

Functional Analysis and PDE

Mounir HAJLI

Email address: hajli@sjtu.edu.cn

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ABSTRACT. Preliminary Lecture notes on

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CHAPTER 1

Banach spaces, Hilbert spaces

1.1. The Banach fixed-point theorem and the iteration method

- \mathbb{K} complete field, e.g. $\mathbb{R}, \mathbb{C}, \dots$
- \checkmark Banach space over \mathbb{K} .
- $A : M \rightarrow M$ an operator.
- M sub-space of X .

THEOREM 1.1. We assume that:

- (a) M is a closed nonempty set of X over \mathbb{K} ,
- (b) $A : M \rightarrow M$ is k -contractive, i.e.

$$\|Au - Av\| \leq k\|u - v\| \quad \forall u, v \in M,$$

k fixed, $0 \leq k < 1$.

Then,


- (1) The equation

$$(1.1) \quad Au = u \quad u \in M$$

has a unique solution. The solution of (1.1) is called *fixed point* of the operator A ,

- (2) The sequence

$$u_{n+1} = Au_n \quad \forall n \in \mathbb{N}$$

where $u_0 \in M$ converges to the fixed point of (1.1). 

@We consider the the variational problem

$$(1.2) \quad \min_{v \in V} J(v) = ?$$

where $J(v) = \frac{1}{2}a(v, v) - L(v)$, and the corresponding the variational equation

$$(1.3) \quad a(u, v) = L(v) \quad \text{for fixed } u \in V, \forall v \in V.$$

Boundary value problem \implies variational equation \Leftrightarrow variational problem.

1.2. Quadratic variational problems

@

We consider the minimum problem

$$(1.4) \quad \min_{v \in V} \left(\frac{1}{2} a(v, v) - L(v) \right) = ?$$

where

(i)

$$a : V \times V \longrightarrow \mathbb{R}$$

is a bounded, symmetric, strongly positive bilinear form, i.e. there exist positive constants μ, M such that

$$(1.5) \quad |a(w, v)| \leq M \|w\| \|v\| \quad a(v, v) \geq \mu \|v\|^2 \quad v, w \in V.$$

$$u \mapsto (u, v), \quad v \mapsto a(u, v)$$

are linear functionals.

(ii) The functional $L : V \rightarrow \mathbb{R}$ is linear and continuous.

Then,

(1) The variational problem (1.5) has a unique solution.

(2) The problem (1.5) is equivalent to the following so-called variational equation:

$$(1.6) \quad a(u, v) = L(v) \quad \text{for fixed } u \in V, \text{ and all } v \in V.$$

.

SKETCH OF THE PROOF. Let $u, v \in V$. We set

$$\phi(t) := \frac{1}{2} a(u + tv, u + tv) - L(u + tv) \quad t \in \mathbb{R}$$

Then,

$$(1.7) \quad \phi(t) = 2^{-1} t^2 a(v, v) + t(a(u, v) - L(v)) + 2^{-1} a(u, u) - L(u).$$

(a) u is a solution to (1.4) then $\phi(t) \geq \phi(0) \quad \forall t$. That is ϕ has a minimum at $t = 0$, i.e.

$$\phi'(0) = 0.$$

This equation is identical to

$$a(u, v) = L(v) \quad \forall v \in V.$$

(b) If u is solution of 1.6, then 1.7 becomes

$$\phi(t) = 2^{-1}t^2a(v, v) - 2^{-1}a(u, u).$$

Then,

$$\frac{1}{2}a(w, w) - L(w) \geq -2^{-1}a(u, u) (= \frac{1}{2}a(u, u) - L(u)) \quad \forall w \in V.$$

So, u solution of 1.4. □

1.3. The Dirichlet principle-Dirichlet problem I

@

Quadratic variational problems \implies (the generalized) Dirichlet problem

- We want to study the following **variational problem**:

$$(1.8) \quad \min_u J(u), \quad u = g \text{ on } \partial\Omega,$$

where $J(u) = 2^{-1} \int_{\Omega} \sum_{j=1}^N (\frac{\partial u}{\partial x_j})^2 dx - \int_{\Omega} f u dx$. This problem is called also **Dirichlet problem**.

•

$$(1.9) \quad \min_u J(u), \quad u - g \in H_0^1(\Omega),$$

called also **the generalized Dirichlet problem**.

The Dirichlet problem says that problem 1.8 (or 1.9) has a solution u .

