$$\rho_{0}^{\left\{\frac{\partial(\overline{u}+u')}{\partial t} + \frac{\partial(\overline{u}+u')^{2}}{\partial x} + \frac{\partial(\overline{u}+u')(\overline{v}+v')}{\partial y} + \frac{\partial(\overline{u}+u')(\overline{w}+w')}{\partial z}\right\}} + \rho_{0}(-f\overline{v} + f_{*}\overline{u})$$

$$= \overline{-\frac{\partial(\overline{p}+p')}{\partial x} + \mu\left(\frac{\partial^{2}(\overline{u}+u')}{\partial x^{2}} + \frac{\partial^{2}(\overline{u}+u')}{\partial y^{2}} + \frac{\partial^{2}(\overline{u}+u')}{\partial z^{2}}\right)} = \frac{F^{x} = \frac{\partial \tau^{xx}}{\partial x} + \frac{\partial \tau^{yx}}{\partial y} + \frac{\partial \tau^{zx}}{\partial z}}{= \frac{1}{\partial x}\left(\mu\frac{\partial u}{\partial x}\right) + \frac{1}{\partial y}\left(\mu\frac{\partial u}{\partial y}\right) + \frac{1}{\partial z}\left(\mu\frac{\partial u}{\partial z}\right)}$$

$$= \mu\nabla^{2}u$$

$$\frac{\Delta = \nabla^{2}}{\partial x}$$

$$= \mu\nabla^{2}u$$

$$\frac{\Delta = \nabla^{2}}{\partial x}$$

$$= \rho_{0}\left(\frac{\partial u'u'}{\partial x} + \frac{\partial u'v'}{\partial y} + \frac{\partial u'w'}{\partial z}\right)$$

$$= \rho_{0}\left(\frac{\partial u'u'}{\partial x} + \frac{\partial u'v'}{\partial y} + \frac{\partial u'w'}{\partial z}\right)$$

$$= \rho_{0}\left(\frac{\partial u'u'}{\partial x} - \frac{\partial u'v'}{\partial y} - \frac{\partial u'w'}{\partial z}\right)$$

$$= \nabla^{2}u$$

$$= \rho_{0}\left(\frac{\partial u'u'}{\partial x} - \frac{\partial u'v'}{\partial y} - \frac{\partial u'w'}{\partial z}\right)$$

$$= \rho_{0}\left(\frac{\partial u'u'}{\partial y} - \frac{\partial u'w'}{\partial z}\right)$$

$$= \rho_{0}\left(\frac{\partial u'u'}{\partial y} - \frac{\partial u'w'}{\partial z}\right)$$

$$= \nabla^{2}u$$

$$= \rho_{0}\left(\frac{\partial u'u'}{\partial x} - \frac{\partial u'w'}{\partial y} - \frac{\partial u'w'}{\partial z}\right)$$

Reynolds stress: stress induced by turbulence acting on mean flow

Reynolds-averaged equations

x:
$$\frac{\partial \mathbf{u}}{\partial \mathbf{t}} + \mathbf{u} \frac{\partial \mathbf{u}}{\partial \mathbf{x}} + \mathbf{V} \frac{\partial \mathbf{u}}{\partial \mathbf{y}} + \mathbf{W} \frac{\partial \mathbf{u}}{\partial \mathbf{z}} - f v + f_* w = -\frac{1}{\rho_0} \frac{\partial p}{\partial x} + \frac{\partial}{\partial x} \left(\mathbf{A}_H \frac{\partial \mathbf{u}}{\partial \mathbf{x}} \right) + \frac{\partial}{\partial y} \left(\mathbf{A}_H \frac{\partial \mathbf{u}}{\partial y} \right) + \frac{\partial}{\partial z} \left(\mathbf{A}_V \frac{\partial \mathbf{u}}{\partial z} \right)$$

y:
$$\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} + fu = -\frac{1}{\rho_0} \frac{\partial p}{\partial y} + \frac{\partial}{\partial x} \left(A_H \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial y} \left(A_H \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial z} \left(A_V \frac{\partial v}{\partial z} \right)$$

z:
$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} + f_* u = -\frac{1}{\rho_0} \frac{\partial p}{\partial z} + \frac{\partial}{\partial x} \left(A_H \frac{\partial w}{\partial x} \right) + \frac{\partial}{\partial y} \left(A_H \frac{\partial w}{\partial y} \right) + \frac{\partial}{\partial z} \left(A_V \frac{\partial w}{\partial z} \right) - \frac{\rho}{\rho_0} g$$

