

Jet evolution in a dense QCD medium

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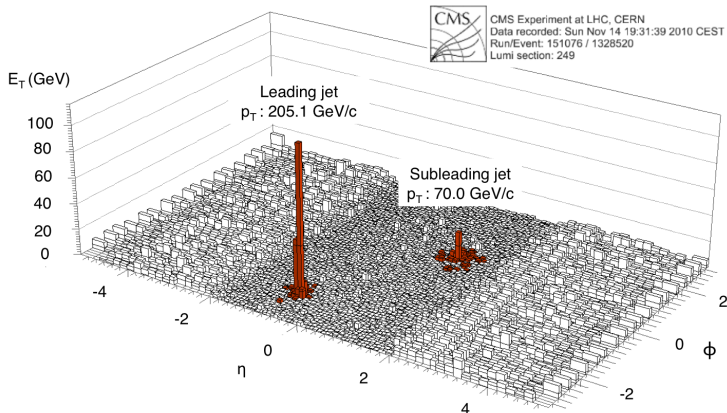
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- *the medium: a quark gluon plasma created in a heavy ion collision*
- *the jet: a collimated spray of particles generated via successive branchings of a parton with high energy produced in the collision*

Structure of presentation:

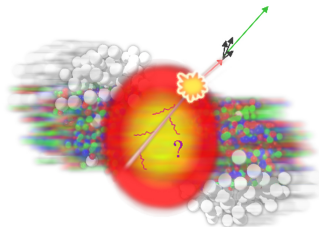
- Context
- Physical picture
- Markov process
- Monte Carlo simulation

Heavy ion collisions, quark gluon plasma and jets, what we observe:

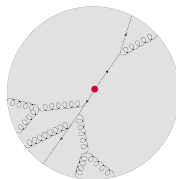
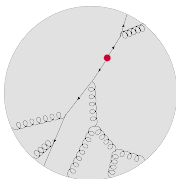


Di-jet asymmetry. Missing energy found among soft hadrons propagating at large angles \Rightarrow different from in vacuum jet evolution.

Heavy ion collisions, quark gluon plasma and jets,
the stereotypical picture:



Implicit assumption: small energy fluctuations.



We need a model that answers:

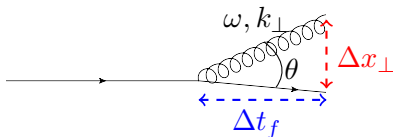
- How can the lost energy end up at such large angles?
- How large are the energy fluctuations?

Physical picture

In a medium, scattering triggers emissions:

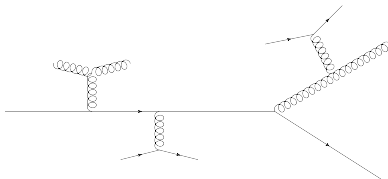
- Destroys quantum coherence
- Gives broadening of transverse momentum

Formation time: $\lambda_{\perp} < \Delta x_{\perp} \Rightarrow \frac{1}{k_{\perp}} < \frac{k_{\perp}}{\omega} \Delta t_f$.



Note: we have mostly gluon emissions so we consider only these.

$\Delta t_f \gg \text{mean free path} \gg \text{Debye length} \Rightarrow \text{multiple scatterings}$
 $\text{lead to one emission and scattering centers are independent} \Rightarrow$
 $\text{random walk} \Rightarrow \langle k_{\perp}^2 \rangle \sim \hat{q} \Delta t.$



All together: $\Delta t_f \sim \sqrt{\frac{\omega}{\hat{q}}}$ and $\theta_f = \left(\frac{\hat{q}}{\omega^3}\right)^{1/4}$
 $\Rightarrow \text{we favour soft gluons at large angles!}$

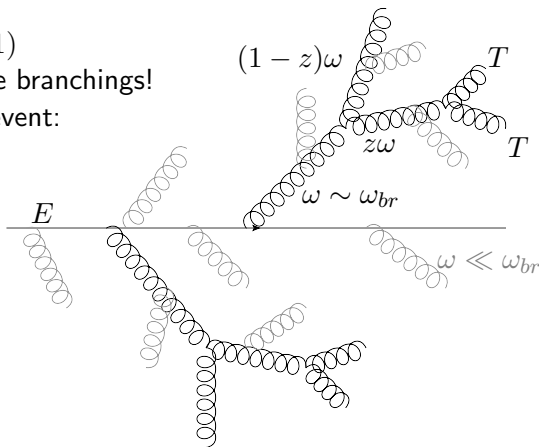
If now $\omega = \text{energy of parent}$: $\Delta P(z, \omega, \Delta t) \sim \alpha_s \Delta t / [\Delta t_f(z\omega)]$