@

- **Boundary-value problem for the Poisson equation**:

$$(1.10) \quad \begin{cases} -\Delta u = f & \text{on } \Omega \\ u = g & \text{on } \partial\Omega, \end{cases}$$

where

$$\Delta = \sum_{j=1}^N \frac{\partial^2}{\partial x_j^2}$$

1.10 is called **the classical Euler-Lagrange equation**.

•

$$\int_{\Omega} \nabla u \cdot \nabla v dx = \int_{\Omega} f v dx \quad \forall v \in \mathcal{C}_0^{\infty}(\Omega).$$

where

$$\nabla(\cdot) = \left(\frac{\partial}{\partial x_1}(\cdot), \dots, \frac{\partial}{\partial x_N}(\cdot) \right).$$

PROPOSITION 1.2. Let Ω be a nonempty bounded open set in \mathbb{R}^N . Let $g : \partial\Omega \rightarrow \mathbb{R}$, $f : \overline{\Omega} \rightarrow \mathbb{R}$ continuous functions. Suppose that $u \in \mathcal{C}^2(\overline{\Omega})$. Then,

Variational problem 1.8 \implies the boundary-value problem 1.10.

PROOF. Let u be a solution of 1.8. Let $v \in \mathcal{C}_0^{\infty}(\Omega)$. Let $t \in \mathbb{R}$. The function

$$w : u + tv$$

is **admissible** for the variational problem 1.8, i.e. $w = g$ on $\partial\Omega$, and $w \in \mathcal{C}^2(\Omega)$.

We set

$$\phi(t) = J(u + tv).$$

Explicitly,

$$\phi(t) = 2^{-1} \int_{\Omega} \sum_j (\partial_j u + t \partial_j v)^2 dx - \int_{\Omega} f \cdot (u + tv) dx.$$

The quadratic function $\phi : \mathbb{R} \rightarrow \mathbb{R}$ has a minimum at $t = 0$. Hence,

$$\phi'(0) = 0$$

That is

$$\int_{\Omega} \nabla u \cdot \nabla v dx - \int_{\Omega} f v dx = 0 \quad \forall v \in \mathcal{C}_0^{\infty}(\Omega)$$

□

@

EXERCISE 1.3. Consider the minimum problem

$$\min_u J(u), \quad u \in \mathcal{C}^1[-1, 1], u(-1) = 0, u(1) = 1.$$

where

$$J(u) = \int_{-1}^1 (xu'(x))^2 dx.$$

Use the sequence

$$u_n(x) = \frac{1}{2} + \frac{1}{2} \frac{\arctan(nx)}{\arctan(n)}, \quad n = 1, 2, \dots$$

to show that this variational problem has no solution.

1.4. Linear Lax-Milgram theory

We consider the **the variational problem**

$$(1.11) \quad \min_{v \in V} J(v) = ?$$

where $J(v) = \frac{1}{2}a(v, v) - L(v)$, and the corresponding **the variational equation**

$$(1.12) \quad a(u, v) = L(v) \quad \text{for fixed } u \in V, \forall v \in V.$$

We assume the following:

- (i) V is a real separable Hilbert space, i.e, there exists a sequence $(V_n)_{n \in \mathbb{N}}$ of finite dimensional linear subspaces V_n of V such that

$$\lim_n \text{dist}_V(u, V_n) = 0 \quad \forall u \in V.$$

- (ii)

$$a : V \times V \longrightarrow \mathbb{R}$$

is a bounded, symmetric, and strongly positive, i.e. there exist positive constants μ, M such that

$$(1.13) \quad |a(w, v)| \leq M \|w\| \|v\| \quad a(v, v) \geq \nu \|v\|^2 \quad v, w \in V.$$

- (iii) The functional $L : V \rightarrow \mathbb{R}$ is linear and continuous.

Let us consider the following problem

$$(1.14) \quad \min_{v \in V_n} \frac{1}{2} a(v, v) - L(v),$$

and

$$(1.15) \quad a(u_n, v) = L(v) \quad \text{for fixed } u_n \in V_n \text{ and for any } v \in V_n.$$

Let $\{e_{n,1}, \dots, e_{n,N_n}\}$ be a basis of V_n ($N_n = \dim V_n$). (1.15) is equivalent to the following system of linear equations:

$$(1.16) \quad \sum_{i=1}^{N_n} \alpha_{n,i} a(e_{n,i}, e_{n,j}) = L(e_{n,j}) \quad j = 1, \dots, N_n.$$

where $u_n = \sum_{i=1}^{N_n} \alpha_{n,i} e_{n,i}$ (we take $v = e_{n,i}$ in (1.15)).

From (1.13), the following matrix

$$(a(e_{n,i}, e_{n,j}))_{1 \leq i, j \leq N_n}$$

is invertible.

