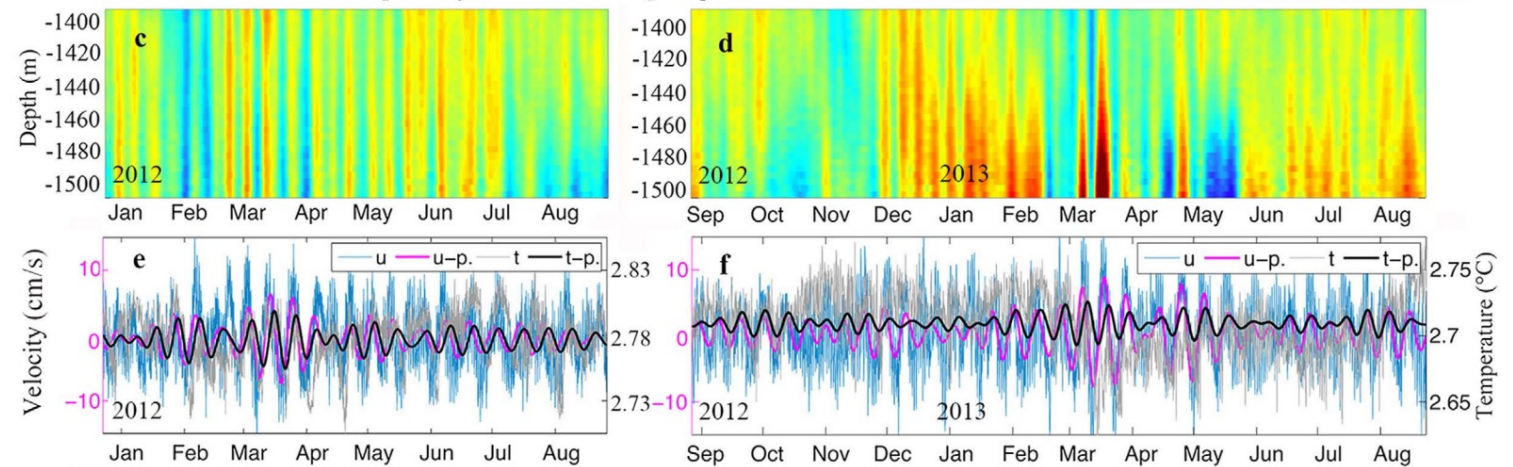


Figure 2. Deep-layer current observations. Observed near-bottom zonal velocity at M1 (unit: cm/s). (a–d) Time series of 3 day low-pass filtered zonal current profiles from the four segments of ADCP observations. Time series of the raw zonal velocity (cyan line) and temperature (grey line) obtained by the Aanderaa current meter at 1730 m (20 m above the bottom) (e) from December 23, 2011 to August 26, 2012 and (f) from August 27, 2012 to August 23, 2013. The solid heavy black and magenta lines represent the 9–14 day band-pass filtered zonal velocity and temperature, respectively. Figures were plotted using MATLAB.