

# Electromagnetism

## 1 | Electrostatics

**seriesCoulomb's law.** Let  $q_1, q_2$  be two point charges at positions  $\mathbf{r}_1, \mathbf{r}_2$ , respectively. Then the force  $\mathbf{F}_2$  experienced by  $q_2$  in the vicinity of  $q_1$  is given by

$$\mathbf{F}_{12} = kq_1q_2 \frac{\mathbf{r}_1 - \mathbf{r}_2}{\|\mathbf{r}_1 - \mathbf{r}_2\|^3},$$

where  $k = \frac{1}{4\pi\epsilon_0}$  and  $\epsilon_0 = 8,854 \text{ F/m}$  is the vacuum permittivity.

**seriesElectric field.** We define the *electric field*  $\mathbf{E}$  as the force per unit of charge. For a point charge, we have that the electric field created by  $q_1$  in the position of  $r_2$  is

$$\mathbf{F}_2 = q_2\mathbf{E}_1(\mathbf{r}_2), \quad \mathbf{E}_1(\mathbf{r}_2) = kq_1 \frac{\mathbf{r}_1 - \mathbf{r}_2}{\|\mathbf{r}_1 - \mathbf{r}_2\|^3}.$$

**seriesSuperposition principle.** **ARREGLAR** Let  $\rho(\mathbf{r}) = \frac{dq}{dV}$  be the volume charge density of an object. Then we have that

$$\mathbf{F} = \int_V \rho(\mathbf{r})\mathbf{E}(\mathbf{r})dV, \quad \mathbf{E}(\mathbf{r}) = k \int_V \rho(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{\|\mathbf{r} - \mathbf{r}'\|^3} dV. \quad ^1$$

**seriesElectric field superposition principle.** Let  $\rho(\mathbf{r}) = dq/dV$ ,  $\sigma(\mathbf{r}) = dq/dA$  and  $\lambda(\mathbf{r}) = dq/d\ell$  be the volume, surface and linear charge densities of an object, respectively. Then the resulting electric field at a point  $\mathbf{r}$  is

$$\mathbf{E}(\mathbf{r}) = k \int_V \rho(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{\|\mathbf{r} - \mathbf{r}'\|^3} dV'.$$

$$\mathbf{E}(\mathbf{r}) = k \int_A \sigma(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{\|\mathbf{r} - \mathbf{r}'\|^3} dA'.$$

$$\mathbf{E}(\mathbf{r}) = k \int_L \lambda(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{\|\mathbf{r} - \mathbf{r}'\|^3} d\ell.$$

**seriesGauß's theorem.**

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} \iff \oint_S \mathbf{E} \cdot \mathbf{n} dS = \frac{1}{\epsilon_0} \int_V \rho(\mathbf{r}) dV = \frac{Q_T}{\epsilon_0}, \quad (1)$$

where  $Q_T$  is the total charge enclosed within  $V$ .

**seriesWork.** The work required to move a point charge  $q$  from  $a$  to  $b$  is

$$W_{b \rightarrow a} = -q \int_a^b \mathbf{E} \cdot d\ell,$$

where the negative sign indicates that the work is done against the field.

**seriesElectric potential.** The electric potential  $\phi$  in a point  $\mathbf{r}$  is defined as:

$$\mathbf{E} = -\nabla\phi, \quad \phi(\mathbf{r}) = k \int \frac{\rho(\mathbf{r}')}{\|\mathbf{r} - \mathbf{r}'\|} dV. \quad (2)$$

<sup>1</sup>Analogously we can define  $\sigma(\mathbf{r}) = \frac{dq}{dS}$  to be the surface charge density and  $\lambda(\mathbf{r}) = \frac{dq}{d\ell}$  to be the linear charge density, and the integrals become as expected.

And then,

$$\phi_a - \phi_b = - \int_b^a \mathbf{E} \cdot d\ell = \int_a^b \mathbf{E} \cdot d\ell.$$

Alternatively we can define the potential from the electric energy. And also if we consider the point  $P$  as a reference point we can define the *electric energy*  $U_a$  as follows

$$U_a - U_b = -q \int_P^a \mathbf{E} \cdot d\ell + q \int_P^b \mathbf{E} \cdot d\ell = -q \int_b^a \mathbf{E} \cdot d\ell = W_{b \rightarrow a}.$$

And finally we get

$$\Delta U(\mathbf{r}) = q\Delta\phi(\mathbf{r}).$$

**seriesPoisson and Laplace equations.** Taking into account formulas 1, 2, we get Poisson's equation

$$\nabla^2\phi = -\frac{\rho}{\epsilon_0}.$$

If  $\rho = 0$ , we obtain Laplace's equation:

$$\nabla^2\phi = 0.$$

**Definition 1.1.** A conductor is a material in which charges can move freely.

**seriesFaraday's cage.** Inside a cavity with no charge of a conductor we have  $\mathbf{E} = 0$  no matter how many charges and the potential are outside. This cavity is known as *Faraday's cage*.

**System of  $N$  conductors**

**FALTA COSA.**

**Definition 1.2 (Capacitor).** A capacitor is formed by two conductors of charges  $\pm q$  and a potential difference  $\Delta\phi$  not depending on the charge of other conductors. As a result, we have the following equality:

$$\Delta\phi = \frac{q}{C},$$

where  $C$  is the *capacitance* of the capacitor and its unit is the *farad* ( $[C] = F$ ).

**seriesCapacitors in series and parallel.** The total capacitance of  $n$  capacitors in series is

$$\frac{1}{C_{\text{total}}} = \sum_{i=1}^n \frac{1}{C_i}.$$

The total capacitance of  $n$  capacitors in parallel is

$$C_{\text{total}} = \sum_{i=1}^n C_i.$$

## Potential energy of a charge distribution

**Discrete charge distribution.** Consider a distribution of  $n$  charges  $q_i$ . If  $\phi_{ij}$  is the potential caused by the charge  $j$  on the point where the charge  $i$  is located, we have that the energy of the distribution  $W$  is

$$W = \sum_{i>j} q_i \phi_{ij} = \frac{1}{2} \sum_{i=1}^n q_i \phi_i,$$

where  $\phi_i = \sum_{i \neq j} \phi_{ij}$ .

**Continuous charge distribution.** Consider a continuous charge distribution of density  $\rho$ . Then

$$W = \frac{1}{2} \int_{\mathcal{V}} \rho \phi d\mathcal{V} = \frac{\varepsilon_0}{2} \int_{\mathbb{R}^3} E^2 d\mathcal{V}.$$

**Definition 1.3.** The radius in which the electrostatic energy equals the rest energy of an electron is called *classical electron radius* and it's equal to:

$$r_0 = \frac{1}{4\pi\varepsilon_0} \frac{e^2}{m_e c^2}.$$