We have

- (a) (1.11) has a unique solution u . This is also the unique solution of the variational equation (1.12). We have **a prior estimate**

$$(1.17) \quad \|u\| \leq \nu^{-1} \|L\|.$$

- (b) u_n is a solution for (1.16), and

$$\|u_n - u\| \leq \frac{M}{\nu} \text{dist}(u, X_n),$$

(approximation of u by u_n .)

1.4.1. Applications to Boundary-Value problems, the method of finite elements, and Elasticity. @ Let us consider the following boundary-value problem:

$$(1.18) \quad \begin{cases} -u''(x) = f(x) & \text{if } \alpha < x \leq \beta \\ u(\alpha) = u(\beta) = 0 & u \in \mathcal{C}^2[\alpha, \beta]. \end{cases}$$

where $-\infty < \alpha < \beta < +\infty$. **In other words, $u \in \mathcal{C}_0^2[\alpha, \beta]$.**

The corresponding variational problem is

$$(1.19) \quad \begin{cases} \min_u J(u), \\ \text{under the conditions } u(\alpha) = u(\beta) = 0 \text{ and } u \in \mathcal{C}^2[\alpha, \beta]. \end{cases}$$

where $J(u) = \int_{\alpha}^{\beta} (\frac{1}{2}(u')^2 - uf)dx$.

- $u(x)$ = **deflection** of a string at the point x under the vertical outer force density $\mathbf{f}(x) = f(x)\mathbf{e}$.
- $J(u)$ = **total potential energy** of the string.
- $\int_{\alpha}^{\beta} \frac{1}{2} u'^2 dx$ elastic energy of the string.
- $L(u) := \int_{\alpha}^{\beta} f u dx$ = work of the outer force density \mathbf{f} with respect to the vertical displacement $\mathbf{u}(x) = u(x)\mathbf{e}$ of the string, which is also $-(\text{potential energy stored by the force density } f)$.

Let us consider the (generalized) variational problem

$$(1.20) \quad \begin{cases} \min_u J(u), \\ \text{under the conditions } u \in H_0^1(\alpha, \beta). \end{cases}$$

Recall (Definition 1.26)

$$(H_0^1(\alpha, \beta), \|\cdot\|_{H_0^1(\alpha, \beta)})$$

is a real Hilbert space, with the associated scalar product

$$(u, v)_{H_0^1(\alpha, \beta)} = \int_{\alpha}^{\beta} (uv + u'v') dx, \quad \|u\|_{H_0^1(\alpha, \beta)} = \left(\int_{\alpha}^{\beta} (u^2 + (u')^2) dx \right)^{\frac{1}{2}}.$$

We set

$$(u, v)_E = \int_{\alpha}^{\beta} u'v' dx, \quad \|u\|_E = (u, v)_E^{\frac{1}{2}},$$

called the **the energetic inner product** and **the energetic norm**¹. Obviously,

$$2^{-1}(u, u)_{H_0^1(\alpha, \beta)} = \text{elastic potential energy with respect to the displacement } u$$

The variational equation reads as follows

$$(1.21) \quad \frac{dJ(u + tv)}{dt} \Big|_{t=0} = \int_{\alpha}^{\beta} (u'v' - fv) dx = 0 \quad \forall v \in H_0^1(\alpha, \beta)$$

This means the following: The time derivative of the energy $J(u + tv)$ at $t = 0$ equals zero.

¹This is, indeed, a norm which moreover is EQUIVALENT to $\|\cdot\|_{H_0^1(\alpha, \beta)}$. This is a consequence, of 1.24.

1.4.2. The Poincaré-Friedrichs Inequality. Let $f \in \mathcal{C}^1[\alpha, \beta]$. We have, for any $\alpha \leq x \leq \beta$,

$$\begin{aligned} f(x)^2 &= f(x)^2 - f(\alpha)^2 + f(\alpha)^2 \\ &= 2 \int_{\alpha}^x f'(y)f(y)dy + f(\alpha)^2 \\ &\leq 2 \left(\int_{\alpha}^x (f'(y))^2 dy \right)^{\frac{1}{2}} \left(\int_{\alpha}^x f(y)^2 dy \right)^{\frac{1}{2}} + f(\alpha)^2 \\ &\leq 2 \left(\int_{\alpha}^{\beta} (f'(y))^2 dy \right)^{\frac{1}{2}} \left(\int_{\alpha}^{\beta} f(y)^2 dy \right)^{\frac{1}{2}} + f(\alpha)^2 \end{aligned}$$

