

TUTORIAL 02 — SOLUTIONS

7CCMCS04

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PROBLEM 2.1

To answer all questions it is useful to find eigenvalues and eigenvectors of \mathbf{Q} first. The eigenvalues are $\lambda_1 = 1$ and $\lambda_2 = 1 - \alpha - \beta$ (as it follows from the fact that $\text{tr}[\mathbf{Q}] = \lambda_1 + \lambda_2 = 2 - \alpha - \beta$). Doing the calculations for the eigenvectors, you will find

$$(1) \quad \begin{aligned} \lambda_1 = 1, \quad |\psi^1\rangle &= \begin{pmatrix} \frac{\beta}{\alpha+\beta} \\ \frac{\alpha}{\alpha+\beta} \end{pmatrix}, & \langle\phi^1| &= (1 \quad 1) \\ \lambda_2 = 1 - \alpha - \beta, \quad |\psi^2\rangle &= \begin{pmatrix} -1 \\ 1 \end{pmatrix}, & \langle\phi^2| &= \left(-\frac{\alpha}{\alpha+\beta} \quad \frac{\beta}{\alpha+\beta}\right). \end{aligned}$$

In the decomposition above, $|\psi^1\rangle \equiv |\mathbf{II}\rangle$. Also, observe that the $(t+1)$ th day correspond to time t (the first day is our $t=0$).

- a. The probability of having rain the second day is simply $Q_{11} = 1 - \alpha$. At the generic $(t+1)$ th day, the probability of having rain is $(\mathbf{Q}^t)_{11} = \sum_a \lambda_a^t \psi_1^a \phi_1^a = \frac{\beta}{\alpha+\beta} + (1 - \alpha - \beta)^t \frac{\alpha}{\alpha+\beta}$.
- b. In this case, using that $|\mathbf{P}(t)\rangle = \mathbf{Q}^t |\mathbf{P}(0)\rangle$ and the fact that here $|\mathbf{P}(0)\rangle = \begin{pmatrix} p \\ 1-p \end{pmatrix}$, then the evolution is

$$\begin{aligned} |\mathbf{P}(t)\rangle &= |\psi^1\rangle + (1 - \alpha - \beta)^t |\psi^2\rangle \langle\phi^2| \mathbf{P}(0)\rangle \\ &= |\psi^1\rangle + \frac{(1 - \alpha - \beta)^t (\beta - p(\alpha + \beta))}{\alpha + \beta} |\psi^2\rangle \end{aligned}$$

so that $P_2(t) = \frac{\alpha}{\alpha+\beta} + \frac{(1-\alpha-\beta)^t (\beta - p(\alpha+\beta))}{\alpha+\beta}$ gives the probability of observing sunshine the second ($t=1$) and third ($t=2$) day.

The question also asks the probability of observing sunshine the first three days,

$$\begin{aligned} \mathbb{P}_3[X_0 = 2, X_1 = 2, X_2 = 2] &= \mathbb{P}_{1|1}[X_2 = 2 | X_1 = 2] \mathbb{P}_{1|1}[X_1 = 2 | X_0 = 2] \mathbb{P}_1[X_0 = 2] \\ &= Q_{22} Q_{22} P_2(0) = (1 - \beta)^2 (1 - p). \end{aligned}$$

Also, the probability of observing sunshine the second and third day is

$$\begin{aligned} \mathbb{P}_2[X_1 = 2, X_2 = 2] &= \mathbb{P}_{1|1}[X_2 = 2 | X_1 = 2] \mathbb{P}_1[X_1 = 2] \\ &= Q_{22} P_2(1) = (1 - \beta)(1 - \beta - p(1 - \alpha - \beta)). \end{aligned}$$

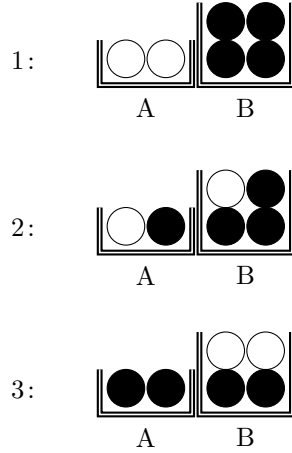
- c. See above.

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- d. It is again $\frac{\beta}{\alpha+\beta}$. Assuming this value, means that we are on the steady state as this is the first entry of $|\psi^1\rangle$ (and the fact that we have 2 states only means that it is enough to specify the entire vector).
- e. We have that $\mathbf{Q} = \sum_{a=1}^2 \lambda_a |\psi^a\rangle \langle \phi^a| = \frac{1}{\alpha+\beta} \begin{pmatrix} \beta & \beta \\ \alpha & \alpha \end{pmatrix} + \frac{1-\alpha-\beta}{\alpha+\beta} \begin{pmatrix} \alpha & -\beta \\ -\alpha & \beta \end{pmatrix}$. Being $\lambda_2 = 1 - \alpha - \beta$, if $|\lambda_2| < 1$ the relaxation time is $\tau = -\frac{1}{\ln|\lambda_2|}$ which depends indeed on $\alpha + \beta$ only.
- f. Yes: we have indeed that $Q_{21}\Pi_1 = \beta \frac{\alpha}{\alpha+\beta} = Q_{12}\Pi_2$.
- g. For $\alpha = \beta = 1$, $\lambda_2 = -1$ and in this case the dynamics does not converge in general to the stationary distribution as there are *two* eigenvalues with modulus equal to 1. Instead for $\alpha = \beta = 0$ we have $\lambda_2 = 1$. Again, there is not a unique asymptotic state: the chain is made of two isolated configurations, and starting in one of them, the system remains in it forever.