Scale analysis

Characteristic scale: a typical scale for a variable

 Table 4.1
 Typical scales of atmospheric and oceanic flows

| Variable | Scale | Unit | Atmospheric value | Oceanic value |
|------------------|---------------|------------------------------------|---|--|
| $\overline{x,y}$ | L | m | $100 \text{ km} = 10^5 \text{ m}$ | $10 \text{ km} = 10^4 \text{ m}$ $10^4 \sim 10^6 \text{ km}$ |
| z | H | m | $1 \text{ km} = 10^3 \text{ m}$ | $100 \text{ m} = 10^2 \text{ m}$ $10^{1} \sim 10^{3} \text{ km}$ |
| t | T | S | $\geq \frac{1}{2} \operatorname{day} \simeq 4 \times 10^4 \operatorname{s}$ | $\geq 1~\mathrm{day} \simeq 9 \times 10^4~\mathrm{s}$ |
| u, v | U | m/s | 10 m/s | 0.1 m/s |
| w | W | m/s | | |
| p | P | ${\rm kg}{\rm m}^{-1}{\rm s}^{-2}$ | variable | |
| ρ | $\Delta \rho$ | kg/m ³ | | |

The continuity equation

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \qquad \qquad \frac{\partial u}{\partial x} = \frac{\Delta u}{\Delta x} \sim \frac{U}{L}$$

scale

$$\frac{U}{L}$$
 $\frac{U}{L}$ $\frac{W}{H}$

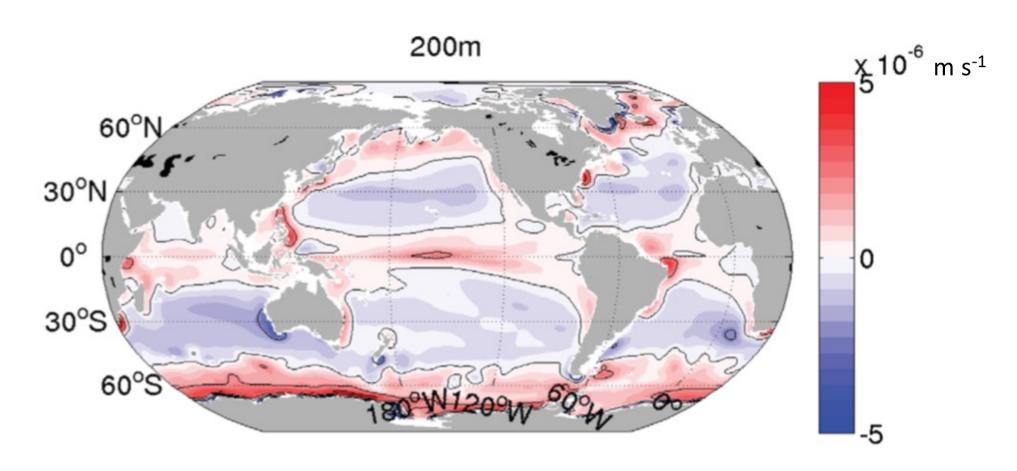
$$\frac{U}{L} \sim \frac{W}{H}$$

$$W \sim U \frac{H}{L} \sim U \delta$$

$$W \sim U \frac{H}{L} \sim U \delta$$
 $\delta = \frac{H}{L}$: aspect ratio (10⁻⁴~10⁻²)

$$W \ll U$$

Vertical velocity in the global ocean



Liang et al. (2017, JGR)

The momentum equation

x direction:
$$\frac{\partial \mathbf{u}}{\partial \mathbf{t}} + \mathbf{u} \frac{\partial \mathbf{u}}{\partial \mathbf{x}} + \mathbf{v} \frac{\partial \mathbf{u}}{\partial \mathbf{y}} + \mathbf{w} \frac{\partial \mathbf{u}}{\partial \mathbf{z}} - f v + f_* w = \frac{U}{T} \frac{U^2}{L} \quad \frac{U^2}{L} \quad \mathbf{w} \frac{U}{H} \quad f U \quad f W$$

$$-\frac{1}{\rho_0}\frac{\partial p}{\partial x} + \frac{\partial}{\partial x}\left(A_H\frac{\partial u}{\partial x}\right) + \frac{\partial}{\partial y}\left(A_H\frac{\partial u}{\partial y}\right) + \frac{\partial}{\partial z}\left(A_V\frac{\partial u}{\partial z}\right)$$

$$\frac{P}{\rho_0 L} \qquad A_H\frac{U}{L^2} \qquad A_H\frac{U}{L^2} \qquad A_V\frac{U}{H^2}$$