That is

$$f(x)^2 \leq 2 \left(\int_{\alpha}^{\beta} (f'(y))^2 dy \right)^{\frac{1}{2}} \left(\int_{\alpha}^{\beta} f(y)^2 dy \right)^{\frac{1}{2}} + f(\alpha)^2.$$

Then

$$(1.22) \quad \int_{\alpha}^{\beta} f(x)^2 dx \leq 2(\beta - \alpha) \left(\int_{\alpha}^{\beta} (f'(y))^2 dy \right)^{\frac{1}{2}} \left(\int_{\alpha}^{\beta} f(y)^2 dy \right)^{\frac{1}{2}} + (\beta - \alpha)f(\alpha)^2.$$

It follows that

$$(1.23) \quad \left(\int_{\alpha}^{\beta} f(x)^2 dx \right)^{\frac{1}{2}} \leq (\beta - \alpha) \left(\int_{\alpha}^{\beta} (f'(y))^2 dy \right)^{\frac{1}{2}} + \sqrt{(\beta - \alpha)^2 \int_{\alpha}^{\beta} (f'(y))^2 dy + (\beta - \alpha)f(\alpha)^2}.$$

Put $X = \left(\int_{\alpha}^{\beta} f(x)^2 dx \right)^{\frac{1}{2}}$, then 1.22 becomes $X^2 \leq 2aX + b..$

We have proved the following lemma

LEMMA 1.4 (The Poincaré-Friedrichs inequality for regular functions in dimension 1).
Let $f \in \mathcal{C}_0^1[\alpha, \beta]$. We have

$$(\beta - \alpha)^{-1} \int_{\alpha}^{\beta} f(x)^2 dx \leq \int_{\alpha}^{\beta} (f'(x))^2 dx$$

Now, let $u \in H_0^1(\alpha, \beta)$, (see Definition 1.26). By definition, $\exists(u_n) \subset \mathcal{C}_0^{\infty}(\alpha, \beta)$ such that

$$u_n \rightarrow u \text{ in } H_0^1(\alpha, \beta),$$

as $n \rightarrow \infty$. We have

$$(\beta - \alpha)^{-1} \int_{\alpha}^{\beta} u_n(x)^2 dx \leq \int_{\alpha}^{\beta} (u'_n(x))^2 dx \quad \forall n \in \mathbb{N}.$$

Letting $n \rightarrow \infty$, we get

$$(1.24) \quad (\beta - \alpha)^{-1} \int_{\alpha}^{\beta} u^2 dx \leq \int_{\alpha}^{\beta} (u')^2 dx \quad \forall u \in H_0^1(\alpha, \beta)$$

@ More generally, we have the Poincaré-Friedrichs inequality:

PROPOSITION 1.5. Let Ω be a nonempty bounded open set in \mathbb{R}^N . There there exists a constant $C > 0$ such that

$$(1.25) \quad C \int_{\Omega} u^2 dx \leq \int_{\Omega} \left(\sum_{j=1}^N \frac{\partial u}{\partial x_j} \right)^2 dx \quad \forall u \in H_0^1(\Omega).$$

Application: With the notations above, we have

$$J(u) = 2^{-1}(u, u)_E - L(u) \quad \forall u \in H_0^1(\alpha, \beta).$$

We see that $a(u, v) := (u, v)_E$ is a bounded, symmetric, strongly positive bilinear form. By Theorem 1.13, the problem 1.20 has a unique solution $u \in H_0^1(\alpha, \beta)$, and hence

$$-\Delta u = f \quad \text{on } H_0^1(\alpha, \beta).$$

EXERCISE 1.6. Let $G(x, y)$ be the function defined on $[\alpha, \beta] \times [\alpha, \beta]$ as follows

$$G(x, y) = \begin{cases} \frac{(\beta-y)(x-\alpha)}{\beta-\alpha} & \text{if } \alpha \leq x \leq y \leq \beta \\ \frac{(\beta-x)(y-\alpha)}{\beta-\alpha} & \text{if } \alpha \leq y < x \leq \beta, \end{cases}$$

Show that the function u given by

$$u(x) = \int_{\alpha}^{\beta} G(x, y) f(y) dy.$$

is a solution of 1.18.