PROBLEM 2.2

Let us denote by 1, 2, 3 respectively, the following:



Then

$$\mathbf{Q} = \begin{pmatrix} 0 & 1/8 & 0 \\ 1 & 1/2 & 1/2 \\ 0 & 3/8 & 1/2 \end{pmatrix}.$$

In principle we can get the eigenvalues by solving the equation for the eigenvalues, $\det(\mathbf{Q} - \lambda \mathbf{I}) = 0$ but it is easier to appeal to the properties of the stochastic matrix: we know indeed that $\lambda_1 = 1$ is always an eigenvalue. Hence from the trace and determinant of \mathbf{Q} we can determine the other two. We know that

$$\text{tr } \mathbf{Q} = \lambda_1 + \lambda_2 + \lambda_3 = 1 \Rightarrow \lambda_2 = -\lambda_3$$

but also

$$\det \mathbf{Q} = \lambda_1 \lambda_2 \lambda_3 = -\frac{1}{16} \Rightarrow \lambda_3 = -\frac{1}{16\lambda_2}$$

so $\lambda_2 = 1/4$ and $\lambda_3 = -1/4$. We know that the left eigenvector $\langle \phi^1 |$ associated to $\lambda_1 = 1$ is $\langle \mathbf{1} |$. The other eigenvectors are found, up to a global factor, by solving the equations for the right and left eigenvectors, respectively, i.e., $\mathbf{Q}|\psi^a\rangle = \lambda_a|\psi^a\rangle$ and $\langle \phi^a|\mathbf{Q} = \langle \phi^a|\lambda_a \forall a = 1, 2, 3$. By requiring the normalization $\langle \phi^a|\psi^b\rangle = \delta_{ab}$, which fixes the factor in front of $|\psi^1\rangle$, we have

$$(2) \quad \langle \phi^1 | = (1 \quad 1 \quad 1), \quad |\psi^1\rangle = \begin{pmatrix} 1/15 \\ 8/15 \\ 2/5 \end{pmatrix},$$

$$(3) \quad \langle \phi^2 | = (4 \quad 1 \quad -2), \quad |\psi^2\rangle = \begin{pmatrix} 1/12 \\ 1/6 \\ -1/4 \end{pmatrix},$$

$$(4) \quad \langle \phi^3 | = (12 \quad -3 \quad 2), \quad |\psi^3\rangle = \begin{pmatrix} 1/20 \\ -1/10 \\ 1/20 \end{pmatrix}.$$

Let us derive these eigenvectors explicitly for the sake of completeness. We know that we *always* have $\langle \phi^1 | = (1 \ 1 \ 1)$. Then to compute $|\psi^1\rangle$, we compute straightforwardly $\mathbf{Q}|\psi^1\rangle = |\psi^1\rangle$, i.e.,

$$\begin{pmatrix} 0 & 1/8 & 0 \\ 1 & 1/2 & 1/2 \\ 0 & 3/8 & 1/2 \end{pmatrix} \begin{pmatrix} \psi_1^1 \\ \psi_2^1 \\ \psi_3^1 \end{pmatrix} = \begin{pmatrix} \frac{1}{8}\psi_2^1 \\ \psi_1^1 + \frac{1}{2}\psi_2^1 + \frac{1}{2}\psi_3^1 \\ \frac{3}{8}\psi_2^1 + \frac{1}{2}\psi_3^1 \end{pmatrix} = \begin{pmatrix} \psi_1^1 \\ \psi_2^1 \\ \psi_3^1 \end{pmatrix}$$

The first equation implies $\psi_2^1 = 8\psi_1^1$, and from the second we get $\psi_1^1 + \frac{1}{2}\psi_2^1 + \frac{1}{2}\psi_3^1 = \psi_1^1 + 4\psi_1^1 + \frac{1}{2}\psi_3^1 = \psi_2^1 = 8\psi_1^1 \Rightarrow \psi_3^1 = 6\psi_1^1$. Note that if we plug what we got in the *third* equation, we will not manage to find ψ_1^1 : this is always the case as eigenvectors are defined up to a multiplicative constant. To fix ψ_1^1 we need to use $\langle 1|\psi^1\rangle = 1 \Rightarrow \psi_1^1 + \psi_2^1 + \psi_3^1 = 15\psi_1^1 = 1 \Rightarrow \psi_1^1 = \frac{1}{15}$, and from this $\psi_2^1 = \frac{8}{15}$ and $\psi_3^1 = \frac{6}{15} = \frac{2}{5}$. Let us now move to the *second* pair of eigenvalues, the one corresponding to $\lambda_2 = 1/4$. Again, to compute $|\psi^2\rangle$, we compute straightforwardly $\mathbf{Q}|\psi^1\rangle = 1/4|\psi^1\rangle$, i.e.,

