## TUTORIAL 06 — SOLUTIONS

## 7CCMCS04 A. ANNIBALE AND G. SICURO

## Problem 6.1

We have

$$\partial_t P_n(t) = \mu P_{n+1}(t) + \mu P_{n-1}(t) - 2\mu P_n(t).$$

Let us introduce the generating function  $F(z,t) = \sum_n P_n(t) z^n$ . Using the master equation

$$\partial_t F(z,t) = \frac{\mu}{z} F(z,t) + \mu z F(z,t) - 2\mu F(z,t)$$

solved by

$$F(z,t) = F(z,0) e^{t\mu(z+1/z-2)}$$
.

For  $P_n(0) = \delta_{n\,0}$ , we have F(z,0) = 1. Expanding in powers of z

$$F(z,t) = e^{t\mu(z+1/z-2)} = e^{-2\mu t} \sum_{k,n=0}^{\infty} \frac{(t\mu)^{k+n} z^{k-n}}{k!n!} = e^{-2\mu t} \sum_{n=-\infty}^{\infty} z^n \sum_{k=0}^{\infty} \frac{(t\mu)^{n+2k}}{(n+k)!k!}$$

so that

$$P_n(t) = e^{-2t\mu} I_n(2t\mu).$$

The properties of  $I_n(x)$  follow from properties of  $P_n(t)$  and from the master equation for  $P_n(t)$ .

Date: February 13, 2025.

7CCMCS04

2

## Problem 6.2

a. Let  $P_n(t)$  be the probability that there are n surviving nuclei at time t. If  $\gamma$  is the decay probability per unit time for one nucleus, then the probability that one nucleus decay during an infinitesimal time interval  $\Delta t$  is  $\gamma \Delta t$  at leading order. Assuming that there are n nuclei, the probability that there is one decay is  $\gamma n \Delta t$  (and the probability that there is no decay  $1 - \gamma n \Delta t$ ) at the leading order. The transition probability from m to n in a short time  $\Delta t$  is therefore

$$\begin{split} \mathbb{P}_{1|1}[\mathsf{N}_{t+\Delta t} = n | \mathsf{N} = m] &= \delta_{m,n} \underbrace{\left(1 - \gamma m \Delta t\right)}_{\mathsf{No \ one \ of } m \ \mathsf{nuclei} \ \mathsf{decay}} \\ &+ \delta_{m-1,n} \underbrace{\gamma m \Delta t}_{\mathsf{One \ of } m \ \mathsf{nuclei} \ \mathsf{decays}} + O(\Delta t^2). \end{split}$$

**b.** Summing over all possible states at time t and weighting with the probability of occurrence of each configuration we get

$$P_n(t+\mathrm{d}\,t) = \sum_m \mathbb{P}_{1|1}[\mathsf{N}_{t+\mathrm{d}\,t} = n | \mathsf{N} = m] P_m(t) = (n+1)\gamma\,\mathrm{d}\,t P_{n+1}(t) + (1-n\gamma\,\mathrm{d}\,t) P_n(t),$$

which leads to the master equation

$$\dot{P}_n(t) = \gamma(n+1)P_{n+1}(t) - \gamma n P_n(t).$$

**c.** In order to get an equation for the generating function  $F(z,t) = \sum_{n=0}^{\infty} P_n(t)z^n$ , we multiply the master equation by  $z^n$  and sum over n

$$\frac{\partial F}{\partial t} = \gamma \sum_{n} (n+1)z^{n} P_{n+1}(t) - \gamma \sum_{n} nz^{n} P_{n}(t)$$
$$= \gamma \frac{\partial}{\partial z} \sum_{n} z^{n+1} P_{n+1}(t) - \gamma z \frac{\partial}{\partial z} \sum_{n} z^{n} P_{n}(t)$$
$$= \gamma (1-z) \frac{\partial F}{\partial z}$$

By using the method of characteristics,

$$\frac{\mathrm{d}\,t}{1} = \frac{\mathrm{d}\,z}{\gamma(z-1)} = \frac{\mathrm{d}\,F}{0}.$$

Two integrals are easily obtained considering the systems t, z and z, F:

$$t = \frac{1}{\gamma} \ln|z - 1| + c \to (z - 1) e^{-\gamma t} = a;$$
  $F(z, t) = \phi \equiv \phi(a).$ 

Then, the general solution is, in terms of a yet unknown function  $\phi$ 

$$F(z,t) = \phi\left((z-1)e^{-\gamma t}\right).$$

From the initial condition  $P_1(n,0) = \delta_{n n_0}$  we get

$$F(z,0) = z^{n_0} = \phi(z-1)$$

and the functional form of  $\phi$  follows as

$$\phi(z) = (z+1)^{n_0}$$

For t > 0 we have therefore

$$F(z,t) = ((z-1)e^{-\gamma t} + 1)^{n_0} = \sum_{m=0}^{n_0} \binom{n_0}{m} z^m e^{-\gamma mt} (1 - e^{-\gamma t})^{n_0 - m}$$

which leads to

$$P_m(t) = \binom{n_0}{m} e^{-\gamma mt} (1 - e^{-\gamma t})^{n_0 - m}$$

**d.** The average number of undecayed nuclei at time t can be obtained using the generating function found in point  $\mathbf{c}$ .

$$F(z,t) = \sum_{n} P_n(t) z^n = ((z-1) e^{-\gamma t} + 1)^{n_0}$$
  

$$\Rightarrow \partial_z F(z,t) = \sum_{n} n P_n(t) z^{n-1} = n_0 e^{-\gamma t} ((z-1) e^{-\gamma t} + 1)^{n_0 - 1}.$$

Indeed

$$\mathbb{E}[\mathsf{N}_t] = \sum_n n P_n(t) = \partial_z F(z, t)|_{z=1} = n_0 \,\mathrm{e}^{-\gamma t}$$

so that the half-life of the process is  $1/\gamma$ . For large t

$$\lim_{t\to\infty} F(z,t) = 1 = P(0,\infty) + \sum_{m\geq 1} z^m P_m(\infty),$$

but  $P_m(\infty) = 0 \ \forall m \geq 1$ . So  $P_0(\infty) = 1$ . Alternatively, one can multiply both sides of the master equation by n and sum over n, thus obtaining

$$\frac{\mathrm{d}\,\mathbb{E}[\mathsf{N}_t]}{\mathrm{d}\,t} = \sum_{n=0}^{\infty} n\dot{P}_n = \gamma \sum_{n=0}^{\infty} n(n+1)P_{n+1} - \gamma \sum_{n=0}^{\infty} n^2 P_n$$
$$= \gamma \sum_{n=0}^{\infty} (n-1)nP_n - \gamma \sum_{n=0}^{\infty} n^2 P_n = -\gamma \sum_{n=0}^{\infty} n P_n = -\gamma \mathbb{E}[\mathsf{N}_t].$$

Solving the above equation for  $\mathbb{E}[\mathsf{N}_0] = n_0$  gives

$$\mathbb{E}[\mathsf{N}_t] = n_0 \,\mathrm{e}^{-\gamma t} \,.$$