Hint: Show that

$$G(x, y) = G(y, x).$$

Let $\phi \in \mathcal{C}_0^{\infty}(\alpha, \beta)$. Compute

$$\int_{\alpha}^{\beta} G(x, y) \phi''(y) dy = ?$$

Use Fubini's theorem to determine

$$\int_{\alpha}^{\beta} u(x) \phi''(x) dx = ?$$

1.5. Generalized derivatives=Weak differentiation

@ Ω open subset of \mathbb{R}^n .

$$L^2(\Omega) = \left\{ u : \Omega \rightarrow \mathbb{R} \cup \{\infty\} \mid u \text{ measurable } \int_{\Omega} |u|^2 dx < \infty \right\}$$

$$\mathcal{C}_0^{\infty}(\Omega),$$

denotes the space of smooth functions on Ω with compact support in Ω . Also denoted by $\mathcal{C}_c^{\infty}(\Omega)$.

DEFINITION 1.7. Let $u, w \in L^2(\Omega)$. Assume that

$$\int_{\Omega} u \partial_j v dx = - \int_{\Omega} w v dx \quad \forall v \in \mathcal{C}_0^{\infty}(\Omega).$$

Then w is called the **generalized derivative** (or **the weak partial derivative**) of the function u on the set Ω of type ∂_j . We write

$$\frac{\partial u}{\partial x_j} := w.$$

EXAMPLE 1.8. Consider the function $u(x) = |x|$ for $x \in (-1, 1)$.

$$u(x) = \begin{cases} u(x) = f(x) & \text{if } -1 < x \leq 0 \\ u(x) = g(x) & \text{if } 0 \leq x < 1, \end{cases}$$

where f and g are differentiable functions such that $f \in L^2(-1, 0)$, $g \in L^2(0, 1)$ and $f(0) = g(0)$. We have, for any $v \in C_0^\infty((-1, 1))$,

$$\begin{aligned} \int_{-1}^1 uv' dx &= \int_{-1}^0 uv' dx + \int_0^1 uv' dx \\ &= - \int_{-1}^0 u'v dx - \int_0^1 u'v dx + [uv]_{-1}^0 + [uv]_0^1 \\ &= - \int_{-1}^0 f'v dx - \int_0^1 g'v dx + f(0)v(0) - g(0)v(0) \\ &= - \int_{-1}^1 wv dx \end{aligned}$$

where w is in $L^2(-1, 1)$.

EXAMPLE 1.9. Recall the definition of G in Exercise 1.6. We have the following equality of currents

$$-\frac{\partial^2}{\partial x^2} G(x, y) = \delta_y$$

equivalently

$$-\Delta G(x, y) = \delta_y$$

where δ_y is the Dirac measure at y , i.e. δ_y is a "generalized function"="distribution"="continuous functional" such that

$$\delta_y(\phi) = \phi(y) \quad \forall \phi \in C_0^\infty(\alpha, \beta)$$

@

DEFINITION 1.10.

(i) The Sobolev space:

$$H^1(\Omega) := \left\{ u \in L^2(\Omega) \mid \forall i = 1, \dots, N, \frac{\partial u}{\partial x_i} \in L^2(\Omega) \right\}$$

also denoted by $W_2^1(\Omega)$. This is a Hilbert space for the scalar product

(a)

$$(u, v) = \int_{\Omega} (uv + \nabla u \cdot \nabla v) dx.$$

(b) The associated norm,

$$\|u\|_{H^1(\Omega)} := \left(\int_{\Omega} (|u|^2 + |\nabla u|^2) dx \right)^{1/2}$$

(ii) The space $H_0^1(\Omega)$.

$$(1.26) \quad \textcolor{red}{H}_0^1(\Omega) := \overline{\mathcal{C}_0^\infty(\Omega)} (\subset H^1(\Omega))$$

the closure wrt the metric of $H^1(\Omega)$.

EXAMPLE 1.11. $\Omega = (a, b) \subset \mathbb{R}$. If $u \in H_0^1(a, b)$, then there exists a unique continuous function $v : [a, b] \rightarrow \mathbb{R}$ such that $u(x) = v(x)$ for almost $x \in (a, b)$ and

$$v(a) = v(b) = 0.$$

1.6. Nonlinear Lax-Milgram theory

@

We want to solve the nonlinear operator equation

$$(1.27) \quad Au = z, \quad u \in V,$$

where

(1) $A : V \rightarrow A$ is an operator on the real Hilbert space V , such that

$$(Au - Av, u - v) \geq \nu \|u - v\|^2 \quad \forall u, v \in V$$

(2)

$$\|Au - Av\| \leq L \|u - v\| \quad \forall u, v \in V.$$

THEOREM 1.12. *For each given $z \in V$, problem (1.27) has a unique solution u .*

