The action is:

$$S = \int d^{4}x \int -g \left[\frac{1}{2} (\nabla \mu a) (\nabla \mu a) - V \left(\frac{a}{\ell} \right) \right] \quad \text{with} \quad V(x) = \frac{\Lambda^{4}}{c_{3}} \left(1 - \sqrt{1 - 4c_{3} \sin^{2} \frac{x}{2}} \right)$$

Metric:
$$d^2 x = (1+2\emptyset) dt^2 - (1-2\emptyset) dr^2 - r^2 d\Omega^2$$

$$\Rightarrow \partial_{\mu\nu} = \begin{pmatrix} 1+2\emptyset \\ -(1-2\emptyset) \\ -n^2 \sin^2 \theta \end{pmatrix}$$

Christoffel symbols (from Mathematica's notebook "christoffel. nb") of the form $\Gamma_{r,r}^{r,\mu}$:

$$\Gamma_{nn}^{n} = \underbrace{g'}_{-1+2g} \simeq -g'(1+2g)$$

$$\int_{00}^{\pi} = \frac{n}{-1+20} = -n \left(1+20\right)$$

$$\Gamma_{ee}^{r} = \frac{\pi \, \text{Nm}^2 \theta}{-1 + 20} \cong -\pi \, \text{Nm}^2 \theta \left(1 + 20\right)$$

$$\int_{++}^{1} = \emptyset' \cong \emptyset' (1+2\emptyset)$$

$$\int_{ee}^{\theta} = - \cosh \theta \sin \theta$$

$$\Gamma_{nn}^{t} = -\frac{\dot{\phi}}{1+2\phi} \approx -\dot{\phi}(1-2\phi)$$

$$\Gamma_{tx}^{t} = \frac{\dot{\phi}}{1+2\phi} \cong \dot{\phi} (1-2\phi)$$

Now, define $\alpha = \frac{a}{fa}$ and make $r \rightarrow \tilde{r} = mr \ \ell \ t \rightarrow \tilde{t} = mt$:

$$\mathcal{L} = \int \frac{d^4 \tilde{x}}{m^4} \sqrt{-g} \left[\int_{z}^{2} m^2 \left(\tilde{\nabla} \mu_{\alpha} \right) \left(\tilde{\nabla} \mu_{\alpha} \right) - V(\alpha) \right]$$

$$\Rightarrow \mathcal{L} = \int \frac{d^4 x}{m^4} \sqrt{-g} \Lambda^4 \left[\frac{1}{2} (\nabla \mu_{\alpha}) (\nabla \mu_{\alpha}) - \sqrt{(\alpha)} \right] \frac{\nabla(\alpha)}{\Lambda^4}$$

Equation of motion for scalar in curved spacetime:

$$\square = -\frac{9\tilde{V}}{9\alpha}, \text{ with } \square = g^{\mu\nu} \nabla_{\mu} \nabla_{\nu}$$

We have

$$\Box \alpha = g^{\mu\nu} \nabla_{\mu} \partial_{\nu} \alpha = g^{\mu\nu} (\partial_{\mu} \partial_{\nu} \alpha - \Gamma_{\mu\nu} \partial_{\beta} \alpha)$$

$$= g^{tt} (\ddot{\alpha} - \Gamma_{tt}^{\beta} \partial_{\beta} \alpha) + g^{rr} (\alpha'' - \Gamma_{rr}^{\beta} \partial_{\beta} \alpha) - (g^{\theta\theta} \Gamma_{\theta\theta}^{\beta} + g^{\theta\theta} \Gamma_{\theta\theta}^{\beta}) \partial_{\beta} \alpha$$
examing spherically symmetrical $\alpha = \alpha(t, r)$

$$\simeq (1-2\emptyset)[\dot{a}-(1-2\emptyset)\dot{\partial}\dot{\alpha}-(1+2\emptyset)\emptyset'\alpha']$$

$$-\frac{1}{r^2}$$
 $\times (1+2\emptyset) \alpha' - \frac{\cos^2\theta}{r^2} \times \sin^2\theta (1+2\emptyset) \alpha'$

$$= (1-2\emptyset)\ddot{\alpha} - (1+2\emptyset) \sim "-2\ddot{\phi}\dot{\alpha} - 2\ddot{\phi}'\alpha' - (1+2\emptyset)\frac{2\alpha'}{2} + \sigma(\emptyset^2)$$

E. o. m becomes:

$$(1-20)\ddot{\varkappa} - (1+20)\alpha'' - (1+20)\frac{2\alpha'}{2} - 2\dot{\sigma}\dot{\varkappa} - 2\dot{\sigma}'\alpha' = -\frac{9\ddot{V}}{9\alpha}$$

$$\Rightarrow \ddot{\varkappa} = (1+40)\left(\alpha'' + \frac{2\alpha'}{2}\right) + 2\dot{\sigma}\dot{\varkappa} + 20'\alpha' - (1+20)\frac{9\ddot{V}}{9\alpha}$$

The action is:

$$S = \int d^{4}x \int -g \left[\frac{1}{2} (\nabla \mu a) (\nabla \mu a) - V \left(\frac{a}{\ell} \right) \right] \quad \text{with} \quad V(x) = \frac{\Lambda^{4}}{c_{3}} \left(1 - \sqrt{1 - 4c_{3} \sin^{2} x} \right)$$