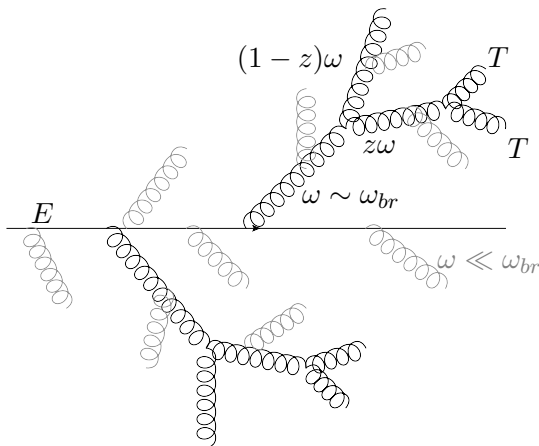
- For $z\omega \sim \omega_{br}$, $\Delta P(z, \omega) \sim \mathcal{O}(1)$ at $\Delta t = L$
- Next, $P(z, \omega_{br}) \sim \mathcal{O}(1)$ for $\Delta t < L$, even for $z \sim 1/2$

$$\Delta P \sim \mathcal{O}(1)$$

\Rightarrow multiple branchings!

A typical event:





- A number of $\mathcal{O}(1)$ of primary gluons with $\omega \sim \omega_{br}$ emitted from leading particle.
- They then branch democratically, transporting away all energy.
- Hence both energy loss and its fluctuations on the scale ω_{br} .
- Time between branchings $>$ formation time \Rightarrow branchings are independent \Rightarrow Markov process

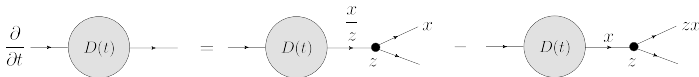
Markov process

Branching rate from BDMPS-Z spectrum:

$$\frac{d^2 P(z, \tau)}{dz d\tau} = \frac{K(z)}{2\sqrt{x}} \equiv K(x, z).$$

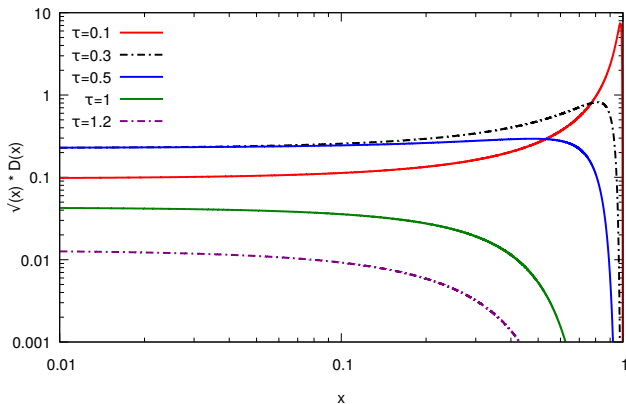
- τ = time scaled by the time needed for the leading particle to branch democratically.
- Splitting kernel: $K(z) = \frac{[1-z(1-z)]^{5/2}}{[z(1-z)]^{3/2}}$.
Simplified: $K(z) = \frac{1}{[z(1-z)]^{3/2}}$.

Get equation for the one point function $D(x, \tau) \equiv xn(x, \tau)$



and similarly (but longer) for the two-point function
 $D^{(2)}(x, x', \tau) \equiv xx'n^{(2)}(x, x', \tau)$.