$$\begin{pmatrix} 0 & 1/8 & 0 \\ 1 & 1/2 & 1/2 \\ 0 & 3/8 & 1/2 \end{pmatrix} \begin{pmatrix} \psi_1^2 \\ \psi_2^2 \\ \psi_3^2 \end{pmatrix} = \begin{pmatrix} \frac{1}{8}\psi_2^2 \\ \psi_1^2 + \frac{1}{2}\psi_2^2 + \frac{1}{2}\psi_3^2 \\ \frac{3}{8}\psi_2^2 + \frac{1}{2}\psi_3^2 \end{pmatrix} = \frac{1}{4} \begin{pmatrix} \psi_1^2 \\ \psi_2^2 \\ \psi_3^2 \end{pmatrix}$$

The first equation implies $\psi_2^2 = 2\psi_1^2$, and from the second we get $\psi_1^2 + \frac{1}{2}\psi_2^2 + \frac{1}{2}\psi_3^2 = \psi_1^2 + \psi_1^2 + \frac{1}{2}\psi_3^2 = \frac{1}{4}\psi_2^2 = \frac{1}{2}\psi_1^2 \Rightarrow \psi_3^2 = -3\psi_1^2$. Once again, the last equation does not give us information on the value of ψ_1^2 . Let us compute $\langle \phi^2 |$, by using

$$(\phi_1^2 \ \phi_2^2 \ \phi_3^2) \begin{pmatrix} 0 & 1/8 & 0 \\ 1 & 1/2 & 1/2 \\ 0 & 3/8 & 1/2 \end{pmatrix} = (\phi_2^2 \ \frac{\phi_1^2 + 4\phi_2^2 + 3\phi_3^2}{8} \ \frac{\phi_2^2 + \phi_3^2}{2}) = \frac{1}{4} (\phi_1^2 \ \phi_2^2 \ \phi_3^2).$$

We immediately obtain $\phi_2^2 = \frac{1}{4}\phi_1^2$ and $\phi_3^2 = -\frac{1}{2}\phi_1^2$. Again, the last equation does not tell us anything about ϕ_1^2 . If we use $\langle \phi^2 | \psi^2 \rangle = 1$ we obtain $\phi_1^2 \psi_1^2 + \frac{1}{2}\phi_1^2 \psi_1^2 + \frac{3}{2}\phi_1^2 \psi_1^2 = 1$, i.e., $\psi_1^2 = \frac{1}{3\phi_1^2}$, so we expressed one constant in term of the other. However, ϕ_1^2 is left arbitrary and indeed *we can pick it at our choice*: the solution presented above has, for example, $\phi_1^2 = 4$, but a different choice (e.g., $\phi_1^2 = 1$) is acceptable as well. Repeating the arguments for the third eigenvalue we can identify the last pair: the same ambiguity will appear, and we will have to fix an arbitrary constant. Note that this ambiguity would not be there if \mathbf{Q} were symmetric: in this case the constant is fixed by the condition that *left and right eigenvectors have the same entries*. Once we have eigenvectors and eigenvalues of \mathbf{Q} , we can decompose it as

$$\mathbf{Q} = \sum_{a=1}^3 \lambda_a |\psi^a\rangle \langle \phi^a| = |\psi^1\rangle \langle \phi^1| + \frac{1}{4} |\psi^2\rangle \langle \phi^2| - \frac{1}{4} |\psi^3\rangle \langle \phi^3|$$

so that, for example, $Q_{ij} = \sum_{a=1}^3 \lambda_a \psi_i^a \phi_j^a$. Note that $\sum_i \psi_i^a = \delta_{a,1}$.

Starting now from the definition of the Markov Chain, $|\mathbf{P}(t)\rangle = \mathbf{Q}^t |\mathbf{P}(0)\rangle$, assuming $P_i(0) = \delta_{i,1}$, then $P_i(t) = (\mathbf{Q}^t)_{i1}$. Therefore

$$P_3(3) = (\mathbf{Q}^3)_{31} = \sum_a \lambda_a^3 \psi_3^a \phi_1^a = 1 \cdot 1 \cdot \frac{2}{5} + \left(\frac{1}{4}\right)^3 \cdot 4 \cdot \left(-\frac{1}{4}\right) + \left(-\frac{1}{4}\right)^3 \cdot 12 \cdot \frac{1}{20} = \frac{3}{8}.$$

Moreover

$$\lim_{t \rightarrow \infty} P_3(t) = \psi_3^1 = \frac{2}{5}.$$