Non-dimensional numbers

Rossby number (R_o) = scale of nonlinear advection term / scale of Coriolis term

=
$$\frac{U^2}{L}/fU = U/fL$$
 measure of the role of Earth's rotation in motions

 $R_o \ll 1$: rotation is important for the motion

Ekman number (E_k) = scale of vertical viscosity term / scale of Coriolis term

$$=A_V \frac{U}{H^2}/fU = \frac{A_V}{fH^2}$$

measure of the importance of frictional force

 $E_k \ll 1$: friction can be neglected

Reynolds number (Re) = scale of the nonlinear term / scale of viscous term

$$= \frac{U^2}{L} / \nu \, \frac{U}{L^2} = \frac{UL}{\nu}$$

small Re: viscous flow

large Re: inviscous, turbulent flow

The momentum equation From the x-direction equation:

z direction:
$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} + f_* u = \frac{\frac{P}{\rho_0 L} \sim fU \sim \frac{U^2}{L} \sim A_H \frac{U}{L^2} \sim A_V \frac{U}{H^2}}{As H \ll L, \quad W \ll U}$$

$$\frac{W}{T} = \frac{UW}{L} = \frac{UW}{L} = \frac{W}{W} = \frac{P}{\rho_0 H} \gg fU$$

$$\frac{P}{\rho_0 H} \gg \frac{UW}{L}$$

$$-\frac{1}{\rho_0} \frac{\partial p}{\partial z} + \frac{\partial}{\partial x} (A_H \frac{\partial w}{\partial x}) + \frac{\partial}{\partial y} (A_H \frac{\partial w}{\partial y}) + \frac{\partial}{\partial z} (A_V \frac{\partial w}{\partial z}) - \frac{\rho}{\rho_0} g = \frac{P}{\rho_0 H} \gg A_H \frac{W}{L^2}$$

$$\frac{P}{\rho_0 H} = A_H \frac{W}{L^2} = A_V \frac{W}{L^2} = A_V \frac{W}{H^2} = \frac{\Delta \rho}{\rho_0} g = \frac{P}{\rho_0 H} \gg A_V \frac{W}{H^2}$$

$$\frac{\partial p}{\partial z} + \rho g = 0$$
 Hydrostatic equation (balance)

Hydrostatic balance: $\frac{\partial p}{\partial z} + \rho g = 0$

$$p = \tilde{p}(z) + p'(x, y, z, t) \qquad \tilde{p}(z) = P_0 - \rho_0 g z$$

$$\frac{\partial \tilde{p}}{\partial z} + \rho_0 g = 0 \qquad \frac{\partial p'}{\partial z} + \Delta \rho g = 0$$

$$\frac{P}{H} g \Delta \rho$$

$$\frac{P}{\rho_0 L} \sim \text{fU} \qquad \frac{gH\Delta\rho}{P} = \frac{gH\Delta\rho}{\rho_0 fLU} = \frac{U}{fL} \cdot \frac{gH\Delta\rho}{\rho_0 U^2}$$

$$= R_o \frac{N^2}{(U/H)^2} \quad \text{Richardson number}$$

The momentum equations

y direction:
$$\frac{\partial \mathbf{v}}{\partial \mathbf{t}} + \mathbf{u} \frac{\partial \mathbf{v}}{\partial \mathbf{x}} + \mathbf{v} \frac{\partial \mathbf{v}}{\partial \mathbf{y}} + \mathbf{w} \frac{\partial \mathbf{v}}{\partial \mathbf{z}} + fu = -\frac{1}{\rho_0} \frac{\partial p}{\partial y} + \frac{\partial}{\partial x} \left(A_H \frac{\partial \mathbf{v}}{\partial \mathbf{x}} \right) + \frac{\partial}{\partial y} \left(A_H \frac{\partial \mathbf{v}}{\partial y} \right) + \frac{\partial}{\partial z} \left(A_V \frac{\partial \mathbf{v}}{\partial z} \right) + \frac{\partial}{\partial z} \left(A_V \frac{\partial \mathbf{v}}{\partial z} \right) + \frac{\partial}{\partial z} \left(A_V \frac{\partial \mathbf{v}}{\partial z} \right)$$

z direction: $\frac{\partial p}{\partial z} + \rho g = 0$ Hydrostatic equation

$$\frac{\partial p'}{\partial z} + \Delta \rho g = 0$$

The thermodynamic equation

1st Law of Thermodyamics (energy conservation, and for per unit mass):

$$\frac{dI}{dt} = Q - W$$
 I: internal energy
 Q: rate of heat gain

W: rate of work done by pressure force

onto surrounding fluids

 C_{v} : specific heat capacity $I = C_{v}T$

$$I = C_{v}T$$

Fourier's Law of heat conduction:

$$Q = \frac{k_T}{\rho} \nabla^2 T$$

Rate of work done by pressure force:

specific volume

$$W = p \frac{d\alpha}{dt} \qquad \boxed{\alpha = 1/\rho}$$

$$C_{v} \frac{dT}{dt} = \frac{k_{T}}{\rho} \nabla^{2}T - p \frac{d\alpha}{dt}$$

$$= \frac{k_{T}}{\rho} \nabla^{2}T + \frac{p}{\rho^{2}} \frac{d\rho}{dt}$$

$$\alpha = 1/\rho$$

From the continuity equation:

$$\frac{d\rho}{dt} = -\rho \nabla \cdot \boldsymbol{v}$$

$$C_{\boldsymbol{v}} \frac{dT}{dt} = \frac{k_T}{\rho} \nabla^2 T - \frac{p}{\rho} \nabla \cdot \boldsymbol{v}$$

For incompressible fluids:

 K_T : thermal diffusivity coefficient

$$\frac{dT}{dt} = \frac{k_T}{\rho C_D} \nabla^2 T \qquad \qquad \frac{dS}{dt} = K_S \nabla^2 S$$

Equation of State

Specific volume:

$$\alpha = 1/\rho$$

$$\alpha = \alpha(T, S, P)$$

For small variations of α around a reference value:

Specific volume:
$$\alpha = 1/\rho$$

$$\alpha = \alpha(T, S, P)$$
 Shih et al. (2019)
$$d\alpha = \left(\frac{\partial \alpha}{\partial T}\right)_{S,p} dT + \left(\frac{\partial \alpha}{\partial S}\right)_{T,p} dS + \left(\frac{\partial \alpha}{\partial p}\right)_{T,S} dp = \alpha(\beta_T dT - \beta_S dS - \beta_p dp)$$

$$\beta_T: \text{thermal expansion coefficient}$$

$$\alpha = \alpha_0 \left[1 + \beta_T (T - T_0) - \beta_S (S - S_0) - \beta_p (p - p_0)\right]$$

$$\beta_S: \text{ salt contraction coefficient}$$

 $\rho = \rho_0 \left[1 - \beta_T (T - T_0) + \beta_S (S - S_0) + \beta_p (p - p_0) \right]$

 β_p : compressibility coefficient

nonlinear nonlinear
$$\alpha = \alpha_0 \left[1 + \beta_T (1 + y^* p) (T - T_0) + \frac{\beta_T^*}{2} (T - T_0)^2 - \beta_S (S - S_0) - \beta_p (p - p_0) \right]$$

TEOS-10

Thermodynamic Equation Of Seawater - 2010

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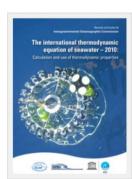
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This site is the official source of information about the Thermodynamic Equation Of Seawater - 2010 (TEOS-10), and the way in which it should be used.

TEOS-10 is based on a Gibbs function formulation from which all thermodynamic properties of seawater (density, enthalpy, entropy sound speed, etc.) can be derived in a thermodynamically consistent manner. <u>TEOS-10 was adopted by the Intergovernmental Oceanographic Commission at its 25th Assembly in June 2009</u> to replace EOS-80 as the official description of seawater and ice properties in marine science.

A significant change compared with past practice is that TEOS-10 uses Absolute Salinity S_A (mass fraction of salt in seawater) as opposed to Practical Salinity S_P (which is essentially a measure of the conductivity of seawater) to describe the salt content of seawater. Ocean salinities now have units of g/kg.

Absolute Salinity (g/kg) is an SI unit of concentration. The thermodynamic properties of seawater, such as density and enthalpy, are now correctly expressed as functions of Absolute Salinity rather than being functions of the conductivity of seawater. Spatial variations of the composition of seawater mean that Absolute Salinity is not simply proportional to Practical Salinity; TEOS-10 contains procedures to correct for these effects.

Summary of the governing equations

Momentum equation:
$$\frac{d\mathbf{v}}{dt} + f\mathbf{k} \times \mathbf{v} = -\frac{1}{\rho_0} \nabla p + \frac{\rho}{\rho_0} \mathbf{g} + \nu_E \nabla^2 \mathbf{v}$$

Continuity equation:
$$\frac{d\rho}{dt} + \rho \nabla \cdot \boldsymbol{v} = 0$$

Tracer equations:
$$\frac{dT}{dt} = K\nabla^2 T$$

$$\frac{dS}{dt} = K\nabla^2 S$$

Equation of state:

$$\rho = \rho(T, S, P)$$