Simple kernel:



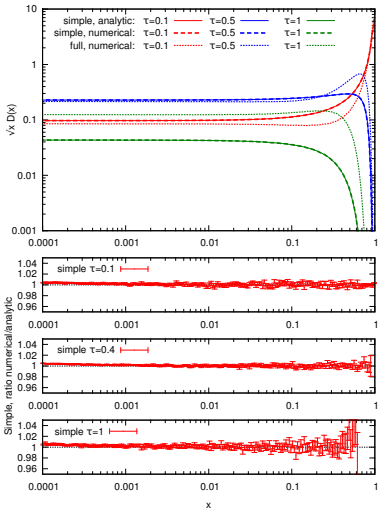
- Energy fraction left in gluon cascade:
 $\int_0^1 dx D(x, \tau) = e^{-\pi\tau^2} \Rightarrow$ decreasing in time.
- Formally: condensate at $x = 0$. Physically: thermalization.
- At small τ , loss $\simeq \pi\omega_{br}$.
- $D^{(2)}(x, x', \tau)$ used to find fluctuations. To order τ^4 :

$$\sigma_\epsilon(\tau) \simeq \langle \epsilon(\tau) \rangle / \sqrt{3}$$

Monte Carlo simulation

Result: full kernel \Rightarrow less efficient branching

$$\sqrt{x}D(x)$$



Simple splitting kernel, numerical versus analytic:

- Good agreement overall
- Bias at small x (pileup)

Corrections from the full splitting kernel:

- Leading peak still present at $\tau = 0.5$
- Less energy lost at $\tau = 1$

Recall:

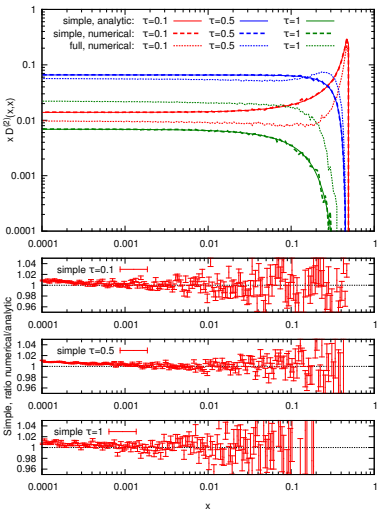
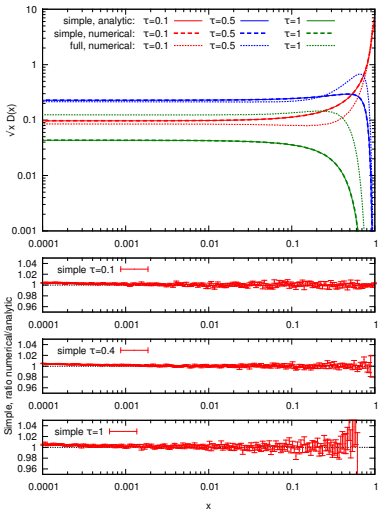
Simple $K(z) = \frac{1}{[z(1-z)]^{3/2}}$
Full $K(z) = \frac{[1-z(1-z)]^{5/2}}{[z(1-z)]^{3/2}}$

Monte Carlo simulation

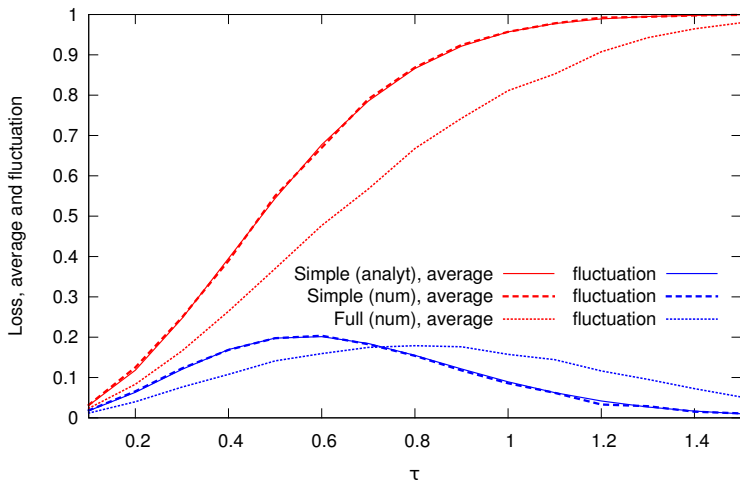
Result: full kernel \Rightarrow less efficient branching

$$\sqrt{x}D(x)$$

$$xD^2(x, x)$$



Energy loss: $1 - \int_0^1 dx D(x, \tau)$.



Full kernel \Rightarrow shifted in τ but same qualitative picture.

Summary:

- Democratic branchings \Rightarrow energy found at large angles
- Prediction: large fluctuations in energy loss
- Possible to make Monte Carlo simulation
- Full kernel \Rightarrow less efficient branching
- Full and simple kernel have same behaviour *qualitatively*

THE END

Questions?