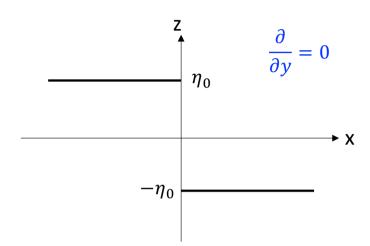
## Barotropic geostrophic adjustment

Initial state (unbalanced):

$$\eta = \begin{cases} \eta_0, & x < 0 \\ -\eta_0, & x > 0 \end{cases} \qquad \eta_0 << H$$

$$u = v = 0$$



For a steady final state (based on shallow water model):

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - fv = -g \frac{\partial \eta}{\partial x}$$

$$\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + fu = -g \frac{\partial \eta}{\partial y}$$

geostrophic flow

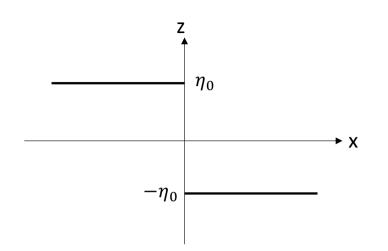
$$\frac{\partial \eta}{\partial t} + \frac{\partial hu}{\partial x} + \frac{\partial hv}{\partial y} = 0$$
 necessary condition: u=0 everywhere

### For PV conservation:

x < 0:

$$\frac{f}{H + \eta_0} = \frac{f + \frac{\partial v}{\partial x}}{H + \eta} = \frac{f + \frac{g}{f} \frac{\partial^2 \eta}{\partial x^2}}{H + \eta}$$

$$\eta_0 << H$$



$$fH + f\eta = fH + f\eta_0 + \frac{gH}{f} \frac{\partial^2 \eta}{\partial x^2} + \frac{g\eta_0}{f} \frac{\partial^2 \eta}{\partial x^2}$$

$$R^{2} \frac{\partial^{2} \eta}{\partial x^{2}} - \eta = -\eta_{0}$$

$$R = \frac{\sqrt{gH}}{f}$$

$$\eta = f(x) + \eta_{0}$$

$$e^{-\frac{1}{R}x} \to \infty \text{ as } x \to -\infty$$

$$f(x) = Ce^{\lambda x} \longrightarrow R^{2} \lambda^{2} - 1 = 0 \longrightarrow \lambda = \pm \frac{1}{R} \longrightarrow \lambda = \frac{1}{R}$$

$$\eta = Ce^{\frac{x}{R}} + \eta_0 \qquad x = 0, \eta = 0, C = -\eta_0$$

$$\eta = -\eta_0 e^{\frac{x}{R}} + \eta_0$$

### For PV conservation:

### x > 0:

$$\frac{f}{H - \eta_0} = \frac{f + \frac{\partial v}{\partial x}}{H + \eta} = \frac{f + \frac{g}{f} \frac{\partial^2 \eta}{\partial x^2}}{H + \eta}$$

$$fH + f\eta = fH - f\eta_0 + \frac{gH}{f} \frac{\partial^2 \eta}{\partial x^2} - \frac{g\eta_0}{f} \frac{\partial^2 \eta}{\partial x^2}$$

$$R^2 \frac{\partial^2 \eta}{\partial x^2} - \eta = \eta_0$$

$$\eta = f(x) - \eta_0$$

$$f(x) = Ce^{\lambda x} \longrightarrow$$

$$R^2\lambda^2 - 1 = 0$$

$$e^{\frac{1}{R}x} \to \infty \ as \ x \to \infty$$

$$f(x) = Ce^{\lambda x} \longrightarrow R^2 \lambda^2 - 1 = 0 \longrightarrow \lambda = \pm \frac{1}{R} \longrightarrow \lambda = -\frac{1}{R}$$

$$\eta = Ce^{-\frac{x}{R}} - \eta_0$$
 $x = 0, \eta = 0, C = \eta_0$ 
 $\eta = \eta_0 e^{-\frac{x}{R}} - \eta_0$ 

$$\eta = \begin{cases} -\eta_0 e^{\frac{x}{R}} + \eta_0, & x < 0 \\ \eta_0 e^{\frac{x}{R}} - \eta_0, & x > 0 \end{cases}$$

$$x \to R, \quad \eta \to -\eta_0$$

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## the adjustment spatial scale is the Rossby deformation radius

$$v = \frac{g}{f} \frac{\partial \eta}{\partial x} = \begin{cases} -\sqrt{\frac{g}{H}} \eta_0 e^{\frac{x}{R}}, & x < 0 \\ -\sqrt{\frac{g}{H}} \eta_0 e^{-\frac{x}{R}}, & x > \frac{1}{2} \end{cases}$$

$$x \to R, \quad v \to 0$$

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# Energetics of geostrophic adjustment