Metric:
$$d^2 x = (1+20) dt^2 - (1-20) (dr^2 + r^2 d\Omega^2)$$
 and in the paper

$$\Rightarrow q_{\mu\nu} = \begin{pmatrix} 1+2\phi \\ -(1-2\phi) \\ -n^2 \sin^2 \theta (1-2\phi) \end{pmatrix}$$

Christoffel symbols (from Mathematica's notebook "christoffel. nb") of the form $\Gamma_{r,s}^{r,n}$, with $\mu=t,r$:

$$\Gamma_{nn}^{n} = \underbrace{\emptyset^{l}}_{-1+2\emptyset} \cong -\emptyset^{l} (1+2\emptyset)$$

$$\Gamma_{00}^{ln} = -\underbrace{n\left(-1+2\phi+n\phi^{l}\right)}_{-1+2\phi} \cong n\left(-1+n\phi^{l}\right)$$

$$\Gamma_{\ell\ell}^{r} = -r \sin^2 \theta \left(-1 + 2\theta + r \theta' \right) \approx r \sin^2 \theta \left(-1 + r \theta' \right)$$

$$-1 + 2\theta$$

$$\Gamma_{tt}^{n} = \emptyset^{1} \cong \emptyset^{1}(1+2\emptyset)$$

$$\Gamma_{nn}^{+} = -\frac{\dot{\phi}}{1+2\phi} \approx -\dot{\phi}(1-2\phi)$$

$$\Gamma_{\theta\theta}^{t} = -\frac{n^{2}\dot{\emptyset}}{1+2\dot{\emptyset}} \stackrel{\sim}{=} -n^{2}\dot{\emptyset} \left(1-2\dot{\emptyset}\right)$$

$$\Gamma_{ee} = -r^2 \sin^2 \theta \ \dot{\varphi} \simeq -r^2 \sin^2 \theta \ \dot{\varphi} \left(1-2\varphi\right)$$

$$\frac{1+2\varphi}{1+2\varphi}$$

$$\int_{1+20}^{1} = \dot{\cancel{g}} (1-20)$$

Equation of motion for scalar in curved spacetime:

$$\square = -\frac{9\tilde{V}}{9\alpha}, \text{ with } \square = g^{\mu\nu} \nabla_{\mu} \nabla_{\nu}$$

We have

$$\Box \alpha = g^{\mu\nu} \nabla_{\mu} \partial_{\nu} \alpha = g^{\mu\nu} (\partial_{\mu} \partial_{\nu} \alpha - \Gamma_{\mu\nu}^{\beta} \partial_{\beta} \alpha)$$

$$= g^{tt} (\ddot{\alpha} - \Gamma_{tt}^{\beta} \partial_{\beta} \alpha) + g^{rr} (\alpha'' - \Gamma_{rr}^{\beta} \partial_{\beta} \alpha) - (g^{\theta\theta} \Gamma_{\theta\theta}^{\beta} + g^{\theta\theta} \Gamma_{\theta\theta}^{\beta}) \partial_{\beta} \alpha$$

arming spherically symmetrical $\alpha \equiv \alpha(t, r)$

$$= (1-20) \left[\ddot{\omega} - \dot{\varphi} \left(1-20 \right) \dot{\omega} - \varphi' \left(1+20 \right) \alpha' \right]$$

$$-\left(1+2\emptyset\right)\left[\infty''+\dot{\emptyset}\left(1-2\emptyset\right)\dot{\alpha}+\emptyset'\left(1+2\emptyset\right)\omega'\right]$$

+
$$(1+2\emptyset)$$
 $\left[-n^{2}\dot{p}(1-2\emptyset)\dot{\alpha}+n(-1+n\emptyset^{1})\alpha^{1}\right]$

+
$$(1+20) \cos^2 \theta \left[-r^2 \sin^2 \theta \dot{\rho} (1-20) \dot{\alpha} + r \sin^2 \theta (-1+r 0) \dot{\alpha} \right]$$

$$= (1-20) \ddot{\alpha} - \dot{\phi} \dot{\alpha} - \dot{\phi} \dot{\alpha} - (1+20) \alpha'' - \dot{\phi} \dot{\alpha} - 0 \alpha'$$

$$-\dot{\phi}\dot{\alpha} - (1+2\phi)\frac{\alpha'}{\pi} + \phi^{\dagger}\dot{\alpha}^{\dagger} - \dot{\phi}\dot{\alpha} - (1+2\phi)\frac{\alpha'}{\pi} + \phi^{\dagger}\dot{\alpha}^{\dagger}$$

E. o.m beames:

$$(1-20)\ddot{\varkappa}-(1+20)\alpha''-(1+20)\frac{2\alpha'}{2}-40\dot{\alpha}=-\frac{0}{2}$$

$$\Rightarrow \ddot{\alpha} = (1+40)(\alpha'' + \frac{2\alpha'}{r}) + 40\dot{\alpha} - (1+20)9\dot{\sqrt{2}}$$
correct!