PROOF. Let $t > 0$, and we consider the operator

$$Bu := u - t(Au - z), \quad u \in V.$$

We have

$$\textcolor{blue}{Bu = u} \Leftrightarrow \textcolor{blue}{Au = z}.$$

For any $u, v \in V$,

$$\begin{aligned} \|Bu - Bv\|^2 &= \|u - v\|^2 - 2t(Au - Av, u - v) + t^2 \|Au - Av\|^2 \\ &\leq (1 - 2\nu t + L^2 t^2) \|u - v\|^2 \\ &= \left(1 - L^2 t \left(\frac{2\nu}{L^2} - t \right) \right) \|u - v\|^2 \end{aligned}$$

So, if $0 < t < \frac{2\nu}{L^2}$, then

$$k := (1 - 2\nu t + L^2 t^2) < 1$$

In this case,

$$\|Bu - Bv\| \leq \sqrt{k}\|u - v\| < \|u - v\| \quad \forall u, v \in V.$$

This means that B is \sqrt{k} -contractive. We conclude by using Theorem 1.1. □

For each given $u_0 \in V$ and $t \in (0, 2\nu/L^2)$, the *iteration method*

$$u_{n+1} = u_n - t(Au_n - z), \quad n = 1, 2, \dots$$

converges to the unique solution u of 1.27. We have the error estimate

$$\|u - u_n\| \leq k^{n/2}(1 - \sqrt{k})^{-1}\|u_1 - u_0\| \quad \forall n \in \mathbb{N}.$$

1.6.1. Application to the Nonlinear Lax-Milgram Theorem. @

Let

- (i) $b : V \rightarrow \mathbb{R}$ linear continuous functional on the real Hilbert space V .
- (ii) $a : V \times V \rightarrow \mathbb{R}$ a function such that, for each $w \in V$,

$$v \mapsto a(w, v)$$

is a linear continuous functional on V .

- (iii) $\exists \nu, L > 0$ such that

$$a(u, u - v) - a(v, u - v) \geq \nu\|u - v\|^2,$$

and

$$|a(u, w) - a(v, w)| \leq L\|u - v\|\|w\|,$$

for all $u, v, w \in V$.

THEOREM 1.13. *The equation*

$$(1.28) \quad a(u, v) = L(v) \quad \forall v \in V$$

has a unique solution u .

PROOF. By Riesz theorem, for each $w \in X$, there is an element called Aw such that

$$a(w, u) = (Aw, u) \quad \forall u \in V.$$

Then (iii) becomes

$$(Au, u - v) - (Av, u - v) \geq \nu\|u - v\|^2,$$

and

$$|(Au - Av, w)| \leq L\|u - v\|\|w\|.$$

Again by Riesz theorem. there is a $z \in V$ such that

$$L(u) = (z, u) \quad \forall u \in V$$

Consequently, (1.28) is equivalent to the operator equation

$$Au = z, \quad u \in V.$$

We conclude by Theorem 1.12. □

1.7. The Dirichlet principle-Dirichlet problem 2

- Let us consider the generalized Dirichlet problem

$$(1.29) \quad \min_u 2^{-1} \int_{\Omega} \sum_{j=1}^N \left(\frac{\partial u}{\partial x_j} \right)^2 dx - \int_{\Omega} f u dx, \quad u - g \in H_0^1(\Omega)^2,$$

- and the generalized boundary-value problem:

$$(1.30) \quad \int_{\Omega} \nabla u \cdot \nabla v dx = \int_{\Omega} f v dx \quad \forall v \in H_0^1(\Omega), u - g \in H_1^0(\Omega).$$

THEOREM 1.14 (Dirichlet principle). *Let Ω be a nonempty bounded open subset of \mathbb{R}^N . Let $f \in L^2(\Omega)$, $g \in H_0^1(\Omega)$. Then the following hold true:*

- (i) *1.29 has unique solution $u \in H^1(\Omega)$.*
- (ii) *This is also the unique solution $u \in H^1(\Omega)$ of 1.30.*

PROOF. Let $V = H^1(\Omega)$, and

$$a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v dx, \quad L(v) = \int_{\Omega} f v dx.$$

□

²This condition is a generalization of the following : $u = g$ on $\partial\Omega$

CHAPTER 2

Introduction to mathematical modeling

2.1. Some classical models; Quelques modèles classiques

Ω is a domain in \mathbb{R}^N .

2.1.1. The convection-diffusion equation.

$$(2.1) \quad \begin{cases} c \frac{\partial \theta}{\partial t} + cV \cdot \nabla \theta - k \Delta \theta = f & \text{in } \Omega \times \mathbb{R}_*^+ \\ \theta = 0 & \text{on } \partial\Omega \times \mathbb{R}_*^+ \\ \theta(\textcolor{red}{t} = \textcolor{red}{0}, x) = \theta_0(x) & \text{in } \Omega. \end{cases}$$

$V(t, x)$ velocity (a vector valued function in \mathbb{R}^N).