For unit length in the y-direction:

### Potential energy:

PE change is limited to  $(-R:R, -\eta_0: -\eta_0)$ 

$$PE_{I} = \int_{0}^{\eta_{0}} \int_{-R}^{0} \rho_{0} gz dx dz + \int_{0}^{-\eta_{0}} \int_{0}^{R} \rho_{0} gz dx dz = \frac{1}{2} \rho_{0} gR \eta_{0}^{2} + \frac{1}{2} \rho_{0} gR \eta_{0}^{2} = \rho_{0} gR \eta_{0}^{2}$$

$$PE_{F} = \int_{0}^{\eta} \int_{-R}^{0} \rho_{0} gz dx dz + \int_{0}^{\eta} \int_{0}^{R} \rho_{0} gz dx dz$$

$$= \frac{1}{2}\rho_0 g \eta_0^2 \int_{-R}^{0} (1 - e^{\frac{x}{R}})^2 dx + \frac{1}{2}\rho_0 g \eta_0^2 \int_{0}^{R} (1 - e^{\frac{x}{R}})^2 dx$$

$$\eta = \begin{cases} -\eta_0 e^{\frac{x}{R}} + \eta_0, & x < 0 \\ \eta_0 e^{-\frac{x}{R}} - \eta_0, & x > 0 \end{cases}$$

$$\Delta PE = -\frac{3}{2}\rho_0 gR\eta_0^2$$

$$=\frac{1}{2}\rho_0g\eta_0^2\left(x|_{-R}^0-2Re^{\frac{x}{R}}|_{-R}^0+\frac{R}{2}e^{\frac{2x}{R}}|_{-R}^0+x|_0^R+2Re^{\frac{x}{R}}|_0^R-\frac{R}{2}e^{\frac{2x}{R}}|_0^R\right)=-\frac{1}{2}\rho_0gR\eta_0^2$$

### Kinetic energy:

$$KE_I = 0$$

$$KE_{F} = \int_{-H}^{0} \int_{-R}^{R} \frac{1}{2} \rho_{0} v^{2} dx dz$$

$$\int_{-H}^{0} \int_{-R}^{0} \frac{1}{2} \rho_{0} v^{2} dx dz$$

$$= \int_{-H}^{0} \int_{-R}^{0} \frac{1}{2} \rho_0 \eta_0^2 \frac{g}{H} e^{\frac{2x}{R}} dx dz + \int_{-H}^{0} \int_{0}^{R} \frac{1}{2} \rho_0 \eta_0^2 \frac{g}{H} e^{\frac{-2x}{R}} dx dz$$

$$= \frac{1}{2}\rho_0 \eta_0^2 g(\int_{-R}^0 e^{\frac{2x}{R}} dx + \int_{-R}^0 e^{\frac{-2x}{R}} dx)$$

$$= \frac{1}{2}\rho_0 \eta_0^2 g(\frac{R}{2}e^{\frac{2x}{R}}|_{-R}^0 - \frac{R}{2}e^{-\frac{2x}{R}}|_0^R)$$

$$=\frac{1}{2}\rho_0 gR\eta_0^2$$

$$v = \begin{cases} -\sqrt{\frac{g}{H}} \eta_0 e^{\frac{x}{R}}, & x < 0 \\ -\sqrt{\frac{g}{H}} \eta_0 e^{-\frac{x}{R}}, & x > 0 \end{cases}$$

$$\Delta PE = -\frac{3}{2}\rho_0 gR{\eta_0}^2$$

$$\frac{\Delta KE}{|\Delta PE|} = 1/3$$

Only 1/3 of released PE is converted into KE, and the rest of lost PE is radiated away by waves