

I. CORRELATION FUNCTION (Problem 2.16)

Evaluate the correlation function

$$C(t) = \langle x(t)x(0) \rangle,$$

for the ground state of the harmonic oscillator $|0\rangle$. How does it decay with time?

Solution. From Equation 2.3.45a,

$$x(t) = x(0) \cos \omega t + \frac{p(0)}{m\omega} \sin \omega t.$$

Therefore,

$$C(t) = \langle 0|x(t)x(0)|0\rangle = \cos \omega t \langle 0|(x(0))^2|0\rangle + \frac{\sin \omega t}{m\omega} \langle 0|p(0)x(0)|0\rangle.$$

Clearly, the time dependence is controlled by the factors outside the expectation values; and we only need to determine the expectation values at $t = 0$. Thus, in order to simplify the notation we let $x(0) \rightarrow x$ and $p(0) \rightarrow p$. Recall from Equation 2.3.24 that

$$x = \sqrt{\frac{\hbar}{2m\omega}}(a + a^\dagger), \quad p = i\sqrt{\frac{m\hbar\omega}{2}}(-a + a^\dagger).$$

Therefore,

$$\begin{aligned} C(t) &= \frac{\hbar \cos \omega t}{2m\omega} \langle 0|a^2 + aa^\dagger + a^\dagger a + a^{\dagger 2}|0\rangle + i\frac{\hbar \sin \omega t}{2m\omega} \langle 0|-a^2 - aa^\dagger + a^\dagger a + a^{\dagger 2}|0\rangle \\ &= \frac{\hbar \cos \omega t}{2m\omega} \langle 0|aa^\dagger|0\rangle - i\frac{\hbar \sin \omega t}{2m\omega} \langle 0|aa^\dagger|0\rangle \\ &= \frac{\hbar}{2m\omega} (\cos \omega t - i \sin \omega t). \end{aligned}$$

The correlation function oscillates and does not decay in time. □

II. DISPLACED OSCILLATOR (Problem 2.12)

At $t = 0$ let

$$|\psi(0)\rangle = e^{-ip(0)l/\hbar} |0\rangle,$$

where $|0\rangle$ is the ground state of a harmonic oscillator, p is the momentum operator, and l is a fixed quantity with dimensions of length. Evaluate the expectation value of the position operator $\langle x \rangle$ as a function of time.

Solution. From Equation 2.3.45,

$$\begin{aligned} x(t) &= x(0) \cos \omega t + \frac{p(0)}{m\omega} \sin \omega t, \\ p(t) &= -m\omega x(0) \sin \omega t + p(0) \cos \omega t. \end{aligned}$$

Therefore, in the Heisenberg picture,

$$\begin{aligned} \langle x(t) \rangle &= \langle \psi(0) | x(t) | \psi(0) \rangle = \cos \omega t \langle \psi(0) | x(0) | \psi(0) \rangle + \frac{\sin \omega t}{m\omega} \langle \psi(0) | p(0) | \psi(0) \rangle \\ &= \cos \omega t \langle 0 | e^{ip(0)l/\hbar} x(0) e^{-ip(0)l/\hbar} | 0 \rangle + \frac{\sin \omega t}{m\omega} \langle 0 | e^{ip(0)l/\hbar} p(0) e^{-ip(0)l/\hbar} | 0 \rangle. \end{aligned}$$

Using the Baker-Hausdorff lemma, for Hermitian operators $x(0)$ and $p(0)$, we have

$$\begin{aligned} e^{ip(0)l/\hbar} x(0) e^{-ip(0)l/\hbar} &= x(0) + \left(\frac{il}{\hbar} \right) [p(0), x(0)] + \left(\frac{i^2 l^2}{2! \hbar^2} \right) [p(0), [p(0), x(0)]] + \dots \\ &= x(0) + \left(\frac{il}{\hbar} \right) [p(0), x(0)] \\ &= x(0) + l, \end{aligned}$$

and

$$e^{ip(0)l/\hbar} p(0) e^{-ip(0)l/\hbar} = p(0) + \left(\frac{il}{\hbar} \right) [p(0), p(0)] + \left(\frac{i^2 l^2}{2! \hbar^2} \right) [p(0), [p(0), p(0)]] + \dots = p(0).$$

Thus,

$$\langle x(t) \rangle = \cos \omega t \langle 0 | x(0) | 0 \rangle + \cos \omega t \langle 0 | l | 0 \rangle + \frac{\sin \omega t}{m\omega} \langle 0 | p(0) | 0 \rangle = l \cos \omega t.$$

□

III. ELECTRON IN A MAGNETIC FIELD (Problem 2.39)

Let an electron move in a uniform magnetic field $\mathbf{B} = B\hat{z}$.

- (a) Evaluate the commutator between the components of the kinematical momentum $[\Pi_x, \Pi_y]$.
- (b) Obtain the energy spectrum

$$E_{kn} = \frac{\hbar^2 k^2}{2m} + \frac{|eB|\hbar}{mc} \left(n + \frac{1}{2} \right),$$

where $\hbar k$ is the eigenvalue of the p_z operator and n is a nonnegative integer.

Solution.

- (a) For an electron in a magnetic field $\mathbf{\Pi} = \mathbf{p} - e\mathbf{A}/c$, where we take $\mathbf{A} = (-By, 0, 0)$ (from Equation 2.7.41). Therefore,

$$\begin{aligned} [\Pi_x, \Pi_y] &= \left[p_x - \frac{eA_x}{c}, p_y - \frac{eA_y}{c} \right] = [p_x, p_y] - \frac{e}{c} ([A_x, p_y] + [p_x, A_y]) + \frac{e^2}{c^2} [A_x, A_y] \\ &= \frac{e}{c} ([p_y, A_x] - [p_x, A_y]) = -\frac{i\hbar e}{c} (\partial_y A_x - \partial_x A_y) = \frac{i\hbar e}{c} (\nabla \times \mathbf{A})_z = \frac{i\hbar e}{c} B_z = \frac{i\hbar e}{c} B. \end{aligned}$$

- (b) Consider the Hamiltonian

$$H = \frac{\mathbf{\Pi}^2}{2m} = \frac{\Pi_x^2 + \Pi_y^2 + \Pi_z^2}{2m}.$$

From the commutation relation, we can let $\Pi_x = y e B / c$ and $\Pi_y = p_y$, where $[y, p_y]$. Also, it is clear that $\Pi_z = p_z$. Therefore,

$$H = \frac{e^2 B^2 y^2}{2mc^2} + \frac{p_y^2}{2m} + \frac{p_z^2}{2m} = H_y + H_z,$$

where H_y corresponds to the Hamiltonian of a harmonic oscillator with $\omega = |eB|/\hbar mc$. Therefore,

$$E = \frac{\hbar^2 k^2}{2m} + \hbar \omega \left(n + \frac{1}{2} \right) = \frac{\hbar^2 k^2}{2m} + \frac{\hbar |eB|}{mc} \left(n + \frac{1}{2} \right).$$

□