2.1.1.1. *The convection-diffusion equation* *The case $N = 1, f \equiv 0$.*

$$(2.2) \quad \begin{cases} \frac{\partial \theta}{\partial t} + V \frac{\partial \theta}{\partial x} - \nu \frac{\partial^2 \theta}{\partial x^2} = 0 & \text{in } \mathbb{R} \times \mathbb{R}_*^+ \\ \theta(t = 0, x) = \theta_0(x) & \text{in } \mathbb{R}. \end{cases}$$

with $\nu = k/c$. 2.2 has the following solution

$$(2.3) \quad \theta(t, x) = \frac{1}{\sqrt{4\pi\nu t}} \int_{-\infty}^{\infty} \theta_0(y) \exp\left(-\frac{(x - Vt - y)^2}{4\nu t}\right) dy.$$

2.1.2. The advection equation. $k = 0$ in 2.1.

$$(2.4) \quad \begin{cases} c \frac{\partial \theta}{\partial t} + cV \cdot \nabla \theta = f & \text{in } \Omega \times \mathbb{R}_*^+ \\ \theta = 0 & \text{on } \partial\Omega \times \mathbb{R}_*^+ \text{ if } V(x) \cdot n(x) < 0 \\ \theta(t = 0, x) = \theta_0(x) & \text{in } \Omega. \end{cases}$$

2.1.2.1. *The advection equation, The case of $N = 1, f \equiv 0$.*

$$(2.5) \quad \begin{cases} \frac{\partial \theta}{\partial t} + V \frac{\partial \theta}{\partial x} = 0 & \text{in } \mathbb{R} \times \mathbb{R}_*^+ \\ \theta(t = 0, x) = \theta_0(x) & \text{in } \mathbb{R}. \end{cases}$$

$$\theta(t, x) = \theta_0(x - Vt)$$

is a solution of (2.5).

EXERCISE 2.1. Let φ be a smooth function with support in $[-1, 1]$. Show that

$$\lim_{\mu \rightarrow 0^+} \frac{1}{\sqrt{4\pi\mu}} \int_{-\infty}^{\infty} \varphi(y) \exp\left(-\frac{y^2}{4\mu}\right) dy = \varphi(0).$$

PROPOSITION 2.2 (The maximum principle). Let $\theta(t, x)$ be solution of (2.2) or (2.5). We have

$$(2.6) \quad \min_{x \in \mathbb{R}} \theta_0(x) \leq \theta(t, x) \leq \max_{x \in \mathbb{R}} \theta_0(x) \quad \forall (x, t) \in \mathbb{R} \times \mathbb{R}^+.$$

2.1.3. The heat flow equation.

$$(2.7) \quad \begin{cases} \frac{\partial u}{\partial t} - \Delta u = f & \text{in } \Omega \times \mathbb{R}_*^+ \\ u = 0 & \text{on } \partial\Omega \times \mathbb{R}_*^+ \\ u(0, x) = u_0(x) & \text{in } \Omega. \end{cases}$$

We have

$$(2.8) \quad \frac{1}{2} \frac{\partial}{\partial t} \left(\int_{\Omega} u^2(t, x) dx \right) = \int_{\Omega} u(t, x) f(x) dx + \int_{\Omega} u \Delta u dx$$

$$\begin{aligned} \frac{1}{2} \frac{\partial}{\partial t} \left(\int_{\Omega} u^2(t, x) dx \right) &= \int_{\Omega} \frac{1}{2} \frac{\partial}{\partial t} u^2(t, x) dx \\ &= \int_{\Omega} u(t, x) \frac{\partial}{\partial t} u(t, x) dx \\ &= \int_{\Omega} u(t, x) (f(x) + \Delta u) dx \\ &= \int_{\Omega} u(t, x) f(x) dx + \int_{\Omega} u \Delta u dx \end{aligned}$$

When $\Omega = (0, 1)$ and $f = 0$, this equation becomes

$$(2.9) \quad \begin{cases} \frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = 0 & \text{in } \mathbb{R}_*^+ \times \Omega \\ u = 0 & \text{on } \partial\Omega \times \mathbb{R}_*^+ \\ u(t = 0, x) = u_0(x) & \text{in } \Omega. \end{cases}$$

In this case (2.8) becomes

$$\frac{1}{2} \frac{\partial}{\partial t} \left(\int_{\Omega} u^2(t, x) dx \right) = \int_{\Omega} u \frac{\partial^2 u}{\partial x^2} dx = \left[u \frac{\partial u}{\partial x} \right]_{\partial\Omega} - \int_{\Omega} \left(\frac{\partial u}{\partial x} \right)^2 dx = - \int_{\Omega} \left(\frac{\partial u}{\partial x}(t, x) \right)^2 dx.$$

EXERCISE 2.3. Let $\Omega = (0, 1)$ Let $v \in \mathcal{C}^1[0, 1]$ such that $v(0) = 0$. Show that

$$\int_{\Omega} v^2(x) dx \leq \int_{\Omega} \left| \frac{dv}{dx} \right|^2 dx$$

PROOF. Let $x \in \Omega$.

$$\begin{aligned}
 v(x)^2 &= v(x)^2 - v(0)^2 \\
 &= 2 \int_0^x v'(y)v(y)dy \\
 &\leq 2 \left(\int_0^x (v'(y))^2 dy \right)^{\frac{1}{2}} \left(\int_0^x v(y)^2 dy \right)^{\frac{1}{2}} \quad \text{Cauchy-Schwartz} \\
 &\leq 2 \left(\int_0^1 (v'(y))^2 dy \right)^{\frac{1}{2}} \left(\int_0^1 v(y)^2 dy \right)^{\frac{1}{2}}.
 \end{aligned}$$

□

2.1.4. The wave equation.

2.1.5. The Laplacian.

2.1.6. Schrödinger's equation.

2.1.7. The Lamé equation.

2.1.8. The Stokes [system](#).

2.1.9. The plate equations. Let u

2.2. Numerical calculation by finite differences

@

Let $(E, \|\cdot\|_E), (F, \|\cdot\|_F)$ be two Banach spaces, e.g. $E = \mathbb{R}^k, F = \mathbb{R}^l$. Let Ω be a nonempty open subset of E , and $f : \Omega \rightarrow F$ a \mathcal{C}^p -function. We have the Taylor's expansion of f at x ,

$$f(x+h) = f(x) + f'(x) \cdot h + \frac{f''(x)}{2!} \cdot h^{(2)} + \dots + \frac{f^{(p-1)}(x)}{(p-1)!} h^{(p-1)} + \int_0^1 \frac{(1-\xi)^{p-1}}{(p-1)!} f^{(p)}(x+\xi t) \cdot h^{(p)} dt$$

where $h^{(j)}$ stands for (h, \dots, h) . From this, we get

$$\left\| f(x+h) - \left(f(x) + f'(x) \cdot h + \frac{f''(x)}{2!} \cdot h^{(2)} + \dots + \frac{f^{(p)}(x)}{p!} h^{(p)} \right) \right\|_F \leq \epsilon \|h\|_E^p,$$

for $\epsilon > 0$ and h sufficiently small. Let $\epsilon > 0$. The continuity of $f^{(p)}$ implies of the continuity of $t \mapsto f(x+\xi t)$ for any $t \in \mathbb{R}$ such that $x+\xi t \in \Omega$. Let $\epsilon > 0$, there exists a $\eta > 0$ such that

$$|t| < \eta \implies \|f(x+\xi t) - f(x)\|_F \leq \epsilon$$

CHAPTER 3

Sobolev spaces

Ⓐ

Ω open subset of \mathbb{R}^n , $s \in \mathbb{N}$.

$$(3.1) \quad H^s(\Omega) = \{u \in L^2(\Omega) \mid D^\alpha u \in L^2(\Omega) \ \forall |\alpha| \leq s\}$$

is called **the Sobolev space**, where D^α is the derivative of u in the sense of distributions.

$$H^s(\Omega) \subset L^2(\Omega), \quad H^0(\Omega) = L^2(\Omega).$$

$$(3.2) \quad \|u\|_s = \left(\sum_{|\alpha| \leq s} \|D^\alpha u\|_0^2 \right)^{1/2}$$

where $\|\cdot\|$ denotes the norm on $L^2(\Omega)$.

$$(u, v)_s = \sum_{|\alpha| \leq s} (D^\alpha u, D^\alpha v)_0$$

where (\cdot, \cdot) is the scalar product on L^2 .

PROPOSITION 3.1. Let $s \in \mathbb{N}$, $H^s(\Omega)$ is a Hilbert space.

CHAPTER 4

La formule de Taylor

@

Soit f une fonction de classe \mathcal{C}^k . Soit $j \leq k$

$$f(x+h) = f(x) + f'(x)h + \frac{1}{2}f''(x)h^2 + \cdots + \frac{1}{j!}f^{(j)}(x)h^j + R_j(h).$$

$$f(x+h) = \sum$$

